

Overlap Hypercube Fermions in QCD with Light Quarks

- I. Construction of the overlap HF
Locality, rotation symmetry and condition number
- II. Applications in the p -regime :
 m_q **vs.** m_π , m_ρ , m_{PCAC} , Z_A and F_π
- III. Applications in the ϵ -regime : F_π and Σ
Topological charges and susceptibility
Zero-mode contributions to $\langle PP \rangle$ in χ limit

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I. Construction of the Overlap HF

For free fermions, the **perfect lattice action** is known analytically (W.B./Wiese '95). Dirac operator:

$$D_{x,y} = \gamma_\mu \rho_\mu(x - y) + \lambda(x - y)$$

with closed expressions for $\rho_\mu(p)$, $\lambda(p)$.

Based on iterated RG transformations \rightarrow no lattice artifacts

Range of $D_{x,y}$ is infinite \rightarrow *optimise* the RGT for *locality*, then truncate by periodic b.c. to a 3^4 hypercube

\rightarrow $\text{supp} [\rho_\mu(x - y), \lambda(x - y)] \subset \{|x_\nu - y_\nu| \leq 1\}$, $\nu = 1 \dots 4$

“Hypercube Fermion” (HF), still excellent scaling (W.B./Brower/Chandrasekharan/Wiese, '96).

Gauging: sum over *shortest lattice paths* , plus **fat links**

$U_\mu(x) \rightarrow (1 - \alpha)U_\mu(x) + \frac{\alpha}{6} \sum \text{staples}$; **link amplification** \rightarrow criticality

Truncation and imperfect gauging \rightarrow **scaling** and **chirality** somewhat **distorted**. **Chirality** can be **corrected** again by inserting the HF into the **overlap formula** (H. Neuberger, '97) :

$$D_{\text{ov}} = \rho(1 + A/\sqrt{A^\dagger A}) , \quad A := D_0 - \rho , \quad \rho \gtrsim 1$$

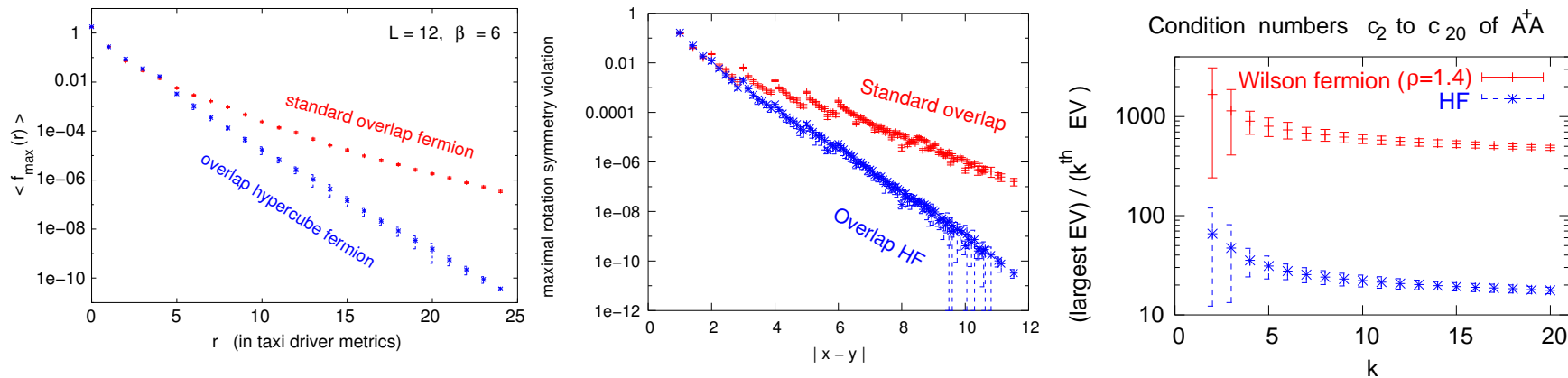
where D_0 is some lattice Dirac operator (with $D_0 = \gamma_5 D_0^\dagger \gamma_5$).

- **Standard overlap fermion: based on $D_0 = D_{\text{Wilson}}$**
Drastic change : $D_{\text{Wilson}} \rightarrow D_{\text{ov-W}}$
- **Overlap HF: $D_0 = D_{\text{HF}}$**
 D_{HF} is approx. chiral already \rightarrow modest modification : $D_{\text{HF}} \rightarrow D_{\text{ov-HF}}$

Both are Ginsparg-Wilson operators \rightarrow *exact χ sym.* (P. Hasenfratz, '98, M. Lüscher, '98)

⇒ Virtues of the HF are essentially inherited by $D_{\text{OV-HF}}$: high degree of **locality**, **approx. rotation symmetry**, **fast convergence**

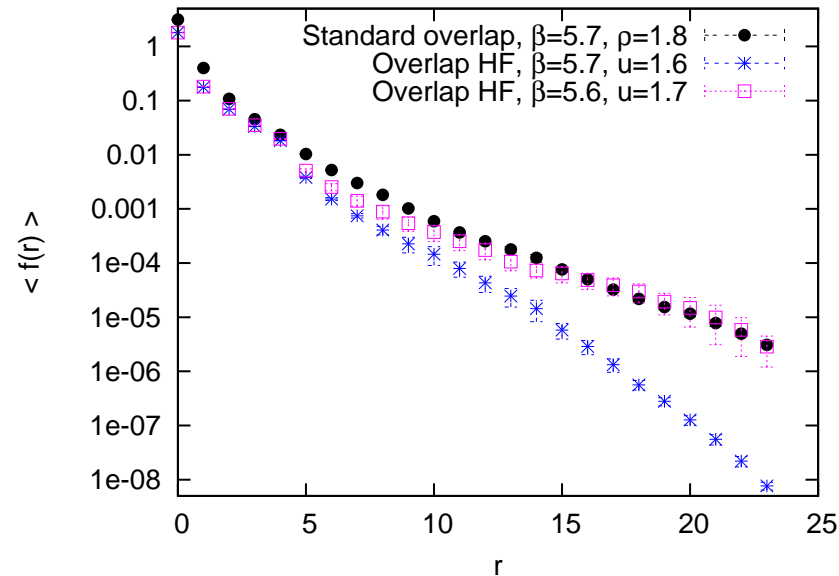
Examples at $\beta = 6$, $V = 12^4$, Wilson gauge action, quenched :



Overlap HF vs. standard overlap at $\rho = 1.4$

- $f_{\max}(r)$ “maximal correlation” between unit source $\bar{\psi}_x$ and ψ_y with $r = \|x - y\|_1$ (Hernández et al. '99)
- Measure for violation of rot. sym. : $f_{\max}(|x - y|) - f_{\min}(|x - y|)$
- Condition numbers c_k :
 (largest EV of $A^\dagger A$) / (smallest EV, after projecting out $k - 1$ modes)
 Gain factor $\approx 25 \rightarrow$ factor 5 in polynomial degree, vs. factor 15 of HF overhead

Locality at strong coupling:



- $\beta = 5.7$ ($a \simeq 0.17$ fm) : $D_{\text{ov-W}}$ with optimal ρ barely local
 $D_{\text{ov-HF}}$ with optimal link amplification has powerful locality,
 stronger than $D_{\text{ov-W}}$ at $\beta = 6$, $\rho = 1.4$
- $\beta = 5.6$: for $D_{\text{ov-W}}$ locality collapses, no valid Dirac operator
 but $D_{\text{ov-HF}}$ is still local

The overlap-HF formulation provides chiral fermions on coarser lattices.

II. Applications in the p -Regime

p -regime : p expansion of χ PT is applicable, box length $\mathbf{L} \gg \frac{1}{m_\pi}$

We consider $\beta = 5.85$ ($a \simeq 0.123$ fm)

$V = 12^3 \times 24$, $V_{\text{phys}} = (1.48 \text{ fm})^3 \times 2.96 \text{ fm}$

Bare quark masses : $am_q = 0.01, 0.02, 0.04, 0.06, 0.08, 0.1$

$\triangleq m_q \simeq 16.1 \text{ MeV} \dots 161 \text{ MeV}$, $D_{\text{ov}}(m_q) = \left(1 - \frac{m_q}{2\rho}\right) D_{\text{ov}}^{(0)} + m_q$

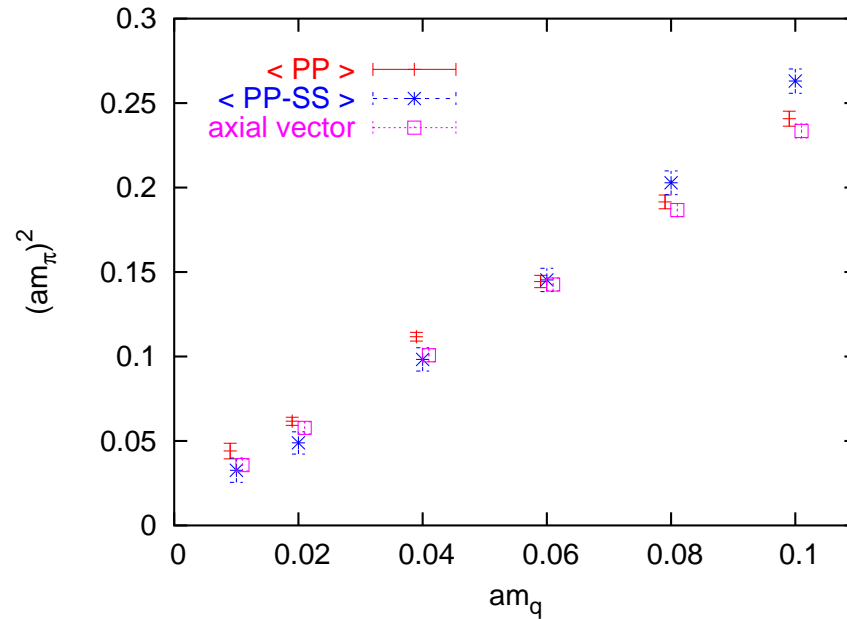
{ *At smallest m_q close to ϵ -regime* }

Statistics : 100 propagators

We evaluate the pion mass in three ways:

- Pseudoscalar correlator $\langle PP \rangle$, $P = \bar{\psi}\gamma_5\psi$
- $\langle A_4 A_4 \rangle$, $A_4 = \bar{\psi}\gamma_5\gamma_4\psi$
- $\langle PP - SS \rangle$, $S = \bar{\psi}\psi$: subtraction useful at small m_q ,
avoids contamination by zero modes, which plagues quenched results

Pion mass

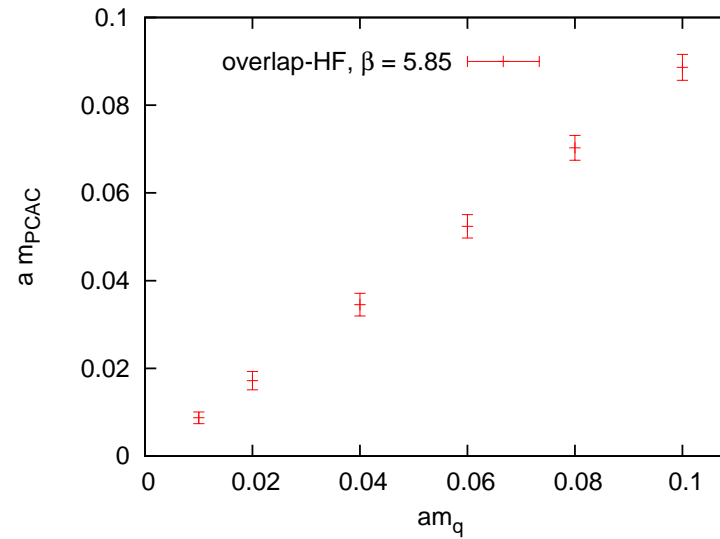
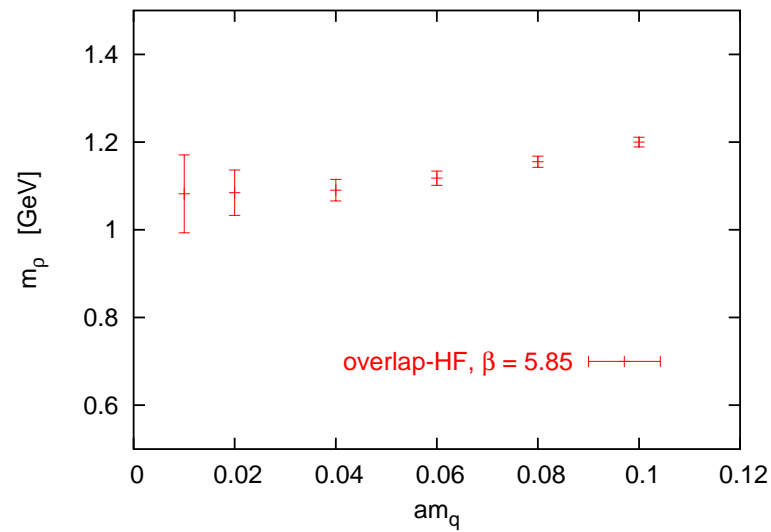


$$m_\pi^2 \propto m_q ; \quad \text{Intercept for } m_{\pi,PP-SS}^2(m_q = 0) = -0.001(15)$$

Hierarchy at small m_q : $m_{\pi,PP} > m_{\pi,AA} > m_{\pi,PP-SS}$
 (agrees with P. Hasenfratz et al. '02)

Smallest pion mass : $m_{\pi,PP-SS}(am_q = 0.01) \simeq (279 \pm 32) \text{ MeV}$
 $\rightarrow \frac{L}{\xi} \approx 2$, edge of p -regime

ρ -meson mass and quark mass according to the axial Ward identity:



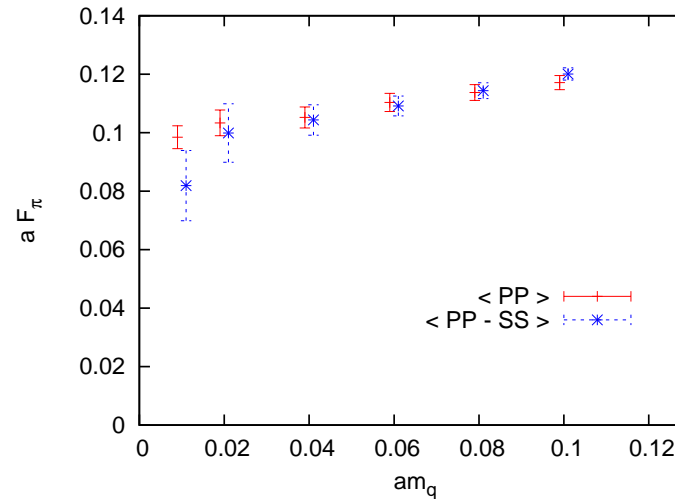
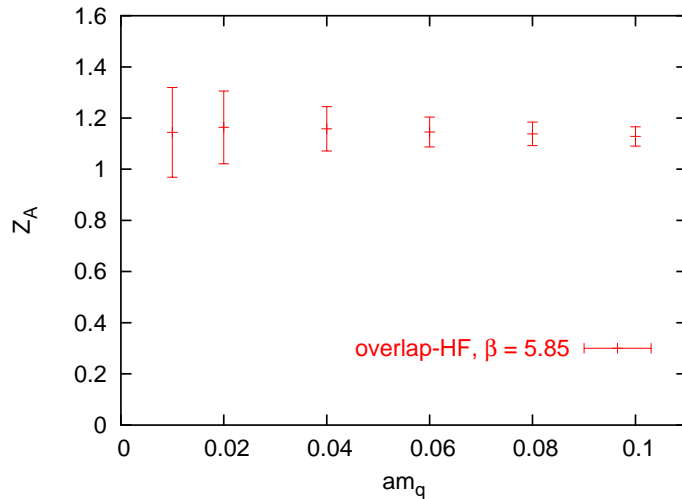
Chiral extrapolation :

$$m_\rho \in [978 \text{ MeV}, 1057 \text{ MeV}] \quad , \quad m_{PCAC} \in [-0.00094, 0.00035]$$

$$m_{PCAC} = \sum_{\vec{x}} \langle \partial_4 A_4^\dagger(x) P(0) \rangle / 2 \sum_{\vec{x}} \langle P^\dagger(x) P(0) \rangle \quad (\text{for } D_{ov-w}, \text{ see } \chi\text{LF, '04})$$

This yields the renormalisation constant $Z_A = \frac{m_q}{m_{\text{PCAC}}}$,

and the pion decay constant $F_\pi = \frac{2m_q}{m_\pi^2} |\langle 0|P|\pi\rangle|$



Chiral extrapolation :

- $Z_A = 1.17(2)$; much closer to 1 than Z_A for standard overlap at same β
($Z_A \simeq 1.45$ at $\rho = 1.6$, χ LF '04; $Z_A \simeq 1.55$ at $\beta = 6$, $\rho = 1.4$, Berruto et al. '03)
- $F_{\pi,PP} \in [109 \text{ MeV}, 114 \text{ MeV}]$, $F_{\pi,PP-SS} \in [95 \text{ MeV}, 113 \text{ MeV}]$
above theor. value in the χ limit: 86 MeV (Colangelo/Dürr, '04)
→ to be reconsidered in the ϵ -regime at smaller $m_q \dots$

III. Applications in the ϵ -Regime

ϵ -regime : $\frac{1}{m_\pi} > L$

In χ PT p -expansion fails due to dominant 0-modes

But: analytical treatment of zero modes with **collective variables** (Gasser/Leutwyler, '87),
higher modes captured by **ϵ -expansion**

Motivation: unphysical setting, but physical values of the **Low Energy Constants (LEC, free parameters in χ PT, $F_\pi, \Sigma \dots$)** can be evaluated in small volumes

→ LEC at $V = \infty$.

Challenge for simulations with Ginsparg-Wilson fermions.

However: quenching → log. finite size effects (Damgaard, '01)

Peculiarity of ϵ -regime: observables depend strongly on the *topological sector* (Leutwyler/Smilga, '92)

Evaluation of Σ from the Dirac Spectrum

χ **Random Matrix Theory** conjectures the densities of the **low lying Dirac eigenvalues** λ in the ϵ -regime (Damgaard/Nishigaki, '98) :

$$\rho_n^{(\nu)}(\mathbf{z})$$

$z := \lambda \Sigma V$ (dim'less)

$n = 1, 2, 3 \dots$ (lowest eigenvalues, excluding zeros)

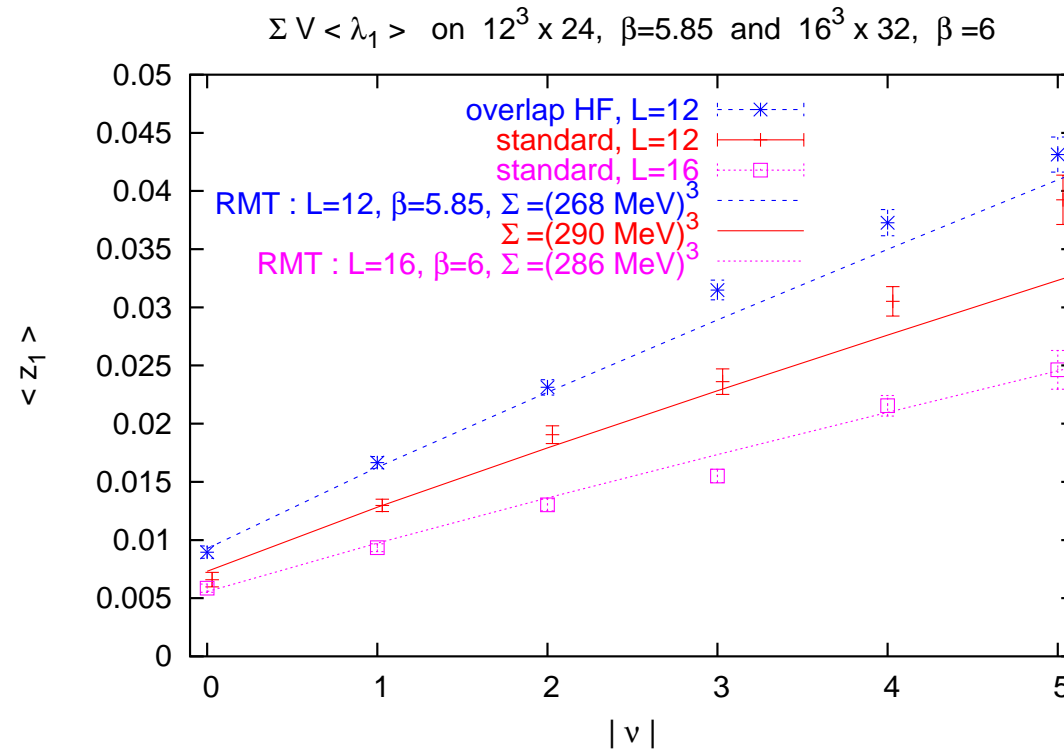
$\nu := \text{index} \equiv \text{top. charge}$

W.B./Jansen/Shcheredin, Giusti/Lüscher/Weisz/Wittig, Galletly et al., '03 :

χ RMT predictions hold for leading EVs if the volume is not too small ($L \gtrsim 1.1$ fm)

Scalar condensate Σ : only free parameter, determined by the fit.

$V = (1.48 \text{ fm})^3 \times 2.96 \text{ fm}$ at $a \simeq 0.12 \text{ fm}$ ($\beta = 5.85$), and $a \simeq 0.093 \text{ fm}$ ($\beta = 6$)



Prediction works well in particular for $|\nu| = 0, 1, 2$

For the overlap-HF the fits yields : $\Sigma = (268(2) \text{ MeV})^3$

Preliminary result for F_π from the axial-vector correlator

Correlators to first order in **quenched** χ PT (q χ PT) (Damgaard et al., '02, '03)

$\langle \text{Vector} - \text{Vector} \rangle = 0$ (all orders)

$\langle \text{Scalar} - \text{Scalar} \rangle$ and $\langle \text{Pseudoscalar} - \text{Pseudoscalar} \rangle$

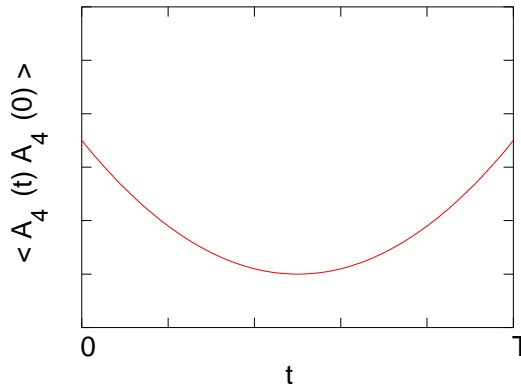
involve additional LEC, specific to quenching.

\Rightarrow Focus on $\langle \text{Axialvector} - \text{Axialvector} \rangle$

in leading order **only** Σ and F_π

Bare axial current at $\vec{p} = \vec{0}$: $A_\mu(t) = \sum_{\vec{x}} \bar{\Psi}(t, \vec{x}) \gamma_5 \gamma_\mu \Psi(t, \vec{x})$

$\langle A_4(t) A_4(0) \rangle$ in a volume $L^3 \times T$:



Parabola: min. at $t = T/2$; F_π^2/T : additive, $\Sigma \rightarrow$ curvature

W.B./Chiarappa/Jansen/Nagai/Shcheredin, '03 :

Simulations at $m_q = 21.3$ MeV, $\beta = 6$ on $V = 10^3 \times 24$ and 12^4

- First lattice failed; consistent with χ RMT study: **L \simeq 0.98 fm is too small**
- MC history in $\nu = 0$ has **strong spikes** , smoother at $\nu \neq 0$,
Spikes exactly for conf's with very **small EV** , most frequent at $\nu = 0$
{ Remedy: "Low Mode Averaging" , Giusti et al. '04 }

Decent agreement with q χ PT in $V = (1.12 \text{ fm})^4$, $|\nu| = 1$:

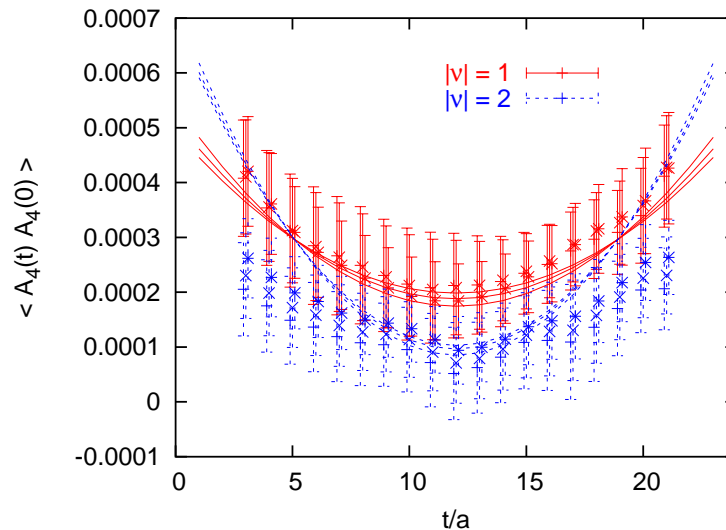
$$F_\pi = (86.7 \pm 4.0) \text{ MeV}$$

Renormalisation with $Z_A \simeq 1.55 \rightarrow F_\pi^r \approx 130 \text{ MeV} > 93 \text{ MeV}$.

Σ cannot be extracted (curves hardly sensitive to $\Sigma = 0$ or $\Sigma = (250 \text{ MeV})^3$).

{ See also : Fukaya/Hashimoto/Ogawa, '05 }

Preliminary result for overlap-HF ($12^3 \times 24$, $\beta = 5.85$) 10 propagators in each of the sectors $|\nu| = 1$, $|\nu| = 2$; $m_q = 1.6$ MeV, 4.8 MeV, 8 MeV

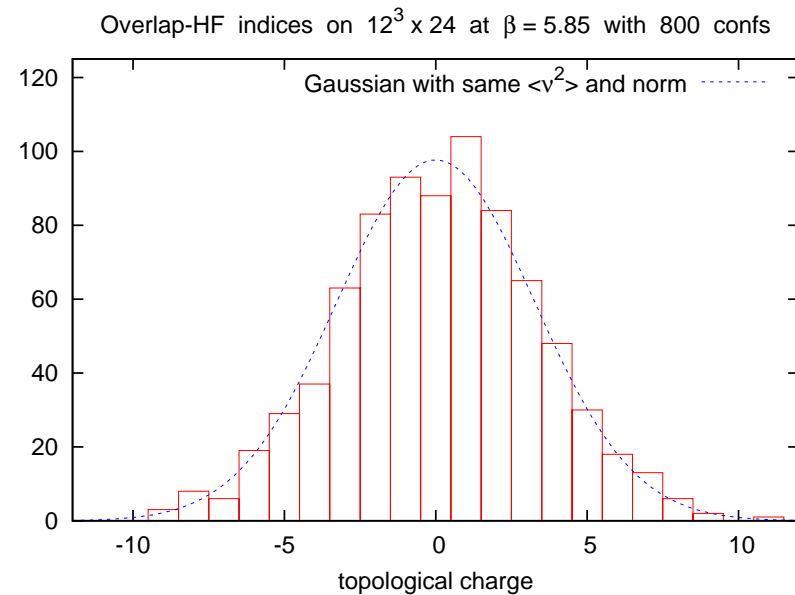
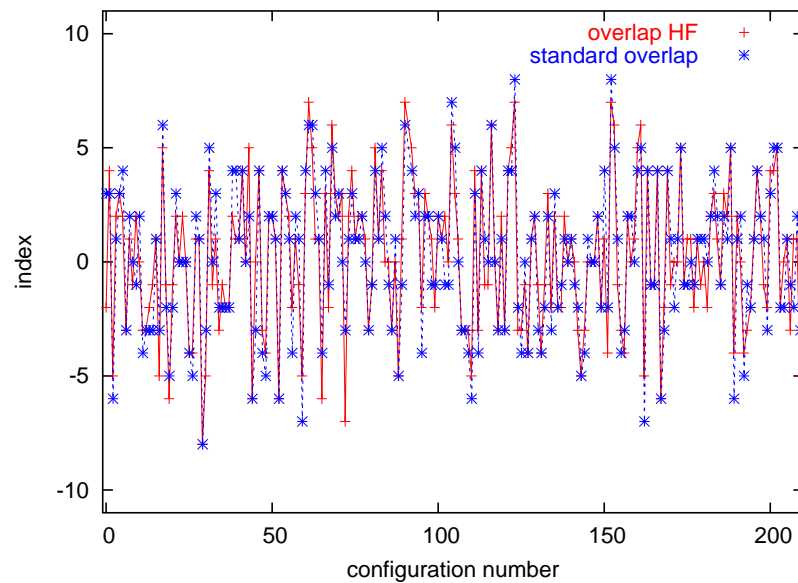


Global fit with $\Sigma = (268 \text{ MeV})^3$: $F_\pi = (96 \pm 10) \text{ MeV}$

Renormalised : $F_\pi^r = (104 \pm 9) \text{ MeV}$ (agrees with quenched literature)

Alternative method based on 0-mode contributions to meson correlators in the χ limit:
see below.

Topological Charges :

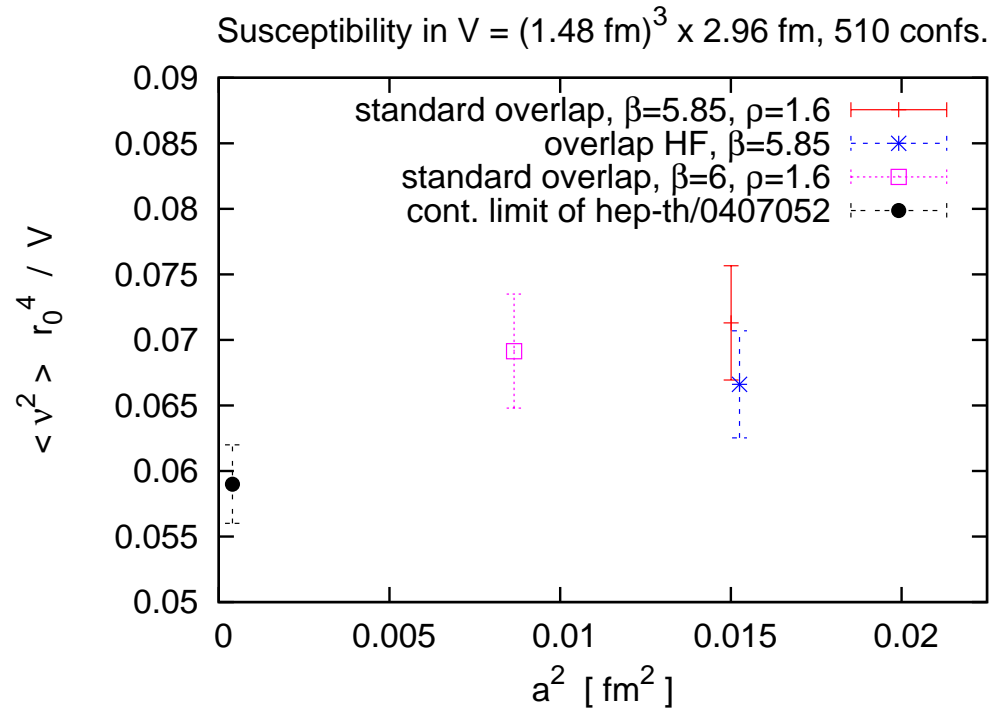


Left: part of **index histories** for **standard overlap** and **overlap-HF**, same conf's,
 $\langle |\nu_{\text{ov-W}} - \nu_{\text{ov-HF}}| \rangle \approx 0.8$, max. deviation: 5

Right: charge histogram for overlap-HF, compatible with a **Gaussian**, parity (?)

(c.f. Alles/D'Elia/DiGiacomo, '05)

Topological susceptibility $\chi = \frac{1}{V} \langle \nu^2 \rangle$ (relevant for mass of η')



Comparison to the continuum limit by Del Debbio/Giusti/Pica '04 (for $\langle \nu^2 \rangle r_0^4 / V$).

Considering finite a and $\rho = 1.6$: agreement within the errors

Overlap HF data closer to cont. limit (same confs at $\beta = 5.85$)

Giusti/Hernández/Laine/Weisz/Wittig '04: qχPT predictions for the

ZERO-MODE CONTRIBUTIONS to $\langle P P \rangle_{|\nu|}$

⇒ from the zero-modes alone we can evaluate F_π and α

α : Low Energy Constant, **specific to quenching**,
enters leading order of $\langle P P \rangle$.

$\mathcal{L}_{\text{q}\chi\text{PT}}^{(2)}$ is formulated with **aux. scalar** ϕ_0 , which supplements the quenching effect

K : couples $\phi_0 U$, plus $\frac{1}{2}\alpha_0(\partial_\mu\phi_0)^2 + \frac{1}{2}m_0^2\phi_0^2$

Count $m_0 = \mathcal{O}(\epsilon)$ and define $\alpha = \alpha_0 - 4N_c^2 K F_\pi / \Sigma$

⇒ **first order observables only involve F_π and α (dim'less)**

Pseudoscalar density $P = \bar{\psi}\gamma_5\psi$

$$\langle P(x)P(y) \rangle = N_f P_1(x, y) - N_f^2 P_2(x, y)$$

$$P_1(x, y) = \text{Tr}[i\gamma_5(D + m_q)^{-1}(x, y) \cdot i\gamma_5(D + m_q)^{-1}(y, x)]$$

$$P_2(x, y) = \text{Tr}[i\gamma_5(D + m_q)^{-1}(x, x)] \cdot \text{Tr}[i\gamma_5(D + m_q)^{-1}(y, y)]$$

Spectral decomposition of propagators \rightarrow residuum given by zero modes:

$$\lim_{m_q \rightarrow 0} (m_q V)^2 \langle P(x) P(0) \rangle_\nu = N_f C_{|\nu|}^{(1)}(x) + N_f^2 C_{|\nu|}^{(2)}(x)$$

$$\text{connected : } C_{|\nu|}^{(1)}(x) = -\langle v_j^\dagger(x) v_k(x) \cdot v_k^\dagger(0) v_j(0) \rangle_\nu$$

$$\text{disconnected : } C_{|\nu|}^{(2)}(x) = \langle v_j^\dagger(x) v_j(x) \cdot v_k^\dagger(0) v_k(0) \rangle_\nu$$

summed over (exact) **zero modes**, $Dv_j = 0$.

Spatial average: $\frac{1}{V} \int d^3x P(x) P(0) \rightarrow C_{|\nu|}^{(i)}(t)$

Fits for $C_{|\nu|}^{(i)}(t)$ are troublesome; more handy:

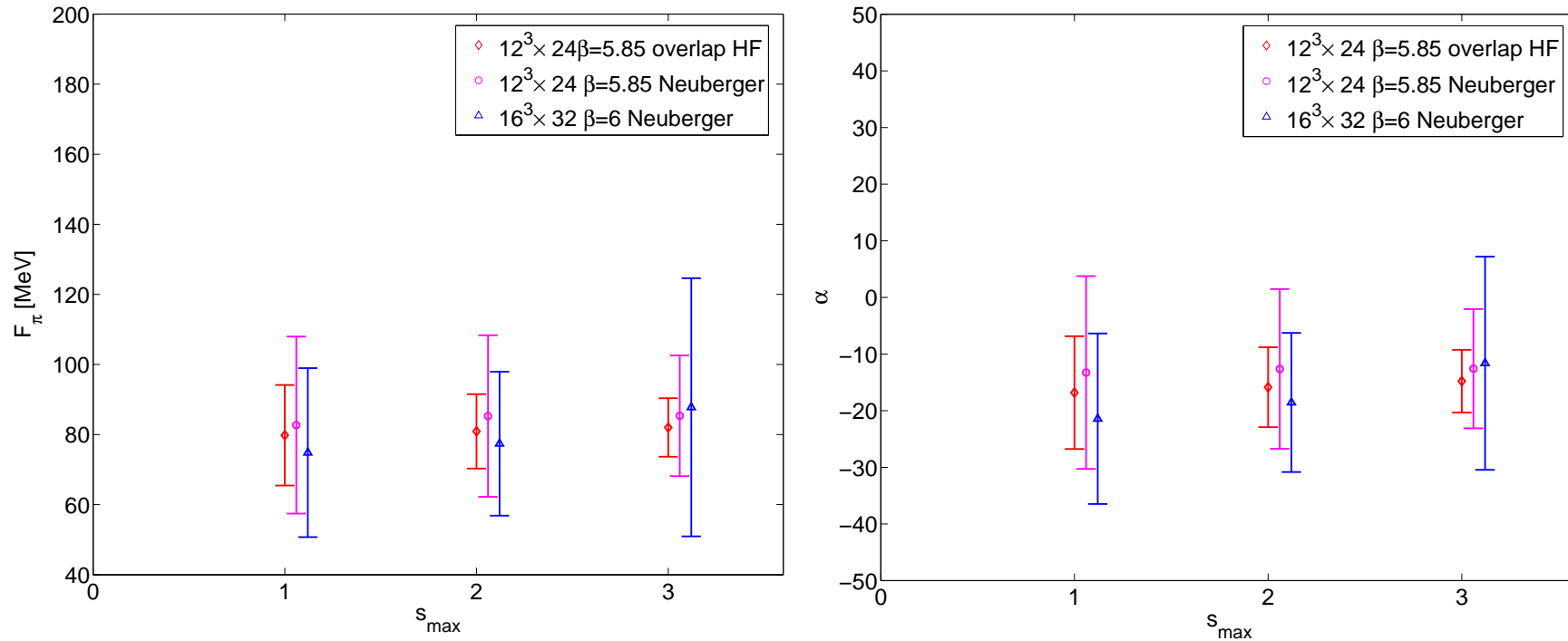
$$\frac{1}{L^2} \frac{d}{dt} C_{|\nu|}^{(i)}(t) |_{t=T/2} = D_{|\nu|}^{(i)} \cdot s + \mathcal{O}(s^3), \quad s = t - \frac{T}{2}$$

Measure slopes $D_{|\nu|}^{(i)}$ up to $s = \pm 1, \pm 2, \dots$, fits yield values for \mathbf{F}_π, α

Combined fit for F_π and α for zero-mode pseudoscalar correlators of **standard overlap operator** and the **overlap-HF operator** in $V = (1.48 \text{ fm})^3 \times 2.96 \text{ fm}$, $\langle \nu^2 \rangle$ as measured in each case.

Statistics :

Dirac operator	β	lattice size		$ \nu = 1$	$ \nu = 2$
standard overlap	5.85	$12^3 \times 24$		132	115
overlap HF	5.85	$12^3 \times 24$		221	192
standard overlap	6	$16^3 \times 32$		115	94



s : fitting range around $T/2$.

In particular the **overlap HF** yields a **neat plateau**.

$F_\pi = (80 \pm 14) \text{ MeV}$, $\alpha = -17 \pm 10$. Standard overlap data consistent.

F_π moves down compared to $\langle AA \rangle$ (theory: 86 MeV at $m_q = 0$).

Giusti et al. in $V \simeq (1.49 \text{ fm})^4$: $F_\pi \simeq (117 \pm 16) \text{ MeV}$, $\alpha \in [-1.8, 7.8]$

Conclusions

- $D_{\text{ov-HF}}$ provides **better locality** than $D_{\text{ov-W}}$ \rightarrow applicable on **coarser lattices**
- **Rotation symmetry** improved
Scaling promising (toy models, preliminary QCD results), under investigation
- p -regime : results for m_π , m_ρ , m_{PCAC} , F_π vs. m_q similar to standard overlap \rightarrow independent confirmation, Z_A much closer to 1
- ϵ -regime : $\Sigma = (268(2) \text{ MeV})^3$ from Dirac spectrum
 $F_\pi \simeq 104(\pm 9) \text{ MeV}$ from $\langle A_4 A_4 \rangle$ (*preliminary*)
agrees exactly with χ extrapolation from p -regime
0-mode result at $m_q = 0$: $F_\pi = (80 \pm 14) \text{ MeV}$, $\alpha = -17 \pm 10$
Topology conserving gauge action may be helpful

Thanks to M. Papinutto and C. Urbach for numerical tools

Modified lattice gauge action to “freeze” the topol. sector :

Hernández/Jansen/Lüscher '99, Neuberger '00

Theoretical condition for topology conservation:

$$S_P := 1 - \frac{1}{3} \text{Re Tr}(U_P) < \varepsilon \simeq 1/20.5$$

U_P : any plaquette

Implementation (Lüscher '01, Fukaya/Onogi '03)

$$S_\varepsilon(U_P) = \begin{cases} \frac{S_P(U_P)}{1 - S_P(U_P)/\varepsilon} & S_P(U_P) < \varepsilon \\ +\infty & \text{otherwise} \end{cases}$$

W.B./Jansen/Nagai/Necco/Scorzato/Shcheredin '04

see also Fukaya/Hashimoto/Hirohashi/Ogawa/Onogi, '05

Relax ε , decrease β for \approx const. scale, (local HMC)

$1/\varepsilon$	β	r_0/a	β_W	$\tau_{\text{aut}}^{\text{top}} / \tau_{\text{aut}}^{\text{plaq}}$
0	6.18	7.14(3)	6.18	$:= 1$
1	1.5	6.6(2)	6.13	~ 29
1.18	1	7.2(2)	6.18	~ 80
1.52	0.3	7.3(4)	6.19	~ 241

Approx. const. scale r_0/a on a 16^4 lattice

β_W : corresponding β value for the Wilson gauge action

τ_{aut} : autocorrelation time with respect to Q_{top} , and to the plaquettes

Allows us to sample a specific top. sector, desired in the ε -regime.

Also condition number for D_{ov} decreases.