## Stochastic Processes

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## Tutorial 5

Brownian Motion
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## Exercise 1 Brownian motion as a Markov process

For reference in this exercise the cumulants for Brownian motion are given in the Appendix to this exercise sheet. In another exercise we will show that Brownian motion is a Gaussian process, i.e., the probability distribution for Brownian motion is completely specified by the first two cumulants.
(a) Show that the probabiltiy distribution $P\left(v, t \mid v_{0}, 0\right)$ to find a velocity $v$ at time $t$ when it was $v_{0}$ at time 0 takes the form

$$
P\left(v, t \mid v_{0}, 0\right)=N(t) \exp \left[-\frac{\left(v-v_{0} \mathrm{e}^{-\lambda t}\right)^{2}}{\frac{2 d}{\lambda}\left(1-\mathrm{e}^{-2 \lambda t}\right)}\right]
$$

where $N(t)$ is an appropriate normalization of this conditional probability.
What is the appropriate normalization $N(t)$ ?
Hint: This expression assumes $\mathcal{C}_{0}(v, v)=0$ and $\left\langle v_{0}\right\rangle=v_{0}$. Why is this justified?
(b) Under which condition on $\lambda$ will Brownian motion become a Markov process for the velocities? What is special about the resulting conditional probability?
(c) Adopt the limit $\lambda \rightarrow \infty$ at a fixed diffusion coefficient $D=d / \lambda^{2}$. Show that in this limit the probabiltiy distribution $P\left(x, t \mid x_{0}, 0\right)$ to find the Brownian particle at position $x$ at time $t$ when it was at $x_{0}$ at time 0 takes the form

$$
\begin{equation*}
P\left(x, t \mid x_{0}, 0\right)=(4 \pi D t)^{-1 / 2} \exp \left[-\frac{\left(x-x_{0}\right)^{2}}{4 D t}\right] \tag{1}
\end{equation*}
$$

(d) Brownian motion is Markovian iff Eq. (1) satisfies the Chapman-Kolmogorov-criterion that for any set of times $t_{1}<t_{2}<t_{3}$ and positions $x_{1}, x_{2}, x_{3}$ one must have

$$
P\left(x_{3}, t_{3} \mid x_{1}, t_{1}\right)=\int \mathrm{d} x_{2} P\left(x_{3}, t_{3} \mid x_{2}, t_{2}\right) P\left(x_{2}, t_{2} \mid x_{1}, t_{1}\right)
$$

Verify that it holds!

## Exercise 2 Variance of positions

In this exercise we derive the expression for $\mathcal{C}_{t}(x, x)$ that is given in the appendix. We start from

$$
x(t)=x_{0}+\frac{v_{0}}{\lambda}\left(1-\mathrm{e}^{-\lambda t}\right)+\int_{0}^{t} \mathrm{~d} s_{1} \mathrm{e}^{-\lambda s_{1}} \int_{0}^{s_{1}} \mathrm{~d} s_{2} A\left(s_{2}\right) \mathrm{e}^{\lambda s_{2}}
$$

(a) Introduce the function $W\left(s_{1}\right)=\int_{0}^{s_{1}} \mathrm{~d} s_{2} A\left(s_{2}\right) \mathrm{e}^{\lambda s_{2}}$ and use integration by parts to show that

$$
\begin{equation*}
x(t)=x_{0}+\frac{v_{0}}{\lambda}\left(1-\mathrm{e}^{-\lambda t}\right)+\lambda^{-1} \int_{0}^{t} \mathrm{~d} s\left(1-\mathrm{e}^{(s-t)}\right) A(s) . \tag{2a}
\end{equation*}
$$

(b) Use Eq. (2a) to evaluate

$$
\mathcal{C}_{t}(x, x)=\left\langle(x(t)-\langle x(t)\rangle)^{2}\right\rangle .
$$

(c) Show that in the limit of long times the resulting expression reduces to

$$
\begin{equation*}
\mathcal{C}_{t}(x, x) \simeq\left(\mathcal{C}_{0}(x, x)-\frac{D}{\lambda}\right)+\lambda^{-2}\left(\mathcal{C}_{0}(v, v)-\frac{d}{\lambda}\right)+2 D t \tag{2b}
\end{equation*}
$$

where we introduced the diffusion coefficient $D=d / \lambda^{2}$.
Provide an interpretation for the three contributions to this expression.
(d) The result Eq. (2b) suggests that $\mathcal{C}_{t}(x, x)$ may take negetive values when $\mathcal{C}_{0}(x, x)=$ $\mathcal{C}_{0}(v, v)=0$ and small $t$.
What is wrong about this argument?
Find the expression that should rather be considered to discuss this special case of $\mathcal{C}_{t}(x, x)$.

## Exercise 3 Covariance of position and velocity

(a) Determine the covariance

$$
\mathcal{C}_{t}(x, v)=\langle(x(t)-\langle x(t)\rangle)(v(t)-\langle v(t)\rangle)\rangle .
$$

(b) Compare the result to the variance $\mathcal{C}_{t}(v, v)$. What do you observe?
(c) For an equilibrated velocity ensemble, where $\mathcal{C}_{0}(v, v)=0$, the asymptotics for large and small times becomes

$$
\mathcal{C}_{t}(x, v) \simeq\left\{\begin{array}{lll}
d t / \lambda & \text { for } & \lambda t \ll 1 \\
d / \lambda^{2} & \text { for } & \lambda t \gg 1
\end{array}\right.
$$

What does this mean physically?
Which interpretation does this suggest for the diffusion coefficient $D=d / \lambda^{2}$ ?

## Exercise 4 Estimating the diffusive displacement

In the lecture I indicated that Nageli $^{1}$ dismissed the role of molecular collisions as origin of Brownian motion. In this exercise we revisit his argument that is based on his estimate of the speed, $U_{B} \simeq 1 \mu \mathrm{~m} / \mathrm{s}$, of a Brownian particles with a diameter of about $R_{B} \simeq 2 \mu \mathrm{~m}$.
(a) According to Stoke's law the friction force on a solid spherical particle is

$$
F_{S}=6 \pi R_{B} \rho_{s} \nu_{s} U_{B}
$$

where $\rho_{s}$ and $\nu_{s}$ are the density and the kinematic viscosity of the surrounding fluid, respectively. For water they take the values $\rho_{s} \simeq 10^{3} \mathrm{~kg} / \mathrm{m}^{3}$ and $\nu_{s} \simeq 10^{-6} \mathrm{~m}^{2} / \mathrm{s}$. Show that for these parameters the damping takes the value $\lambda \simeq 10^{6} / \mathrm{s}$.
Bonus: Note that smaller particles have a larger damping. Which radius will result in the damping $\lambda \simeq 10^{7} / \mathrm{s}$ that was quoted in the lecture?
(b) When the particle is at thermal equilibrium it should have a velocity $U_{E}$

$$
\frac{1}{2} \frac{4 \pi \rho_{B} R_{B}^{3}}{3} U_{E}^{2}=\frac{3}{2} k_{B} T
$$

Estimate $U_{E}$ for a particle that has roughly the same density as water.
(c) Assume that water molecules have an effective radius of about $R_{w} \simeq 4 \times 10^{-10} \mathrm{~m}$. What would the momentum balance

$$
M_{B} U_{B} \simeq M_{w} U_{w}
$$

imply about typical verlocity $U_{B}$ for our Brownian particle when it collides with water molecules in thermal equilibrium?
(d) Show that the diffusion coefficient takes a value of the order to $D \simeq 10^{-13} \mathrm{~m}^{2} / \mathrm{s}$, and calculate the diffusive displacement $\Delta X(t)=2 D t$ for time intervals $t=0.1,1.0,10,100 \mathrm{~s}$.
(e) Compare now Nageli's estimate of $U=10^{-6} \mathrm{~m} / \mathrm{s}$ to the velocity $U_{t}=\Delta X(t) / t$.

What does this imply about the time and space resolution of Nageli's observation?
Observe also that $U_{E}>U_{t}>U_{P}$. Why would one expect this relation?

[^0]
## Appendix: Derivation of the Cumulants

## Velocity

$$
v(t)=v_{0} \mathrm{e}^{-\lambda t}+\int_{0}^{t} \mathrm{~d} s A(s) \mathrm{e}^{\lambda(s-t)}
$$

with velocity $v(t)$ at time $t$
$v_{0}$ at initial time 0
relaxation rate $\lambda$
random forces $A(t)$

$$
\begin{array}{ll}
\text { where } & \langle A(t)\rangle=0 \\
\text { where } & \left\langle A\left(t_{1}\right) A\left(t_{2}\right)\right\rangle=2 d \delta\left(t_{1}-t_{2}\right)
\end{array}
$$

## Expectation

$$
\mathcal{C}_{t}(v)=\langle v(t)\rangle=\left\langle v_{0}\right\rangle \mathrm{e}^{-\lambda t}+\int_{0}^{t} \mathrm{~d} s\langle A(s)\rangle \mathrm{e}^{\lambda(s-t)}=\left\langle v_{0}\right\rangle \mathrm{e}^{-\lambda t}
$$

Starting from its initial value $\left\langle v_{0}\right\rangle$ the expectation decays exponentially to zero.

## Variance

$$
\begin{aligned}
\mathcal{C}_{t}(v, v) & =\left\langle(v(t)-\langle v(t)\rangle)^{2}\right\rangle=\left\langle\left(\left(v_{0}-\left\langle v_{0}\right\rangle\right) \mathrm{e}^{-\lambda t}+\int_{0}^{t} \mathrm{~d} s A(s) \mathrm{e}^{\lambda(s-t)}\right)^{2}\right\rangle \\
& =\left\langle\left(v_{0}-\left\langle v_{0}\right\rangle\right)^{2}\right\rangle \mathrm{e}^{-2 \lambda t}+2 \mathrm{e}^{-\lambda t} \int_{0}^{t} \mathrm{~d} s\left\langle\left(v_{0}-\left\langle v_{0}\right\rangle\right) A(s)\right\rangle \mathrm{e}^{\lambda(s-t)}+\left\langle\left(\int_{0}^{t} \mathrm{~d} s A(s) \mathrm{e}^{\lambda(s-t)}\right)^{2}\right\rangle \\
& =\mathcal{C}_{0}(v, v) \mathrm{e}^{-2 \lambda t}+\int_{0}^{t} \mathrm{~d} s_{1} \mathrm{e}^{\lambda\left(s_{1}-t\right)} \int_{0}^{t} \mathrm{~d} s_{2} \mathrm{e}^{\lambda\left(s_{2}-t\right)}\left\langle A\left(s_{1}\right) A\left(s_{2}\right)\right\rangle \\
& =\mathcal{C}_{0}(v, v) \mathrm{e}^{-2 \lambda t}+2 d \int_{0}^{t} \mathrm{~d} s_{2} \mathrm{e}^{2 \lambda\left(s_{2}-t\right)} \\
& =\left(\mathcal{C}_{0}(v, v)-\frac{d}{\lambda}\right) \mathrm{e}^{-2 \lambda t}+\frac{d}{\lambda}
\end{aligned}
$$

Starting from ints initial value $\mathcal{C}_{0}(v, v)$ the variance decays exponentially to the value $d / \lambda$.

## Position

$$
x(t)=x_{0}+\int_{0}^{t} \mathrm{~d} s v(s)=x_{0}+\frac{v_{0}}{\lambda}\left(1-\mathrm{e}^{-\lambda t}\right)+\int_{0}^{t} \mathrm{~d} s_{1} \int_{0}^{s_{1}} \mathrm{~d} s_{2} A\left(s_{2}\right) \mathrm{e}^{\lambda\left(s_{2}-s_{1}\right)}
$$

## Expectation

$$
\mathcal{C}_{t}(x)=\langle x(t)\rangle=\left\langle x_{0}\right\rangle+\frac{\left\langle v_{0}\right\rangle}{\lambda}\left(1-\mathrm{e}^{-\lambda t}\right)
$$

When the expectation of the velocity in the initial ensemble vanishes, $\left\langle v_{0}\right\rangle=0$, the expectation of the position remains constant at $\left\langle x_{0}\right\rangle$. Otherwise, it decays exponentially to its asymptotic value $\left\langle x_{0}\right\rangle+\left\langle v_{0}\right\rangle / \lambda$.

## Variance

$$
\begin{aligned}
\mathcal{C}_{t}(x, x) & =\left\langle(x(t)-\langle x(t)\rangle)^{2}\right\rangle \\
& =\mathcal{C}_{0}(x, x)+\frac{\left(1-\mathrm{e}^{-\lambda t}\right)^{2}}{\lambda^{2}}\left(\mathcal{C}_{0}(v, v)-\frac{d}{\lambda}\right)+\frac{2 d}{\lambda^{3}}\left(\lambda t-\left(1-\mathrm{e}^{-\lambda t}\right)\right)
\end{aligned}
$$

The derivation and interpretation is given as an exercise.


[^0]:    ${ }^{1}$ K. von Nageli, Sitzungsberichte der Königlich Bayrischen Akademie der Wissenschaften München, Mathematisch-physikalische Klasse 9 (1879) 389-453.

