

Appendix A

Curvilinear coordinates

A.1 Lamé coefficients

Consider set of equations

$$\xi_i = \xi_i(x_1, x_2, x_3), \quad i = 1, 2, 3$$

where ξ_1, ξ_2, ξ_3 independent, single-valued and continuous

(x_1, x_2, x_3) : coordinates of point P in Cartesian system with radius-vector \mathbf{x}

(ξ_1, ξ_2, ξ_3) : coordinates of point P in curvilinear system

$\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3$: base unit vectors in Cartesian coordinates

$\mathbf{e}_{\xi_1}, \mathbf{e}_{\xi_2}, \mathbf{e}_{\xi_3}$: base unit vectors in curvilinear system

we restrict ourselves to orthogonal curvilinear coordinates systems

$$\mathbf{e}_{\xi_i} \cdot \mathbf{e}_{\xi_j} = \delta_{ij}$$

Consider $\mathbf{x} = \mathbf{x}(x_i(\xi_j))$ and take the total differential (summation over double appearing indices understood)

$$d\mathbf{x} = \frac{\partial \mathbf{x}}{\partial \xi_j} d\xi_j$$

$\frac{\partial \mathbf{x}}{\partial \xi_1}$ partial derivative keeping ξ_2, ξ_3 fixed: vector tangential to coordinate line ξ_1

$$\longrightarrow \frac{\partial \mathbf{x}}{\partial \xi_1} = h_1 \mathbf{e}_{\xi_1}$$

in general

$$d\mathbf{x} = h_1 d\xi_1 \mathbf{e}_{\xi_1} + h_2 d\xi_2 \mathbf{e}_{\xi_2} + h_3 d\xi_3 \mathbf{e}_{\xi_3}$$

h_i Lamé coefficients

Those coefficients (might) depend on the coordinates ξ_i

$$h_i = h_i(\xi_1, \xi_2, \xi_3)$$

$h_i d\xi_i$ (no summation) components of length element along \mathbf{e}_{ξ_i}

Without restriction to orthogonal systems the length element $ds = |d\mathbf{x}|$ is defined as

$$ds^2 = d\mathbf{x} \cdot d\mathbf{x} = dx_j dx_j = \frac{\partial x_j}{\partial \xi_l} \frac{\partial x_j}{\partial \xi_m} d\xi_l d\xi_m = g_{lm} d\xi_l d\xi_m$$

Coefficients

$$g_{lm} = \frac{\partial x_j}{\partial \xi_l} \frac{\partial x_j}{\partial \xi_m}$$

are called the elements of the metric tensor

For an orthogonal curvilinear system the base vectors $\mathbf{e}_{\xi_1}, \mathbf{e}_{\xi_2}, \mathbf{e}_{\xi_3}$ are mutually perpendicular and form a right-handed system

We get the length element squared

$$ds^2 = h_1^2 (d\xi_1)^2 + h_2^2 (d\xi_2)^2 + h_3^2 (d\xi_3)^2$$

the Lamé coefficients

$$g_{11} = h_1^2, \quad g_{22} = h_2^2, \quad g_{33} = h_3^2, \quad g_{ik} = 0 \text{ for } i \neq k$$

the volume element (elementary rectangular box)

$$d^3x \equiv dV = h_1 h_2 h_3 d\xi_1 d\xi_2 d\xi_3$$

A.2 Some special orthogonal coordinates

- Cartesian coordinates ($-\infty < x_1 < \infty$, $-\infty < x_2 < \infty$, $-\infty < x_3 < \infty$)

$$\begin{aligned} h_1 &= h_2 = h_3 = 1 \\ ds^2 &= dx_1^2 + dx_2^2 + dx_3^2 \\ d^3x &= dx_1 dx_2 dx_3 \end{aligned}$$

- Cylindrical coordinates ($0 \leq \rho < \infty$, $0 \leq \varphi < 2\pi$, $-\infty < z < \infty$)

$$x_1 = \rho \cos \varphi, \quad x_2 = \rho \sin \varphi, \quad x_3 = z$$

$$\begin{aligned} h_\rho &= h_z = 1, \quad h_\varphi = \rho \\ ds^2 &= d\rho^2 + \rho^2 d\varphi^2 + dz^2 \\ d^3x &= \rho d\rho d\varphi dz \end{aligned}$$

- Spherical coordinates ($0 \leq r < \infty$, $0 \leq \theta \leq \pi$, $0 \leq \varphi < 2\pi$)

$$x_1 = r \sin \theta \cos \varphi, \quad x_2 = r \sin \theta \sin \varphi, \quad x_3 = r \cos \theta$$

$$\begin{aligned} h_r &= 1, \quad h_\theta = r, \quad h_\varphi = r \sin \theta \\ ds^2 &= dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\varphi^2 \\ d^3x &= r^2 dr \sin \theta d\theta d\varphi = r^2 dr d\cos \theta d\varphi \equiv r^2 dr d\Omega \end{aligned}$$

- Parabolic cylindrical coordinates
parametrization in Mathematica ($0 \leq u < \infty$, $0 \leq v < \infty$, $-\infty < z < \infty$)

$$x_1 = \frac{1}{2}(u^2 - v^2), \quad x_2 = uv, \quad x_3 = z$$

$$\begin{aligned} h_u &= h_v = \sqrt{u^2 + v^2}, \quad h_z = 1 \\ ds^2 &= (u^2 + v^2)(du^2 + dv^2) + dz^2 \\ d^3x &= (u^2 + v^2) du dv dz \end{aligned}$$

parametrization of Arfken ($0 \leq \xi < \infty$, $0 \leq \eta < \infty$, $-\infty < z < \infty$)

$$x_1 \leftrightarrow x_2, \quad u \rightarrow \eta, \quad v \rightarrow \xi$$

- Parabolic coordinates
parametrization of Arfken, Landau/Lifshitz ($0 \leq \xi < \infty$, $0 \leq \eta < \infty$, $0 \leq \varphi < 2\pi$)

$$x_1 = \sqrt{\xi \eta} \cos \varphi, \quad x_2 = \sqrt{\xi \eta} \sin \varphi, \quad x_3 = \frac{1}{2}(\xi - \eta)$$

$$\begin{aligned}
h_\xi &= \frac{1}{2} \sqrt{\frac{\xi + \eta}{\xi}}, \quad h_\eta = \frac{1}{2} \sqrt{\frac{\xi + \eta}{\eta}}, \quad h_\varphi = \sqrt{\xi \eta} \\
ds^2 &= \frac{1}{4} \frac{\xi + \eta}{\xi} d\xi^2 + \frac{1}{4} \frac{\xi + \eta}{\eta} d\eta^2 + \xi \eta d\varphi^2 \\
d^3x &= \frac{1}{4} (\xi + \eta) d\xi d\eta d\varphi
\end{aligned}$$

parametrization in Mathematica ($0 \leq u < \infty$, $0 \leq v < \infty$, $0 \leq \varphi < 2\pi$)

$$\xi = u^2, \quad \eta = v^2, \quad x_1 = uv \cos \varphi, \quad x_2 = uv \sin \varphi, \quad x_3 = \frac{1}{2}(u^2 - v^2)$$

$$\begin{aligned}
h_u &= h_v = \sqrt{u^2 + v^2}, \quad h_\varphi = uv \\
ds^2 &= (u^2 + v^2) (du^2 + dv^2) + uv d\varphi^2 \\
d^3x &= uv(u^2 + v^2) du dv d\varphi
\end{aligned}$$

A.3 Vector operations in orthogonal coordinates

Gradient

The length element along coordinate line ξ_1 is $h_1 d\xi_1$

$$\rightarrow \mathbf{e}_{\xi_1} \cdot \text{grad}\psi = \mathbf{e}_{\xi_1} \cdot \nabla\psi = \frac{1}{h_1} \frac{\partial\psi}{\partial\xi_1}$$

$$\nabla\psi = \frac{1}{h_1} \frac{\partial\psi}{\partial\xi_1} \mathbf{e}_{\xi_1} + \frac{1}{h_2} \frac{\partial\psi}{\partial\xi_2} \mathbf{e}_{\xi_2} + \frac{1}{h_3} \frac{\partial\psi}{\partial\xi_3} \mathbf{e}_{\xi_3}$$

\Rightarrow form of the nabla operator

$$\nabla = \sum_{i=1}^3 \mathbf{e}_{\xi_i} \frac{1}{h_i} \frac{\partial}{\partial\xi_i}$$

Divergence and Laplacian

Use the divergence definition

$$\nabla \cdot \mathbf{A} = \lim_{\Delta V \rightarrow 0} \frac{1}{\Delta V} \oint_S \mathbf{A} \cdot \mathbf{n} da$$

with

$$\mathbf{A} = A_1(\xi_1, \xi_2, \xi_3) \mathbf{e}_{\xi_1} + A_2(\xi_1, \xi_2, \xi_3) \mathbf{e}_{\xi_2} + A_3(\xi_1, \xi_2, \xi_3) \mathbf{e}_{\xi_3}$$

and choose as volume ΔV an elementary rectangular box of sides $h_1\Delta\xi_1 h_2\Delta\xi_2 h_3\Delta\xi_3$. Analogously to the derivation in Cartesian coordinates (compare Chapter 1.2) we calculate the net outward flux in direction of coordinate line ξ_1 divided by the volume:

$$\begin{aligned} & \lim_{\Delta V \rightarrow 0} \left[A_1(\xi_1 + \Delta\xi_1, \bar{\xi}_2, \bar{\xi}_3) h_2(\xi_1 + \Delta\xi_1, \bar{\xi}_2, \bar{\xi}_3) \Delta\xi_2 h_3(\xi_1 + \Delta\xi_1, \bar{\xi}_2, \bar{\xi}_3) \Delta\xi_3 - \right. \\ & \left. A_1(\xi_1, \bar{\xi}_2, \bar{\xi}_3) h_2(\xi_1, \bar{\xi}_2, \bar{\xi}_3) \Delta\xi_2 h_3(\xi_1, \bar{\xi}_2, \bar{\xi}_3) \Delta\xi_3 \right] / (h_1\Delta\xi_1 h_2\Delta\xi_2 h_3\Delta\xi_3) \\ & \rightarrow \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial\xi_1} (h_2 h_3 A_1) \end{aligned}$$

\Rightarrow we get for the divergence

$$\nabla \cdot \mathbf{A} = \frac{1}{h_1 h_2 h_3} \left\{ \frac{\partial}{\partial\xi_1} (h_2 h_3 A_1) + \frac{\partial}{\partial\xi_2} (h_3 h_1 A_2) + \frac{\partial}{\partial\xi_3} (h_1 h_2 A_3) \right\}$$

The Laplacian operator follows ($\nabla^2 = \nabla \cdot \nabla = \Delta$):

$$\Delta = \frac{1}{h_1 h_2 h_3} \left\{ \frac{\partial}{\partial\xi_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial}{\partial\xi_1} \right) + \frac{\partial}{\partial\xi_2} \left(\frac{h_3 h_1}{h_2} \frac{\partial}{\partial\xi_2} \right) + \frac{\partial}{\partial\xi_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial}{\partial\xi_3} \right) \right\}$$

Curl

Using a similar derivation as in Cartesian coordinates we get for the curl of vector \mathbf{A}

$$\nabla \times \mathbf{A} = \frac{1}{h_2 h_3} \left\{ \frac{\partial}{\partial \xi_2} (h_3 A_3) - \frac{\partial}{\partial \xi_3} (h_2 A_2) \right\} \mathbf{e}_{\xi_1} + \text{cyclic permutations}$$

Representation as determinant:

$$\nabla \times \mathbf{A} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \mathbf{e}_{\xi_1} & h_2 \mathbf{e}_{\xi_2} & h_3 \mathbf{e}_{\xi_3} \\ \frac{\partial}{\partial \xi_1} & \frac{\partial}{\partial \xi_2} & \frac{\partial}{\partial \xi_3} \\ h_1 A_1 & h_2 A_2 & h_3 A_3 \end{vmatrix}$$

A.4 Some explicit forms of vector operations

Cartesian coordinates (x_1, x_2, x_3) with base vectors $\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3$

$$\begin{aligned}\nabla\psi &= \frac{\partial\psi}{\partial x_1}\mathbf{e}_1 + \frac{\partial\psi}{\partial x_2}\mathbf{e}_2 + \frac{\partial\psi}{\partial x_3}\mathbf{e}_3 \\ \nabla\cdot\mathbf{A} &= \frac{\partial A_1}{\partial x_1} + \frac{\partial A_2}{\partial x_2} + \frac{\partial A_3}{\partial x_3} \\ \nabla\times\mathbf{A} &= \left(\frac{\partial A_3}{\partial x_2} - \frac{\partial A_2}{\partial x_3}\right)\mathbf{e}_1 + \left(\frac{\partial A_1}{\partial x_3} - \frac{\partial A_3}{\partial x_1}\right)\mathbf{e}_2 + \left(\frac{\partial A_2}{\partial x_1} - \frac{\partial A_1}{\partial x_2}\right)\mathbf{e}_3 \\ \Delta\psi &= \frac{\partial^2\psi}{\partial x_1^2} + \frac{\partial^2\psi}{\partial x_2^2} + \frac{\partial^2\psi}{\partial x_3^2}\end{aligned}$$

Cylindrical coordinates (ρ, φ, z) with base vectors $\mathbf{e}_\rho, \mathbf{e}_\varphi, \mathbf{e}_z$

$$\begin{aligned}\nabla\psi &= \frac{\partial\psi}{\partial\rho}\mathbf{e}_\rho + \frac{1}{\rho}\frac{\partial\psi}{\partial\varphi}\mathbf{e}_\varphi + \frac{\partial\psi}{\partial z}\mathbf{e}_z \\ \nabla\cdot\mathbf{A} &= \frac{1}{\rho}\frac{\partial}{\partial\rho}(\rho A_\rho) + \frac{1}{\rho}\frac{\partial A_\varphi}{\partial\varphi} + \frac{\partial A_z}{\partial z} \\ \nabla\times\mathbf{A} &= \left(\frac{1}{\rho}\frac{\partial A_z}{\partial\varphi} - \frac{\partial A_\varphi}{\partial z}\right)\mathbf{e}_\rho + \left(\frac{\partial A_\rho}{\partial z} - \frac{\partial A_z}{\partial\rho}\right)\mathbf{e}_\varphi + \frac{1}{\rho}\left(\frac{\partial}{\partial\rho}(\rho A_\varphi) - \frac{\partial A_\rho}{\partial\varphi}\right)\mathbf{e}_z \\ \Delta\psi &= \frac{1}{\rho}\frac{\partial}{\partial\rho}\left(\rho\frac{\partial\psi}{\partial\rho}\right) + \frac{1}{\rho^2}\frac{\partial^2\psi}{\partial\varphi^2} + \frac{\partial^2\psi}{\partial z^2}\end{aligned}$$

Spherical coordinates (r, θ, φ) with base vectors $\mathbf{e}_r, \mathbf{e}_\theta, \mathbf{e}_\varphi$

$$\begin{aligned}\nabla\psi &= \frac{\partial\psi}{\partial r}\mathbf{e}_r + \frac{1}{r}\frac{\partial\psi}{\partial\theta}\mathbf{e}_\theta + \frac{1}{r\sin\theta}\frac{\partial\psi}{\partial\varphi}\mathbf{e}_\varphi \\ \nabla\cdot\mathbf{A} &= \frac{1}{r^2}\frac{\partial}{\partial r}(r^2 A_r) + \frac{1}{r\sin\theta}\frac{\partial}{\partial\theta}(\sin\theta A_\theta) + \frac{1}{r\sin\theta}\frac{\partial A_\varphi}{\partial\varphi} \\ \nabla\times\mathbf{A} &= \frac{1}{r\sin\theta}\left(\frac{\partial}{\partial\theta}(\sin\theta A_\varphi) - \frac{\partial A_\theta}{\partial\varphi}\right)\mathbf{e}_r \\ &\quad + \left(\frac{1}{r\sin\theta}\frac{\partial A_r}{\partial\varphi} - \frac{1}{r}\frac{\partial}{\partial r}(r A_\varphi)\right)\mathbf{e}_\theta + \frac{1}{r}\left(\frac{\partial}{\partial r}(r A_\theta) - \frac{\partial A_r}{\partial\theta}\right)\mathbf{e}_\varphi \\ \Delta\psi &= \frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial\psi}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial\psi}{\partial\theta}\right) + \frac{1}{r^2\sin^2\theta}\frac{\partial^2\psi}{\partial\varphi^2}\end{aligned}$$

Note

$$\frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial\psi}{\partial r}\right) \equiv \frac{1}{r}\frac{\partial^2}{\partial r^2}(r\psi)$$

A.5 Relation between unit base vectors and their time derivatives

Relation between the local unit base vectors of cylindrical coordinates $\mathbf{e}_\rho(t), \mathbf{e}_\varphi(t), \mathbf{e}_z$ and their time derivatives and the global constant unit base vectors of Cartesian coordinates $\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3$

$$\begin{aligned}
 \mathbf{e}_\rho &= \cos \varphi \mathbf{e}_1 + \sin \varphi \mathbf{e}_2 & \mathbf{e}_1 &= \cos \varphi \mathbf{e}_\rho - \sin \varphi \mathbf{e}_\varphi \\
 \mathbf{e}_\varphi &= -\sin \varphi \mathbf{e}_1 + \cos \varphi \mathbf{e}_2 & \mathbf{e}_2 &= \sin \varphi \mathbf{e}_\rho + \cos \varphi \mathbf{e}_\varphi \\
 \mathbf{e}_z &= \mathbf{e}_3 & \mathbf{e}_3 &= \mathbf{e}_z \\
 \\
 \dot{\mathbf{e}}_\rho &= -\sin \varphi \dot{\varphi} \mathbf{e}_1 + \cos \varphi \dot{\varphi} \mathbf{e}_2 = \dot{\varphi} \mathbf{e}_\varphi & \dot{\mathbf{e}}_1 &= 0 \\
 \dot{\mathbf{e}}_\varphi &= -\dot{\varphi} \mathbf{e}_\rho & \dot{\mathbf{e}}_2 &= 0 \\
 \dot{\mathbf{e}}_z &= 0 & \dot{\mathbf{e}}_3 &= 0
 \end{aligned}$$

Same for the local unit base vectors of spherical coordinates $\mathbf{e}_r(t), \mathbf{e}_\theta(t), \mathbf{e}_\varphi(t)$ and their time derivatives

$$\begin{aligned}
 \mathbf{e}_r &= \sin \theta \cos \varphi \mathbf{e}_1 + \sin \theta \sin \varphi \mathbf{e}_2 + \cos \theta \mathbf{e}_3 \\
 \mathbf{e}_\theta &= \cos \theta \cos \varphi \mathbf{e}_1 + \cos \theta \sin \varphi \mathbf{e}_2 - \sin \theta \mathbf{e}_3 \\
 \mathbf{e}_\varphi &= -\sin \varphi \mathbf{e}_1 + \cos \varphi \mathbf{e}_2 \\
 \\
 \mathbf{e}_1 &= \sin \theta \cos \varphi \mathbf{e}_r + \cos \theta \cos \varphi \mathbf{e}_\theta - \sin \varphi \mathbf{e}_\varphi \\
 \mathbf{e}_2 &= \sin \theta \sin \varphi \mathbf{e}_r + \cos \theta \sin \varphi \mathbf{e}_\theta + \cos \varphi \mathbf{e}_\varphi \\
 \mathbf{e}_3 &= \cos \theta \mathbf{e}_r - \sin \theta \mathbf{e}_\theta \\
 \\
 \dot{\mathbf{e}}_r &= \dot{\theta} \mathbf{e}_\theta + \sin \theta \dot{\varphi} \mathbf{e}_\varphi \\
 \dot{\mathbf{e}}_\theta &= -\dot{\theta} \mathbf{e}_r + \cos \theta \dot{\varphi} \mathbf{e}_\varphi \\
 \dot{\mathbf{e}}_\varphi &= -\dot{\varphi} (\sin \theta \mathbf{e}_r + \cos \theta \mathbf{e}_\theta)
 \end{aligned}$$

Relation between the local unit base vectors $\mathbf{e}_\rho(t), \mathbf{e}_\varphi(t), \mathbf{e}_z$ and $\mathbf{e}_r(t), \mathbf{e}_\theta(t), \mathbf{e}_\varphi(t)$

$$\begin{aligned}
 \mathbf{e}_\rho &= \sin \theta \mathbf{e}_r + \cos \theta \mathbf{e}_\theta & \mathbf{e}_r &= \sin \theta \mathbf{e}_\rho + \cos \theta \mathbf{e}_z \\
 \mathbf{e}_\varphi &= \mathbf{e}_\varphi & \mathbf{e}_\theta &= \cos \theta \mathbf{e}_\rho - \sin \theta \mathbf{e}_z \\
 \mathbf{e}_z &= \cos \theta \mathbf{e}_r - \sin \theta \mathbf{e}_\theta & \mathbf{e}_\varphi &= \mathbf{e}_\varphi
 \end{aligned}$$