

# Irreversible worm dynamics

On critical speeding-up for Ising models in high dimensions

# Collaboration

- Jens Grimm (Monash, )
- Lijie Ding (USTC, )
- Abraham Nasrawi (Monash, )
- Timothy Garoni (Monash, )
- Youjin Deng (USTC, )

## Lifted Worm Algorithm for the Ising Model

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We design an irreversible worm algorithm for the zero-field ferromagnetic Ising model by using the lifting technique. We study the dynamic critical behavior of an energy estimator on both the complete graph and toroidal grids, and compare our findings with reversible algorithms such as the Prokof'ev-Svistunov worm algorithm. Our results show that the lifted worm algorithm improves the dynamic exponent of the energy estimator on the complete graph, and leads to a significant constant improvement on toroidal grids.

### I. INTRODUCTION

Markov-chain Monte Carlo (MCMC) algorithms are a powerful and widely-used tool in various areas of physics and other disciplines, such as in machine learning [1] and statistics [2]. In many practical applications MCMC algorithms are constructed via the Metropolis [3] or heat bath update scheme [4]. Such algorithms are necessarily *reversible*.

One important example of a Metropolis algorithm is the Prokof'ev-Svistunov Worm Algorithm (P-S worm algorithm) which has widespread application for both classical and quantum systems [5, 6]. As opposed to cluster algorithms like the Wolff [7] or Swendsen-Wang algorithm [8], the updates of the worm algorithm are purely *local*. On the simple-cubic lattice with periodic bound-

For the Ising model on the complete graph, it was numerically observed that the lifted single-spin flip Metropolis algorithm improves the scaling (with volume) of the rate of decay of the autocorrelation function of the magnetization [14]. Another study [13] proved that a lifted MCMC algorithm for uniformly sampling leaves from a given tree reduces the mixing time. In other examples [16, 20, 22] it was numerically observed that lifting speeds up reversible MCMC algorithms by a possibly large constant factor but does not asymptotically affect the scaling with the system size.

In this work we investigate how lifting affects worm algorithms. More precisely, we design a lifted worm algorithm for the zero-field ferromagnetic Ising model, and numerically study the dynamic critical behavior of an estimator of the energy. Our simulations were performed on both the complete graph and toroidal grids in dimen-

<https://arxiv.org/abs/1711.05346>

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1. From reversible to irreversible worm through Berretti-Sokal & lifting.
2. Non-extensive (*but impressive*) efficiency boosts on high-dim. tori.
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# P-S reversible worm

$$Z = 2^{|V|} \cosh^{|E|}(\beta) \sum_{\omega \in \mathcal{C}_0} \tanh^{|\omega|}(\beta)$$

---

## Algorithm 1 P-S Worm Algorithm

---

**if**  $\omega \in \mathcal{C}_0$  **then**

    Choose a uniformly random vertex  $x$

**else**

    Choose a uniformly random odd vertex  $x$

**end if**

    Choose a uniformly random edge  $xx'$  among the set of edges incident to  $x$ . With probability  $a_{\text{P-S}}(\omega, \omega \Delta xx')$ , let  $\omega \rightarrow \omega \Delta xx'$ . Otherwise  $\omega \rightarrow \omega$

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# P-S reversible worm

- Applications in classical and quantum systems

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# P-S reversible worm

- Applications in classical and quantum systems
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- Rapidly mixing on any graph and any temperature.

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# P-S reversible worm

- Applications in classical and quantum systems
  - Here classical Ising only
- Updates completely local.
- Rapidly mixing on any graph and any temperature.
- On  $Z^3_L$  more efficient than SW for measuring susceptibility & 2nd-moment correlation length.

$$Z = 2^{|V|} \cosh^{|E|}(\beta) \sum_{\omega \in \mathcal{C}_0} \tanh^{|\omega|}(\beta)$$

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## Algorithm 1 P-S Worm Algorithm

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**if**  $\omega \in \mathcal{C}_0$  **then**

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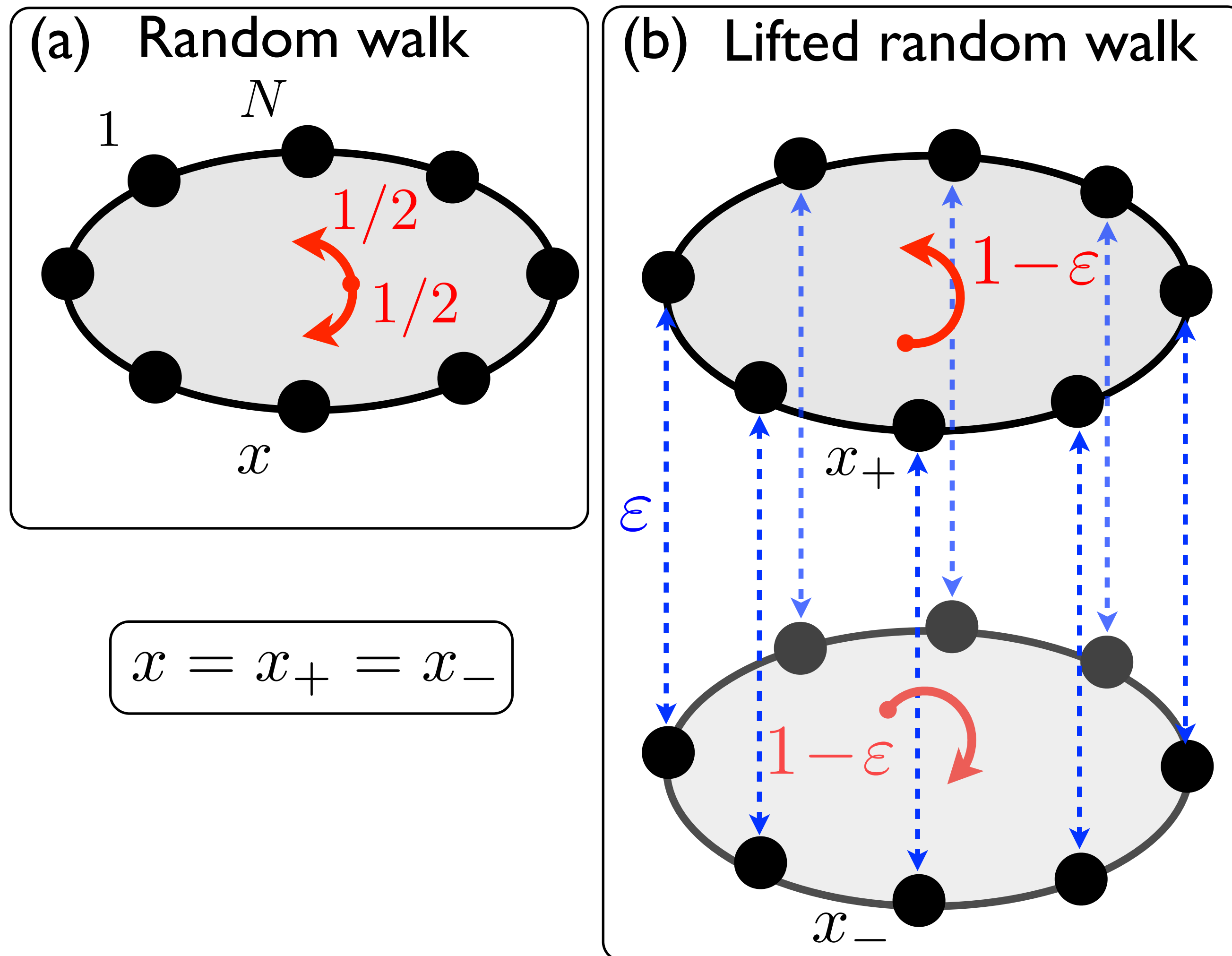
    Choose a uniformly random odd vertex  $x$

**end if**

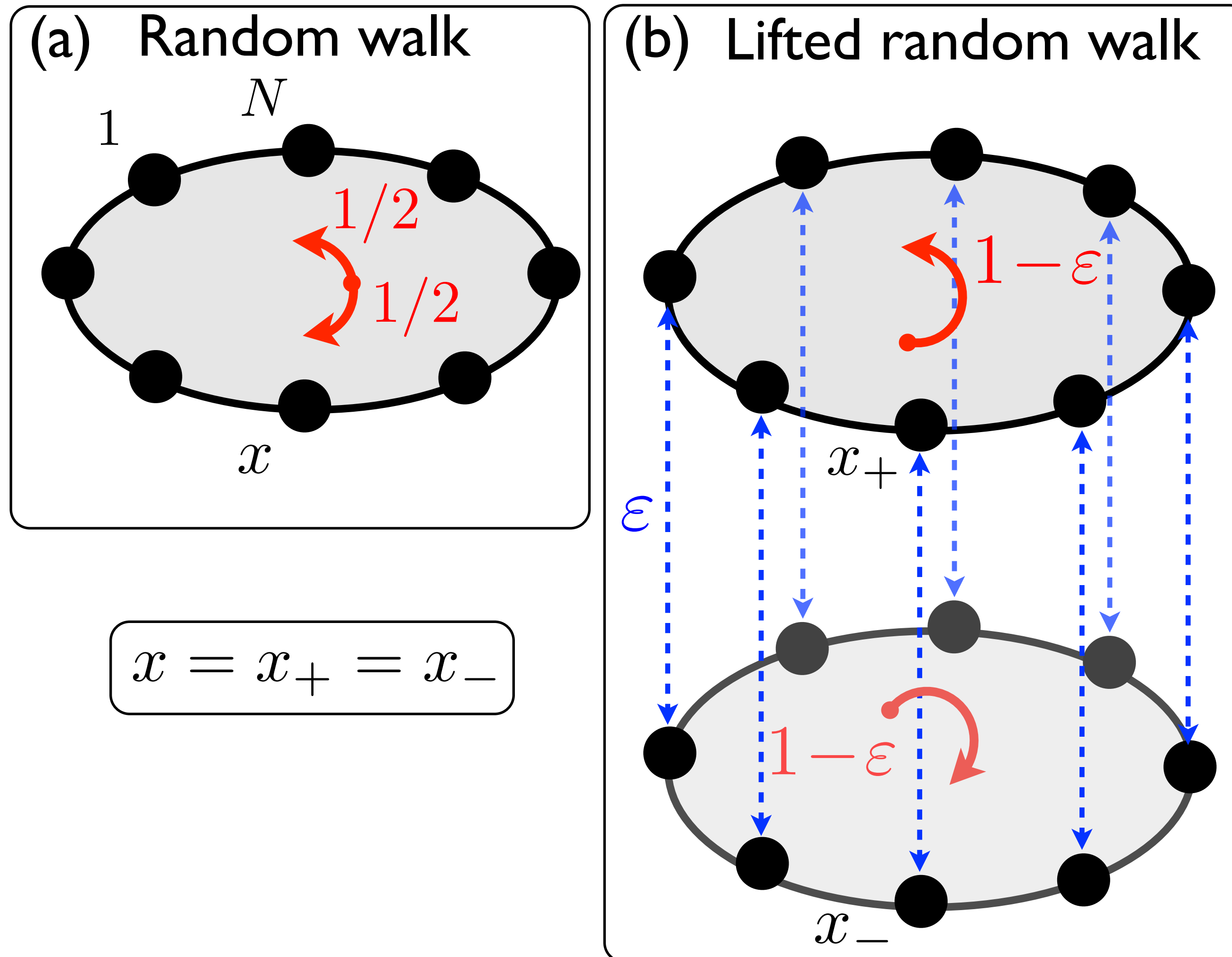
    Choose a uniformly random edge  $xx'$  among the set of edges incident to  $x$ . With probability  $a_{P-S}(\omega, \omega \Delta xx')$ , let  $\omega \rightarrow \omega \Delta xx'$ . Otherwise  $\omega \rightarrow \omega$

---

# The lifting technique

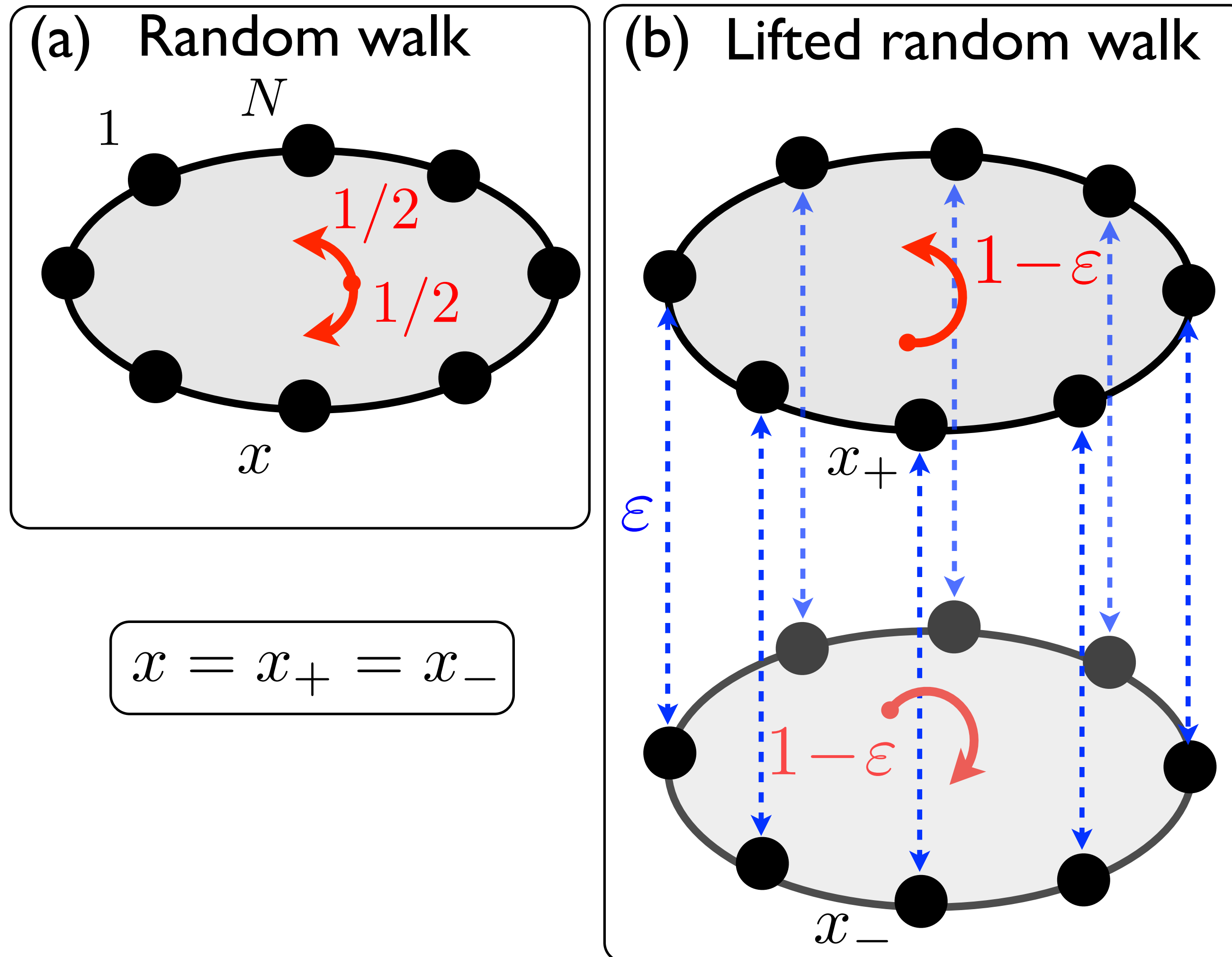


# The lifting technique



- For worm dynamics one expects energy (edge) diffusion to be the slowest “mode” (cause of slowing-down)

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- For worm dynamics one expects energy (edge) diffusion to be the slowest “mode” (cause of slowing-down)
- Idea: Apply lifting along “edge”-direction, but how?

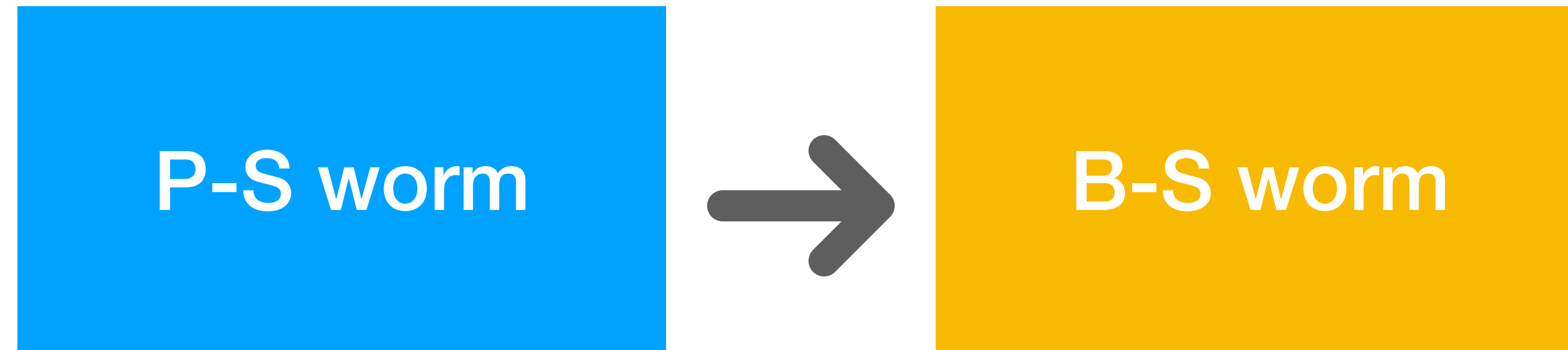


# Our construction

# Our construction

P-S worm

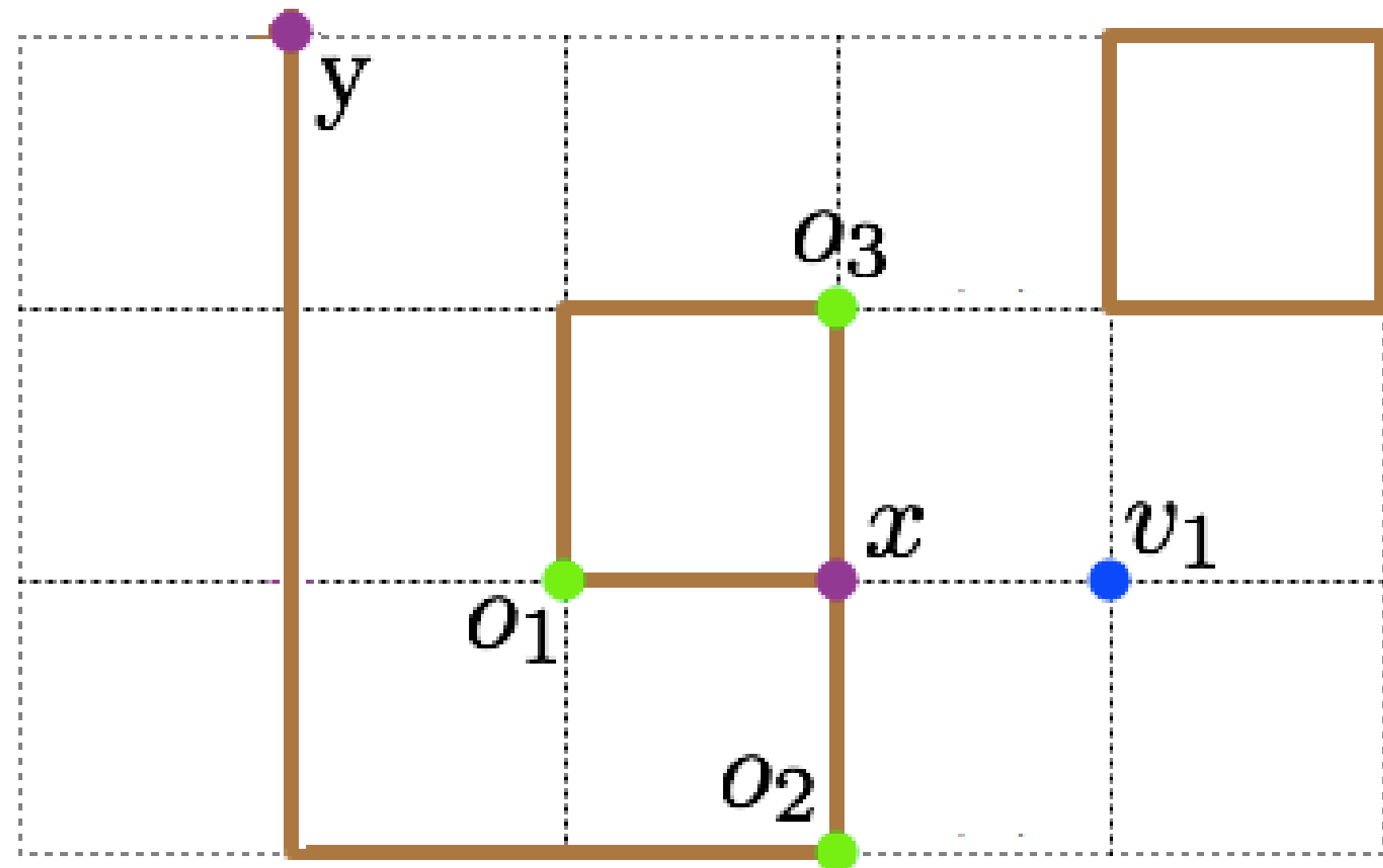
# Our construction



# Our construction



# The Berretti-Sokal worm



---

## Algorithm 2 B-S type Worm Algorithm

---

Choose  $\lambda = \{+, -\}$  uniformly at random

**if**  $\omega \in \mathcal{C}_0$  **then**

    Choose a uniformly random vertex  $x$

**else**

    Choose a uniformly random odd vertex  $x$

**end if**

**if**  $N_\omega(x, \lambda) = \emptyset$  **then**

    Set  $\omega \rightarrow \omega$  and skip all following steps

**else**

    Choose a uniformly random edge  $xx' \in N_\omega(x, \lambda)$ . With probability  $a_{\text{B-S}}(\omega, \omega \Delta xx')$ , let  $\omega \rightarrow \omega \Delta xx'$ . Otherwise  $\omega \rightarrow \omega$ .

**end if**

---

# Lifting the B-S worm

---

## Algorithm 2 B-S type Worm Algorithm

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**end if**

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# Lifting the B-S worm

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## Algorithm 2 B-S type Worm Algorithm

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if  $\omega \in \mathcal{C}_0$  then

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    Choose a uniformly random odd vertex  $x$

end if

if  $N_\omega(x, \lambda) = \emptyset$  then

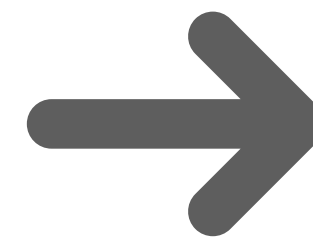
    Set  $\omega \rightarrow \omega$  and skip all following steps

else

    Choose a uniformly random edge  $xx' \in N_\omega(x, \lambda)$ . With probability  $a_{\text{B-S}}(\omega, \omega \Delta xx')$ , let  $\omega \rightarrow \omega \Delta xx'$ . Otherwise  $\omega \rightarrow \omega$ .

end if

---



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## Algorithm 3 Irreversible Worm Algorithm

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if  $\tilde{\omega} = (\omega, \lambda)$  where  $\omega \in \mathcal{C}_0$  then

    Choose a uniformly random vertex  $x$

else

    Choose a uniformly random odd vertex  $x$

end if

if  $N_\omega(x, \lambda) = \emptyset$  then

    Set  $(\omega, \lambda) \rightarrow (\omega, -\lambda)$  and skip all following steps

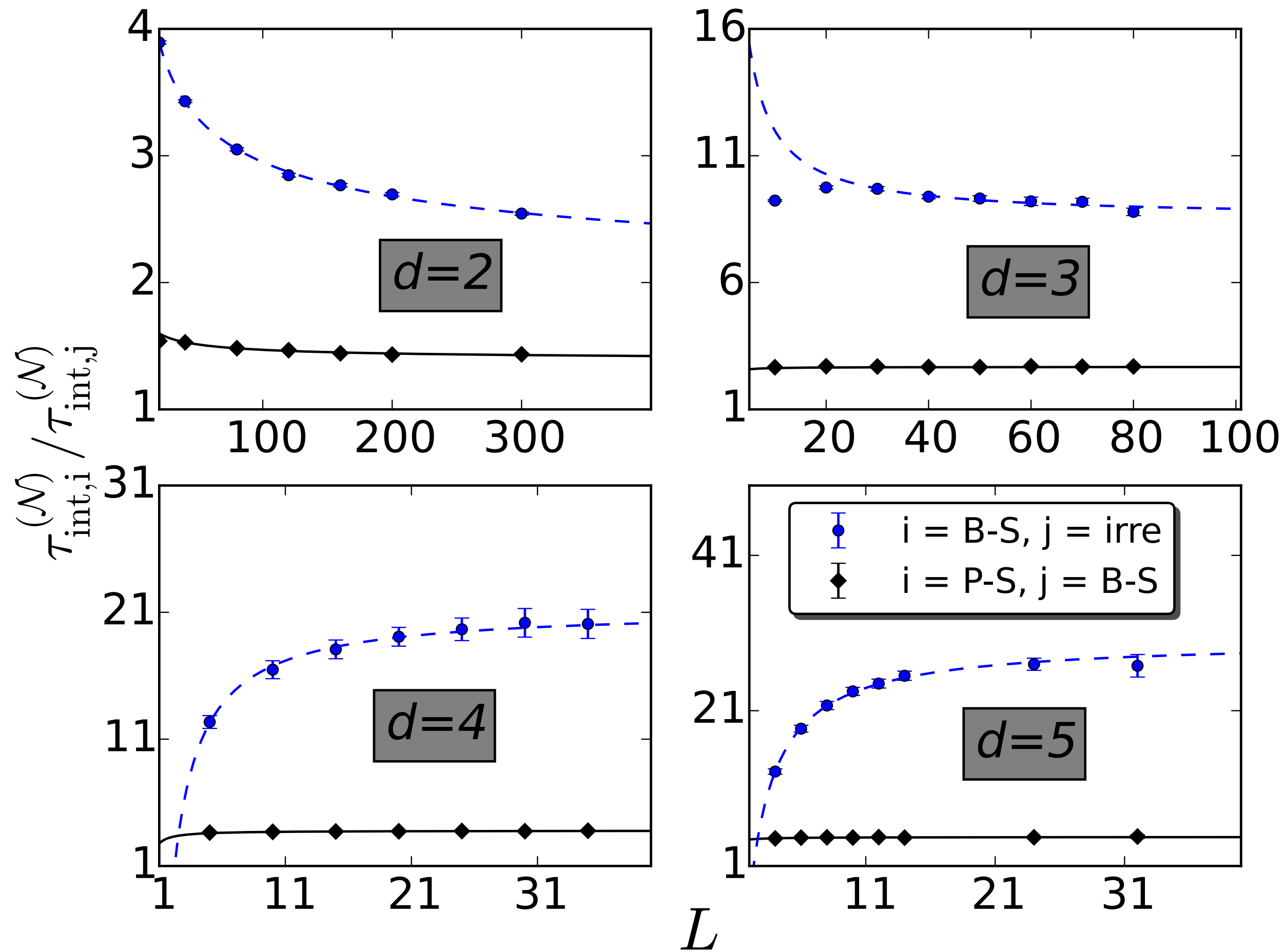
else

    Choose a uniform random edge  $xx' \in N_\omega(x, \lambda)$ . With probability  $a_{\text{B-S}}(\omega, \omega \Delta xx')$ , let  $(\omega, \lambda) \rightarrow (\omega \Delta xx', \lambda)$ . Otherwise  $(\omega, \lambda) \rightarrow (\omega, -\lambda)$

end if

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# Boosts on tori



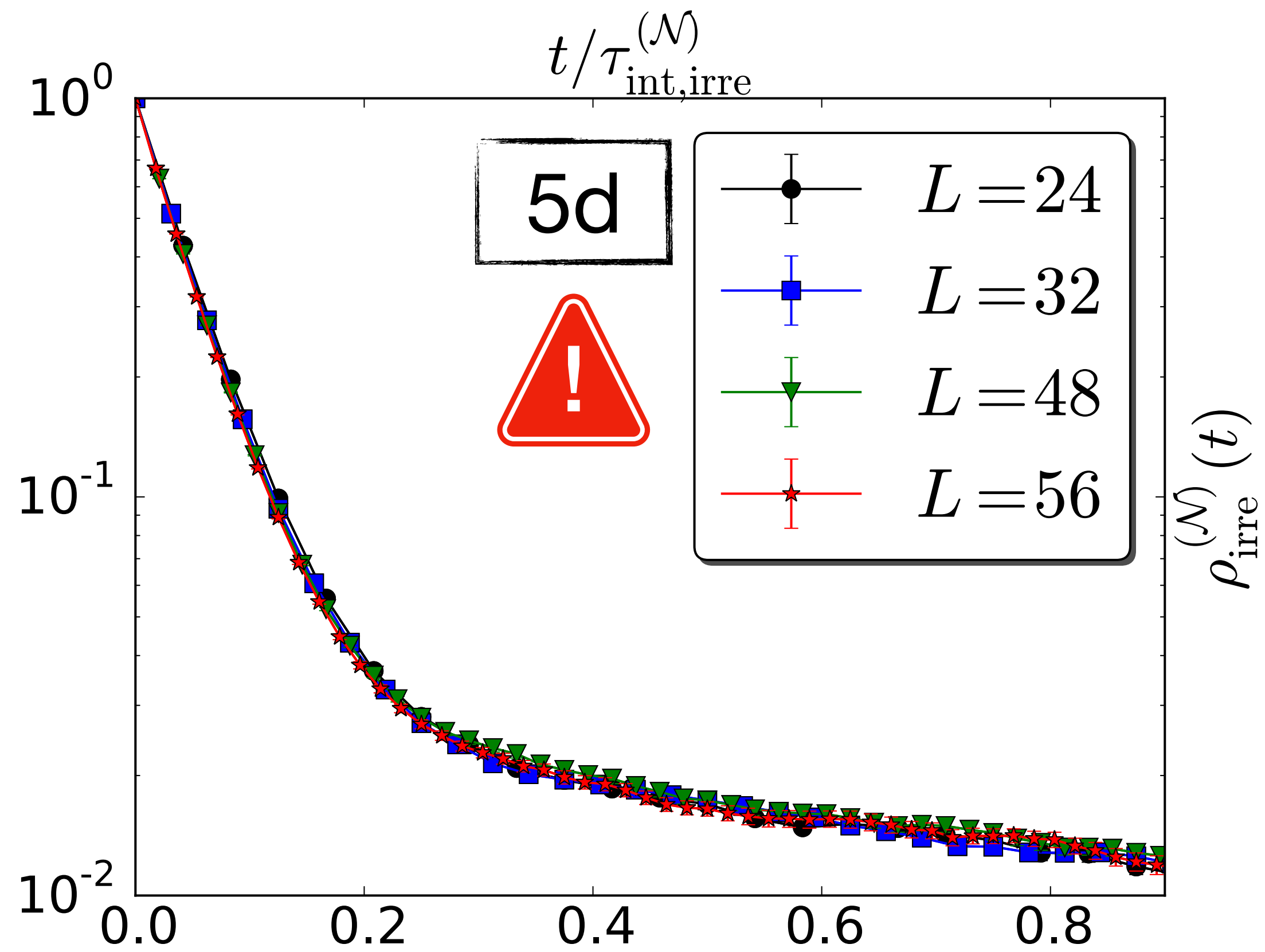
$$\text{Var}(\bar{\mathcal{N}}) \sim 2\tau_{\text{int}}^{(\mathcal{N})} \frac{\text{Var}(\mathcal{N}_0)}{M} \quad M \rightarrow \infty$$

$$\tau_{\text{int}}^{(\mathcal{N})} := \frac{1}{2} + \sum_{t=1}^{\infty} \rho^{(\mathcal{N})}(t).$$

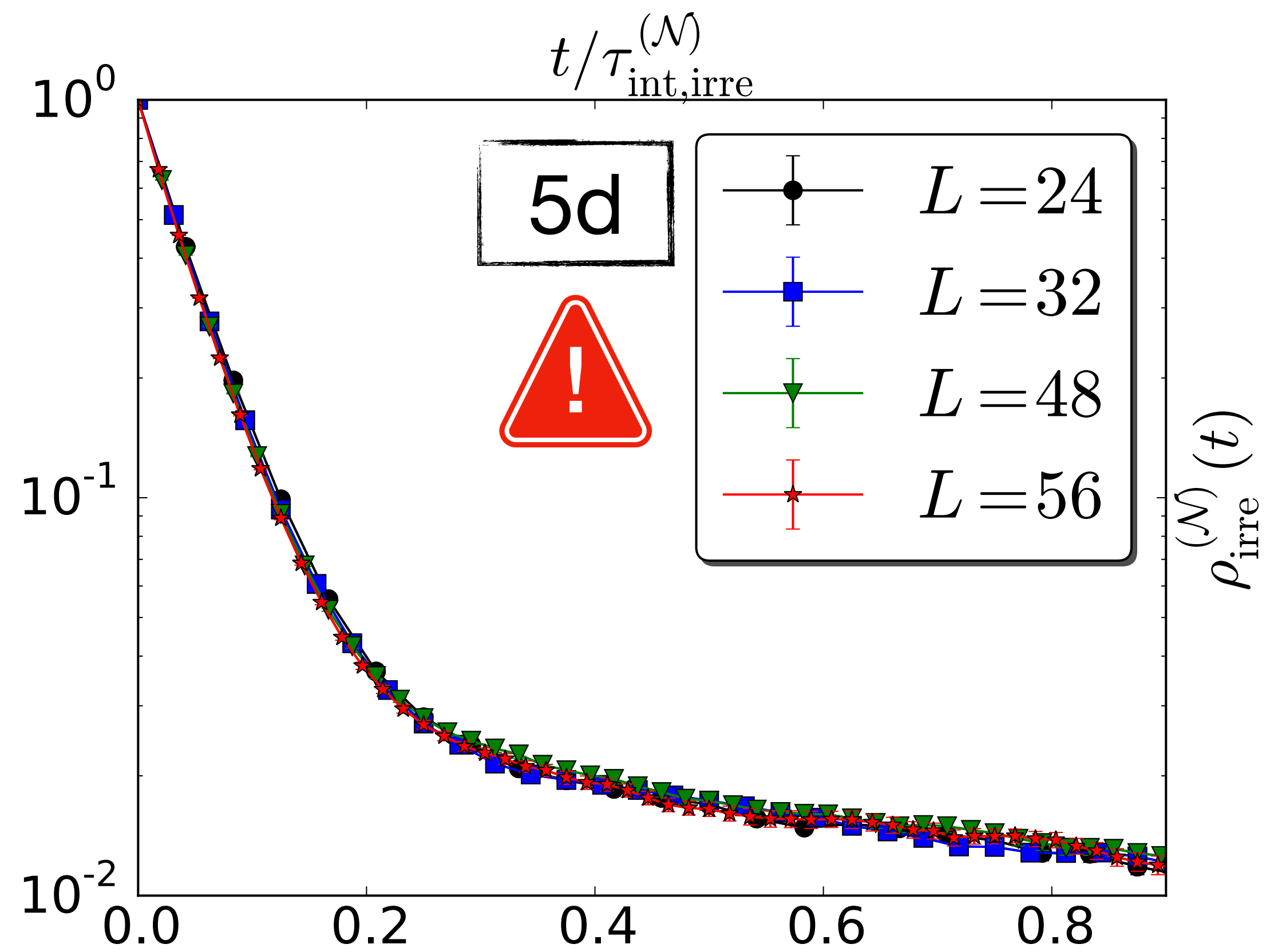
	P-S $\rightarrow$ B-S	B-S $\rightarrow$ irre	P-S $\rightarrow$ irre
$d = 2$	1.4(1)	1.7(2)	2.4(4)
$d = 3$	2.68(9)	8.2(4)	22(2)
$d = 4$	3.79(3)	24(1)	91(4)
$d = 5$	4.7(1)	30(1)	141(6)



# Slow suppressed modes

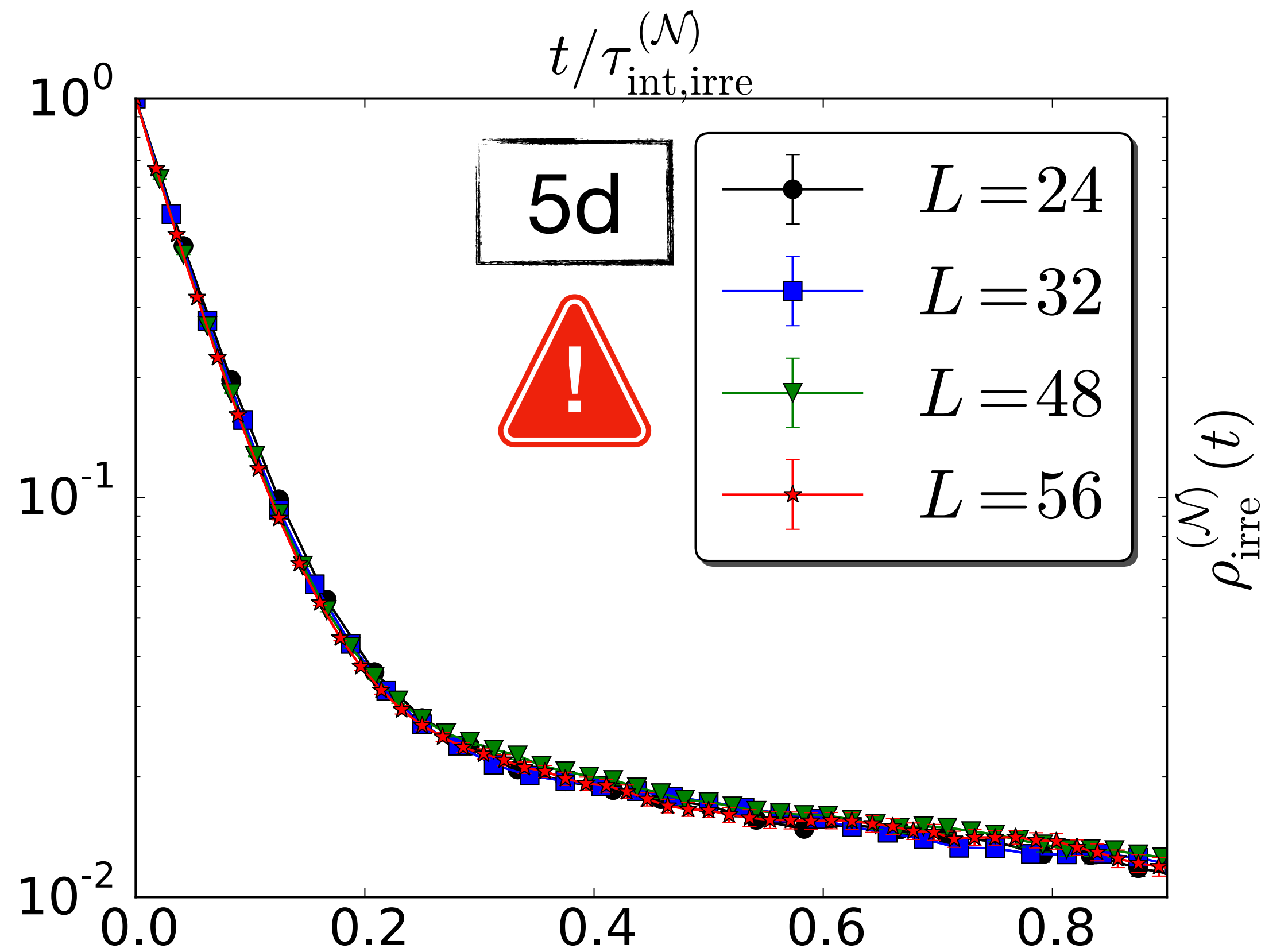


# Slow suppressed modes



Two-time-scale ansatz

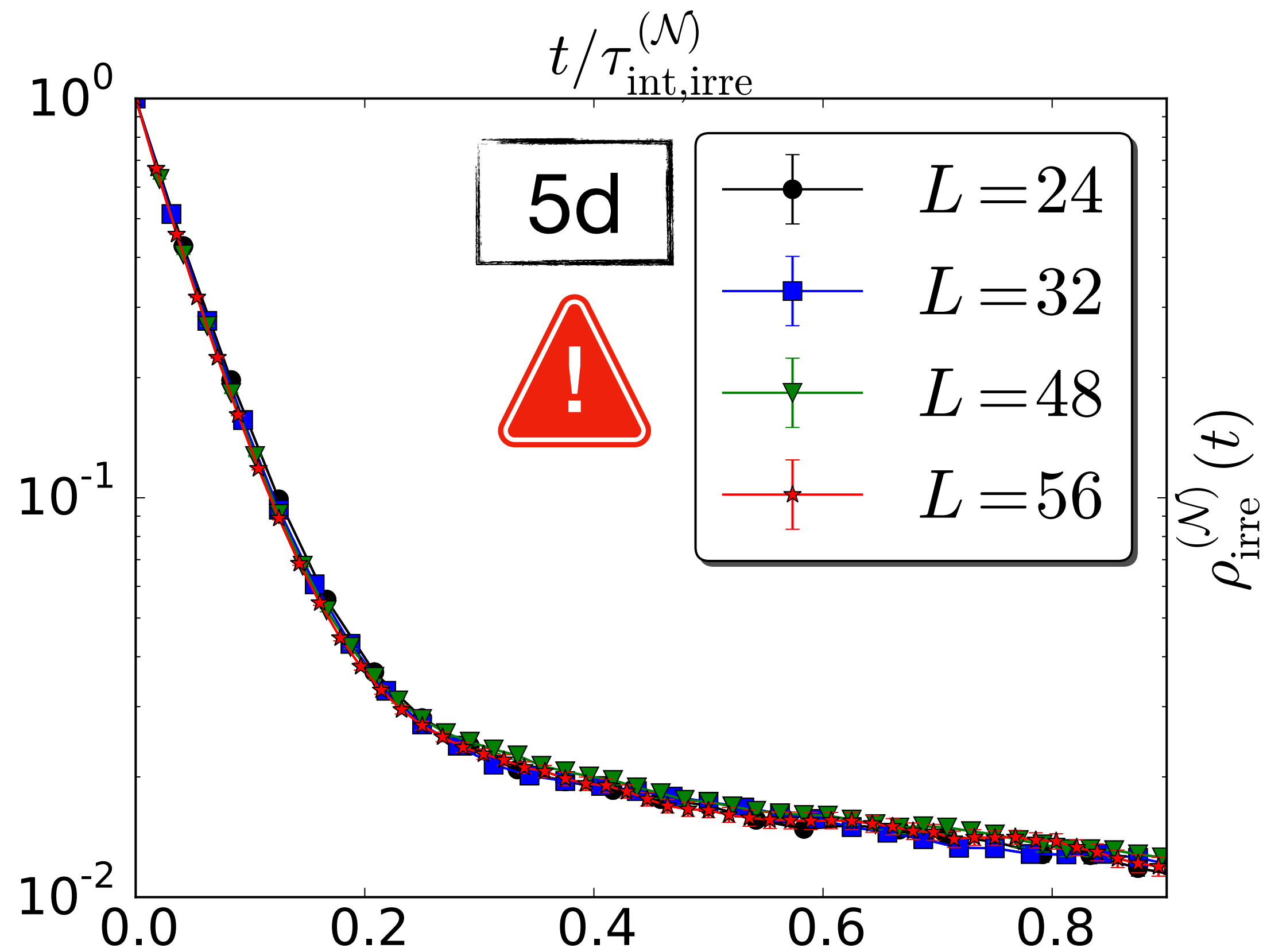
# Slow suppressed modes



Two-time-scale ansatz

$$\rho_{\text{irre}}^{(\mathcal{N})}(t) = \alpha_1 \exp(-t/\tau_1) + \alpha_2 \exp(-t/\tau_2)$$

# Slow suppressed modes

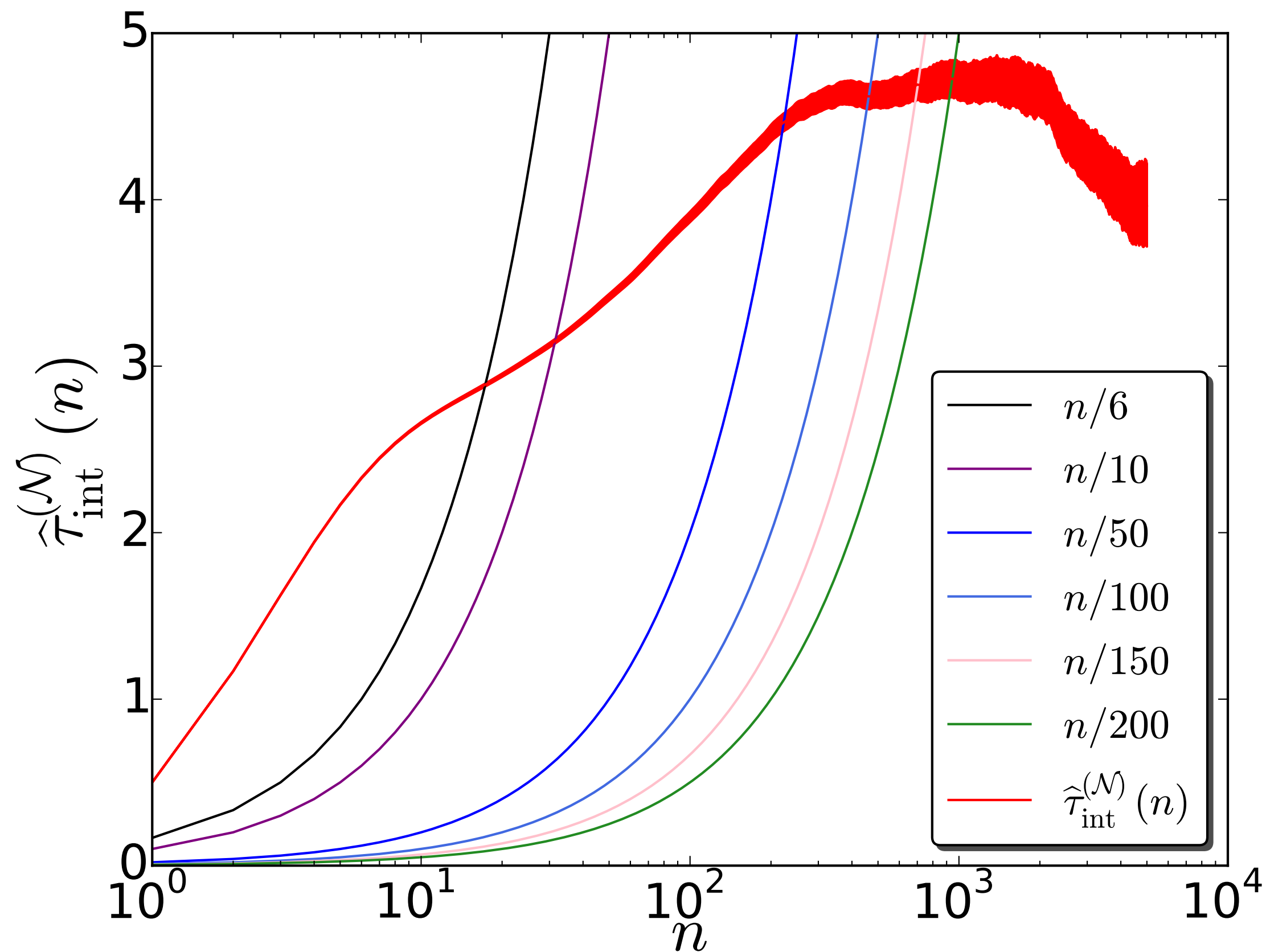


## Two-time-scale ansatz

$$\rho_{\text{irre}}^{(\mathcal{N})}(t) = \alpha_1 \exp(-t/\tau_1) + \alpha_2 \exp(-t/\tau_2)$$

$$\frac{\alpha_2}{\alpha_1} \stackrel{L \rightarrow \infty}{\sim} 40.6(4) \quad \& \quad \frac{\tau_1}{\tau_2} \stackrel{L \rightarrow \infty}{\sim} 32.4(6)$$

# Slow suppressed modes



$$\hat{\tau}_{\text{int}}^{(\mathcal{N})}(n) := \frac{1}{2} + \sum_{t=1}^n \rho^{(\mathcal{N})}(t).$$

## Two-time-scale ansatz

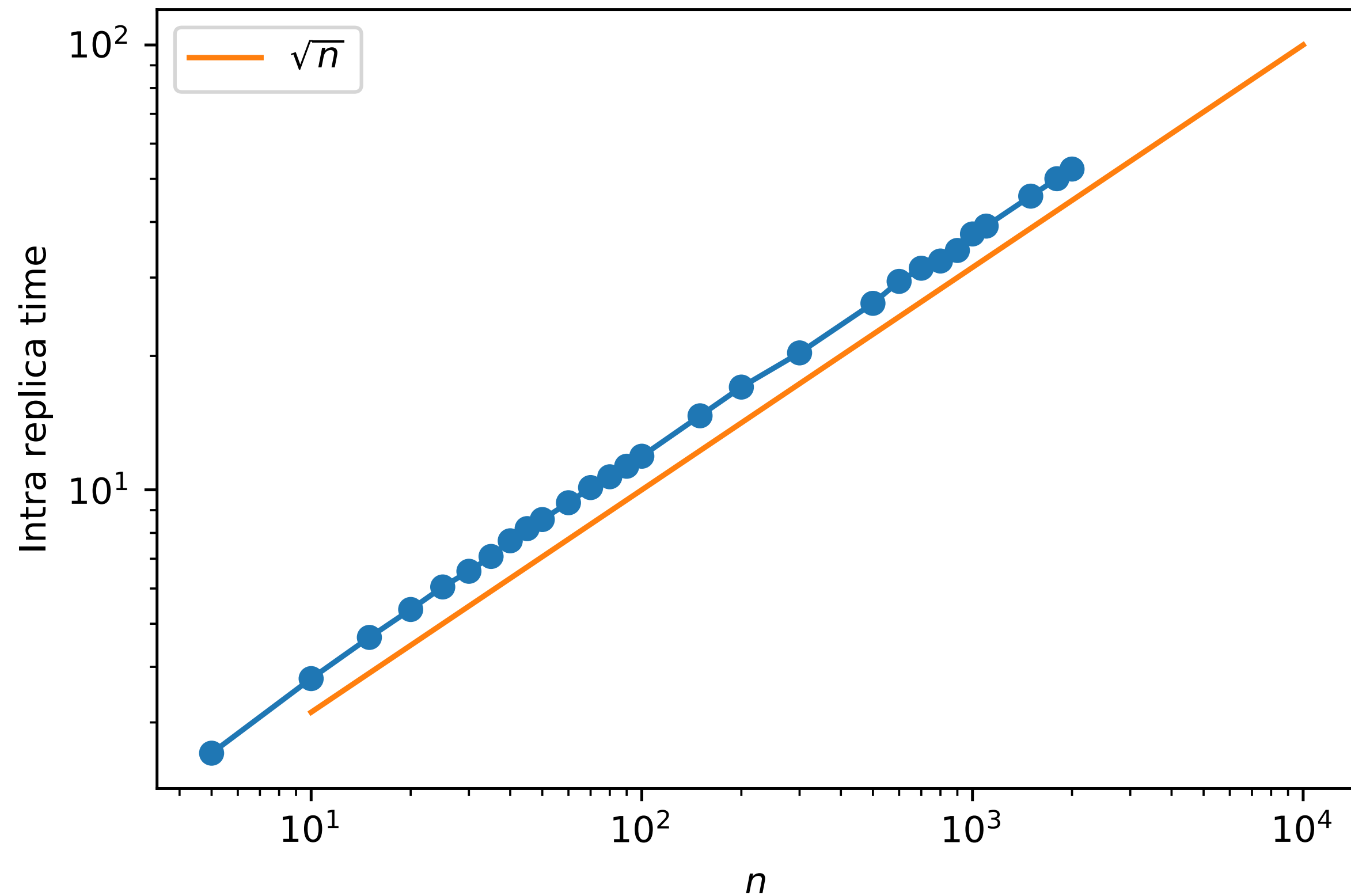
$$\rho_{\text{irre}}^{(\mathcal{N})}(t) = \alpha_1 \exp(-t/\tau_1) + \alpha_2 \exp(-t/\tau_2)$$

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# Critical speeding-up in the mean-field limit

# Extensive ballistic drift

## Some scaling results

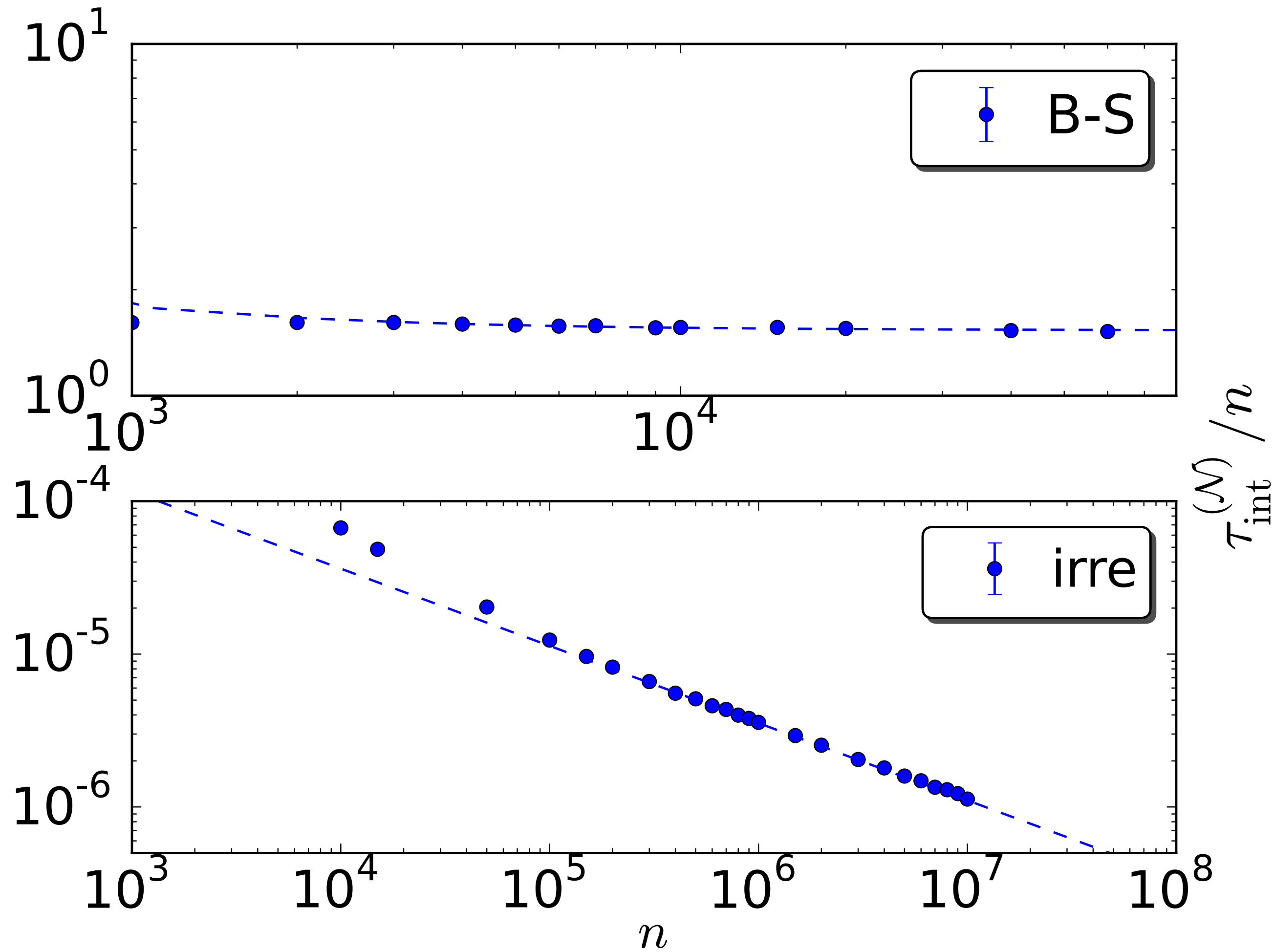


$$\beta = \frac{1}{n} \Rightarrow x \sim \frac{1}{n}$$

$$\mu_N^{(\mathcal{E}_0)} \sim \frac{\Gamma(7/4)}{\sqrt{3}\Gamma(5/4)} n^{\frac{1}{2}} \approx 0.5854143 n^{\frac{1}{2}}$$

$$\sigma_N \sim \sqrt{\frac{9}{4} - \frac{24\Gamma(5/4)^4}{\pi^2}} n^{\frac{1}{2}} \approx 0.7801732 n^{\frac{1}{2}}$$

# Critical speeding-up



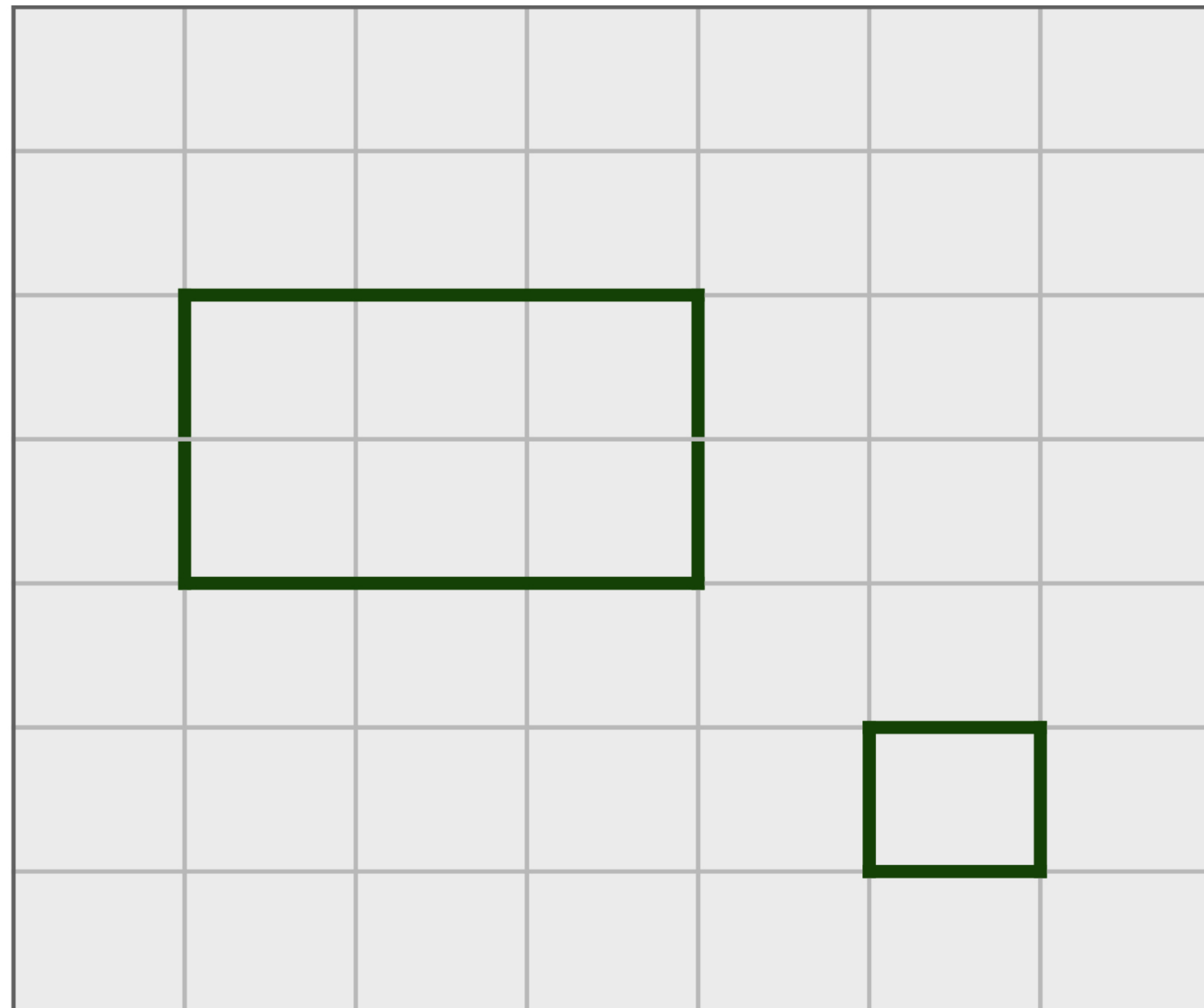
$$\tau_{\text{int}, \mathcal{N}}^{(\text{BS})} \approx n$$

$$\tau_{\text{int}, \mathcal{N}}^{(\text{irr})} \approx n^{\frac{1}{2}}$$

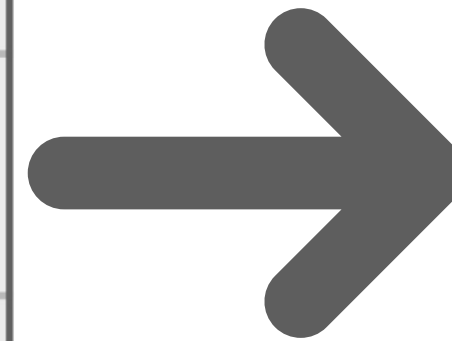


# An irreversible cluster algorithm

Eulerian loop configuration



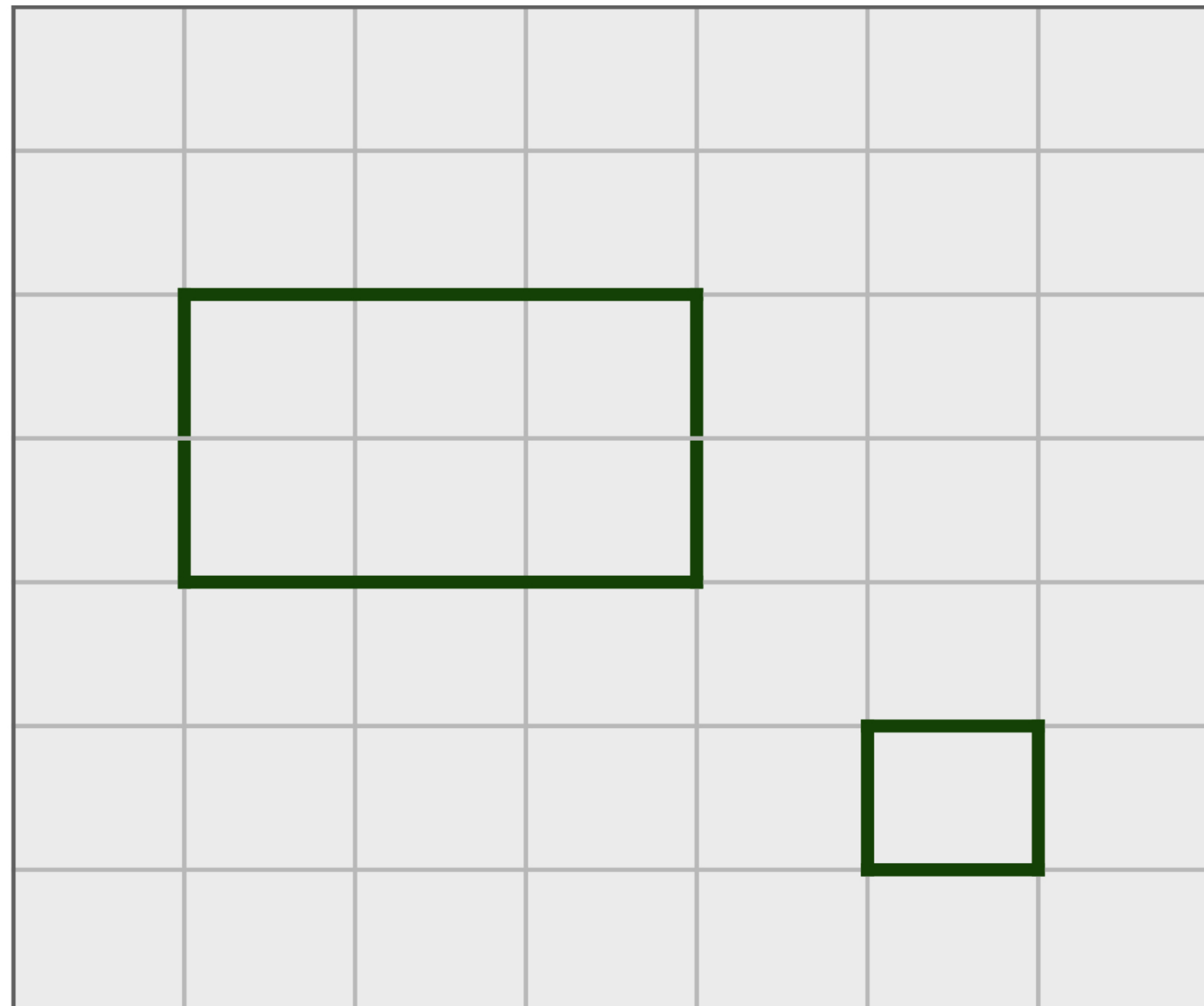
Add vacant  
edges  
with prob.  $x$



$$x \in [0,1]$$

# An irreversible cluster algorithm

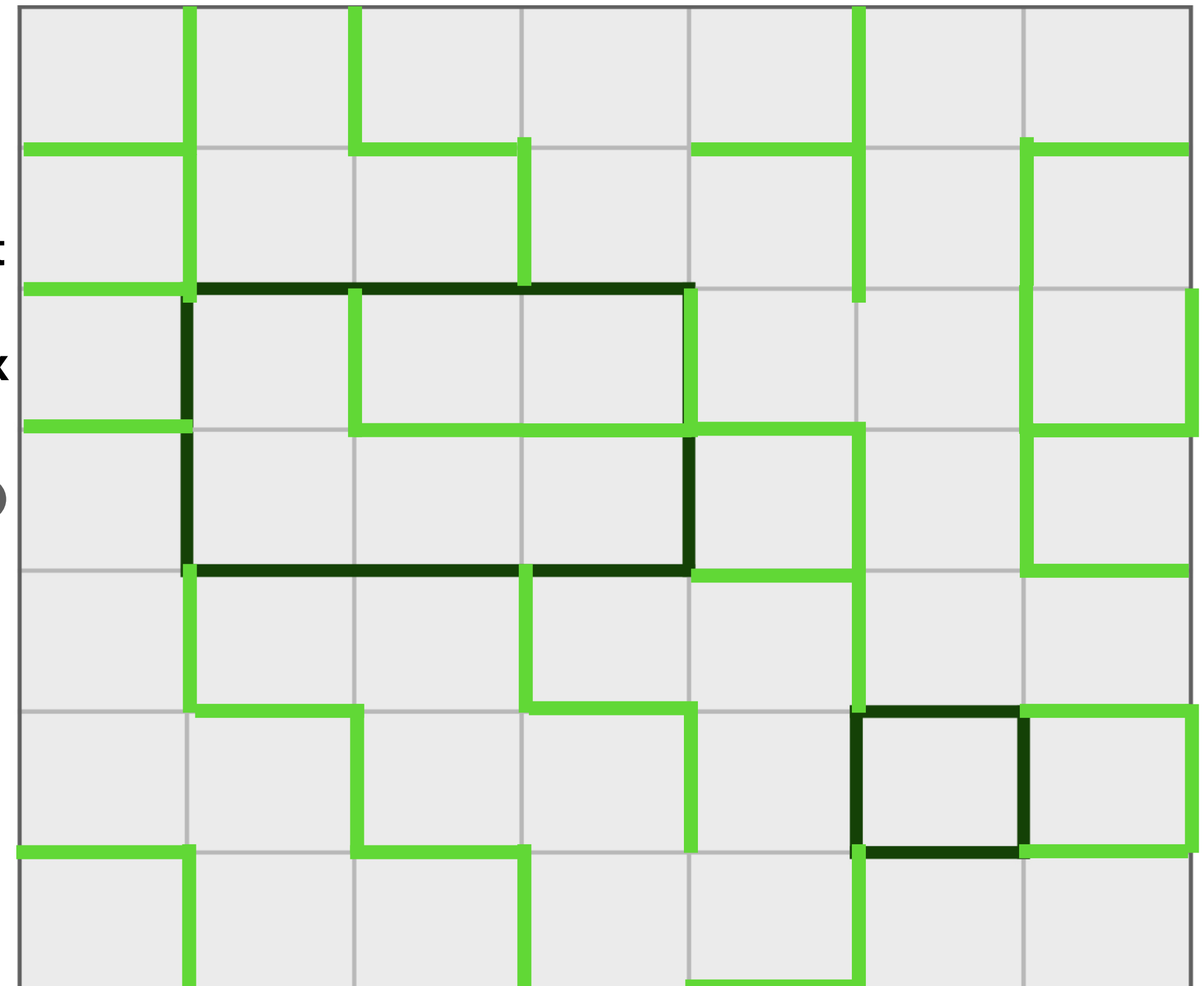
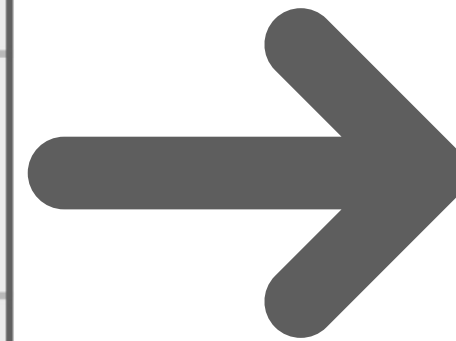
Eulerian loop configuration



$$x \in [0,1]$$

FK Ising configuration

Add vacant  
edges  
with prob.  $x$

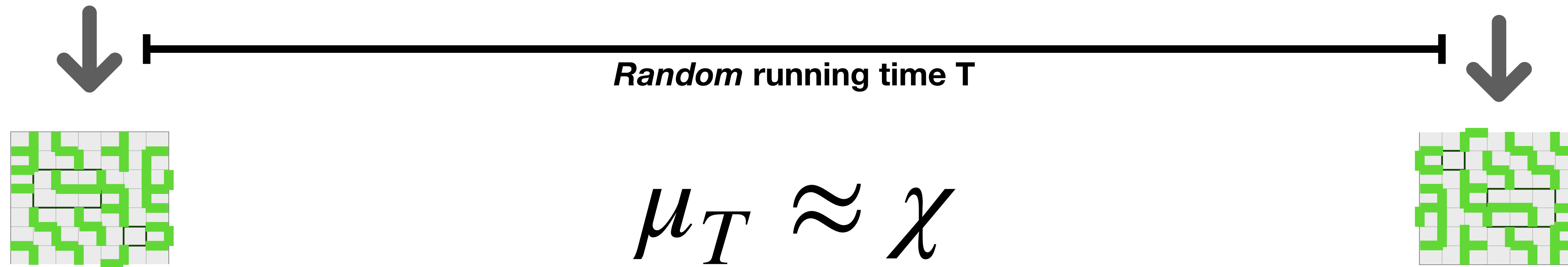


$$p = \frac{2x}{1+x}$$

$$q = 2$$

# An irreversible cluster algorithm

$$\omega_1 \in \mathcal{C}_0 \rightarrow \omega_2 \in \mathcal{C}_2 \rightarrow \omega_3 \in \mathcal{C}_2 \rightarrow \dots \rightarrow \omega_T \in \mathcal{C}_2 \rightarrow \omega_{T+1} \in \mathcal{C}_0$$

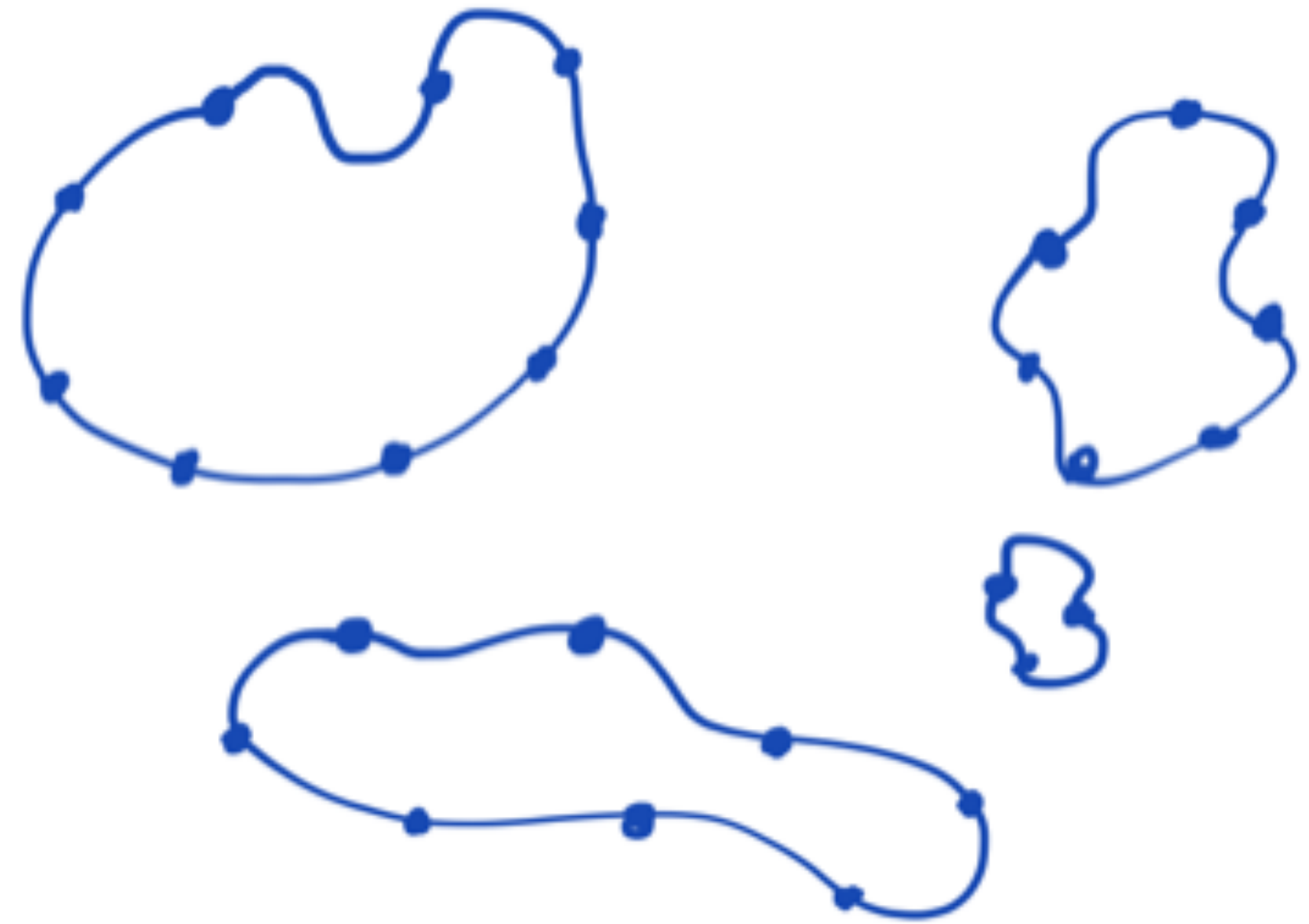


## Preliminary analysis / picture on complete graph

- Scaling  $\chi \approx n^{\frac{1}{2}}$
- Mean of T is dominant time scale (deviations that dominate mean asymptotically are at least exp. unlikely)
- Induced FK cluster algorithm has integrated autocorrelation for edges (energy) that is bounded as n becomes large.
- Compare: SW is expected to have mixing time  $n^{\frac{1}{4}}$

# Ballistic path creation

- Some text



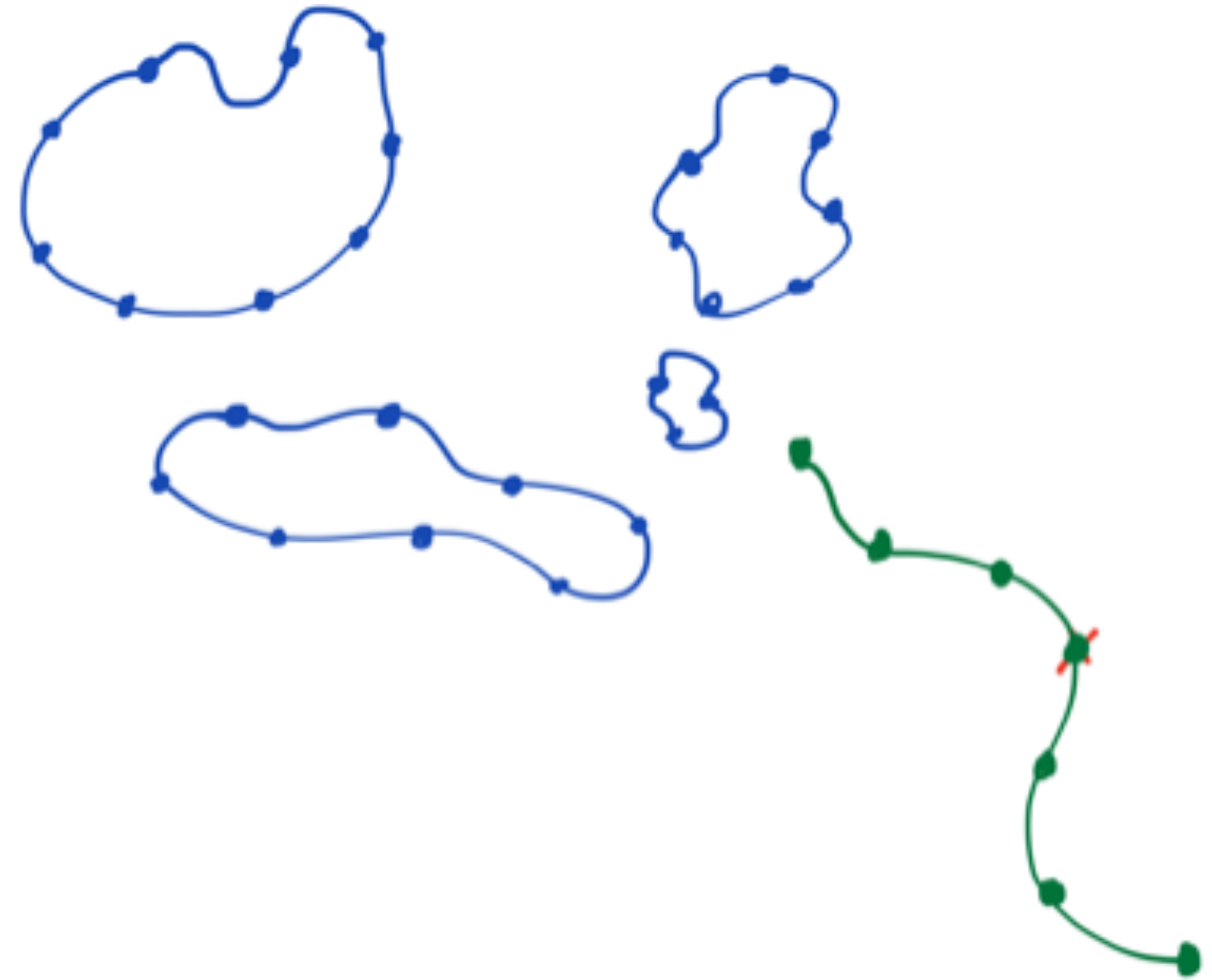
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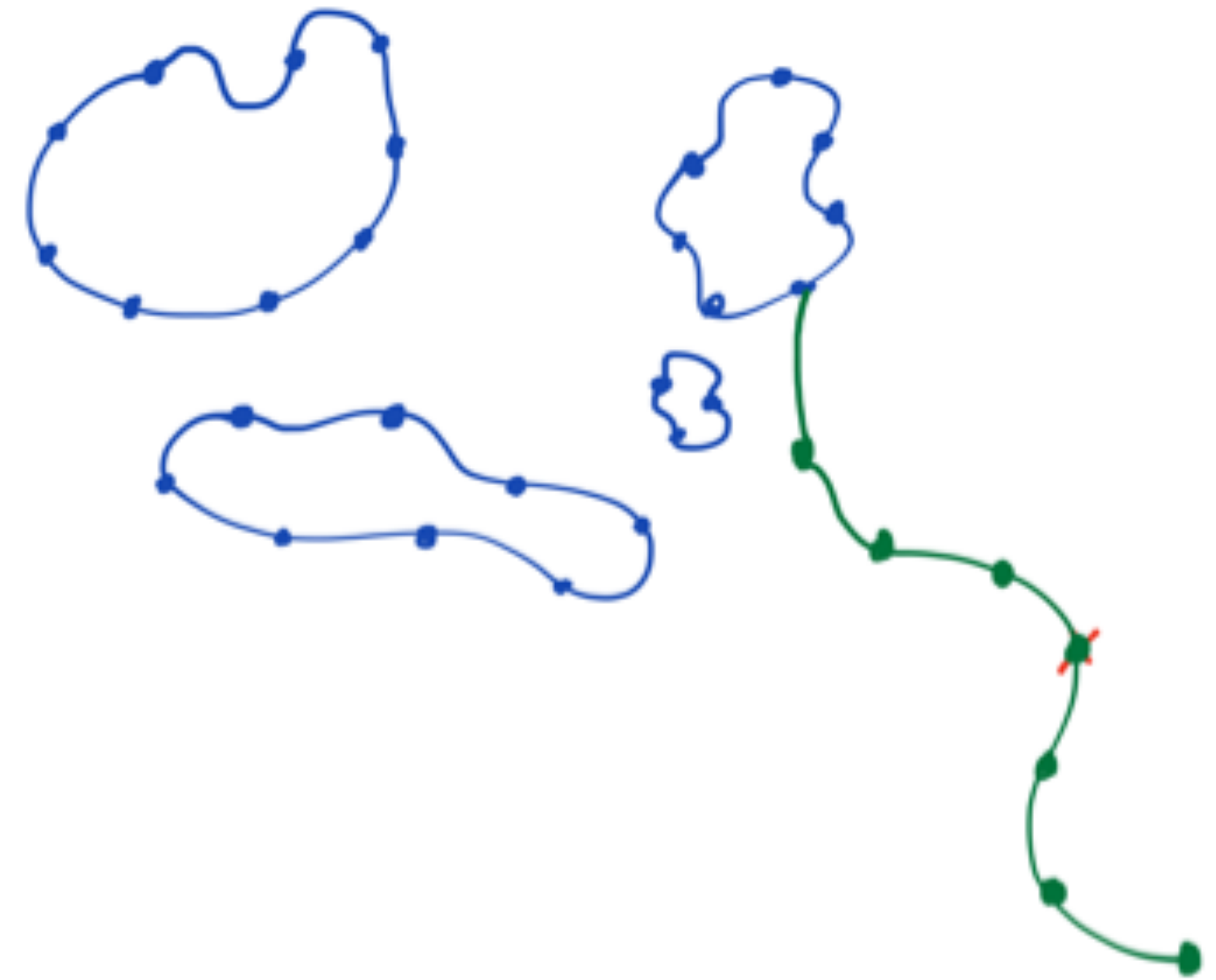
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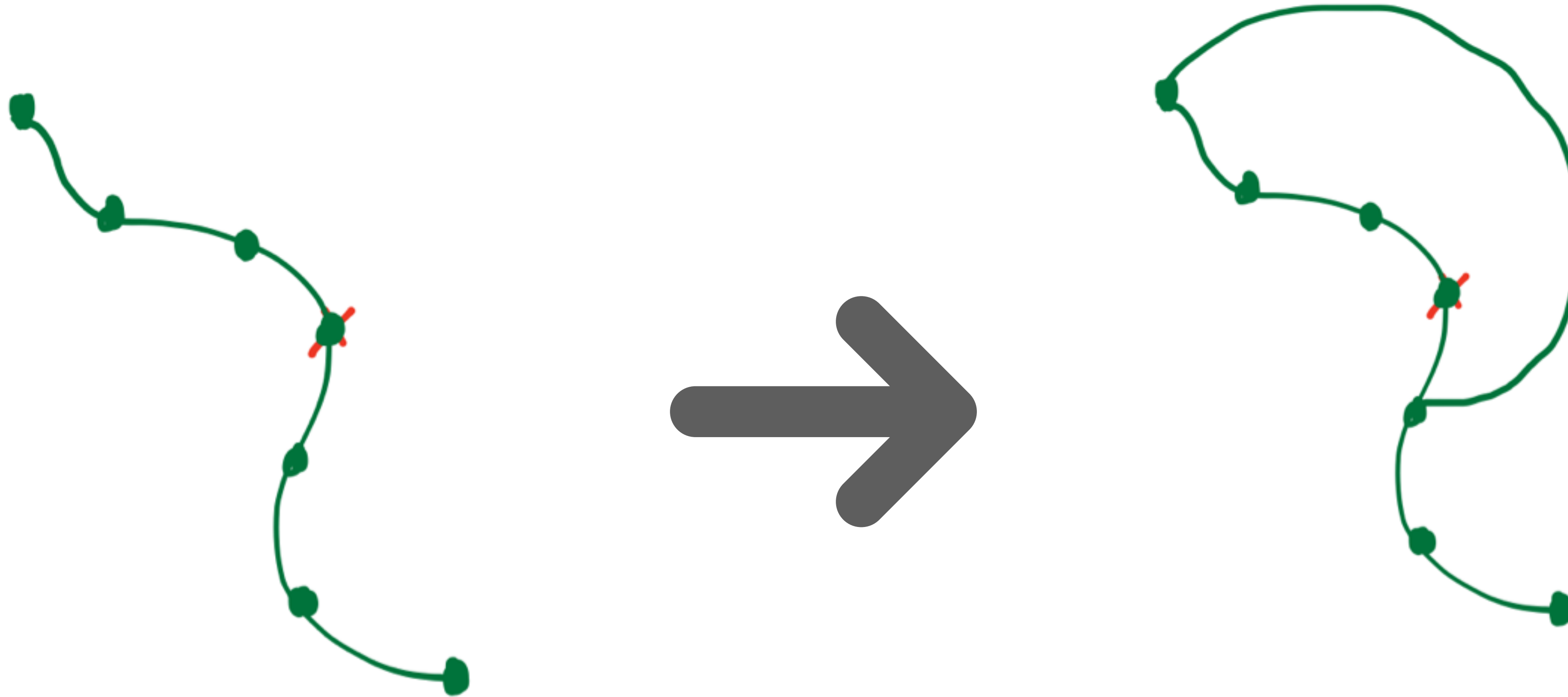


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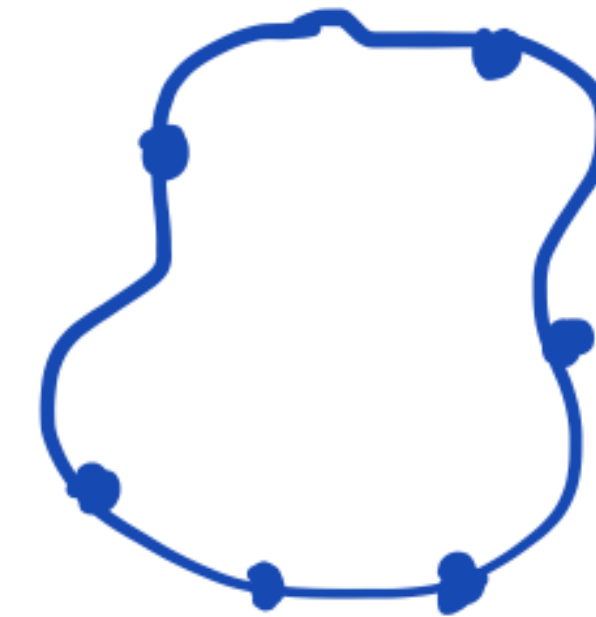
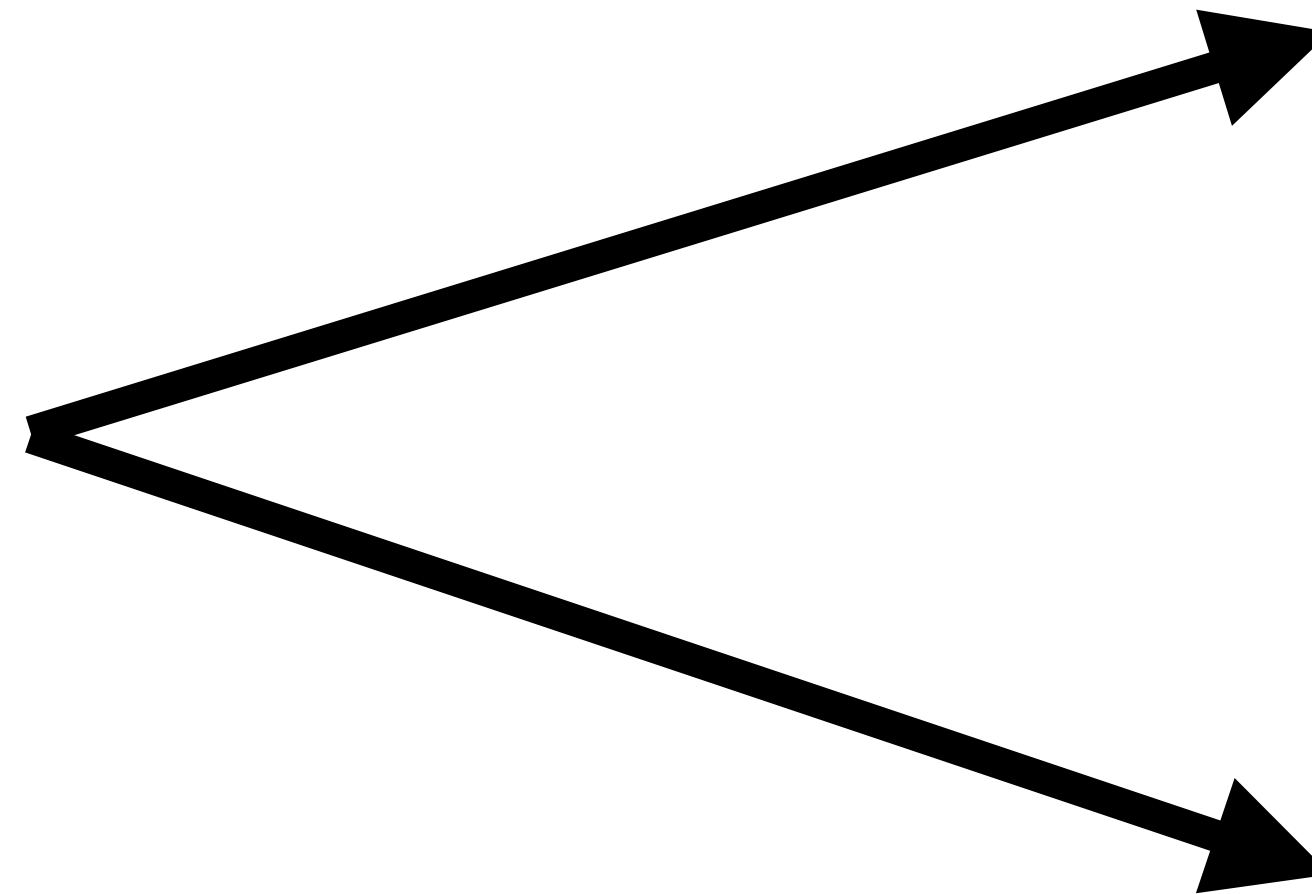
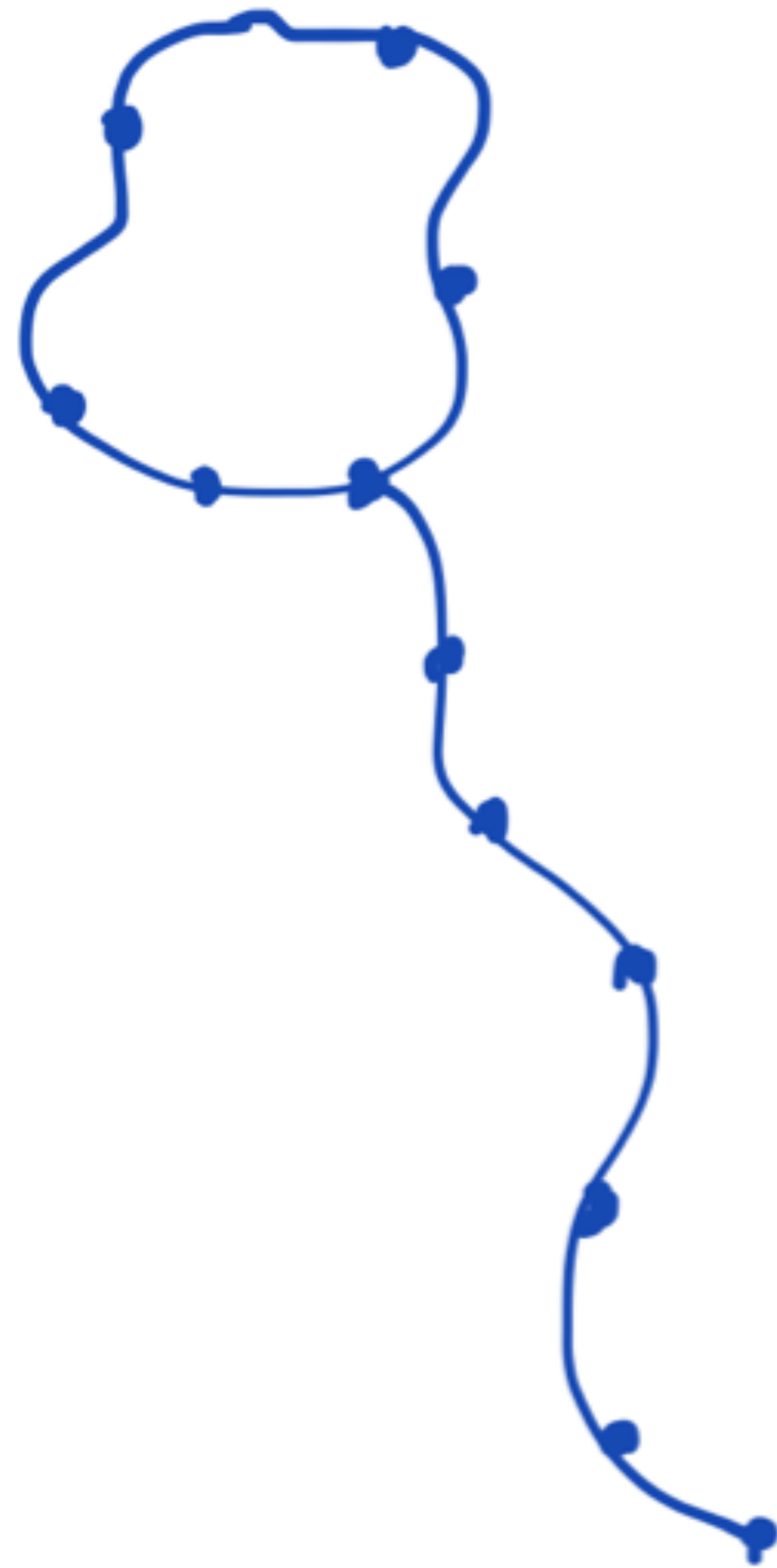


# Ballistic cycle creation





# Ballistic path / cycle erasure



# Facts

$$\sigma_N \sim \sqrt{\frac{9}{4} - \frac{24\Gamma(5/4)^4}{\pi^2}} n^{\frac{1}{2}} \approx 0.7801732 n^{\frac{1}{2}}$$

$$\sigma_N^{(\mathcal{E}_0)} \sim \sqrt{\frac{3}{4} - \frac{6\pi^2}{\Gamma(1/4)^4}} n^{\frac{1}{2}} \approx 0.4072901 n^{\frac{1}{2}}$$

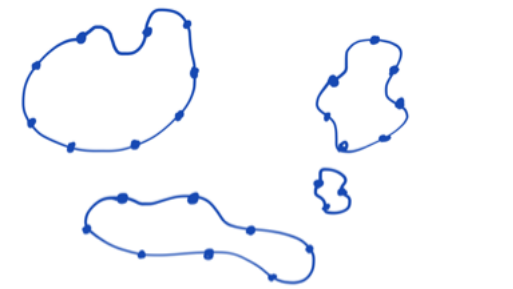
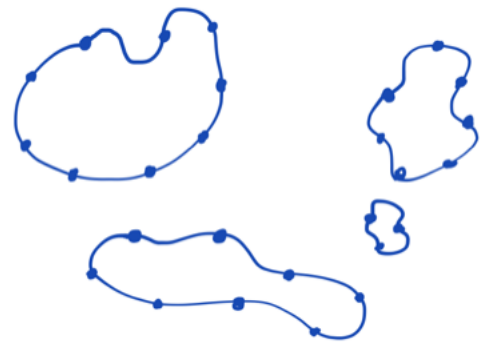
$$\mu_N^{(\mathcal{E}_0)} \sim \frac{\Gamma(7/4)}{\sqrt{3}\Gamma(5/4)} n^{\frac{1}{2}} \approx 0.5854143 n^{\frac{1}{2}}$$

$$\mu_N \sim \frac{2\sqrt{6}\Gamma(5/4)^2}{\pi} n^{\frac{1}{2}} \approx 1.2811439 n^{\frac{1}{2}}$$

$$\beta = \frac{1}{n} \Rightarrow x \sim \frac{1}{n}$$

$$\chi \approx n^{\frac{1}{2}}$$

# Figure pool



x

