

Exercise 1 *The classical Fluctuation-Dissipation Theorem*

In the derivation of the Fluctuation Dissipation Relation we used

$$\langle B(t)A(0) \rangle = \langle A(0) B(t + i\hbar\beta) \rangle.$$

- (a) Proof that the relation holds.
- (b) What is the order of magnitude of $\hbar\beta$? How does it compare to time scales that can be realized in experiments in classical macroscopic systems? Can you find an experimental setting where $\beta\hbar$ will be of order one?
- (c) Assume that $B(t)$ is analytical in t such that it can be expanded in a Taylor expansion,

$$B(t + i\hbar\beta) = \sum_{j=0}^{\infty} \frac{(i\hbar\beta)^j}{j!} B^{(j)}(t).$$

Use this expansion to evaluate the integral in

$$T_{BA}(\omega) - T_{AB}(-\omega) = \int_{\mathbb{R}} dt e^{-i\omega t} \langle [B(t), A(0)] \rangle.$$

Compare your result to the small $\beta\hbar\omega$ limit of the expression given in the lecture.

Exercise 2

— will be inserted —

Exercise 3 *Generalized Fokker-Planck Equation for the Poisson process*

The Poisson process models the statistics of counts where discrete signals arrive at random times with a uniform rate ν . An example is the number of registered events encountered in an experiment at a synchrotron beamline (e.g., the number of Higgs Bosons identified at CERN). The stochastic variable is the number of counts $N(t)$ after beam time t . Now we explore the distribution $P(N, t)$ of the variable $N(t)$.

- (a) Why is it justified to model this problem as a Markov process? What is the set $\mathcal{S}$ of states? Which transmission are admissible and what are the rates t_k^j , $j, k \in \mathcal{S}$?

- (b) Go back to the analysis of $\frac{d}{dt}\langle M(a, t) \rangle$ in the lecture. Since $\mathcal{S}$ is discrete the integrals turn in to sums. Verify that in this case the generalized diffusivities $D_n(j, t)$ take the form

$$D_n(j, t) = \frac{1}{n!} \sum_{k \in \mathcal{S}} (k - j)^n t_k^j.$$

Determine $D_n(j, t)$ for the Poisson process! Write down the generalized Fokker-Planck equation.

Hint: If everything went fine the generalized Fokker-Planck equation should be a Taylor expansion of the Master equation

$$\partial_t P(N, t) = \nu [P(N - 1, t) - P(N, t)].$$

- (c) Demonstrate that the Master equation is solved by the Poisson distribution

$$P(N, t) = \frac{(\mu(t))^N}{N!} e^{-\mu(t)},$$

with an appropriately chosen function $\mu(t)$. Determine $\mu(t)$.

- (d) Determine the expectation $\langle N(t) \rangle$ based on the equation

$$\frac{d}{dt} \langle N \rangle_t = \langle D_1 \rangle_t,$$

and compare the result to the expectation of the Poisson distribution.

- (e) Determine the second moment $\langle N^2 \rangle_t$ based on the equation

$$\frac{d}{dt} \langle N^2 \rangle_t = 2 \langle D_1 N \rangle_t + 2 \langle D_2 \rangle_t,$$

and compare the resulting expression of the variance to the one of the Poisson distribution.

Exercise 4 *Burst Theorem*

The Burst theorem states that if any D_n for $n > 2$ exists then there is an infinite set of D_n . In the lecture I indicated that this can be seen based on the inequality¹

$$D_{m+n}^2 \leq D_{2m} D_{2n}, \quad \text{for } m, n > 0. \quad (4a)$$

- Why do we need the condition $m, n > 0$?
- Show that when $D_{2m} = 0$ then $D_k = 0$ for all $k > m$.
- Show that when $D_m \neq 0$ then $D_k \neq 0$ for all positive k in $\{2 + 2^j(m - 2), j \in \mathbb{N}_0\}$.
- Show that therefore either $D_m = 0$ for all $m > 2$ or they all differ from zero.

¹A proof of the inequality is provided in the appendix.

Appendix. Proof of the inequality Eq. (4a)

Let

$$a = \frac{1}{n!} \lim_{\Delta t \rightarrow 0} \frac{(a(t + \Delta t) - a(t))^n}{\Delta t}$$
$$b = \frac{c}{n!} \lim_{\Delta t \rightarrow 0} \frac{(a(t + \Delta t) - a(t))^m}{\Delta t}$$

with some fixed number $c \in \mathbb{R}$. Then we have

$$D_{m+n} = c \langle a b \rangle, \quad D_{2n} = \langle a^2 \rangle, \quad D_{2m} = c^2 \langle b^2 \rangle$$

such that Eq. (4a) is equivalent to

$$\langle a b \rangle^2 \leq \langle a^2 \rangle \langle b^2 \rangle$$

for some value of c that will be chosen appropriately.

To derive the latter equation we observe

$$0 \leq \langle (a \pm b)^2 \rangle \Leftrightarrow 2|\langle a b \rangle| \leq \langle a^2 \rangle + \langle b^2 \rangle$$
$$\Leftrightarrow 4 \langle a b \rangle^2 \leq 4 \langle a^2 \rangle \langle b^2 \rangle + (\langle a^2 \rangle - \langle b^2 \rangle)^2.$$

However, the second term on the right hand side vanishes when we choose $c = \sqrt{\langle a^2 \rangle / \langle b^2 \rangle}$.

□