

Towards a realistic simulation of disorder in unconventional superconductors

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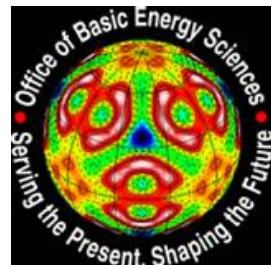
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Wei Ku

Condensed Matter Physics and Materials Science Department, Brookhaven National Laboratory, Upton, New York 11973, USA



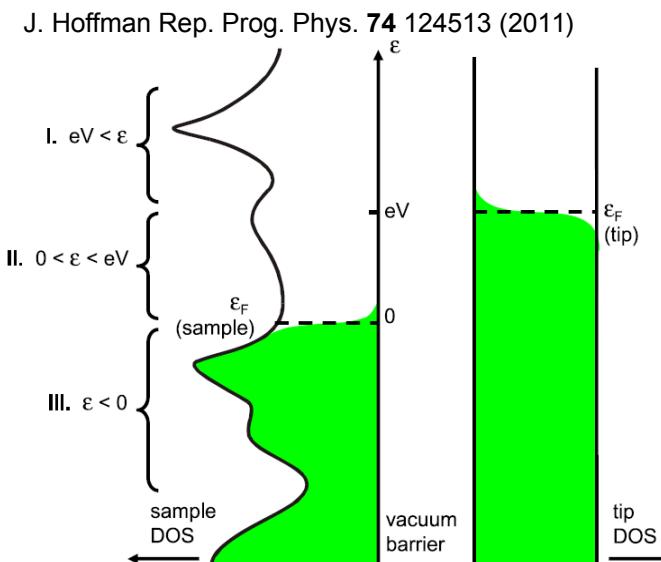
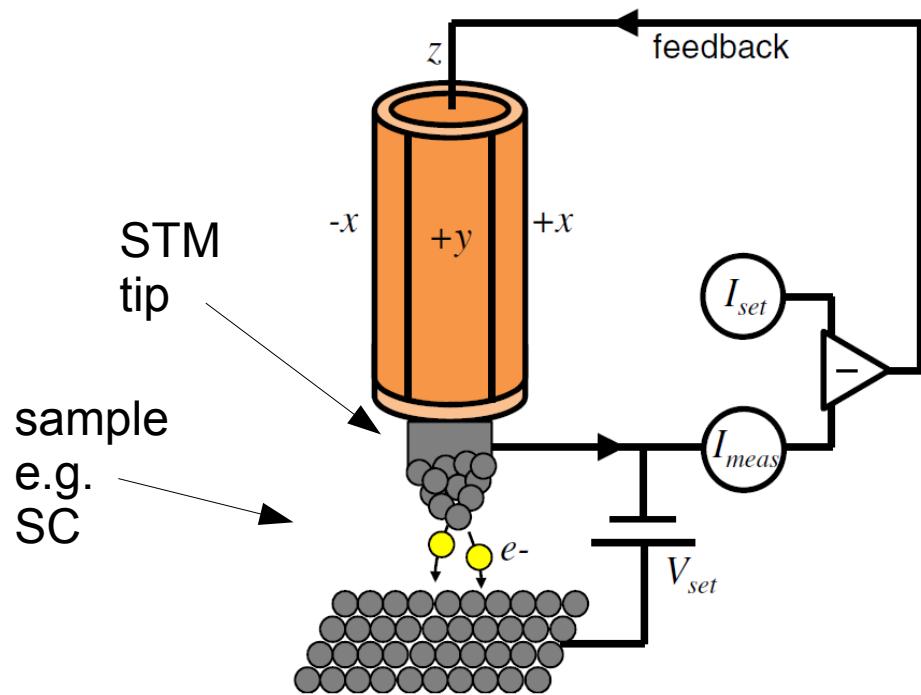
Phys. Rev. B **90**, 134520 (2014)
Phys. Rev. Lett. **114**, 217002 (2015)
arXiv:1607.03192



Outline

- Motivation
 - STM: impurities as probe for electronic structure, order parameter
 - layered superconductors, complications
- Theoretical methods to investigate impurity physics in superconductors
- Using wavefunction information in layered superconductors
- Applications
 - BiSrCaCuO (single band, d-wave)
 - LiFeAs (multiband, s-wave)

Scanning tunnelling microscopy



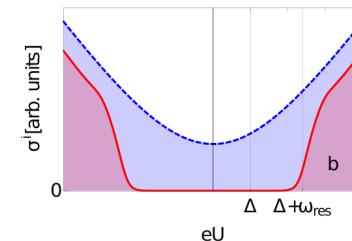
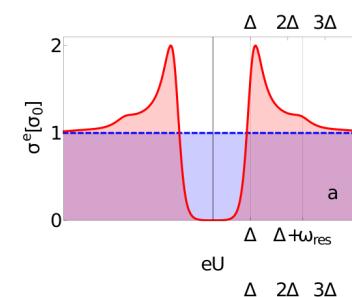
Inelastic tunneling:
coupling to bosonic mode

Tunnelling current:

$$I(V, x, y, z) = -\frac{4\pi e}{\hbar} \rho_t(0) |M|^2 \int_0^{eV} \rho(x, y, z, \epsilon) d\epsilon$$

Local Density Of States (LDOS)
of sample at given energy **at the tip position**

J. Tersoff and D. R. Hamann, PRB **31**, 805 (1985)



superconducting
state

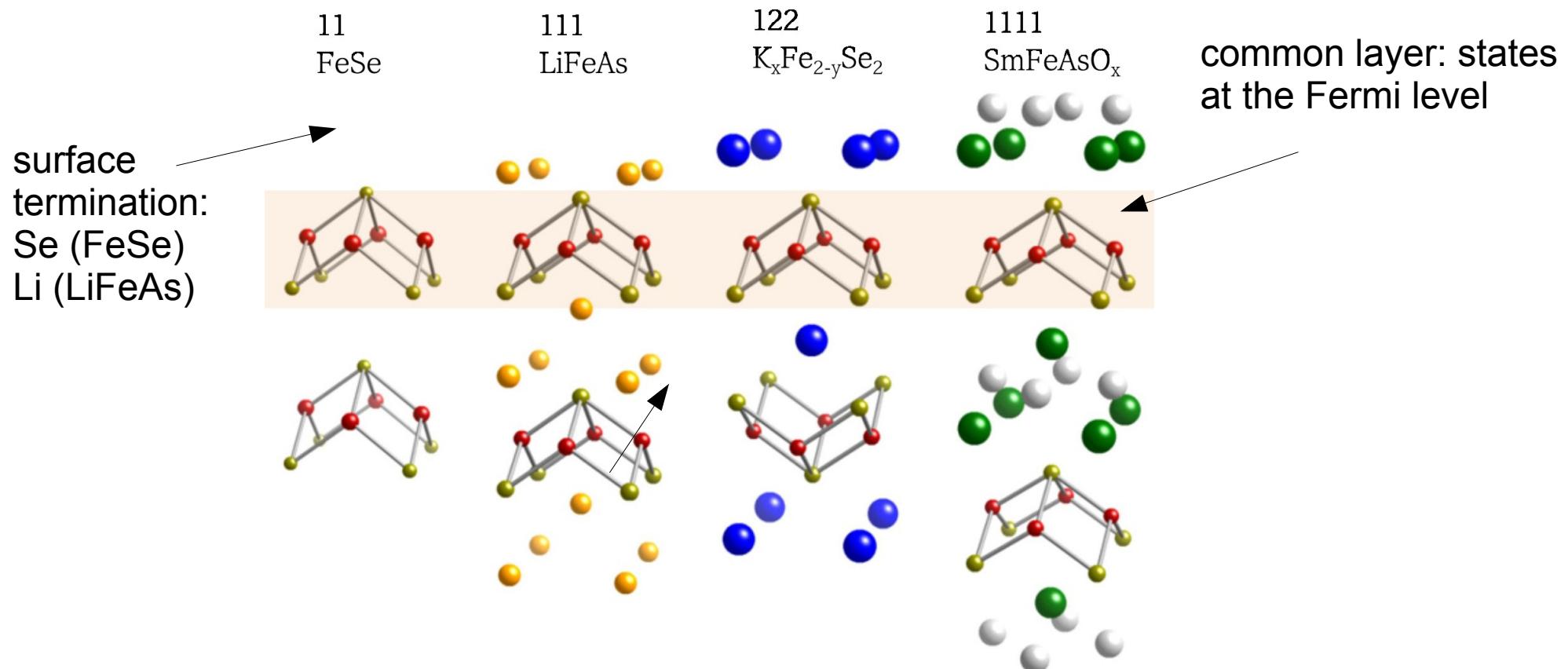
normal state

J. R. Kirtley and D. J. Scalapino, PRL **65**, 798 (1990)
J. R. Kirtley, PRB **47**, 11379 (1993)
Patrik Hlobil, et al., arXiv:1603.05288 (2016)

Layered superconductors

LDOS of sample at given energy **at the tip position**

- Iron based superconductors



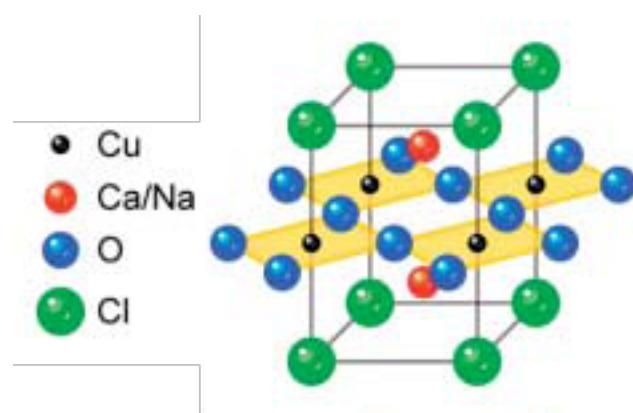
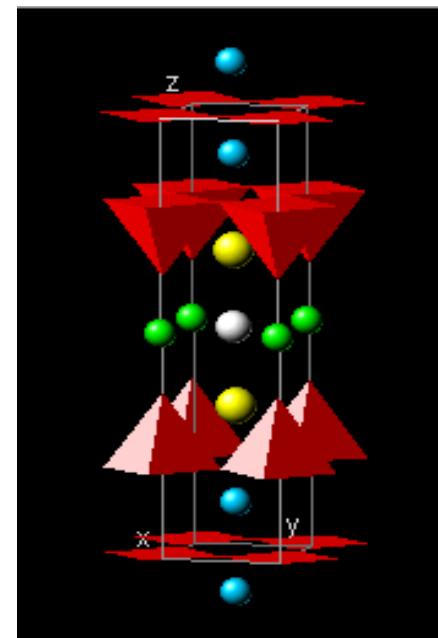
Layered superconductors

- Cuprates

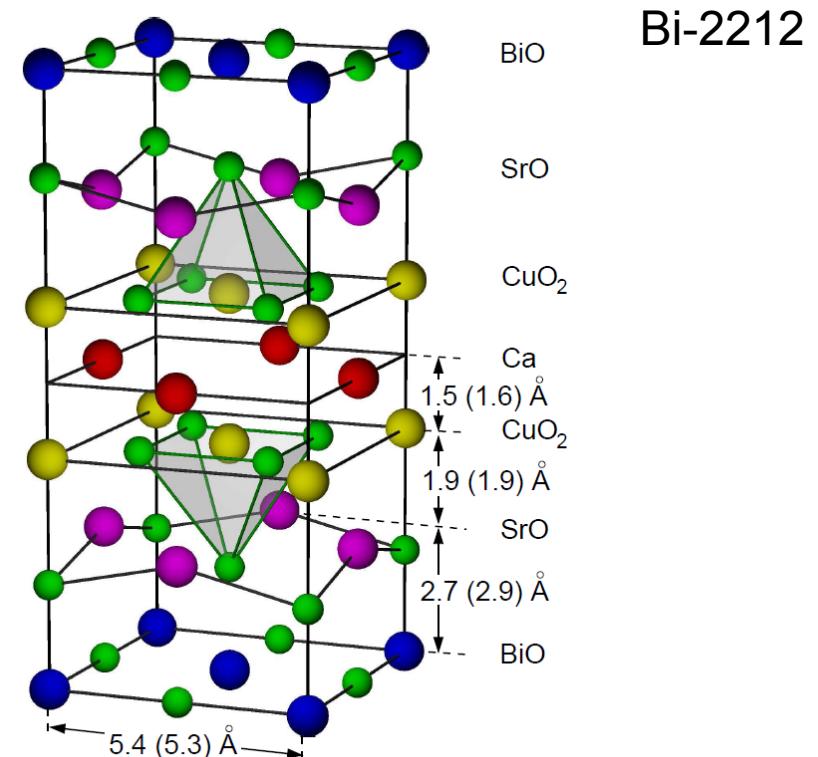
$\text{Hg}_1\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_8$

$T_c = 135 \text{ K}$
under pressure: 153 K

▲ Cu/O



Na-CCOC



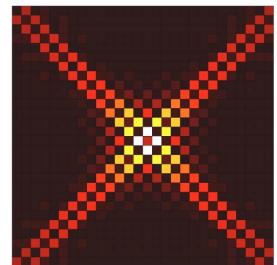
Bi-2212

Bound states of nonmagnetic impurity

- d-wave superconductors
 - local LDOS: 4 fold pattern
 - low energy bound state

$$\Omega \equiv \Omega' + i\Omega'' = \Delta_0 \frac{\pi c/2}{\ln(8/\pi c)} \left[1 + \frac{i\pi}{2} \frac{1}{\ln(8/\pi c)} \right]$$

J. M. Byers, M. E. Flatté, and D. J. Scalapino
Phys. Rev. Lett. **71**, 3363 (1993)

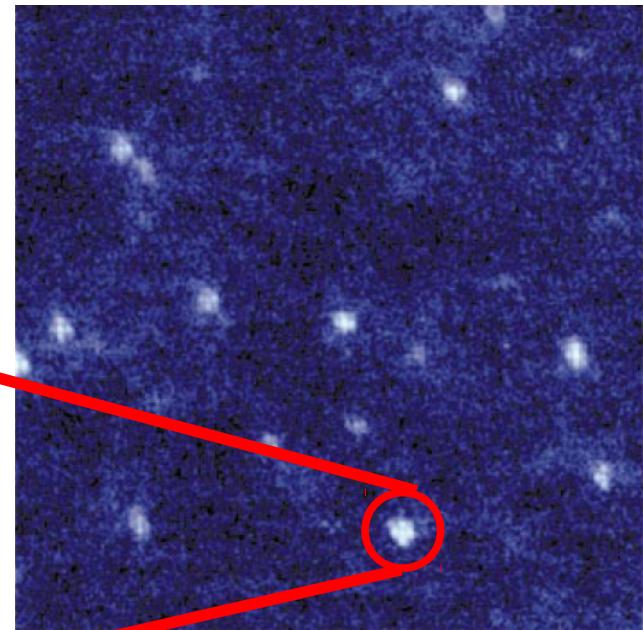
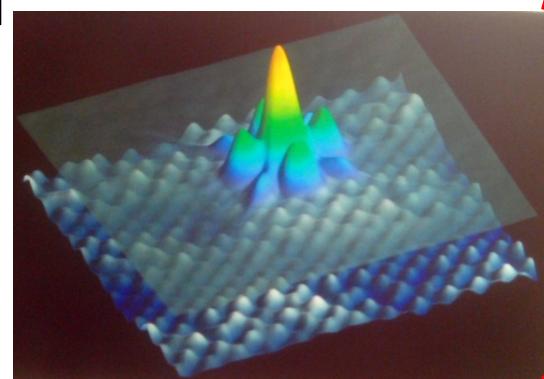


A. V. Balatsky, M. I. Salkola, and A. Rosengren
Phys. Rev. B **51**, 15547 (1995)
Stamp, Journal of Magnetism and Magnetic Materials,
63, 429 - 431 (1987) (p-wave)

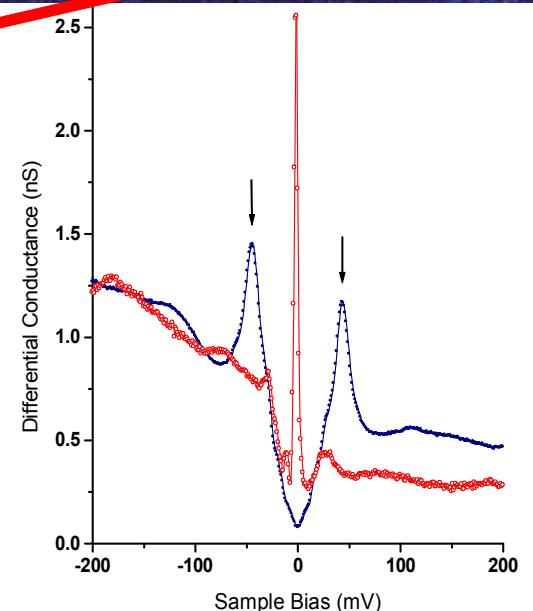
STM experiment on Bi-2212

- LDOS map at -1.5 meV
~20 Zn atoms in field of view
 - 4 fold pattern

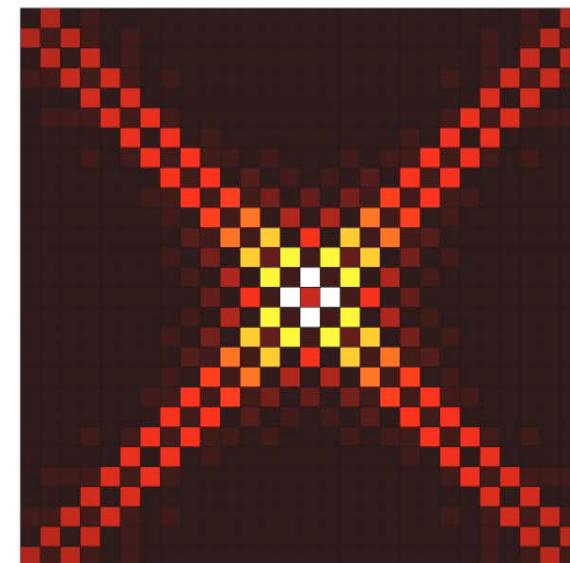
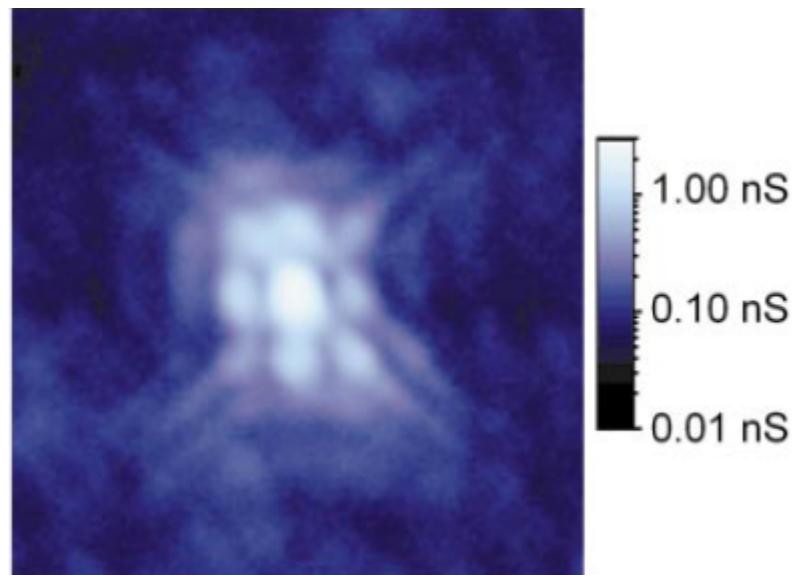
Pan et al., Nature
403, 746 (2000)



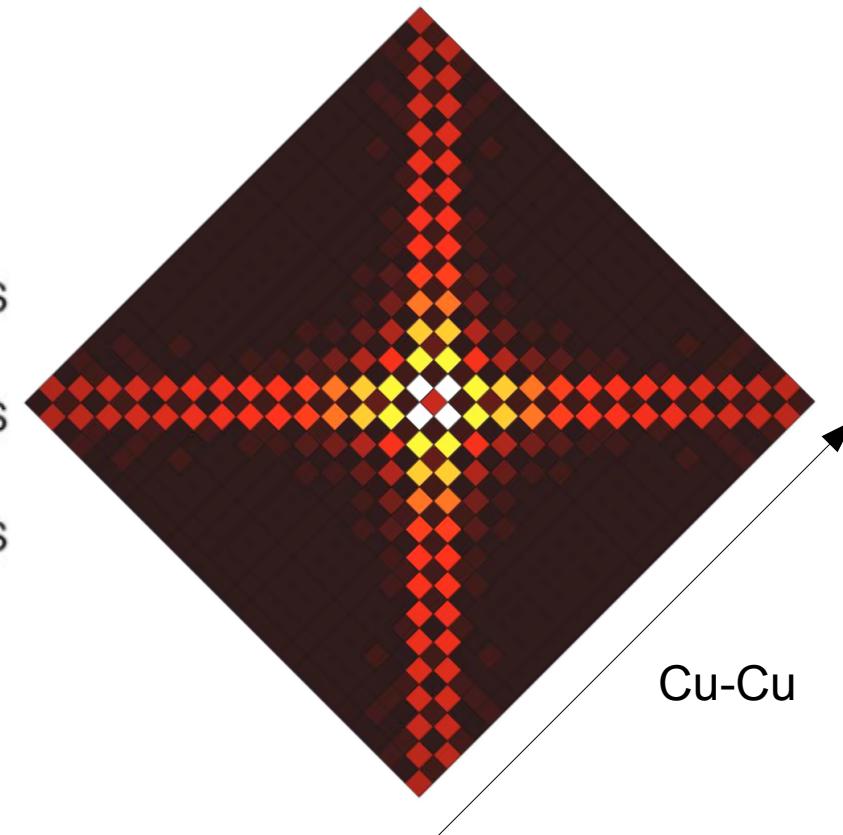
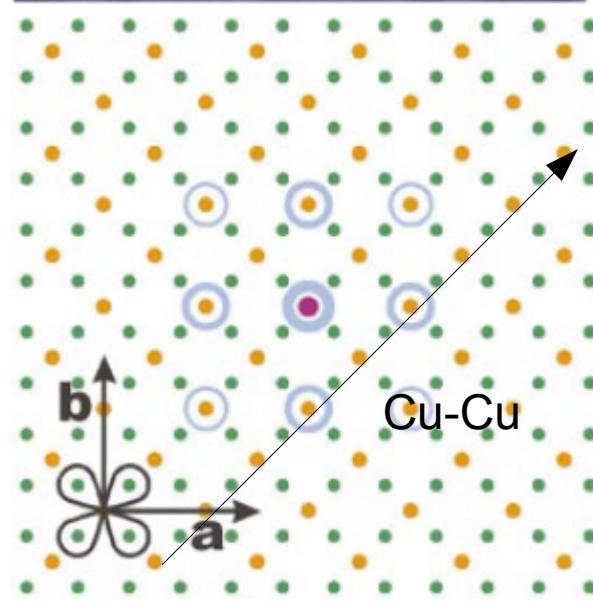
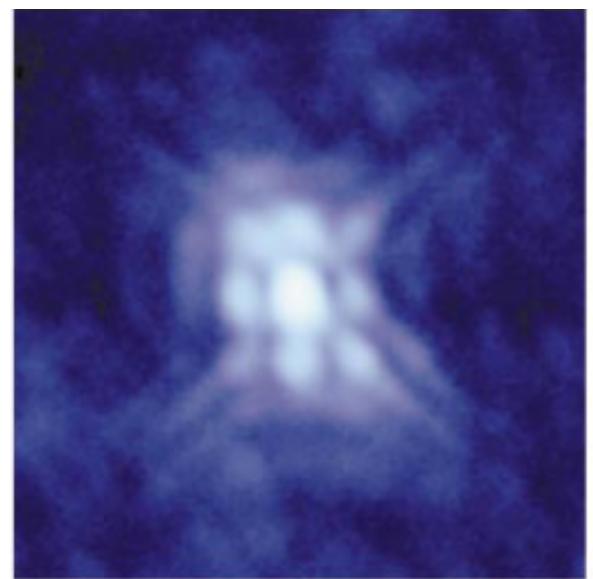
- low energy bound state
in on site spectrum



comparison to theory



comparison to theory

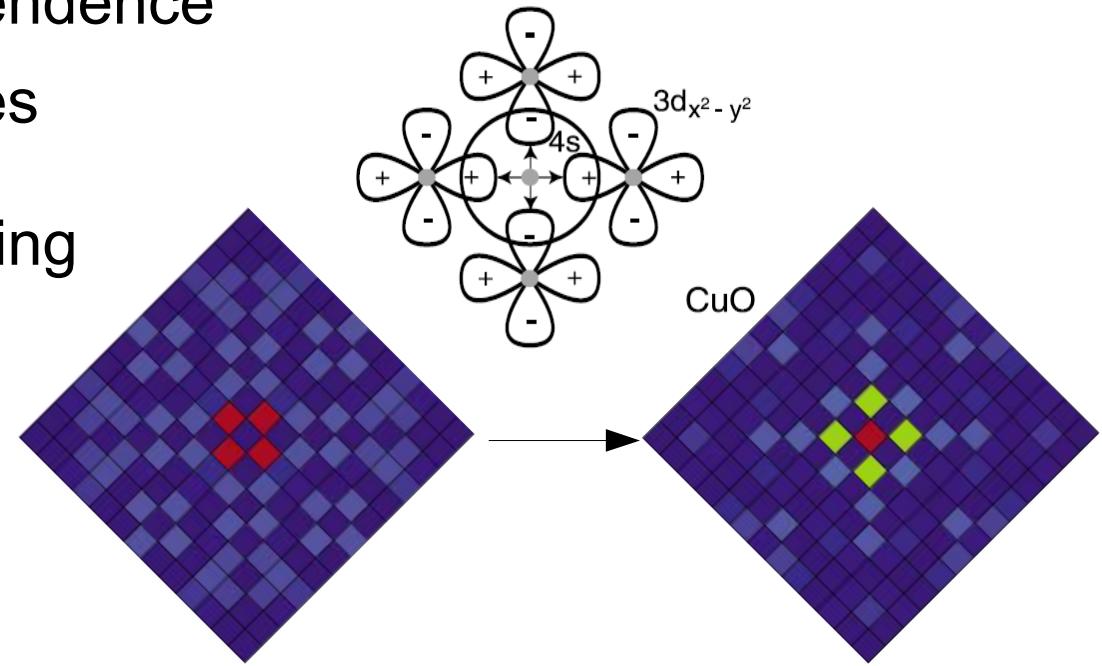
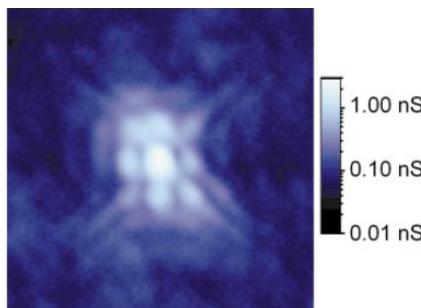


Theories for spatial pattern

- extended impurity potentials
(magnetic Ni impurity)
- Correlations: “Kondo screening”
(magnetic impurity), “modifications of
the theory for the case of a
nonmagnetic impurity” → consistent
with experimental spatial dependence
- Bi-O filter function: STM probes
neighbouring Cu states due to
momentum dependent tunnelling
matrix elements

Jian-Ming Tang and Michael E.
Flatté PRB **66**, 060504(R)
(2002)

Anatoli Polkovnikov PRB **65**, 064503 (2002)



Martin et al., PRL **88**, 097003 (2002)

Theory: State of the art methods

- Hamiltonian

$$H_0 = \sum_{R, R', \sigma} t_{R, R'} c_{R, \sigma}^\dagger c_{R', \sigma} - \mu_0 \sum_{R, \sigma} c_{R, \sigma}^\dagger c_{R, \sigma}$$

$$H = H_0 + H_{\text{BCS}} + H_{\text{imp}}$$

band structure
 kinetic energy superconductivity
 gap function / pairing

$$H_{\text{BCS}} = - \sum_{R, R'} \Delta_{R, R'} c_{R, \uparrow}^\dagger c_{R', \downarrow} + H.c.,$$

$$H_{\text{imp}} = \sum_{\sigma} V_{\text{imp}} c_{R, * \sigma}^\dagger c_{R, * \sigma}$$

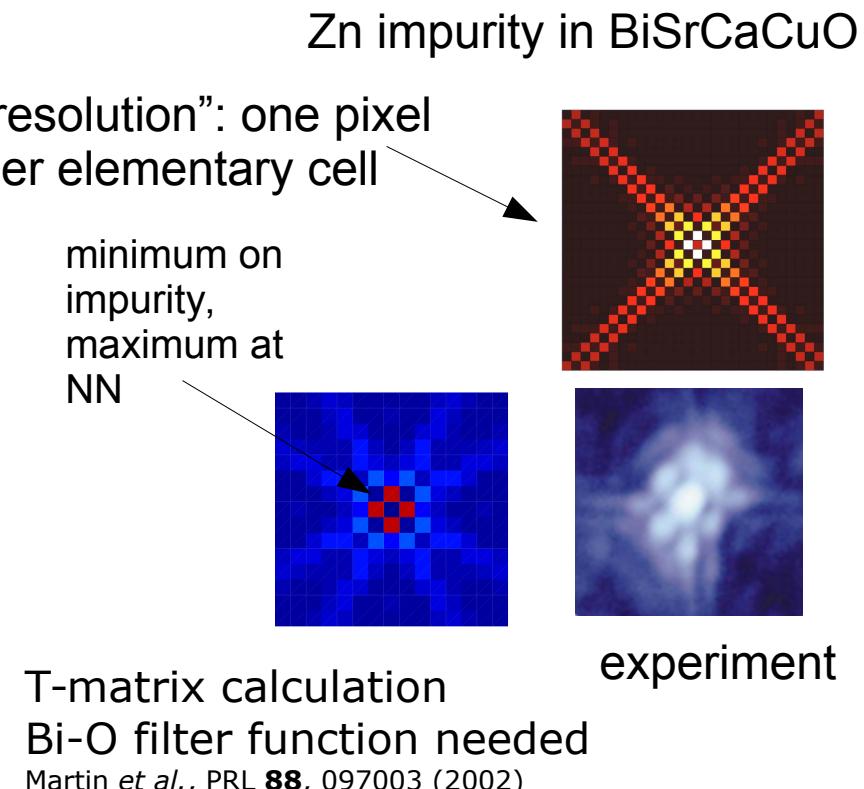
impurity scatterer
 (non)magnetic
 potential / T_c scatterer

- T-matrix calculations
- self-consistent mean field theory (BdG)
- lattice Green function

$$\hat{G}(\mathbf{r}, \mathbf{r}'; \omega) = \hat{G}_0(\mathbf{r} - \mathbf{r}', \omega) + \hat{G}_0(\mathbf{r}, \omega) \hat{T}(\omega) \hat{G}_0(\mathbf{r}', \omega)$$

- Local Density of States (LDOS)

$$N_{\text{imp}}(\mathbf{r}, \omega) = - \frac{1}{\pi} \text{Im} [\hat{G}_0(\mathbf{r}, \omega) \hat{T}(\omega) \hat{G}_0(\mathbf{r}, \omega)]_{11}$$



Theory: State of the art methods

Bogoliubov-de Gennes (BdG)

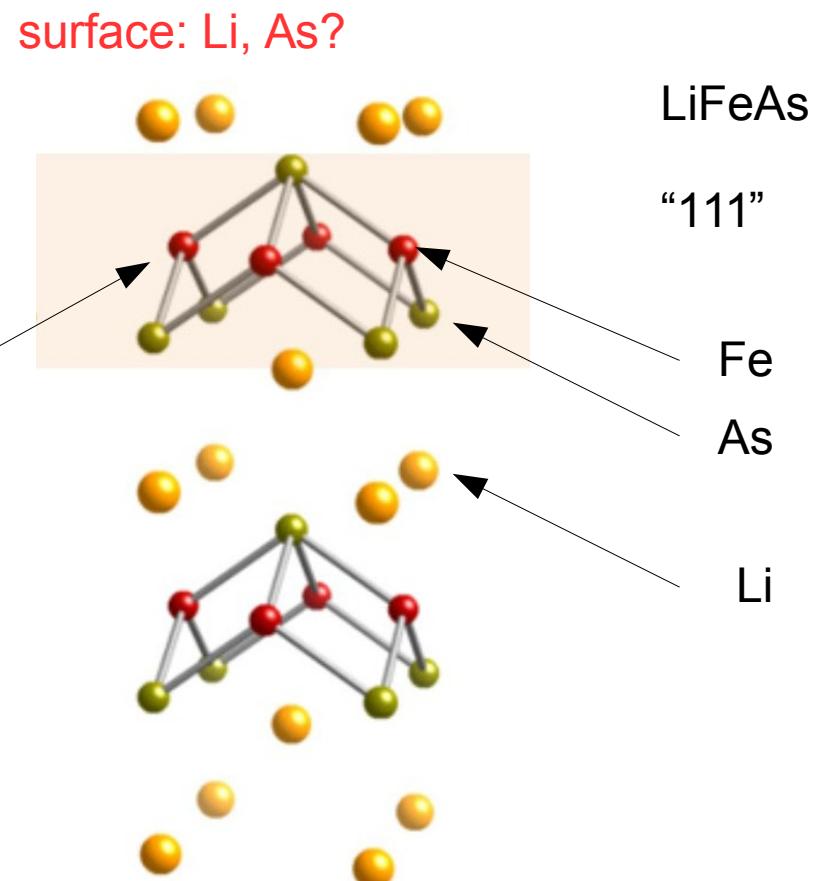
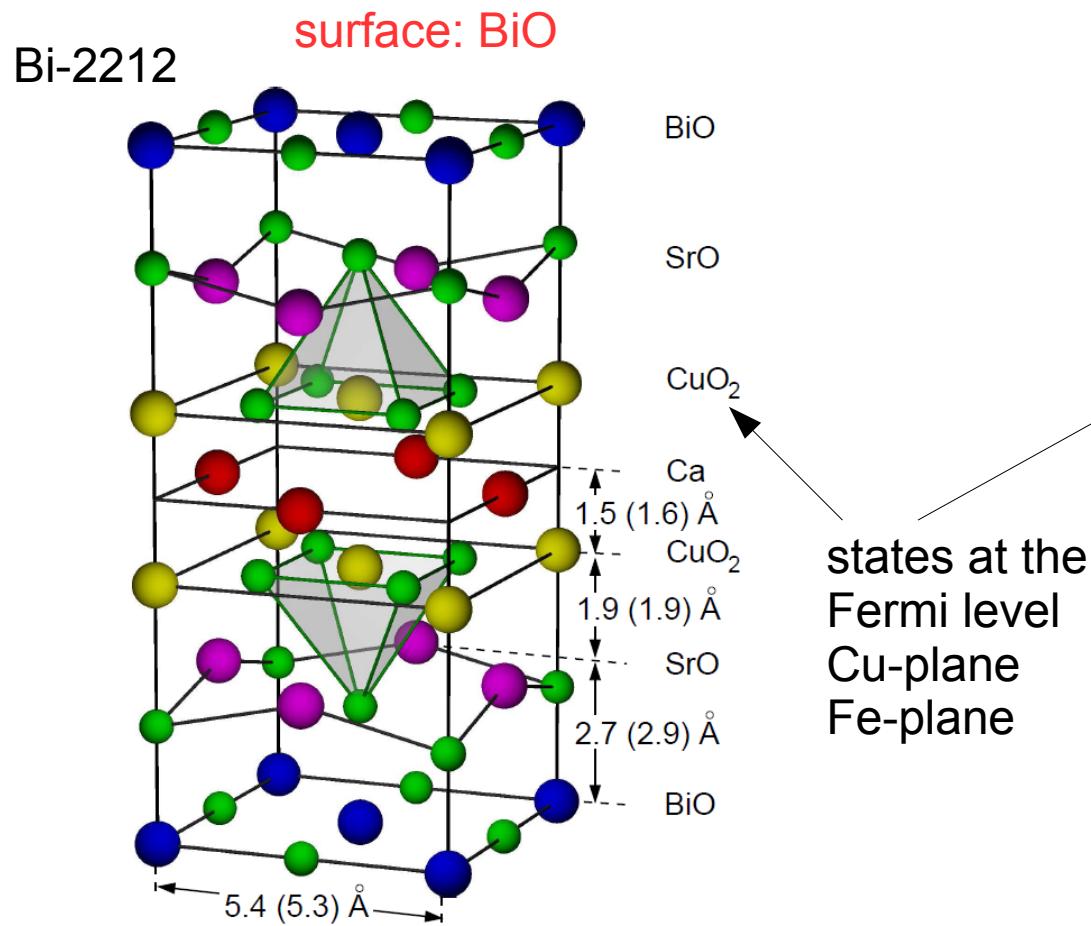
- Hamiltonian $H = H_0 + H_{\text{BCS}} + H_{\text{imp}}$
- self-consistent solution in real space
(NxN grid, determine gaps) $\Delta_{R R'} = \Gamma_{R R'} \langle c_{R'} \downarrow c_R \uparrow \rangle$
- eigenvalues E_n , eigenvectors (u_n, v_n)
- lattice Green function

$$G_\sigma(R, R'; \omega) = \sum_n \left(\frac{u_R^{n\sigma} u_{R'}^{n\sigma*}}{\omega - E_{n\sigma} + i0^+} + \frac{v_R^{n-\sigma} v_{R'}^{n-\sigma*}}{\omega + E_{n-\sigma} + i0^+} \right)$$

Layered superconductors

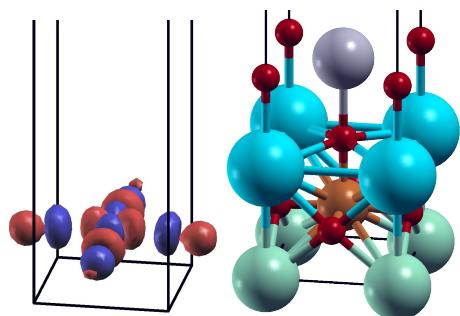
- 2 examples: surface atoms \neq superconducting layer

Cuprates	Iron based superconductors
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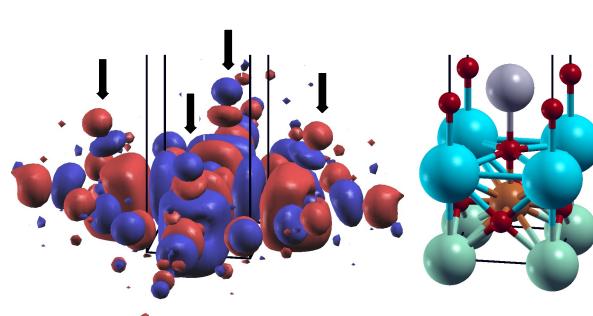


Wannier method

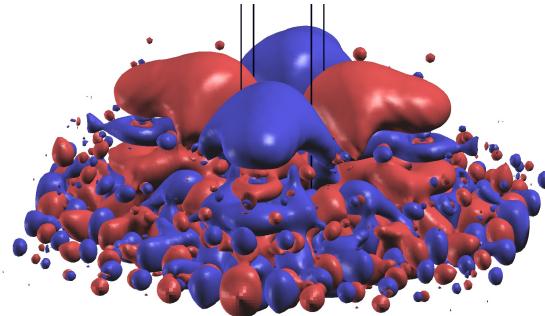
- first principles calculation (surface)
- 1 band tight binding model:
1 Wannier function



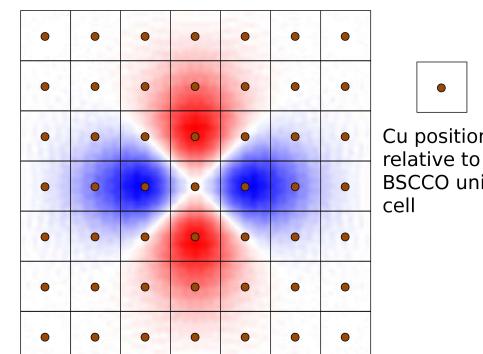
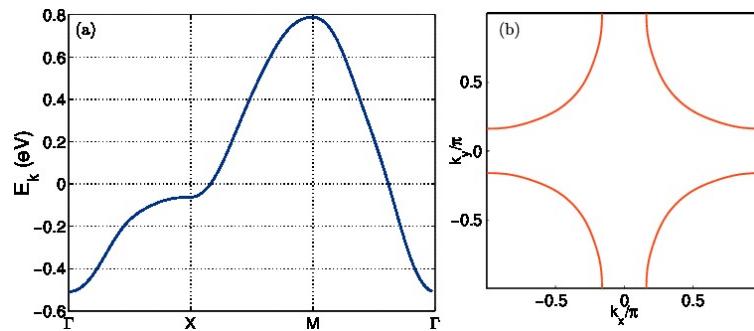
Cu dxy



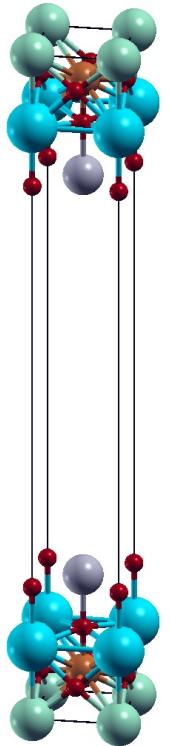
NN apical O tails



at surface: only contributions to NN



Cu position
relative to
BSCCO unit
cell



Superconductivity

- superconducting order parameter (d-wave) (phenomenology or calculation fx. mean-field)

- continuum Green function

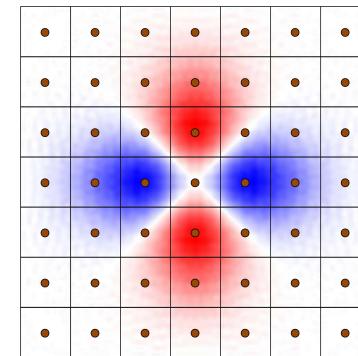
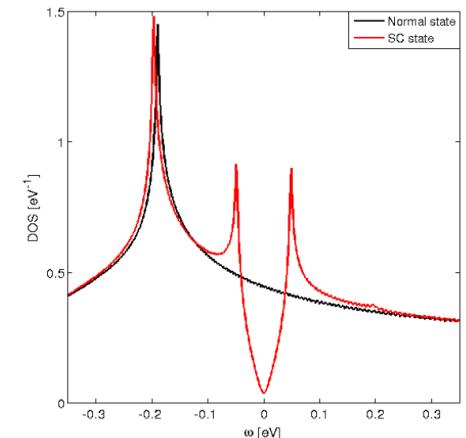
$$\psi_\sigma(\mathbf{r}) = \sum_{\mathbf{R} \mu} c_{\mathbf{R} \mu \sigma} w_{\mathbf{R} \mu}(\mathbf{r})$$

$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_{\mathbf{R}, \mathbf{R}'} G(\mathbf{R}, \mathbf{R}'; \omega) w_{\mathbf{R}}(\mathbf{r}) w_{\mathbf{R}'}^*(\mathbf{r}')$$

lattice Green function

continuum position

nonlocal contributions



local density of states (LDOS)

$$\rho(\mathbf{r}, \omega) \equiv -\frac{1}{\pi} \operatorname{Im} G(\mathbf{r}, \mathbf{r}; \omega)$$

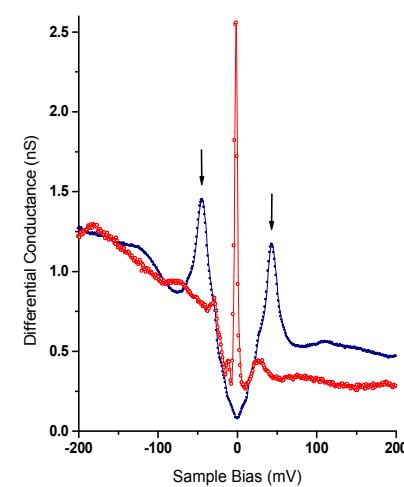
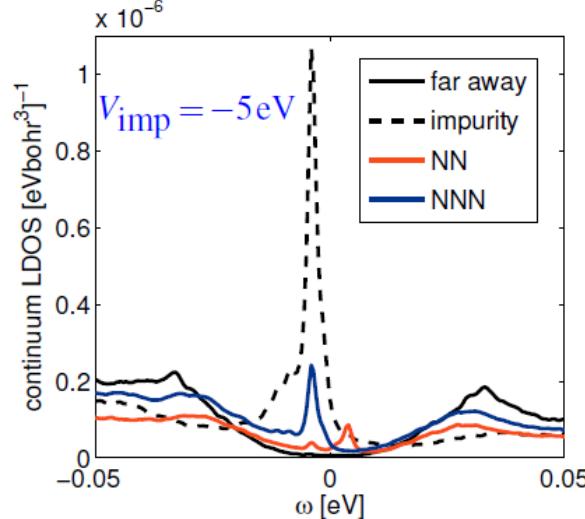
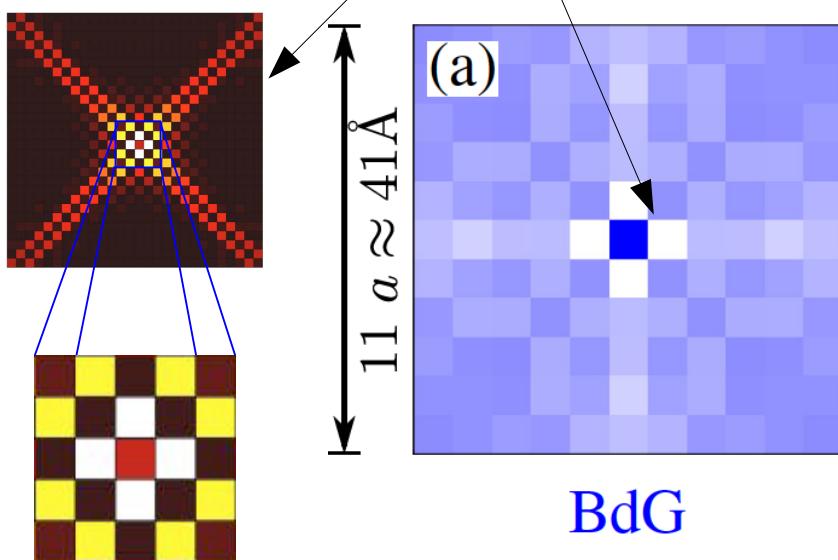
surface Wannier function with phases

BSCCO: Results

STM maps and spectra

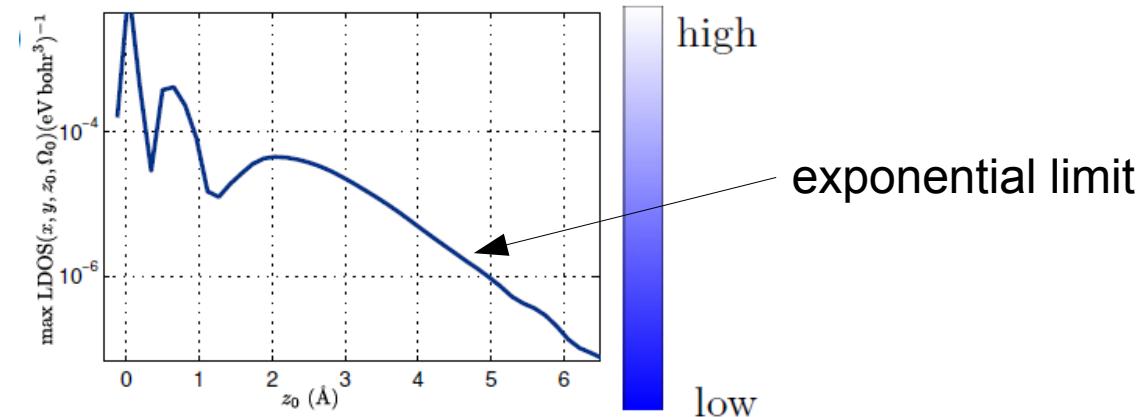
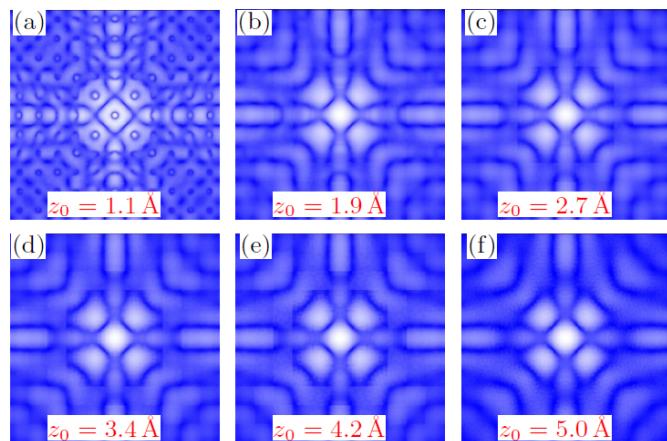
- d-wave order parameter
- Zn impurity:
 $V_{\text{imp}} = -5 \text{ eV}$
 resonance: -3.6 meV

Zhu et al., PRB
67, 094508
 (2003)

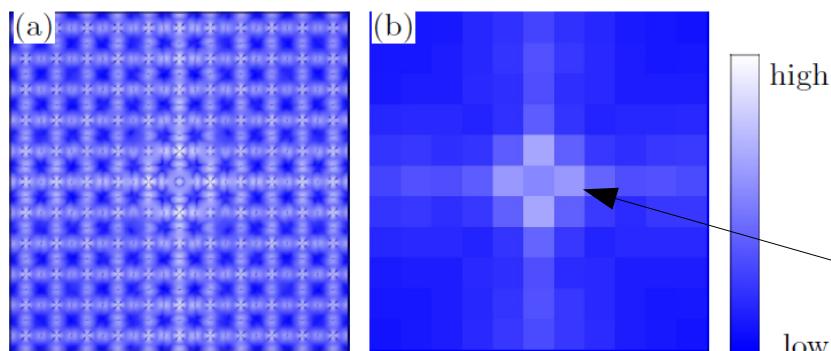


Pan et al., Nature
403, 746 (2000)

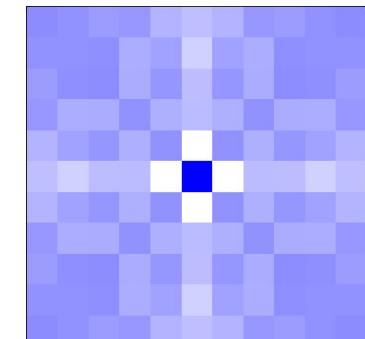
- dependence on tip height



- continuum LDOS in the Cu-plane

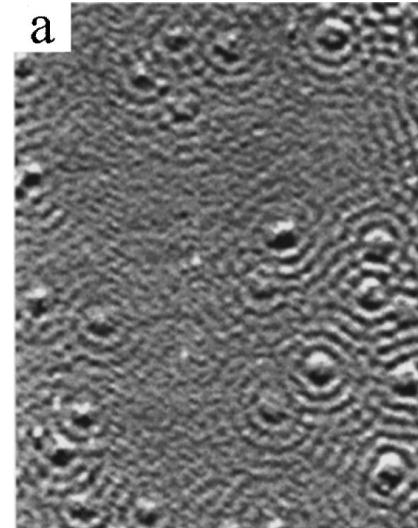


convolution with Gaussian blur
of 1 pixel per elementary cell



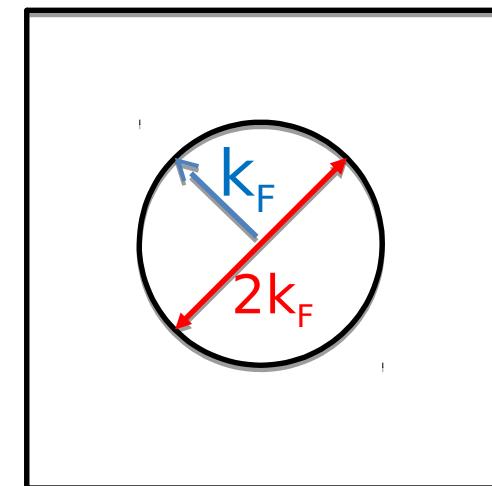
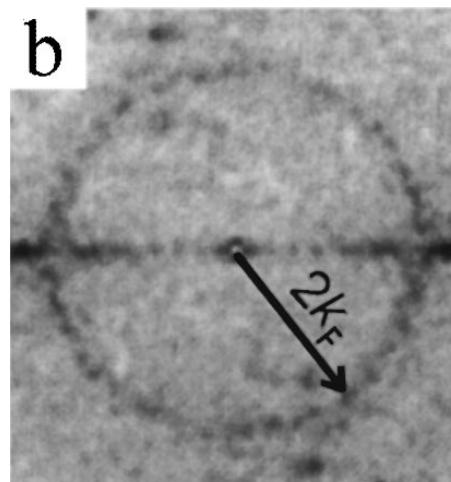
Quasiparticle Interference (QPI)

- STM on normal metal (Cu)
 - impurities
 - Friedel oscillations



L. Petersen, et al.
PRB **57**, R6858(R)
(1998)

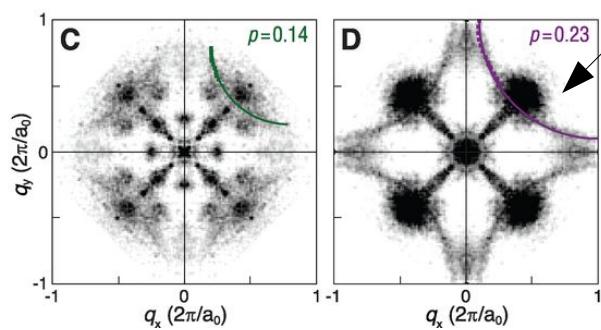
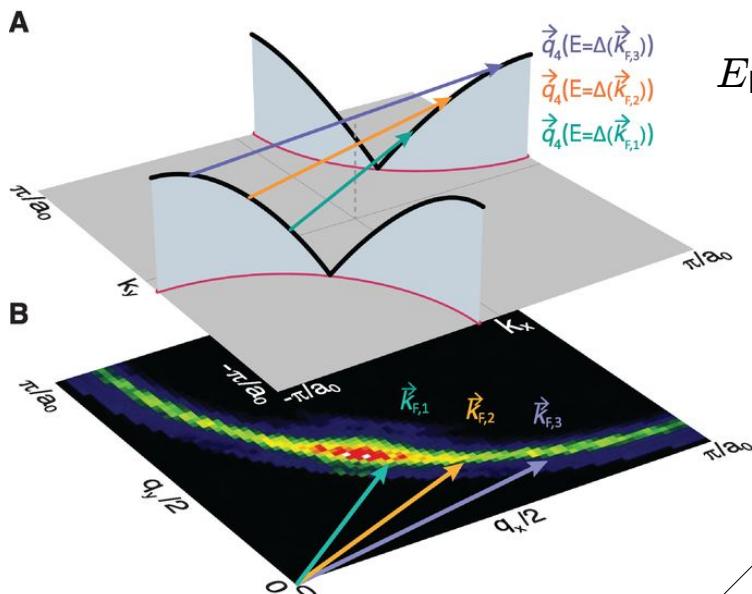
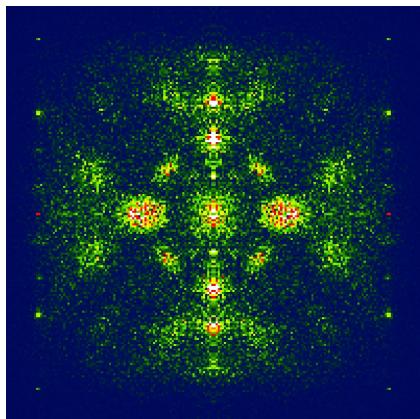
- Fourier transform of conductance map
 - mapping of constant energy contour



QPI in superconductors

- Fourier transform of differential conductance maps

FT of conductance map at +22 meV

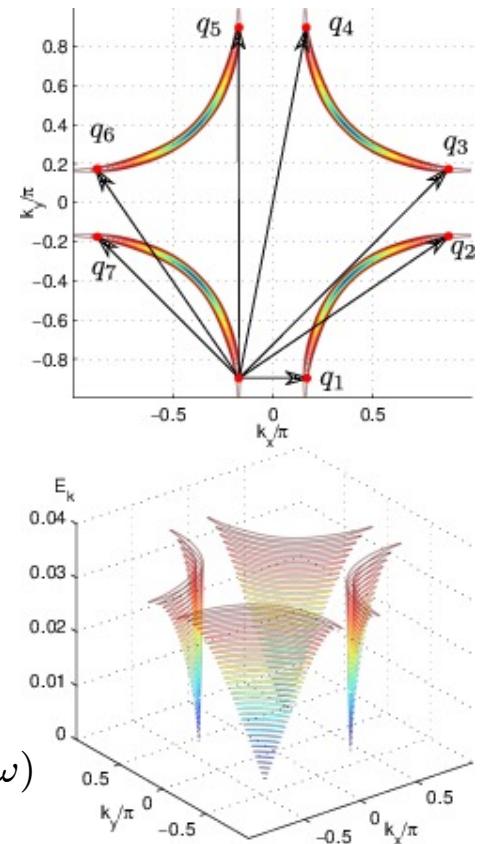


K Fujita et al. Science 344, 612 (2014)

$$E_k = \pm \sqrt{\epsilon_k^2 + \Delta_k^2}$$

energy integrated maps: trace back Fermi surface

$$\Lambda(\mathbf{q}) = \int_0^{\Delta_0} d\omega Z(\mathbf{q}, \omega)$$

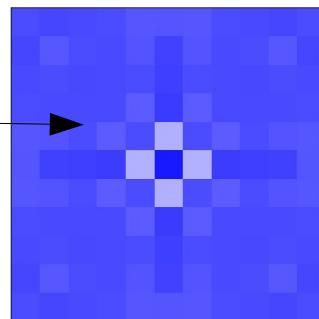


octet model: 7 scattering vectors between regions of high DOS

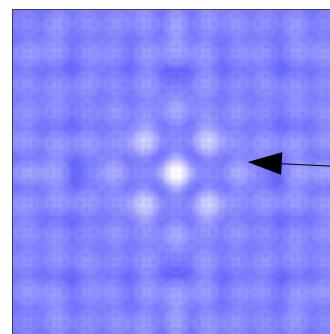
Quasi Particle Interference (QPI)

- Fourier transform of conductance maps
- BSCCO: weak potential scatterer

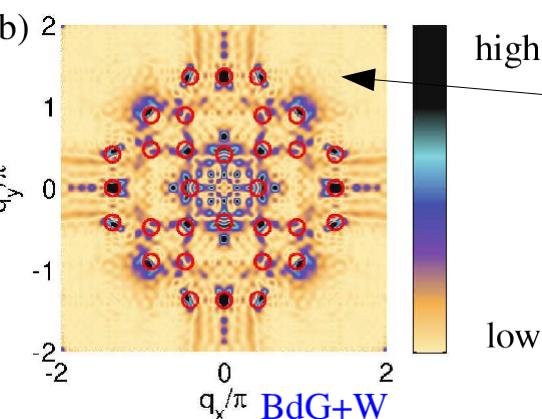
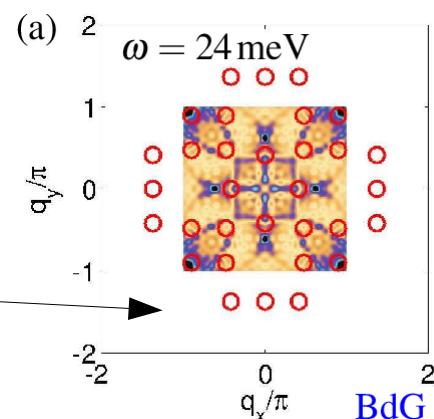
no intra-unitcell information
1 pixel per elementary cell



atomic scale local density of states at STM tip position

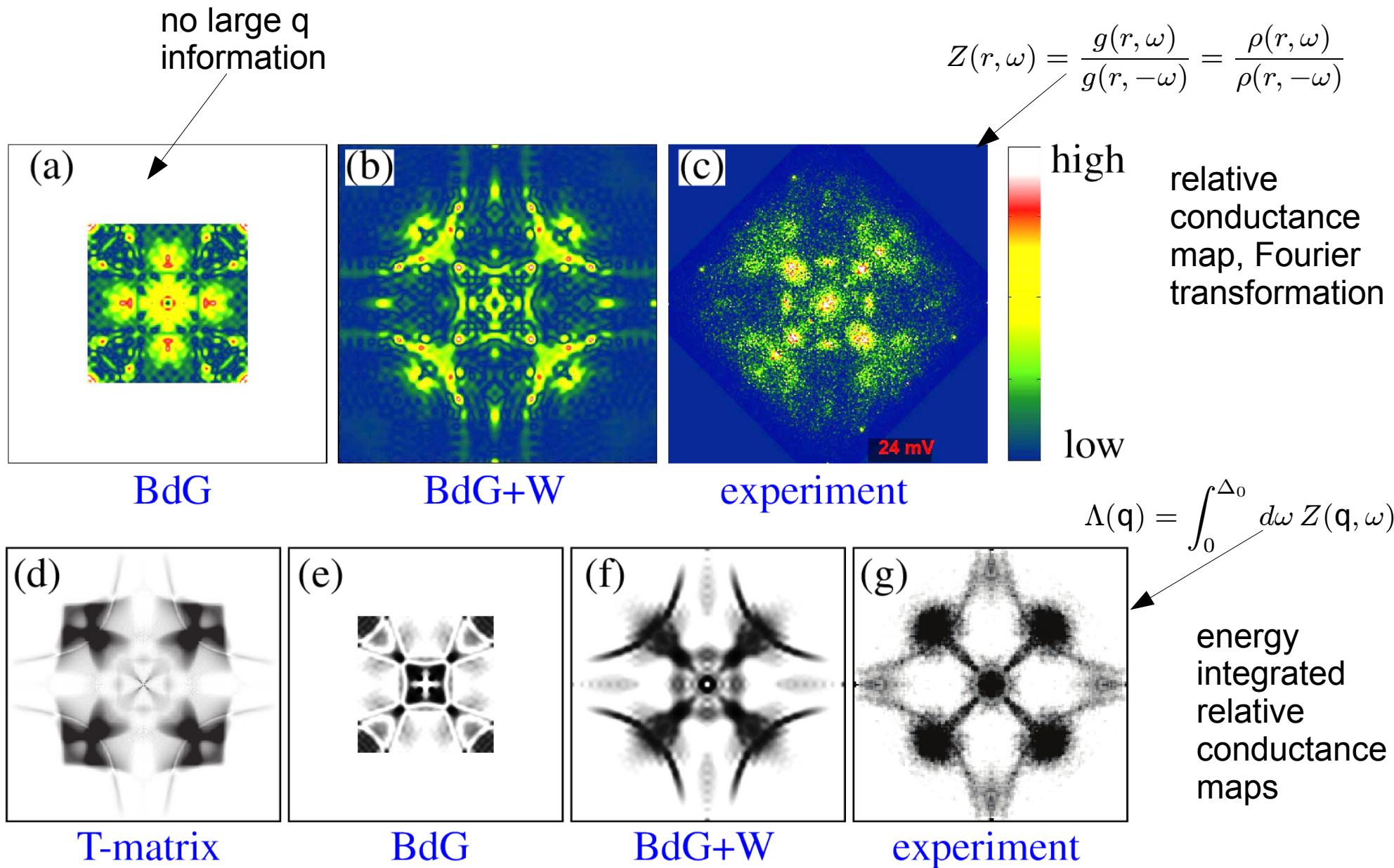


no information beyond first BZ



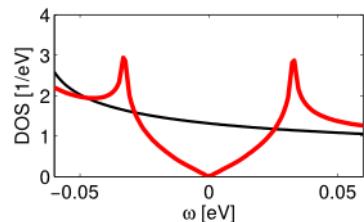
full information for all scattering vectors

Comparison to experiment



Homogeneous superconductor

- phenomenological pairing interactions
similar results from spin-fluctuation pairing



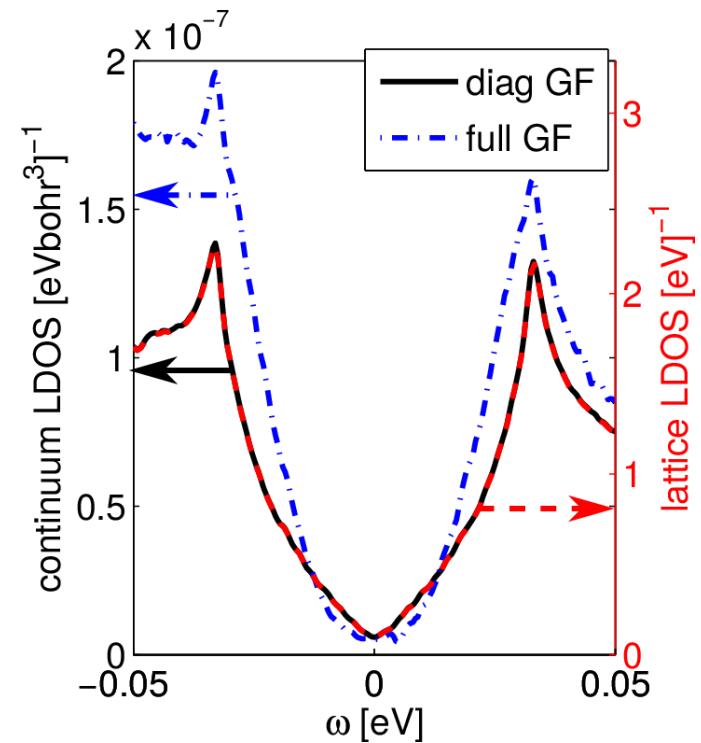
DOS of homogeneous superconductor

- spectra measured at the surface

$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_{\mathbf{R}, \mathbf{R}'} G(\mathbf{R}, \mathbf{R}'; \omega) w_{\mathbf{R}}(\mathbf{r}) w_{\mathbf{R}'}^*(\mathbf{r}')$$

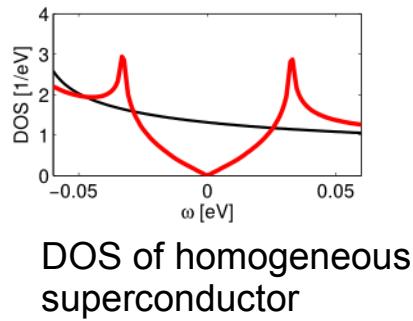
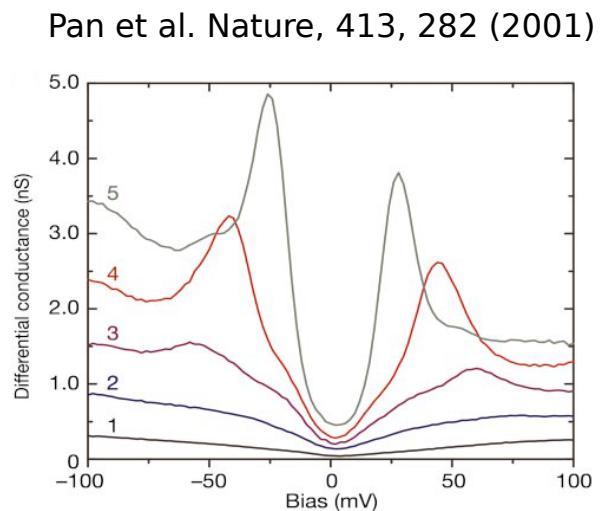
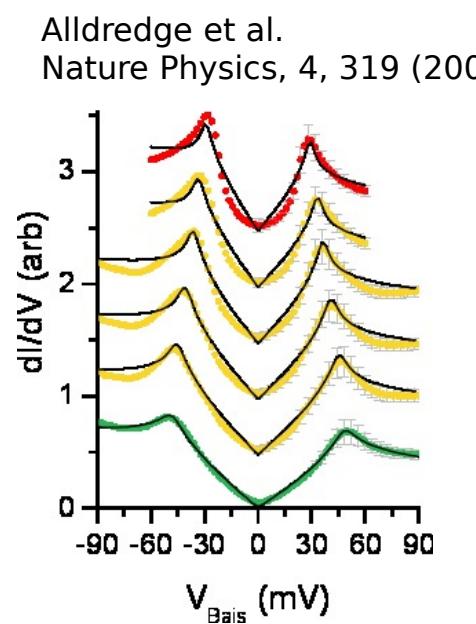
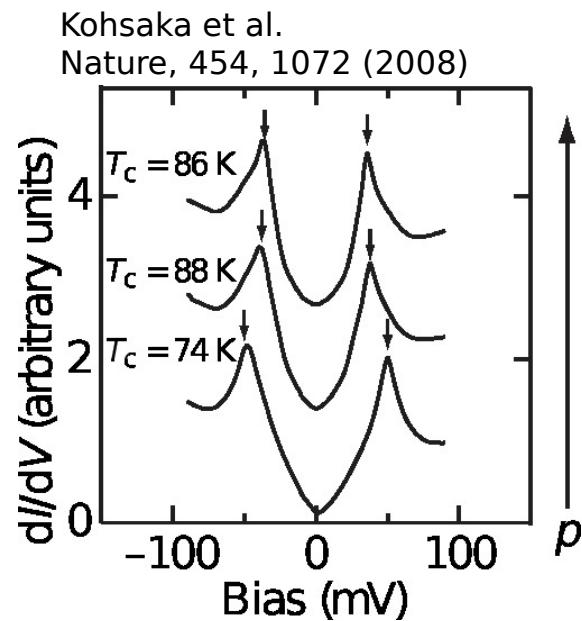
local density of states (LDOS)

$$\rho(\mathbf{r}, \omega) \equiv -\frac{1}{\pi} \text{Im} G(\mathbf{r}, \mathbf{r}; \omega)$$

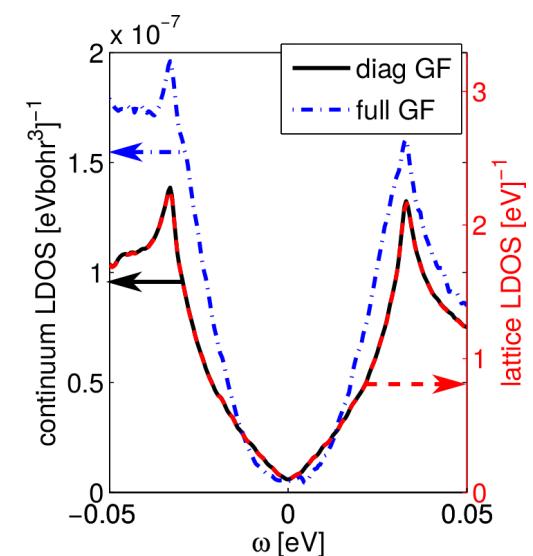


STM Spectra: homogeneous SC

- overdoped: U-shape, lower doping: V-shape

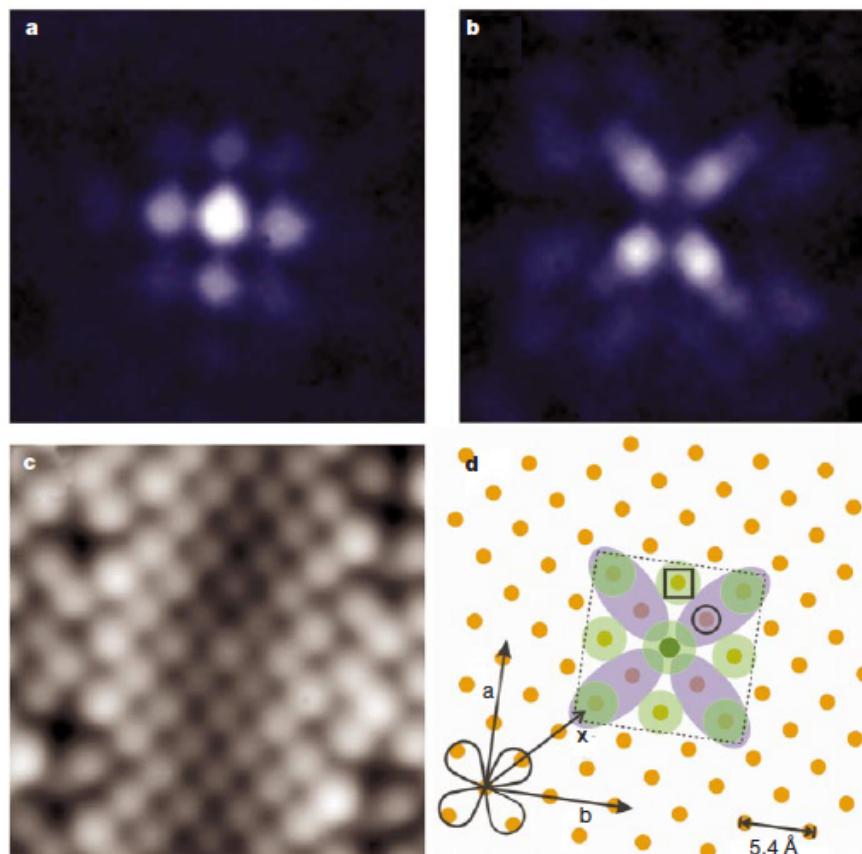


BdG+W: U shape enters naturally within our method, applicable to overdoped regime

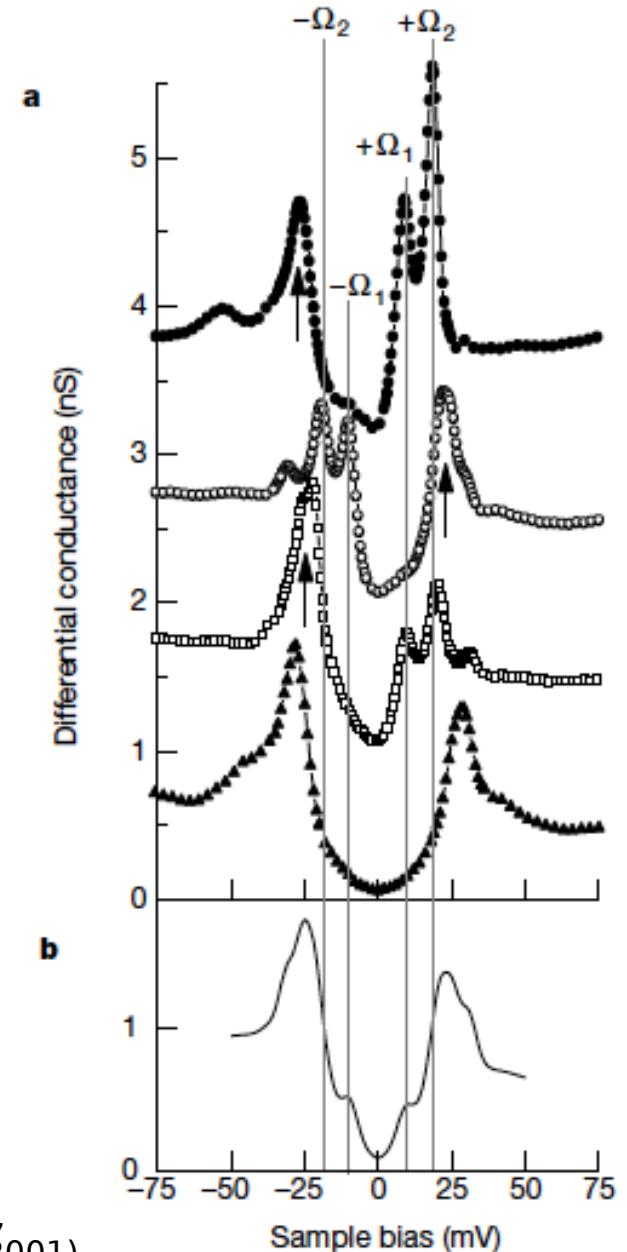


BSCCO: Magnetic impurity

- Ni on BSCCO:
weak magnetic scatterer
- double resonance

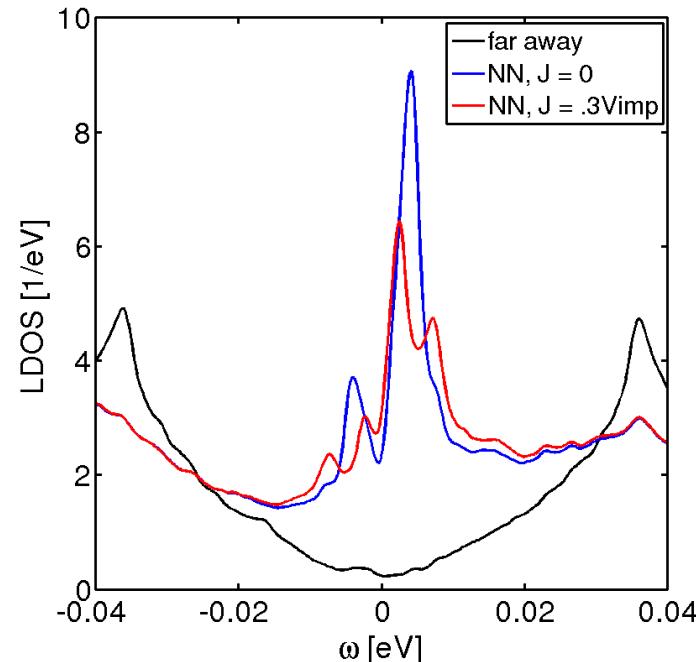


E. W. Hudson et al.,
Nature **411**, 920 (2001)

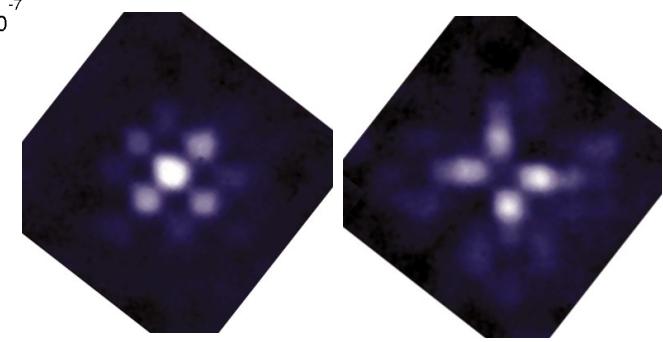
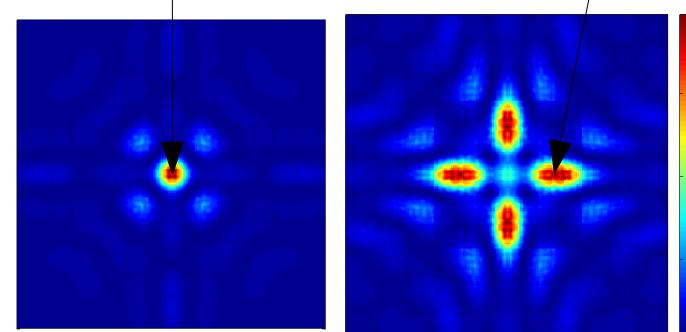
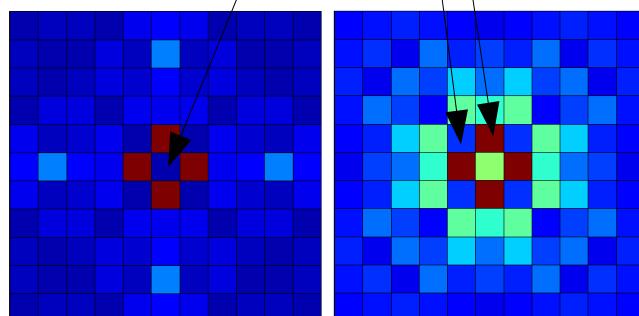
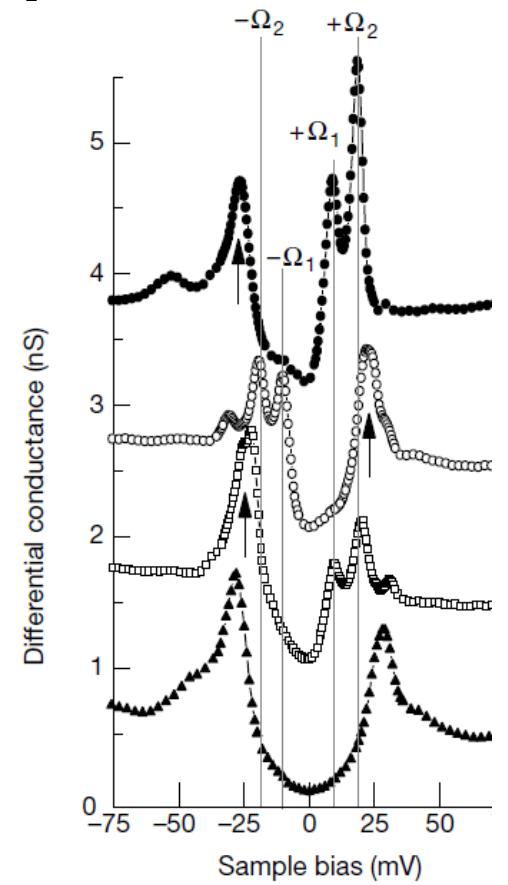
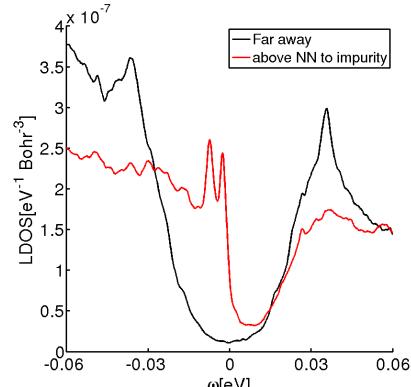
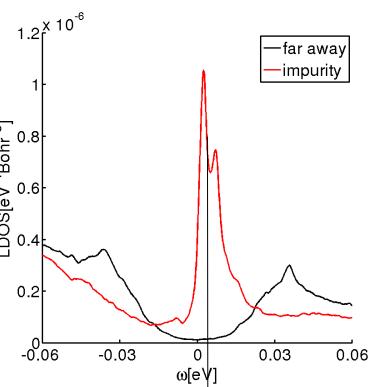
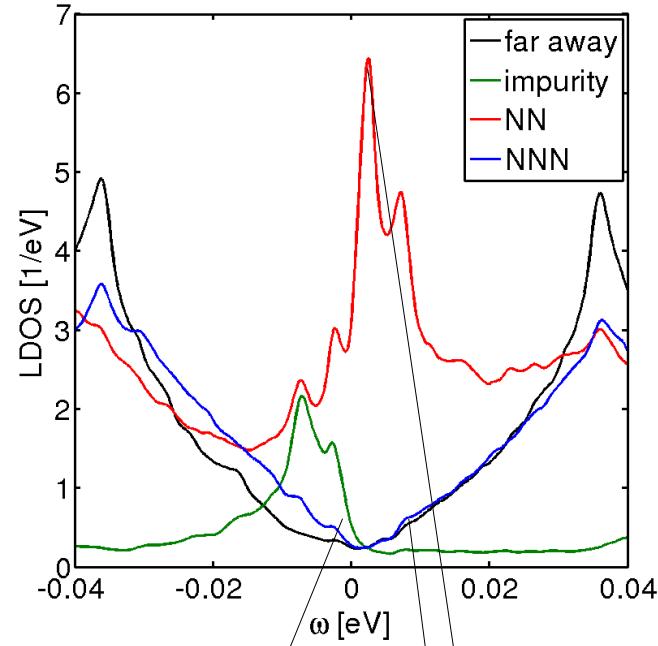


Ni impurity on BSCCO

- chemistry: Ni 3d₈ configuration → magnetic moment with S=1
- Classical spin: additional magnetic potential
$$H_{\text{imp}}^{\text{mag}} = J(n_R * \uparrow - n_R * \downarrow)$$
- resonances at
+/- 2.4 meV (up)
+/- 7.2 meV (down)



Results: spectra, maps



$\Omega = -2.4 \text{ meV}$

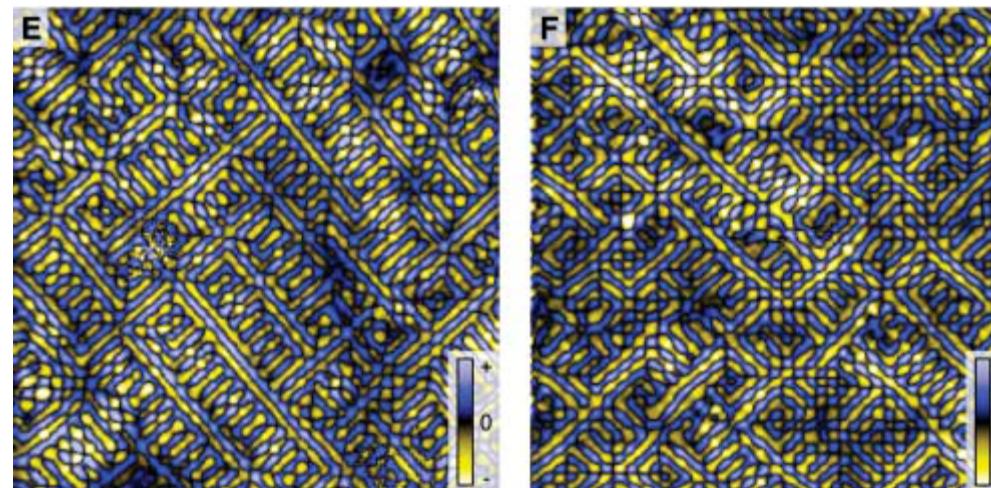
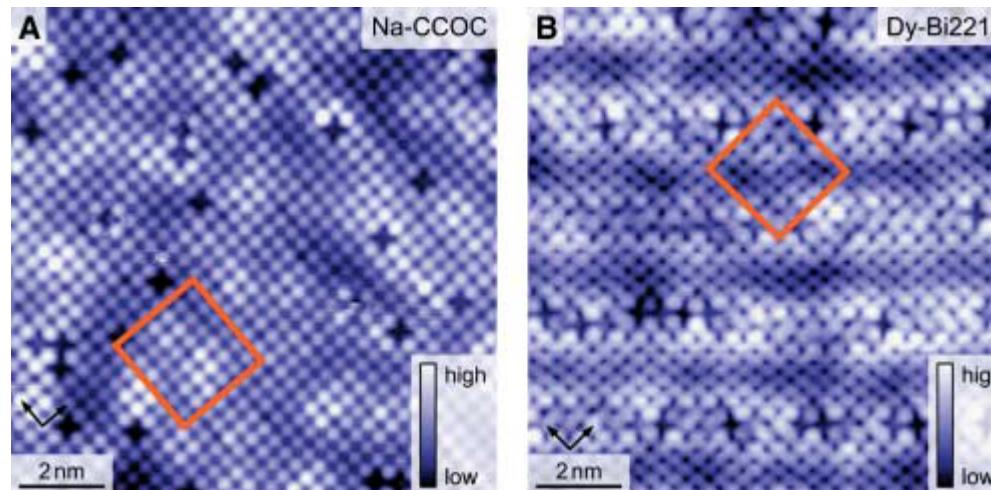
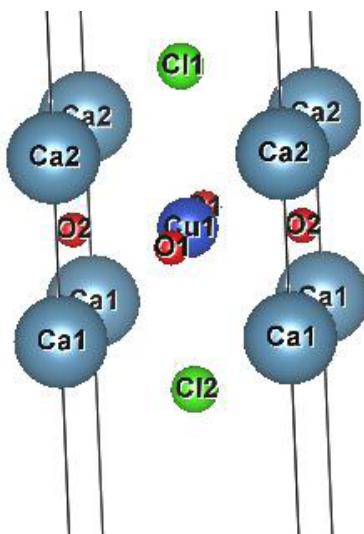
$\Omega = 2.4 \text{ meV}$

Differences between materials?

- Different tunnelling layers, same physics

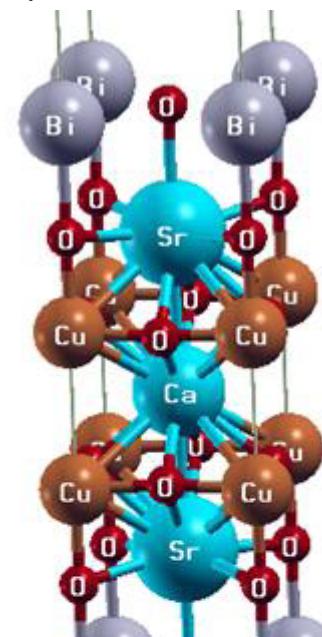
Na-CCOC

Ca-Cl surface

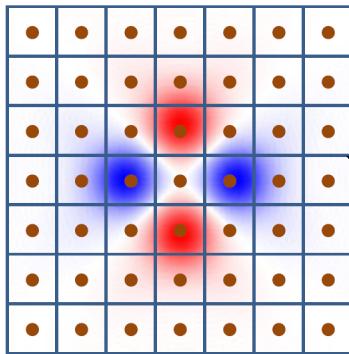


Dy-Bi2212

Bi-O surface
(Sr-O in between)



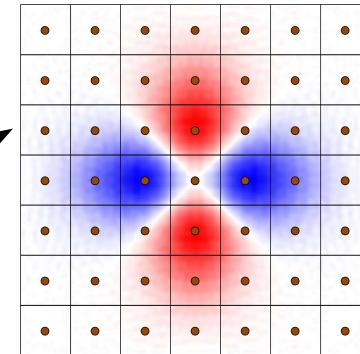
Wannier functions



Na-CCOC

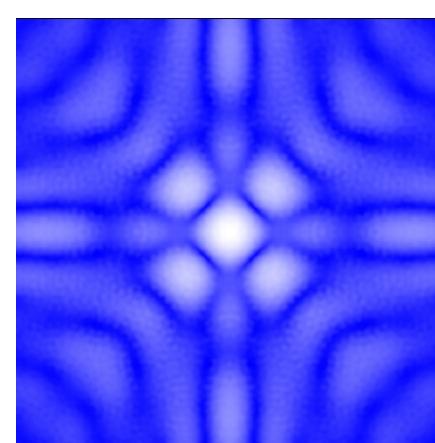
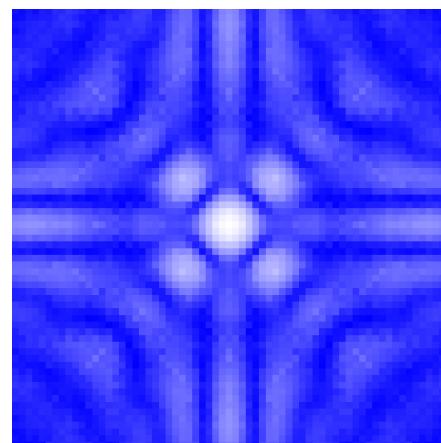
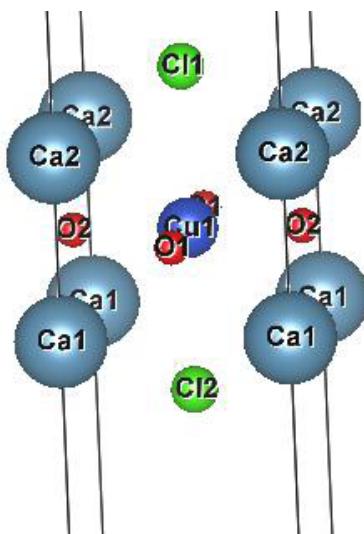
Ca-Cl surface

Wavefunctions very similar due to symmetry constraints on downfolding to dx^2-y^2 band in conjunction with crystal symmetry

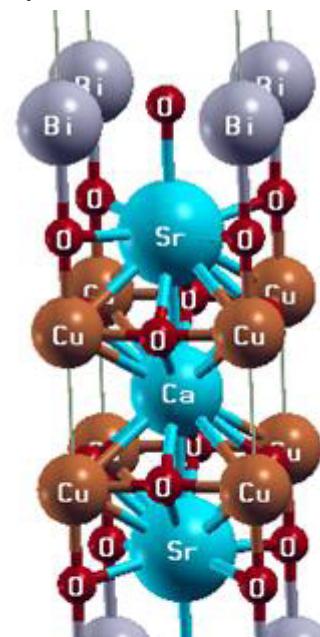


Dy-Bi2212

Bi-O surface
(Sr-O in between)



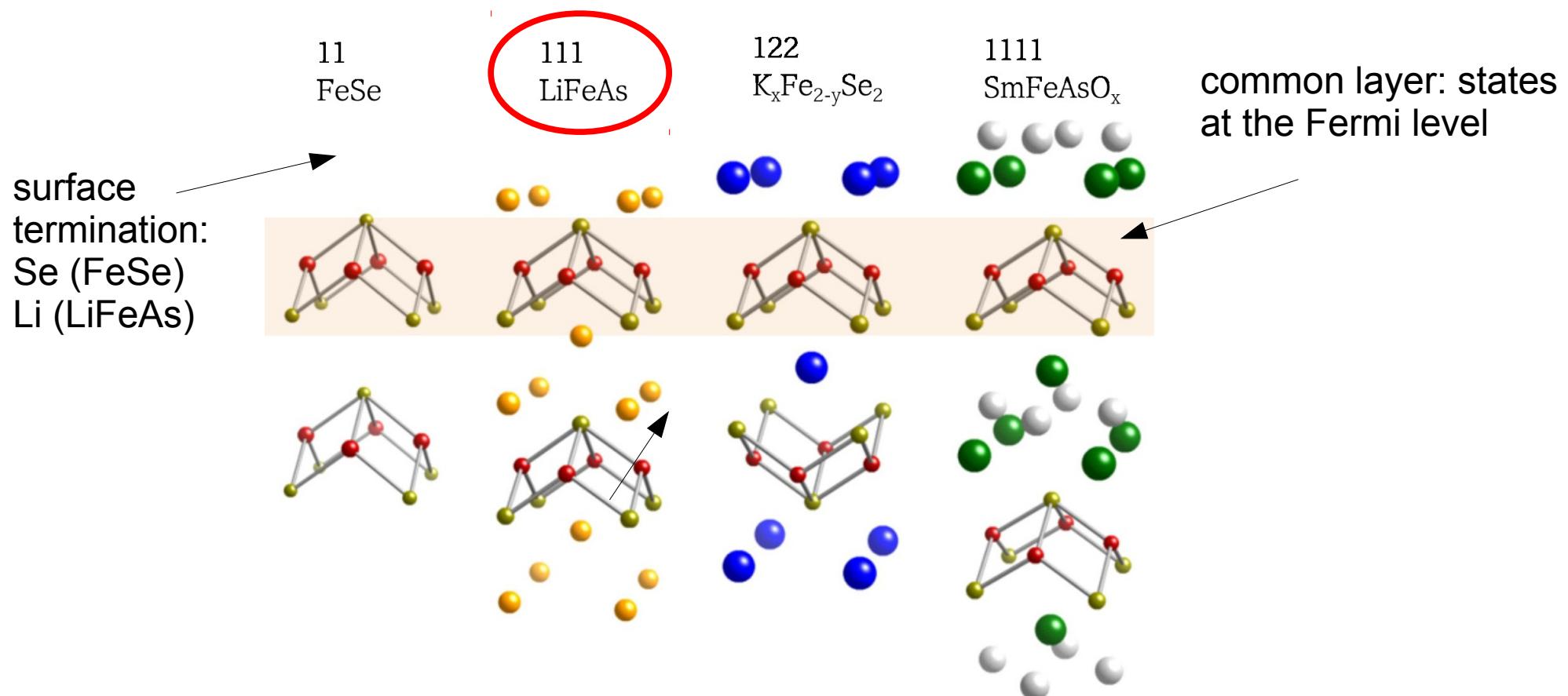
Conductance maps for a strong potential scatterer (in the SC state, d-wave)



Layered superconductors

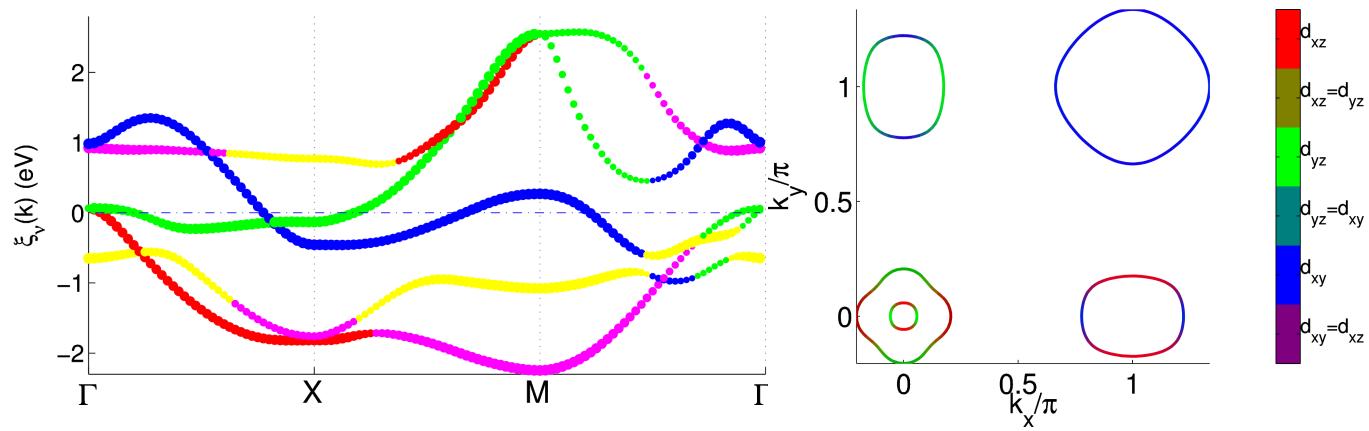
LDOS of sample at given energy **at the tip position**

- Iron based superconductors



Gap symmetries: FeSC

- Fermi surface
5 band model



- Possible order parameters

s_{++}



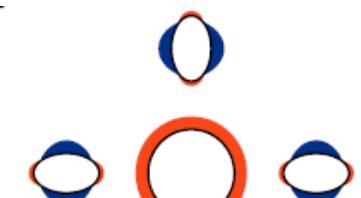
s_{\pm}



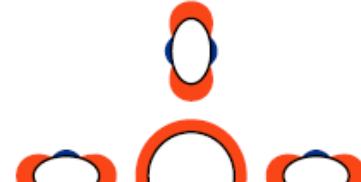
nodal s_{\pm}



d



s_{\pm}



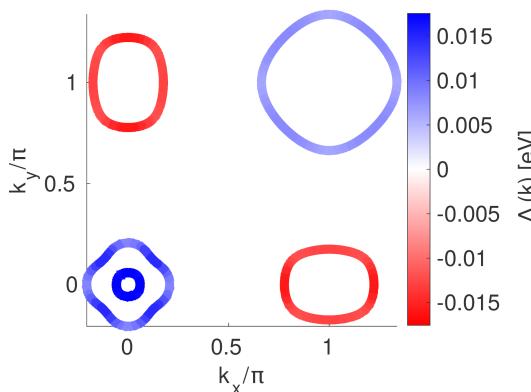
s_{++}

LiFeAs: 18K superconductor

- proposed gap structures

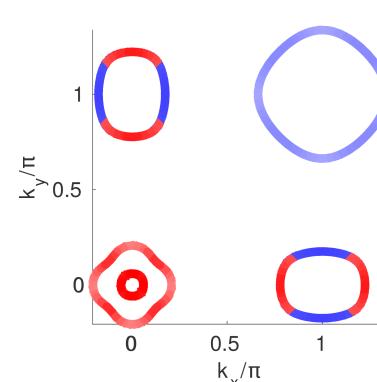
conventional s_{\pm}

Y. Wang, A. Kreisel, et al.
Phys. Rev. B 88, 174516 (2013)



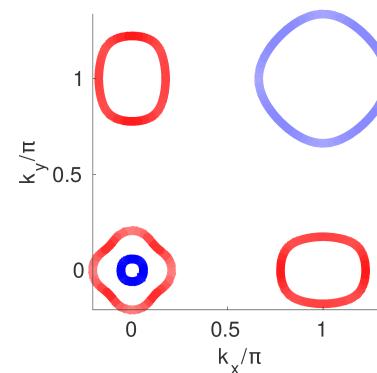
antiphase s_{\pm}

Z. P. Yin, K. Haule, G. Kotliar
Nature Physics 10, 845-850 (2014)



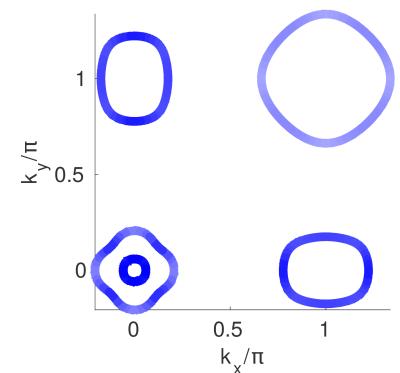
novel s_{\pm}

F. Ahn, et al. Phys. Rev. B 89, 144513 (2014)

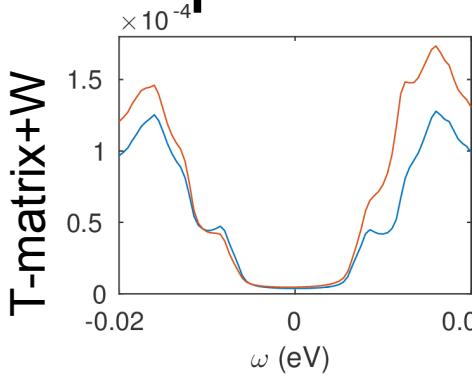


conventional s_{++}

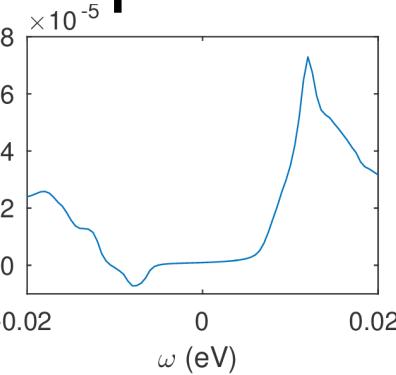
Tetsuro Saito, et al.
Phys. Rev. B 90, 035104 (2014)



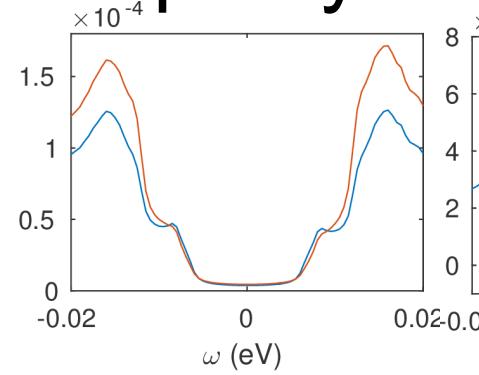
- experimental probes? → impurity bound states



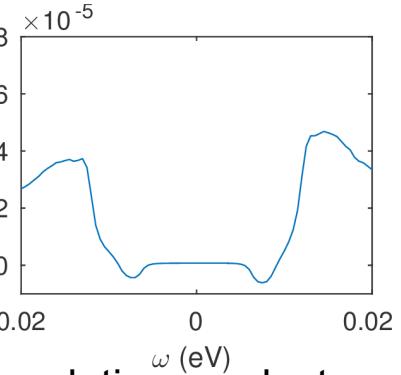
conventional s_{\pm}



relative conductance:
peak-dip



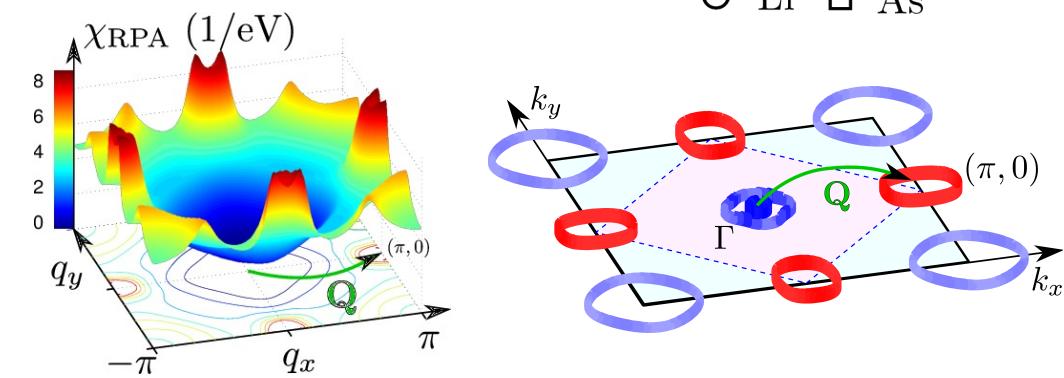
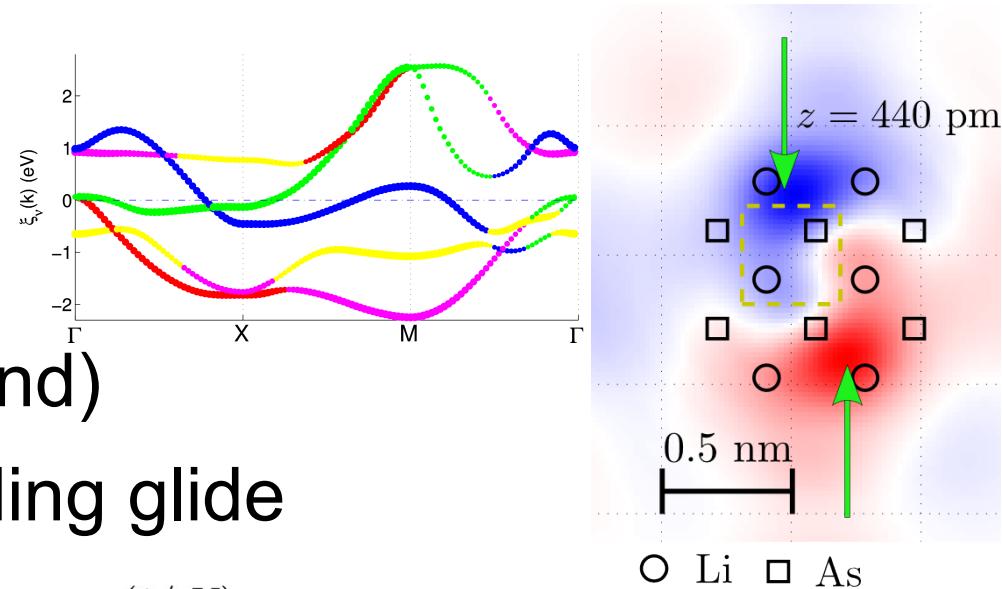
conventional s_{++}



relative conductance:
modification of
normal state DOS

Theory: T-matrix+Wannier

- ab-initio calculation for LiFeAs (surface)
 - tight-binding model (5 band)
 - Wannier functions (including glide plane symmetry)
- superconducting order parameter from spin-fluctuation theory
- ab-initio impurity potentials (Ni, Mn, Co) used in T-matrix calculation



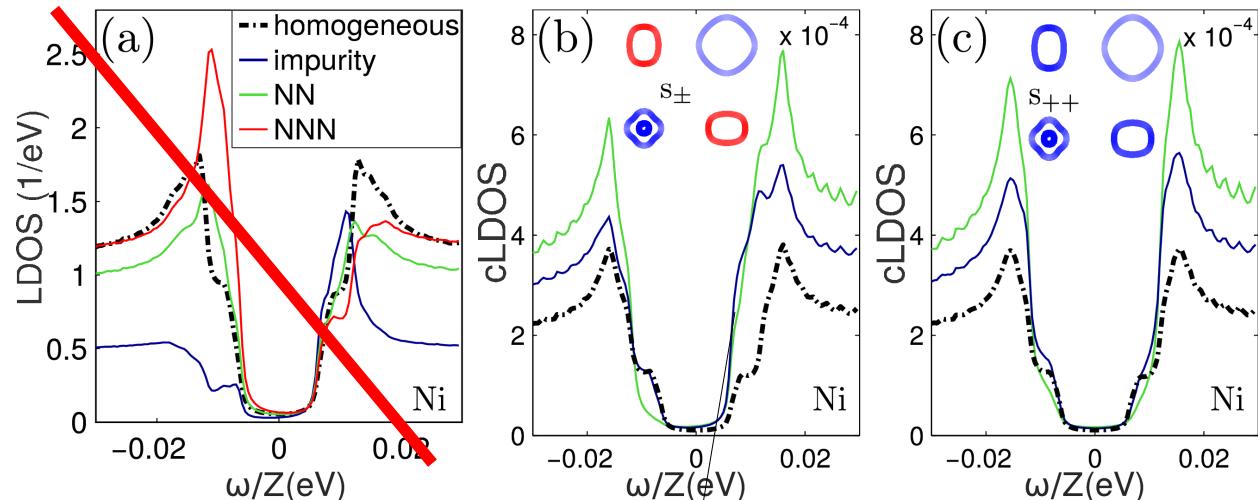
$$\hat{G}(\mathbf{r}, \mathbf{r}'; \omega) = \hat{G}_0(\mathbf{r} - \mathbf{r}', \omega) + \hat{G}_0(\mathbf{r}, \omega) \hat{T}(\omega) \hat{G}_0(\mathbf{r}', \omega)$$

$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_{\mathbf{R}, \mathbf{R}'} G(\mathbf{R}, \mathbf{R}'; \omega) w_{\mathbf{R}}(\mathbf{r}) w_{\mathbf{R}'}^*(\mathbf{r}')$$

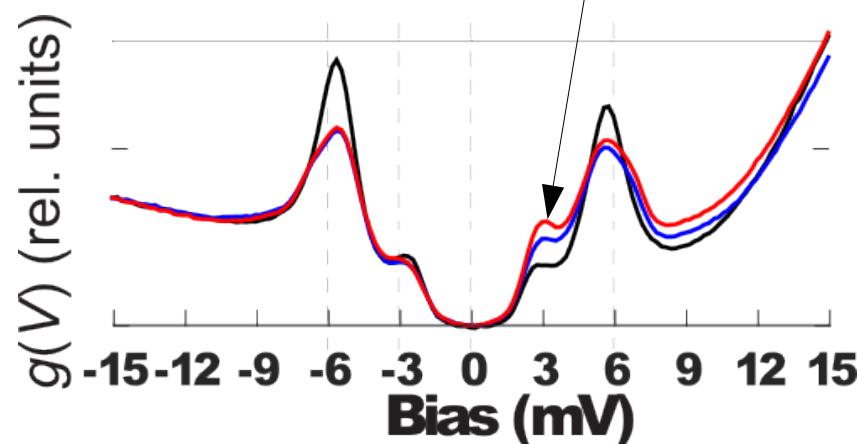
$$\rho(\mathbf{r}, \omega) \equiv -\frac{1}{\pi} \text{Im} G(\mathbf{r}, \mathbf{r}; \omega)$$

LiFeAs: spectra

- evidence for sign-changing order parameter by in-gap state with engineered impurity

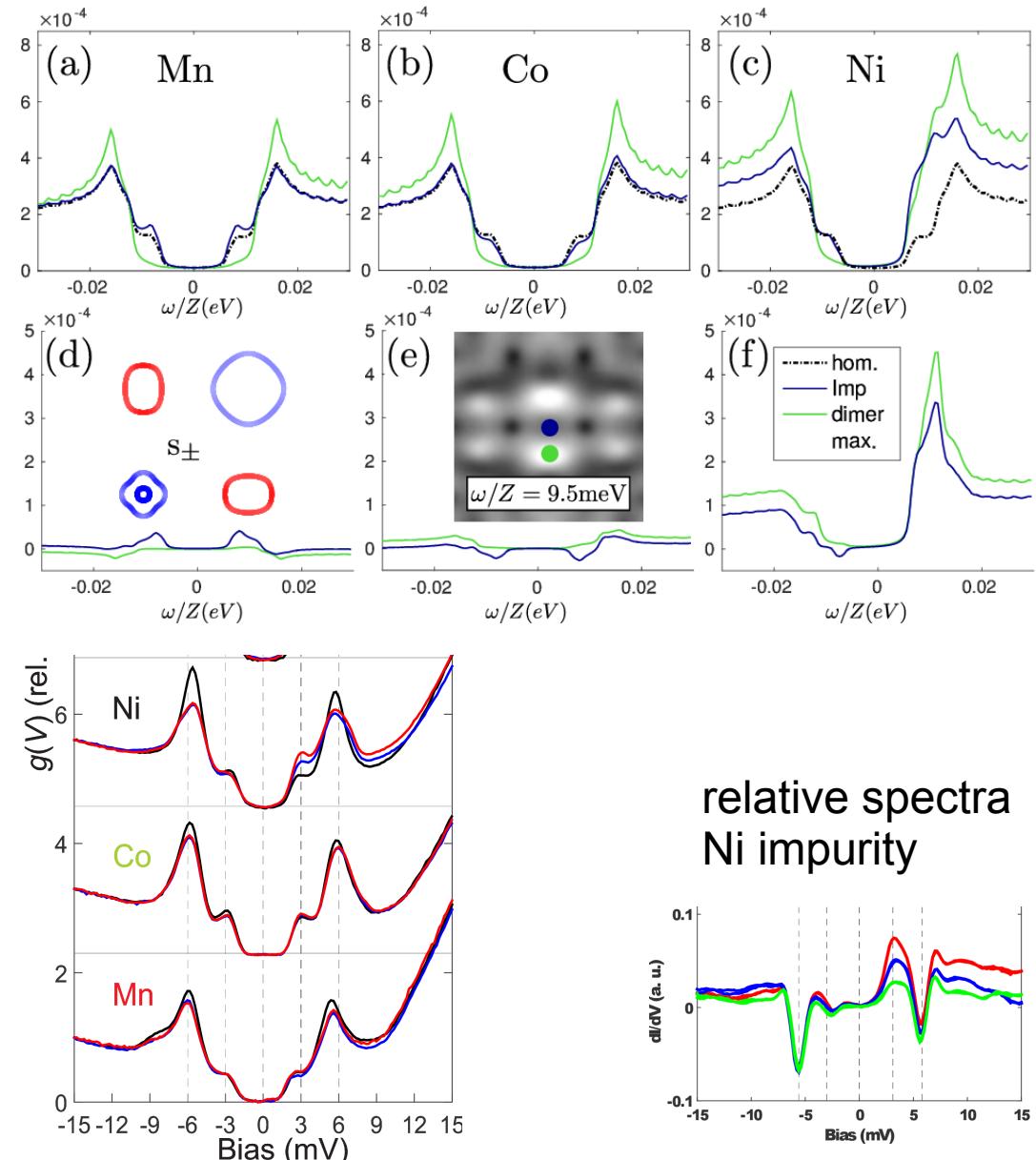


lattice LDOS: strong response at negative bias

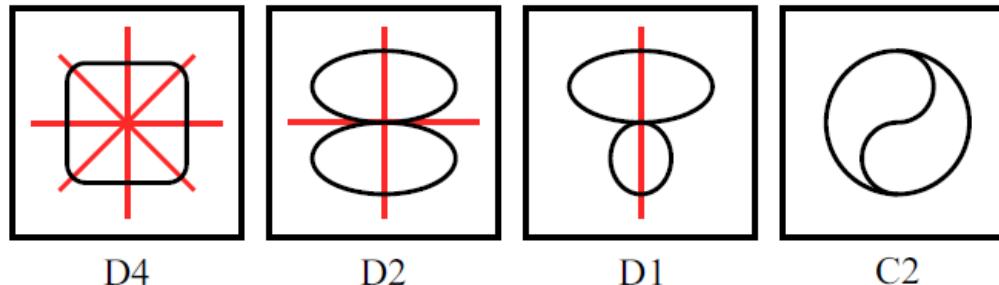


LiFeAs: spectra

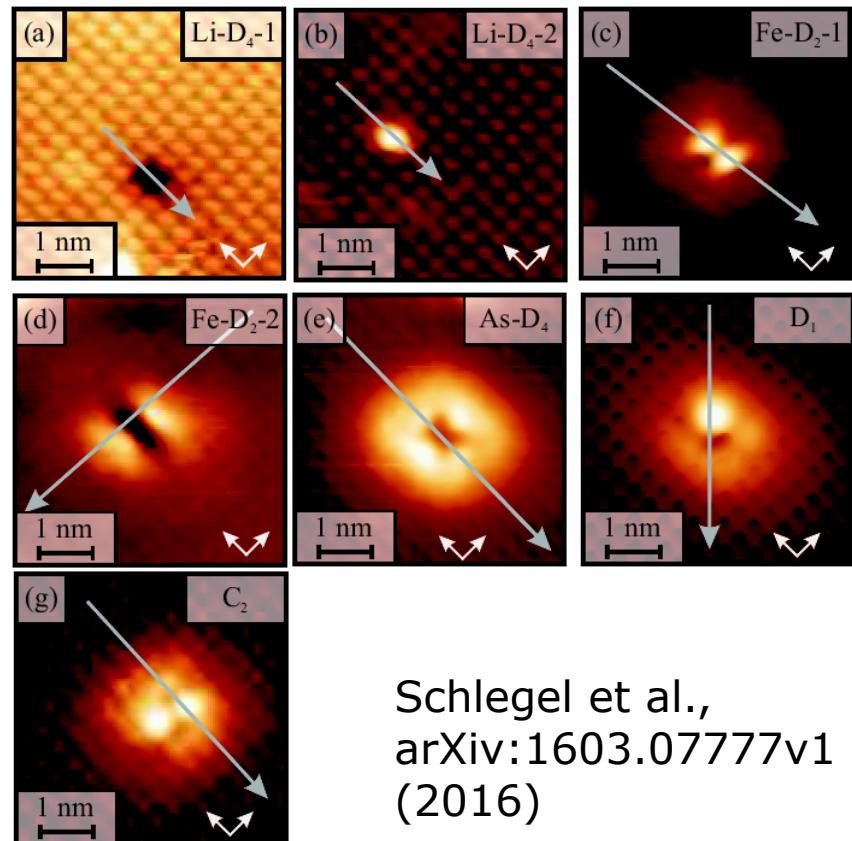
- sequence of impurity potentials from ab-initio calculation correct, but overall renormalization downwards required.



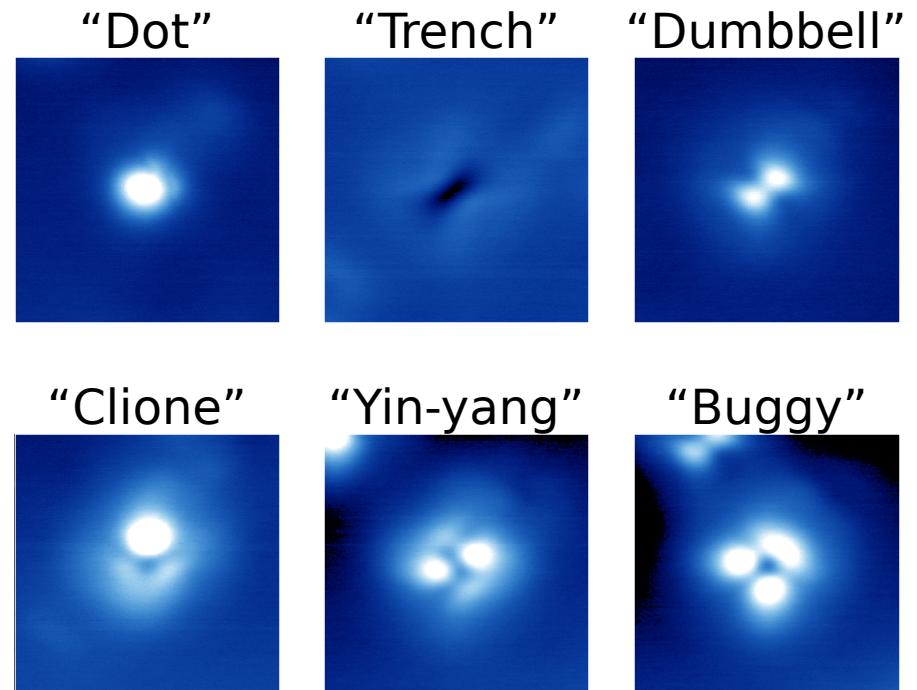
LiFeAs: other native impurities



Schönflies classification
of impurities



Hanaguri, unpublished (KITP 2011)

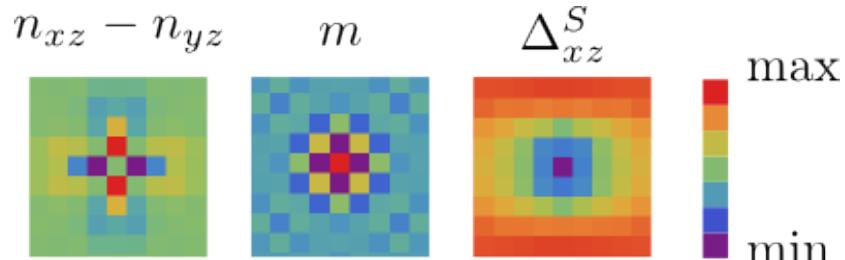


Schlegel et al.,
arXiv:1603.07777v1
(2016)

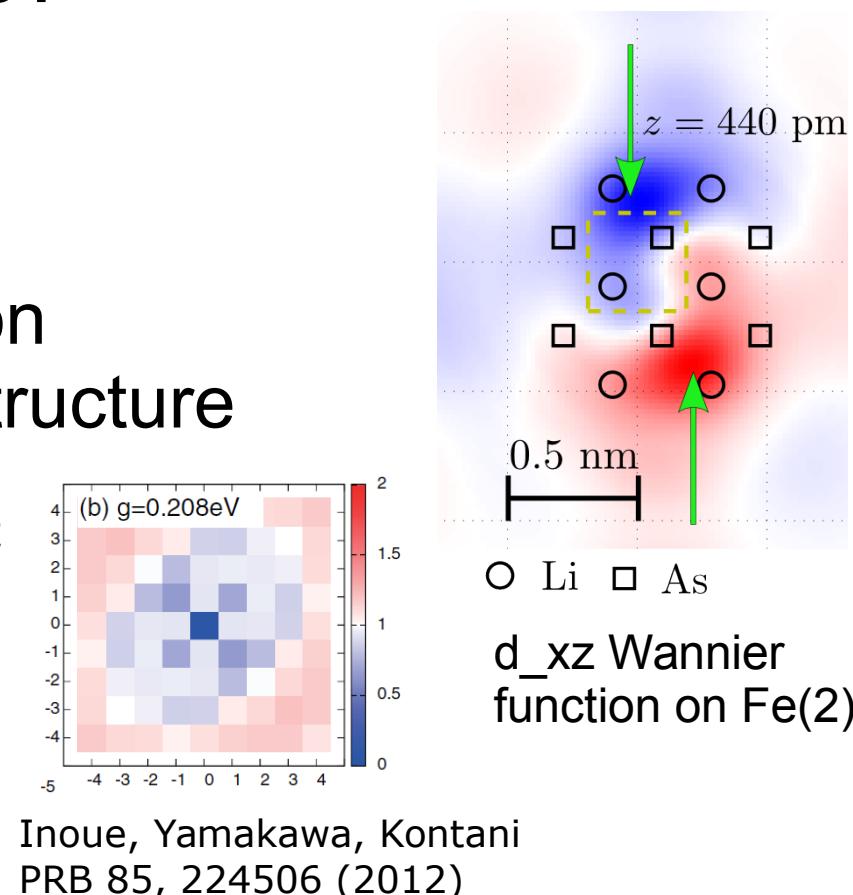
Chiral defects

- from a symmetry perspective not compatible to impurities on any single site in LiFeAs
- multiple impurities?
- local order?

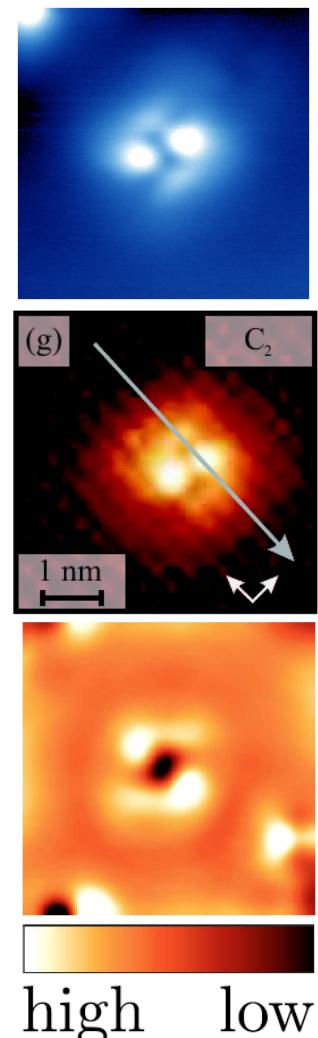
local orbital order
+ Wannier function
→ chiral defect structure



Gastiasoro, Andersen, J. Supercond Nov. Magn., **26**, 2651 (2013)

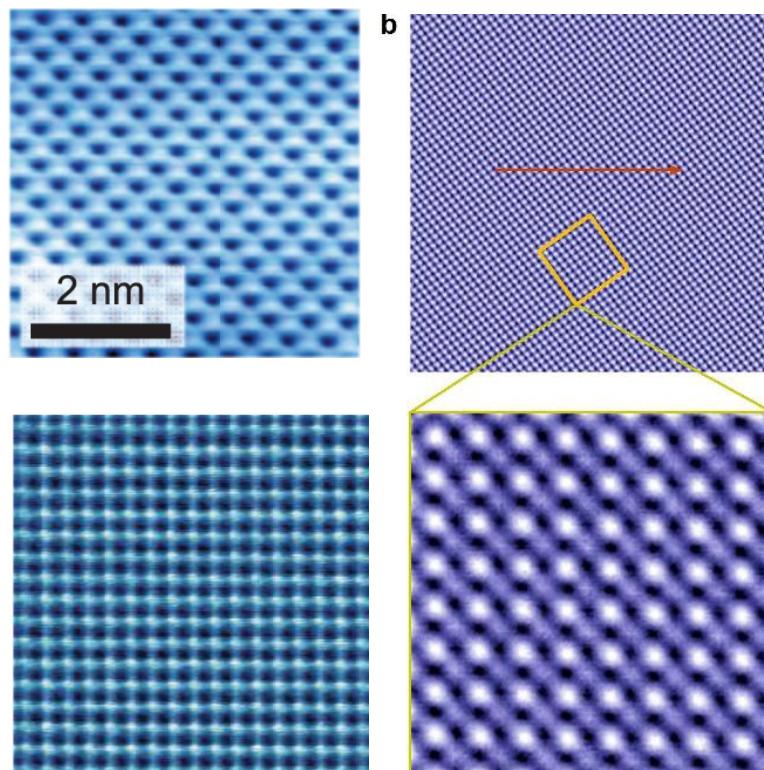


Inoue, Yamakawa, Kontani PRB 85, 224506 (2012)

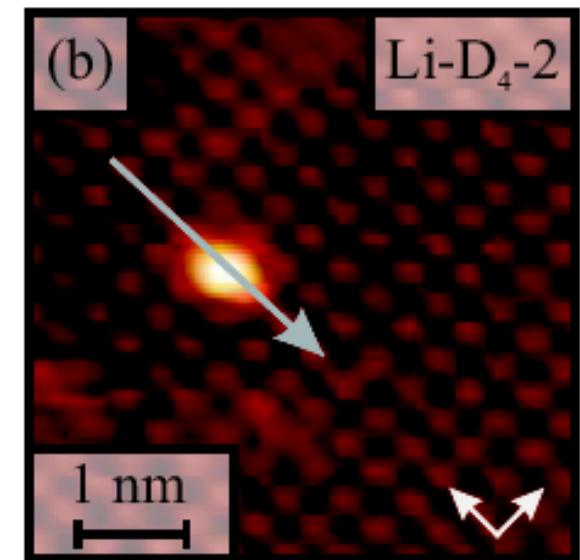


Height and current dependence of topographs

- experiment: Li or As lattice?



height maxima at Li positions!?
counter-intuitive from chemistry point of view

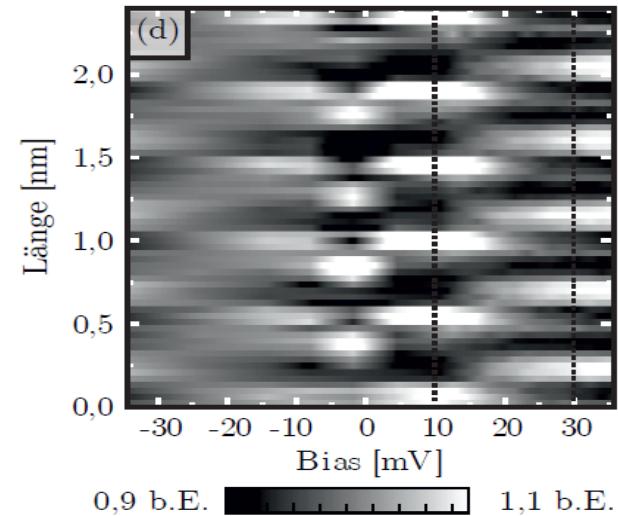
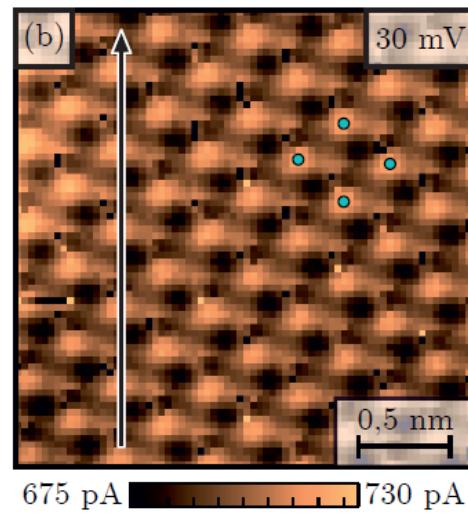
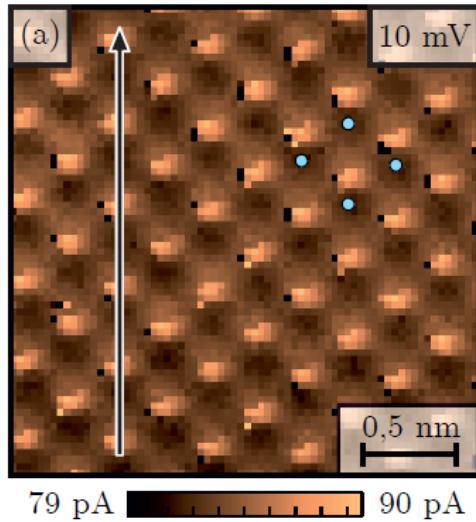


Shun Chi, et al., PRL 109, 087002 (2012)
T. Hanaguri, et al. PRB 85, 214505 (2012)
S. Grothe, et al., PRB 86, 174503 (2012)
J. -X. Yin, et al., arXiv, 1602.04949 (2016)

Schlegel et al.,
arXiv:1603.07777v1
(2016)

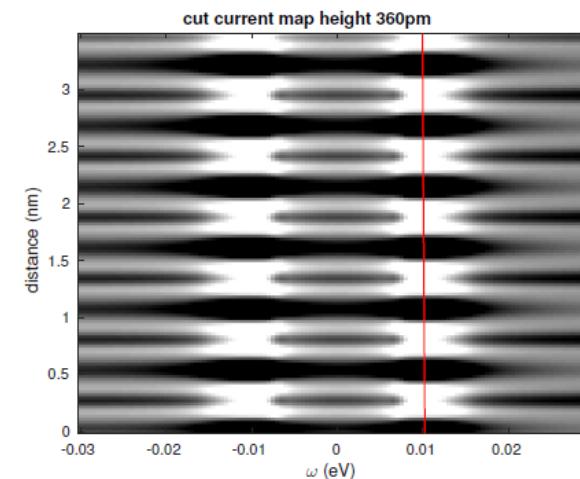
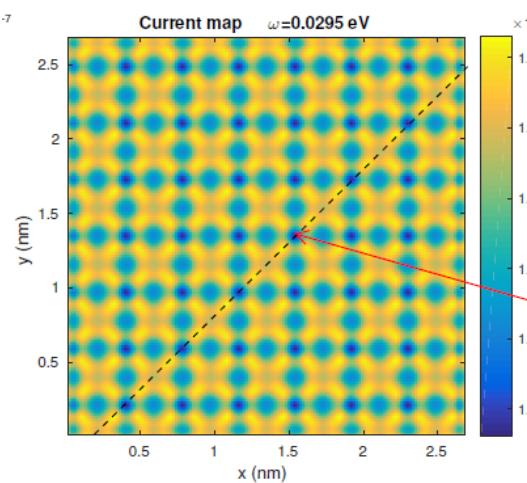
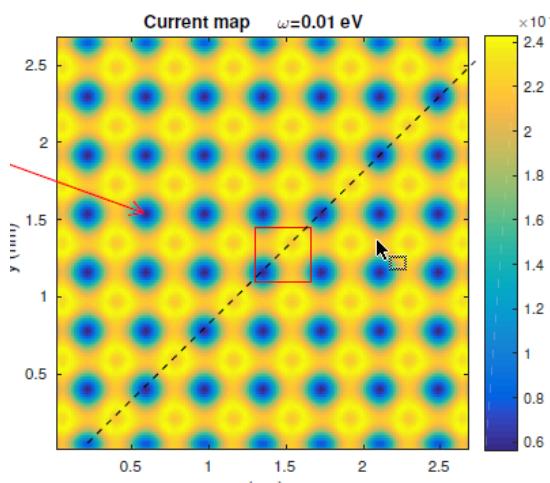
Further experimental evidences?

experiment (current maps)



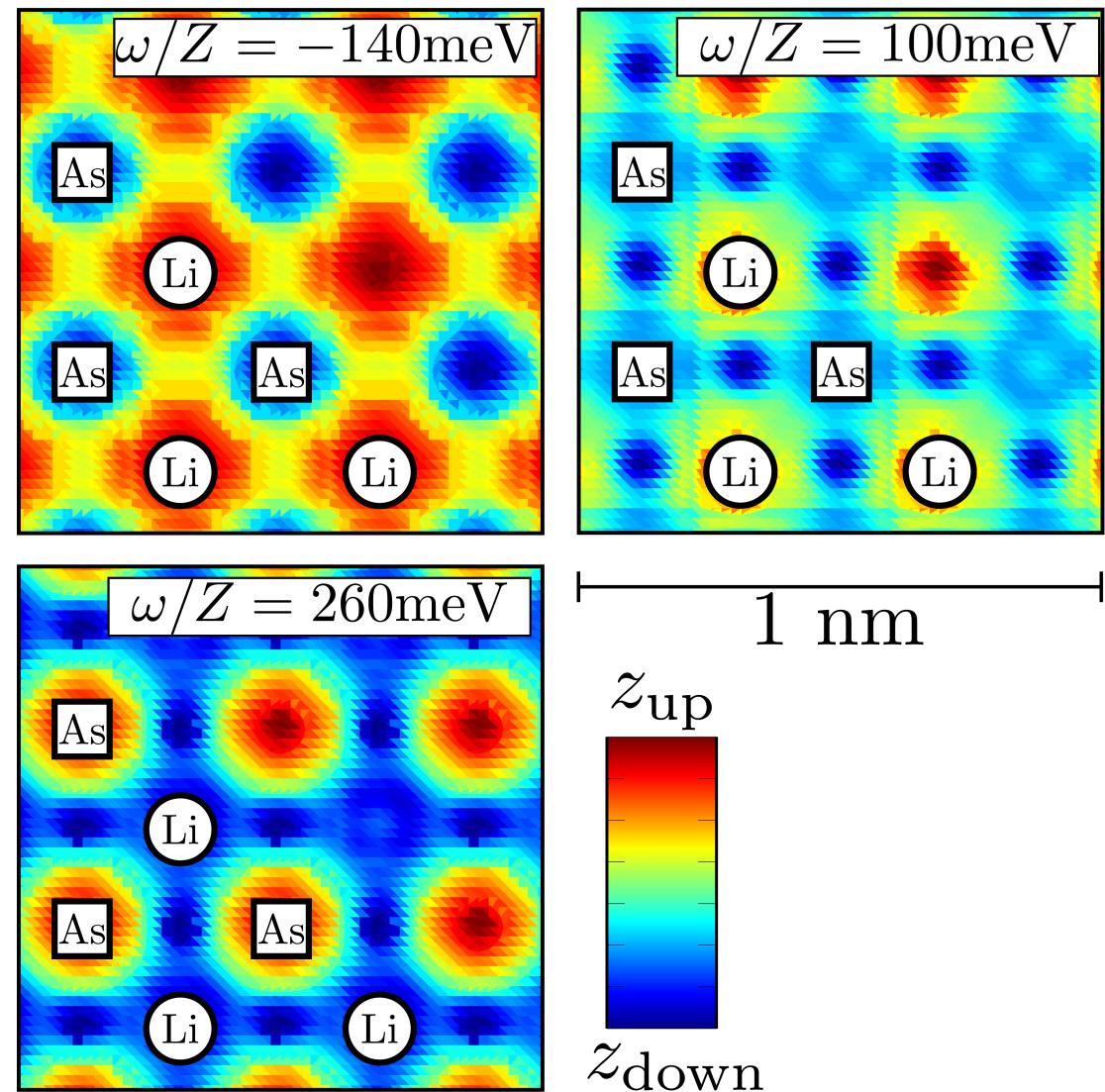
Ronny Schlegel, Dissertation, TU Dresden
(thanks to C. Hess)

theory



Simulation of topographs

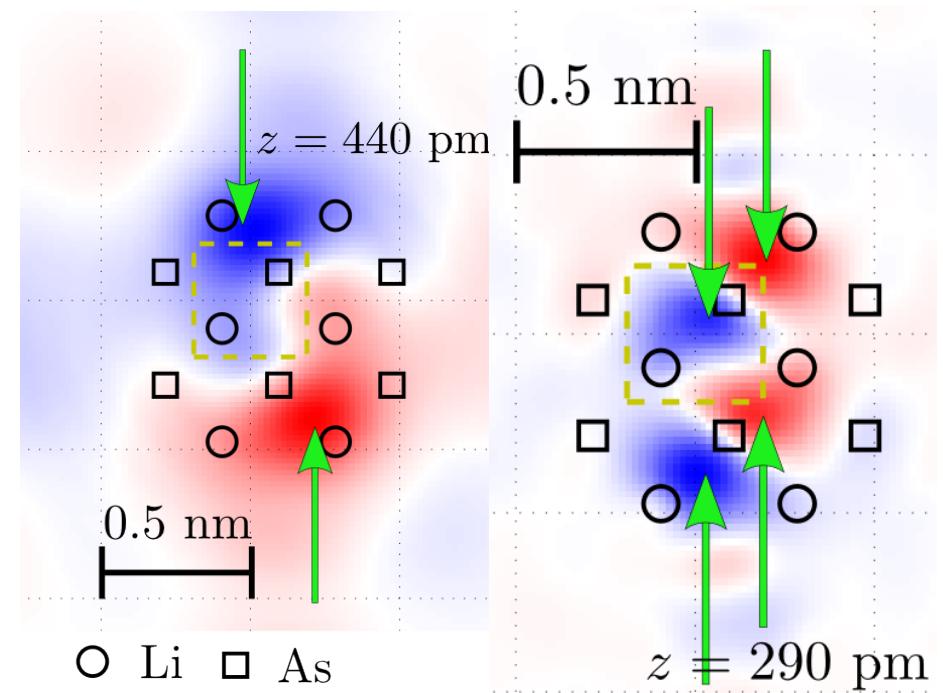
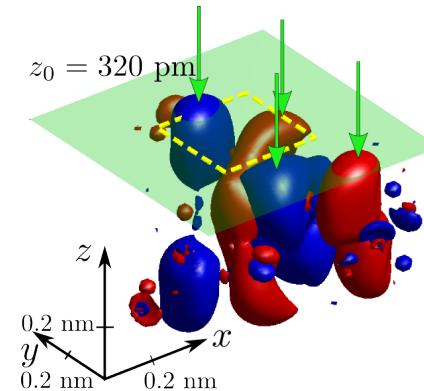
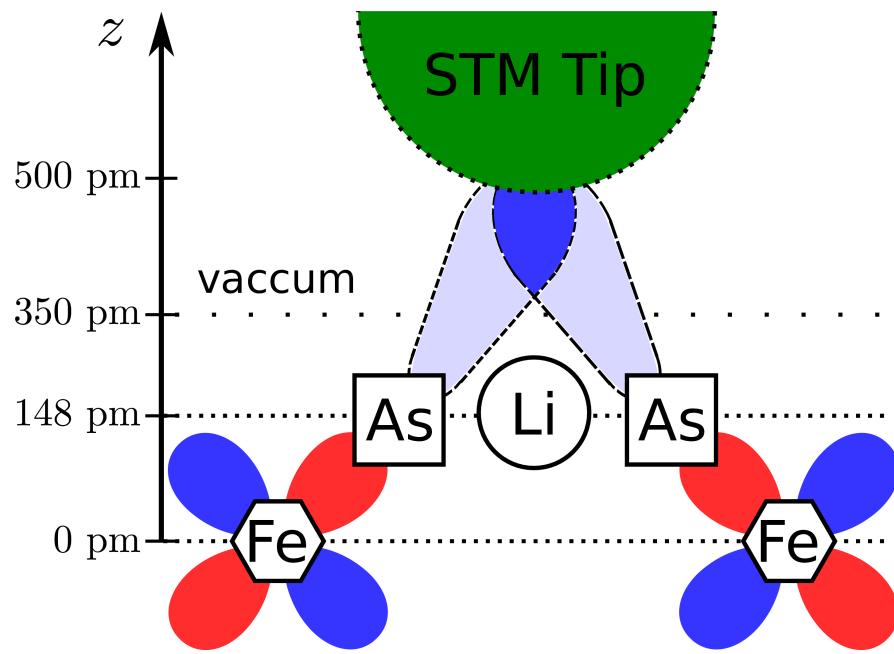
- switching of height maxima as a function of bias voltage



Simulation of topographs

- interplay of interference of wavefunctions and their spatial structure

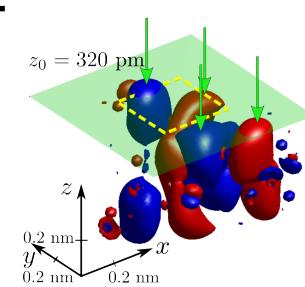
$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_{\mu, \nu, R, R'} G(R, \mu, R', \nu; \omega) w_{R, \mu}(\mathbf{r}) w_{R', \nu}^*(\mathbf{r}')$$



Recapitulation: BdG/T-matrix+W

- **simple:** just a basis transformation of the Green's function
- **powerful** tool for calculation of local density of states at the surface (STM tip position) of superconductors
- takes into account atomic scale information and symmetries of the elementary cell and the contained atoms
- **shown to work**

$$G(\mathbf{r}, \mathbf{r}'; \omega) = \sum_{\mathbf{R}, \mathbf{R}'} G(\mathbf{R}, \mathbf{R}'; \omega) w_{\mathbf{R}}(\mathbf{r}) w_{\mathbf{R}'}^*(\mathbf{r}')$$

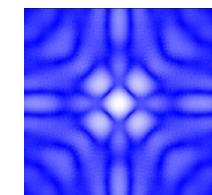
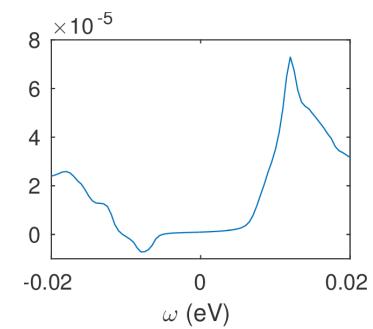
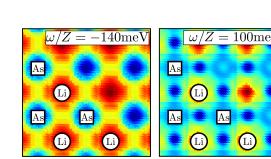
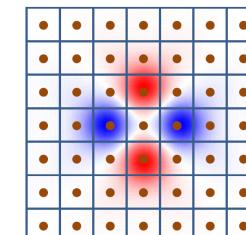
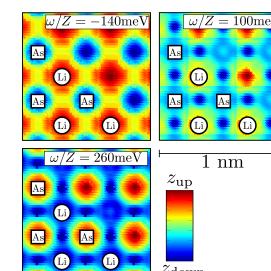


– LiFeAs: metal ion impurities

Shun Chi, (...), A. Kreisel, et al. arXiv:1607.03192

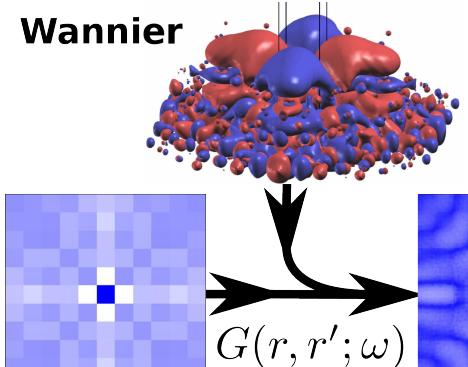
– BiSrCaCuO: Zn impurity, QPI

Kreisel *et al.* PRL 114, 217002 (2015)

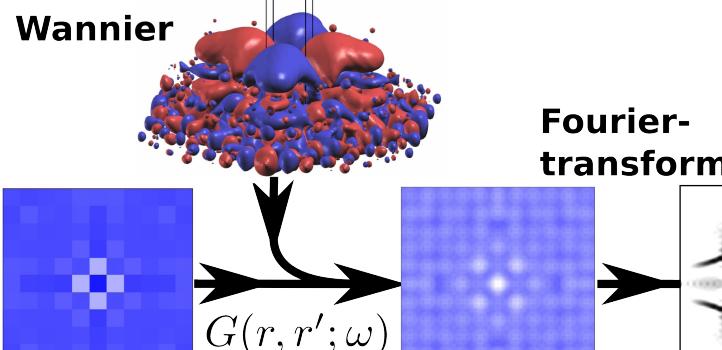


Summary

Kreisel *et al.*
PRL 114, 217002 (2015)



BdG+W experiment

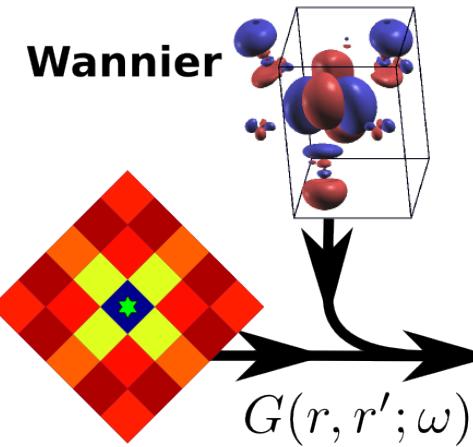


Fourier-transform

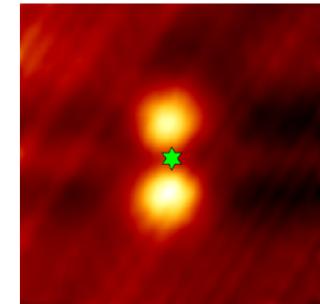
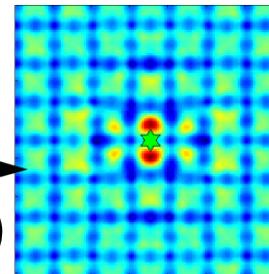
BdG+W

BSCCO
QPI: weak scatterer

experiment



$$G(r, r'; \omega)$$



arXiv:1607.03192

Acknowledgements



University of
St Andrews

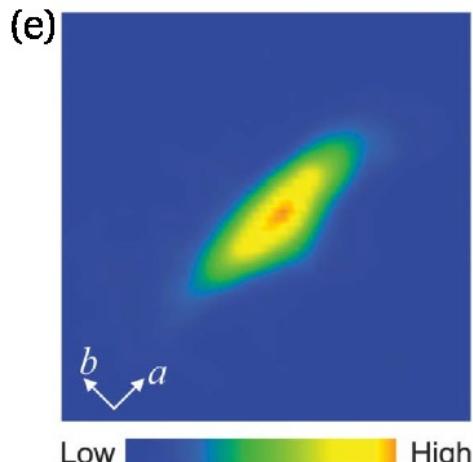
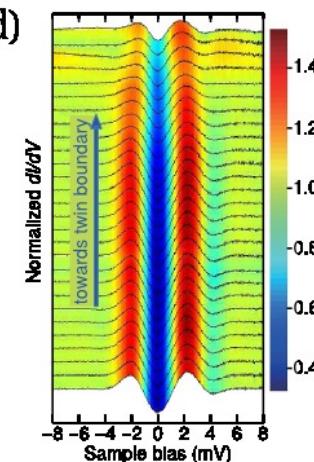
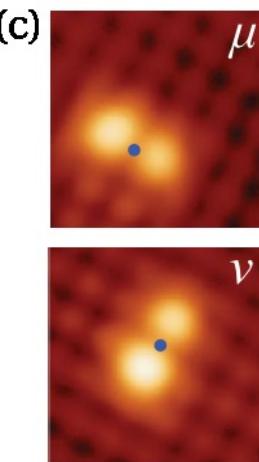
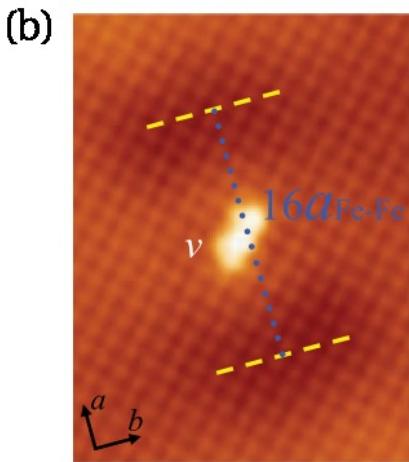
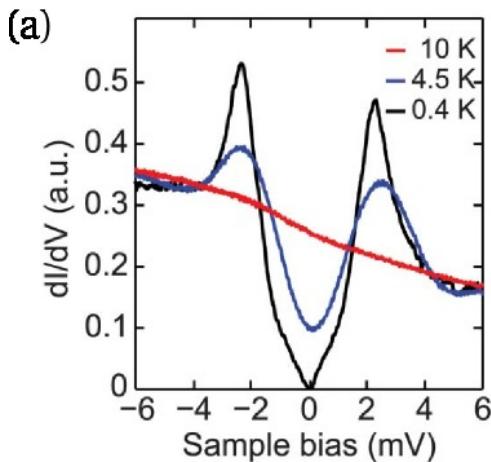
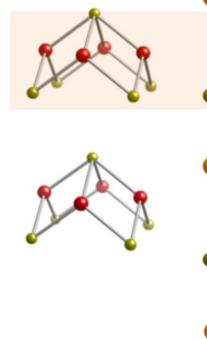
THE UNIVERSITY OF BRITISH COLUMBIA



FeSe: simplest crystal structure

- Tc 8K, under pressure ~40K Medvedev, et al. Nat. Mater. 8, 630 (2009)
- Tc 100K (single layer) Ge et al. Nat. Mater. 14, 285 (2015)
- nematic phase Baek, et al. Nat. Mat. 14, 210 (2015)
no magnetism
- consequences: nodal gapstructure, anisotropy

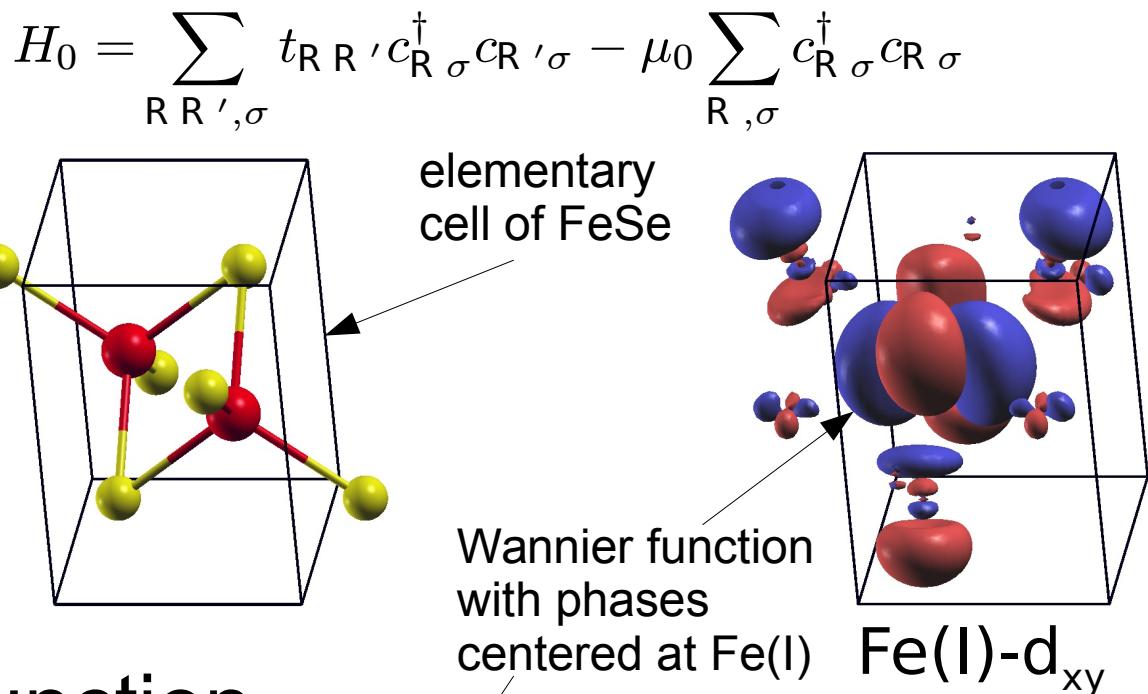
11
FeSe



BdG+Wannier method

- first principles calculation

- band structure
- Wannier functions
wavefunctions in
real space



- continuum Green function

$$G(r, r'; \omega) = \sum_{R, R'} G(R, R'; \omega) w_R(r) w_{R'}^*(r')$$

continuum position

nonlocal contributions

lattice Green function

local density of states (LDOS)

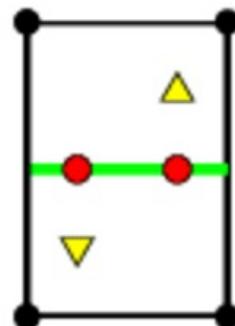
$$\rho(r, \omega) \equiv -\frac{1}{\pi} \text{Im} G(r, r; \omega)$$

FeSe: BdG+W Results

$$I(V, x, y, z) = -\frac{4\pi e}{\hbar} \rho_t(0) |M|^2 \int_0^{eV} \rho(x, y, z, \epsilon) d\epsilon$$

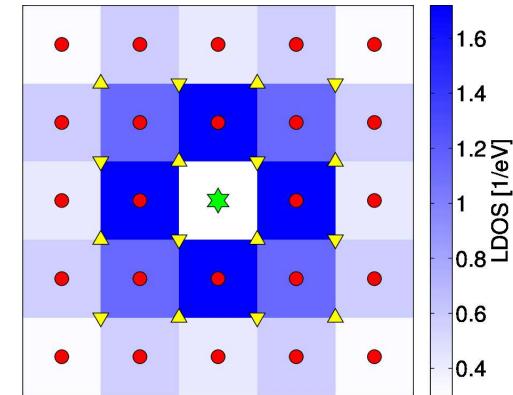
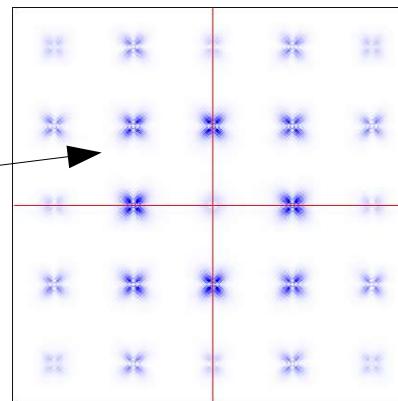
- continuum density of states

- at Fe plane

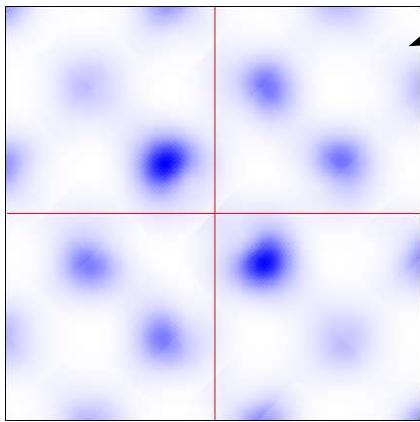
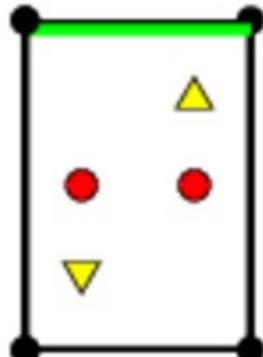


C4 symmetry!

2 meV

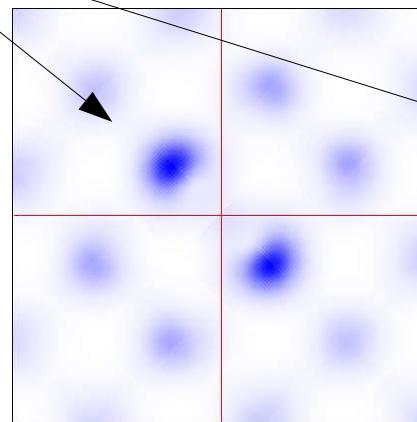
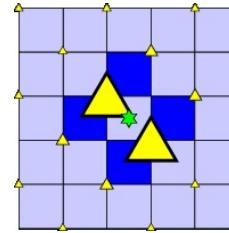


- at STM tip position

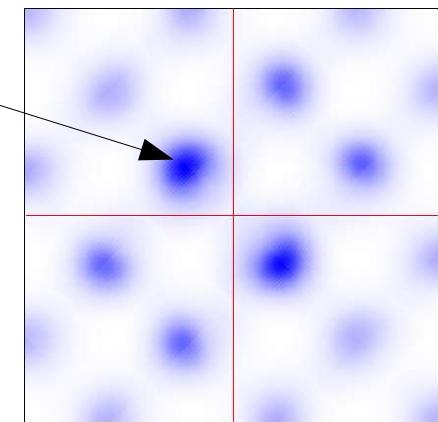


-2 meV

C2 symmetry!



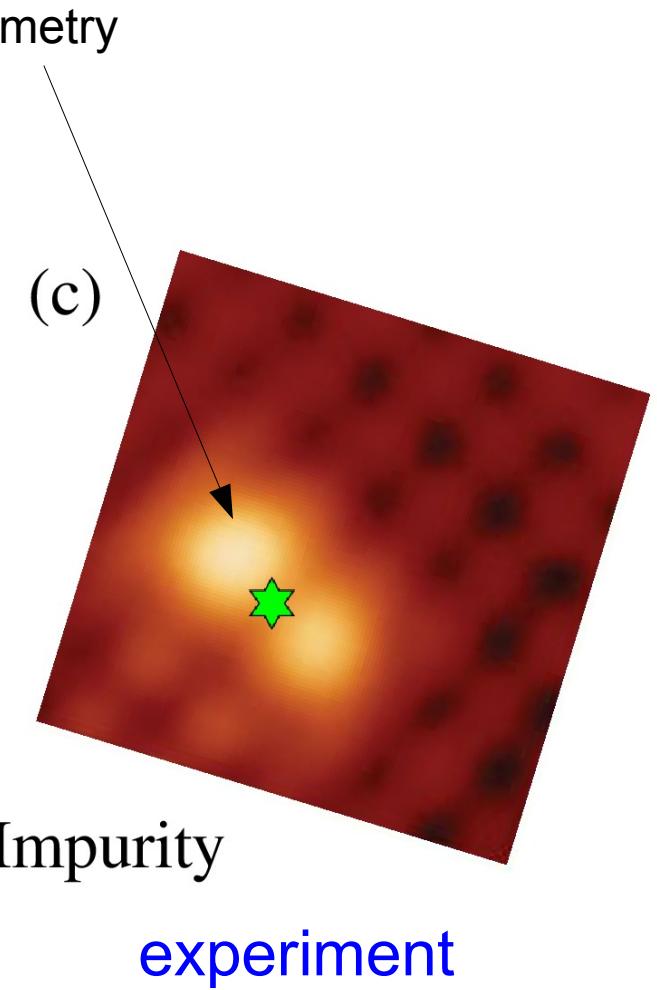
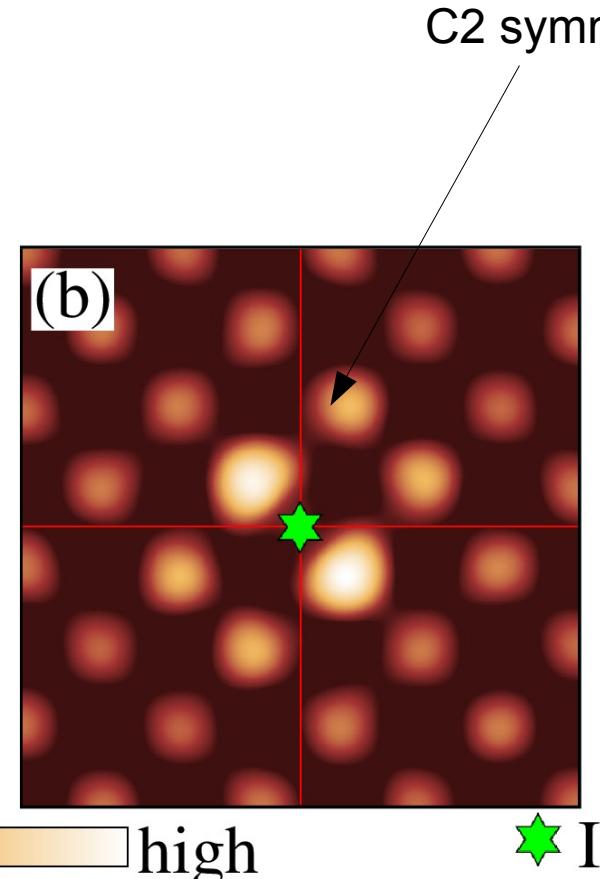
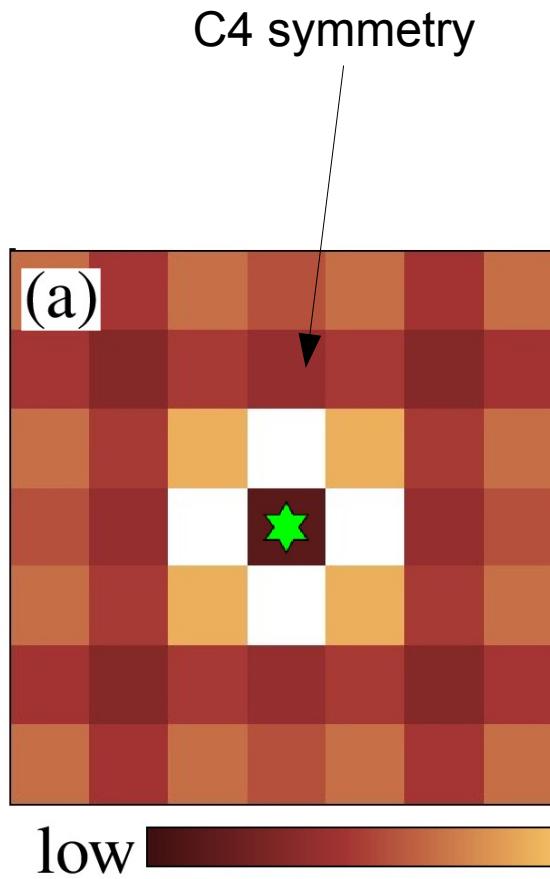
+2 meV



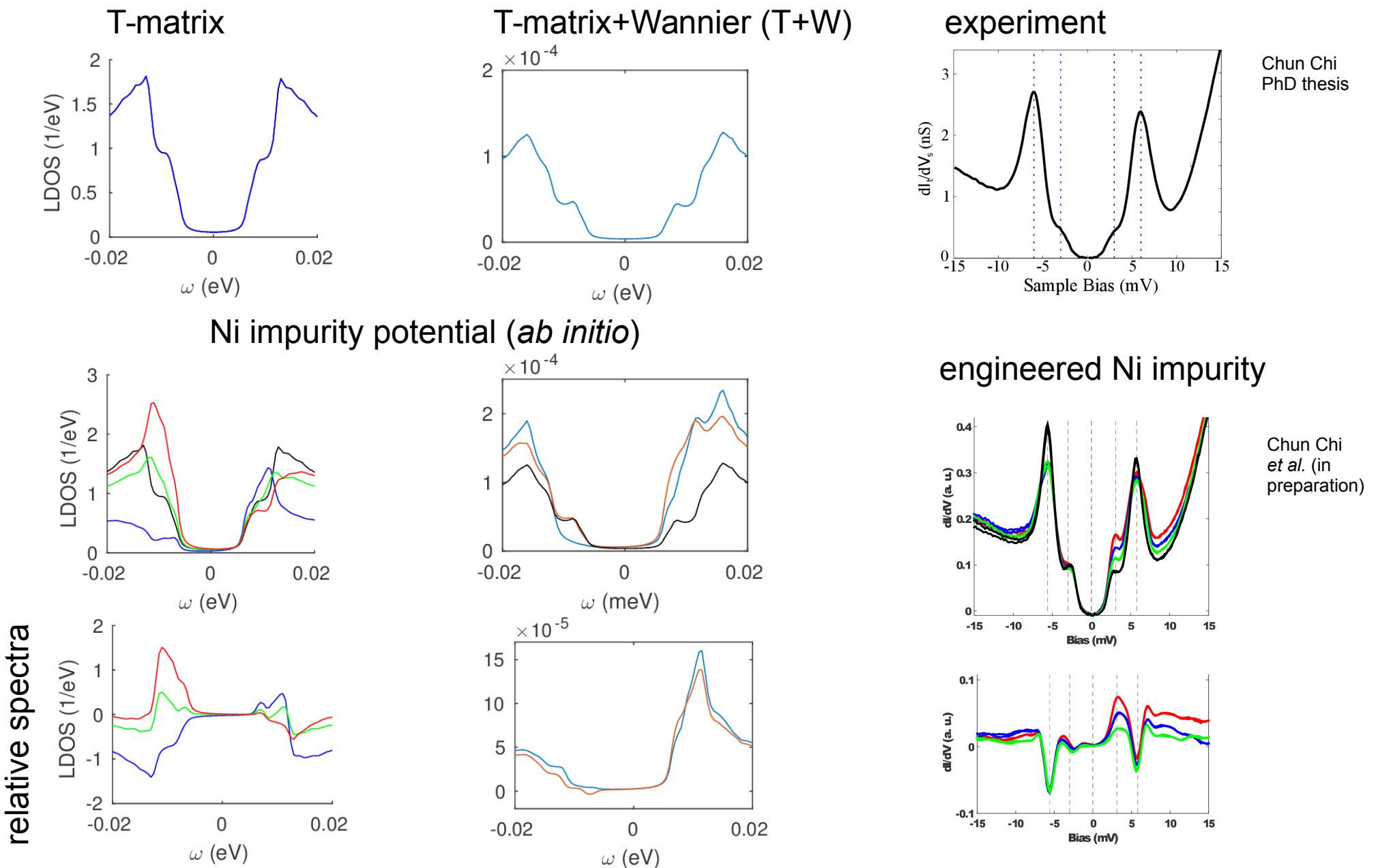
+30 meV

FeSe: Comparison to experiment

STM topography on (bulk) FeSe with Fe-centered impurity

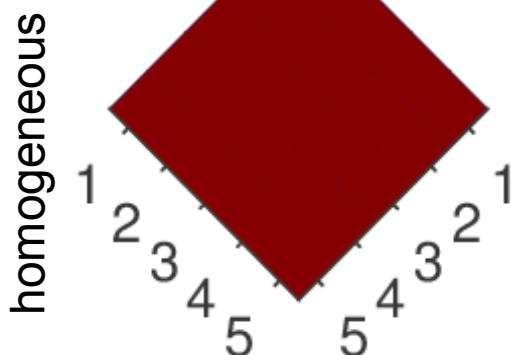


LiFeAs: spectra

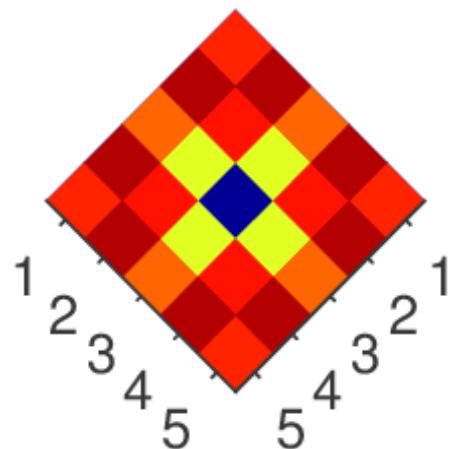


LiFeAs: spatial conductance

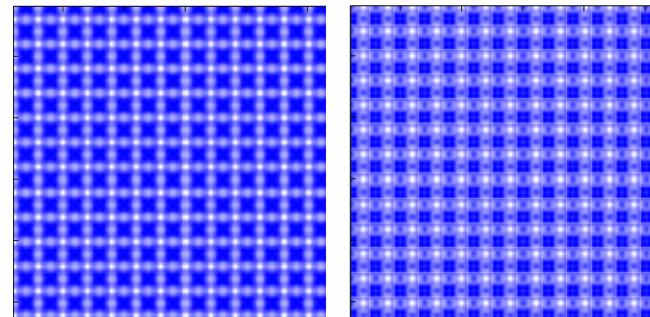
T-matrix



one pixel per Fe atom!

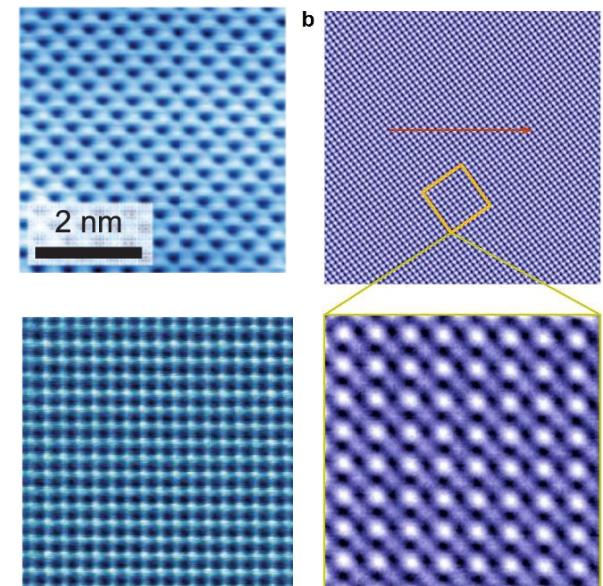


T-matrix+Wannier (T+W)

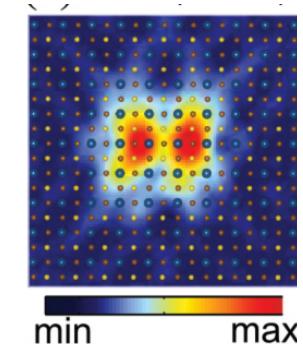


As lattice Li lattice

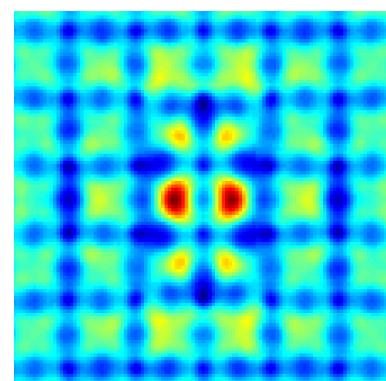
experiment (Li \leftrightarrow As lattice ?)



engineered
Ni defect



Ni impurity:
C2 symmetric impurity state



Shun Chi, et al., PRL 109, 087002 (2012)
T. Hanaguri, et al. PRB 85, 214505 (2012)
S. Grothe, et al., PRB 86, 174503 (2012)
J.-X. Yin, et al., arXiv, 1602.04949 (2016)