

# Diagrams of states for single flexible-semiflexible multiblock copolymer chains in different selective solvents:

## Computer simulations by means of stochastic approximation Monte Carlo

Daria Maltseva<sup>1</sup>, Sergey Zablotskiy<sup>1</sup>, Julia Martemyanova<sup>1</sup>, Viktor Ivanov<sup>1</sup>,  
Timur Shakirov<sup>2</sup>, Wolfgang Paul<sup>2</sup>



<sup>1</sup>Lomonosov Moscow State University, Moscow, Russia

<sup>2</sup>Martin-Luther-University Halle-Wittenberg, Halle, Germany

E-mail: dv.maltseva@physics.msu.ru

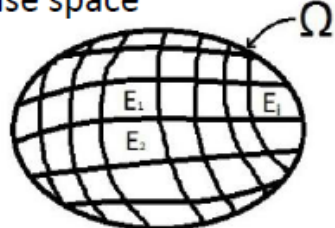


## Method of simulation

### Monte-Carlo algorithms of simulation

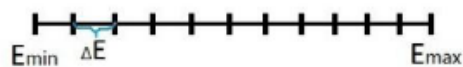
Phase space

Statistical sum of the chain



$$Z_{NVT} = \int_{\Omega} e^{-\frac{E\{\vec{x}\}}{kT}} d\Omega$$

$\{\vec{x}\}$  - conformation of the molecule



$g(E)$  - density of states function

$$\int_{E_{\min}}^{E_{\max}} g(E_i) \cdot e^{-\frac{E_i}{kT}}$$

### Advantages of SAMC

- 1) SAMC vs. Metropolis algorithm
- 2) The calculation with two-dimensional density of states function vs. the calculation with one-dimensional

We accumulate the estimation of density of states function with certain value of monomer-monomer interaction contributions and certain value of stiffness energy contribution

We can calculate average values of observables for any other meanings of parameter  $\epsilon_{st}$ :

$$\langle A(T) \rangle_{\epsilon_{st}, T=\text{fixed}} = \frac{1}{Z(T)} \sum_{E_{nv}, n_{st}} \bar{A}(E_{nv}, n_{st}) \cdot g(E_{nv}, n_{st}) \cdot e^{-\frac{E_{nv} + \epsilon_{st} n_{st}}{kT}}$$

$$\langle Z(T) \rangle_{\epsilon_{st}, T=\text{fixed}} = \sum_{E_{nv}, n_{st}} g(E_{nv}, n_{st}) \cdot e^{-\frac{E_{nv} + \epsilon_{st} n_{st}}{kT}}$$

### Stochastic Approximation Monte-Carlo (SAMC)

Visits all available **macrostates** uniformly

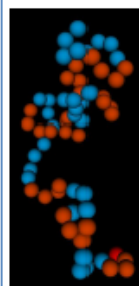
$\{E_{nv}, n_{st}\}$  - **macrostate**, conformation - **microstate**

$$p(\text{old} \rightarrow \text{trial}) = \min \left\{ 1, \frac{g(E_{nv}^{(\text{old})}, n_{st}^{(\text{old})})}{g(E_{nv}^{(\text{trial})}, n_{st}^{(\text{trial})})} \right\} \quad \text{- acceptance probability}$$

$$\ln g(E_{nv}, n_{st}) = \ln g(E_{nv}, n_{st}) + \gamma_t \cdot \left[ \delta_{(E_{nv}, n_{st}), (E_{nv}^{(\text{new})}, n_{st}^{(\text{new})})} - \frac{1}{n} \right],$$

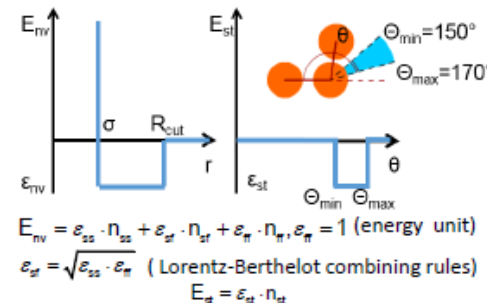
$$\gamma_t = \gamma_0 \cdot \min \left( 1, \frac{t_0}{t} \right)$$

### Model of investigated system



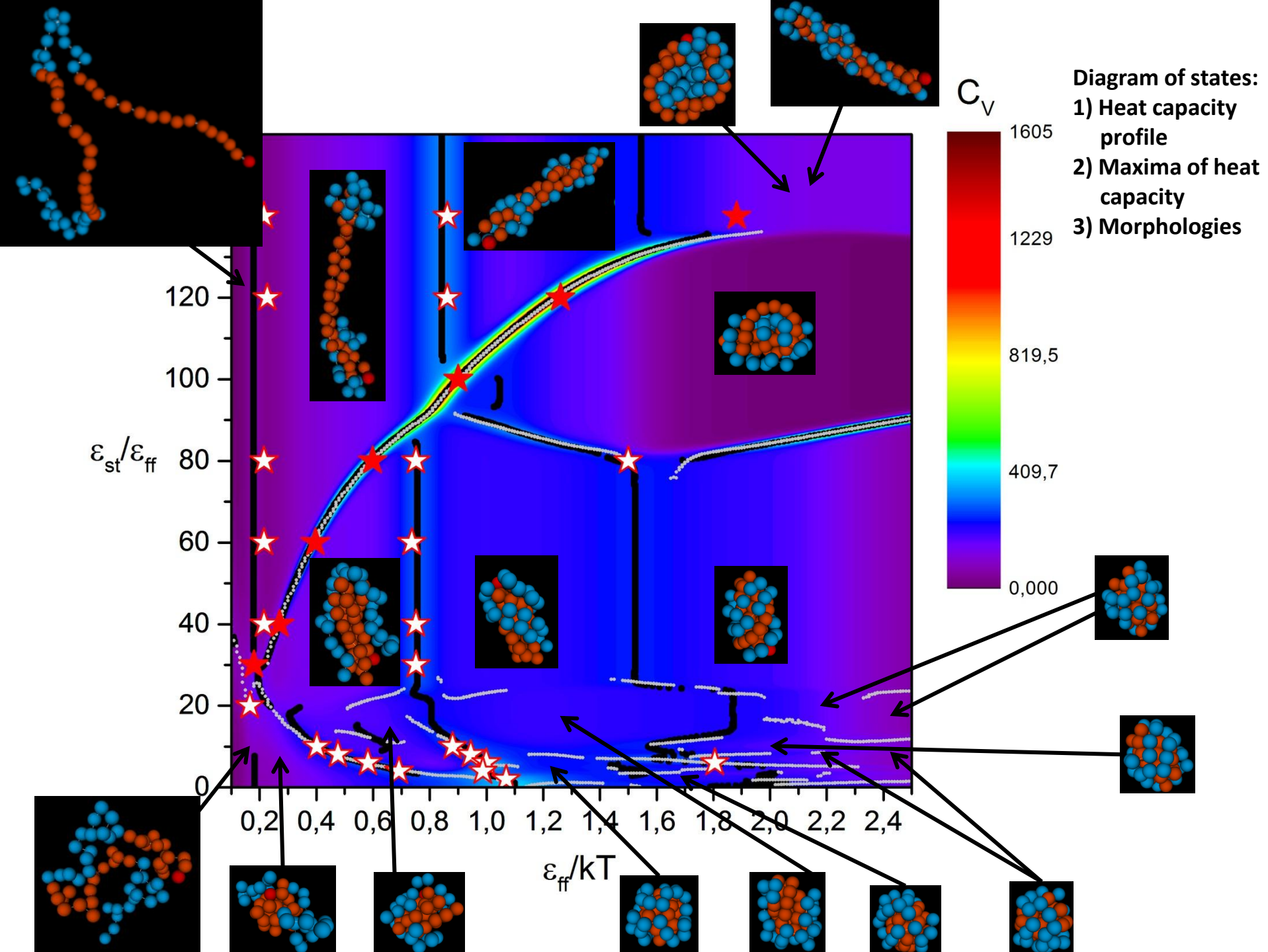
Snapshot for a chain with  $b=4$

- Continuum bead-on-spring
- $N=64$  monomers
- Bond length  $0.8\sigma-1.25\sigma$
- Flexible and semi-flexible blocks
- Block length  $b=4, 8, 16$  or  $32$
- Two energy contributions:
  - 1) non-valence interactions  $E_{nv}$
  - 2) energy of stiffness  $E_{st}$  (for semi-flexible blocks)
- $R_{\text{cut}}=1.5\sigma$
- Two types of trial moves: local moves and "end-cut-and-regrow"



Full configurational energy of a chain:

$$E = \epsilon_{rr} \cdot n_{rr} + \epsilon_{ss} \cdot n_{ss} + \sqrt{\epsilon_{ss} \cdot \epsilon_{rr}} \cdot n_{sr} + \epsilon_{st} \cdot n_{st}$$



## Analysis in microcanonical ensemble:

1) On a base of  $g(U)$

2) On a base of  $g(E)$

