

Orbital-Selective Pairing and Gap Structures of Iron-Based Superconductors

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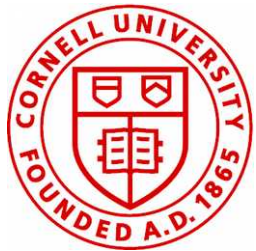
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Peter O. Sprau, ..., A. Kreisel, et al.

arXiv:1611.02134

A. Kreisel, et al.

arXiv:1611.02643



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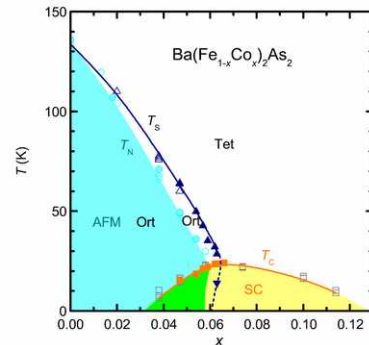
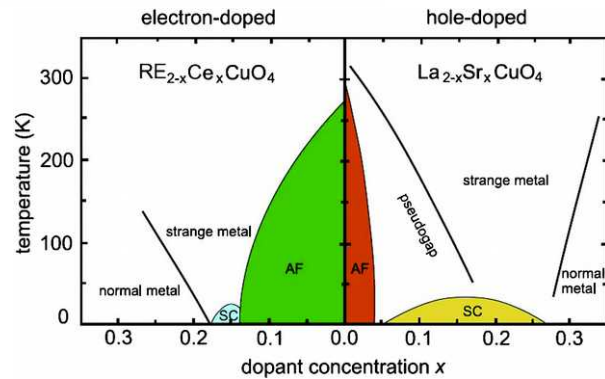
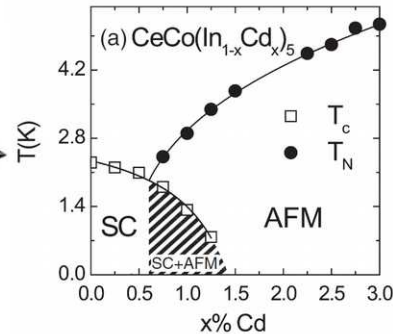
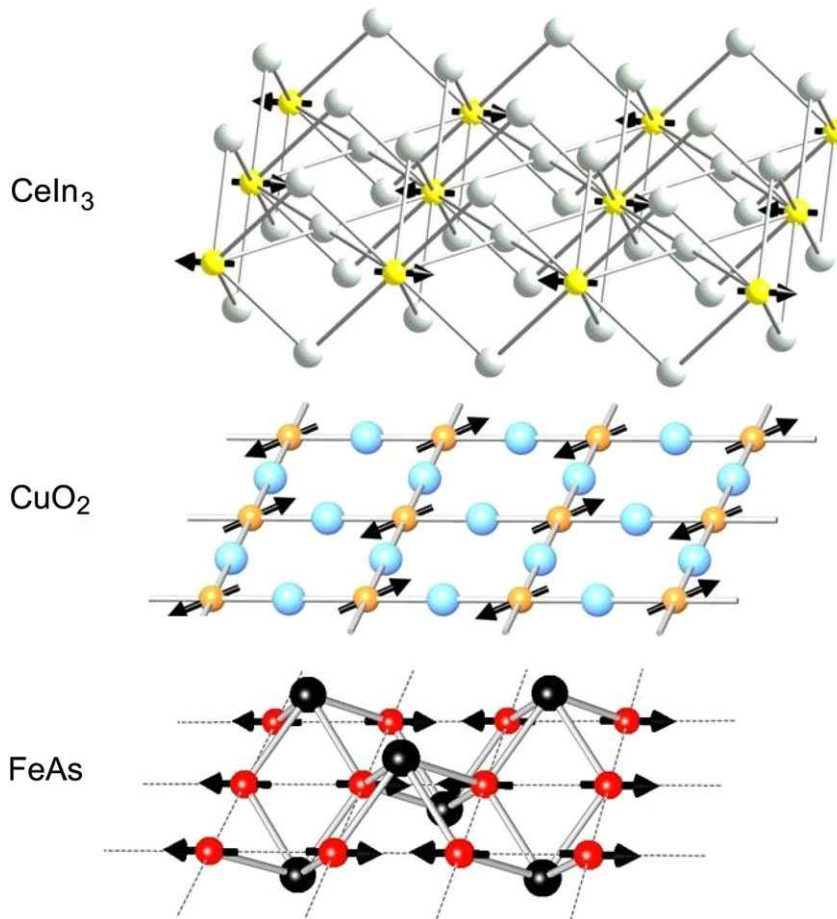
Outline

- Unconventional superconductivity
 - some “big questions”
- Example of Fe-based SC: FeSe (bulk)
 - Why interesting?
 - Nematic order
 - Superconducting gap structure
 - Orbital selectivity
- Other materials: FeSe (monolayer), LiFeAs

Unconventional superconductivity

2d active layers, antiferromagnetic spin orders (undoped compounds)

Phase diagram: close proximity of antiferromagnetism and superconductivity

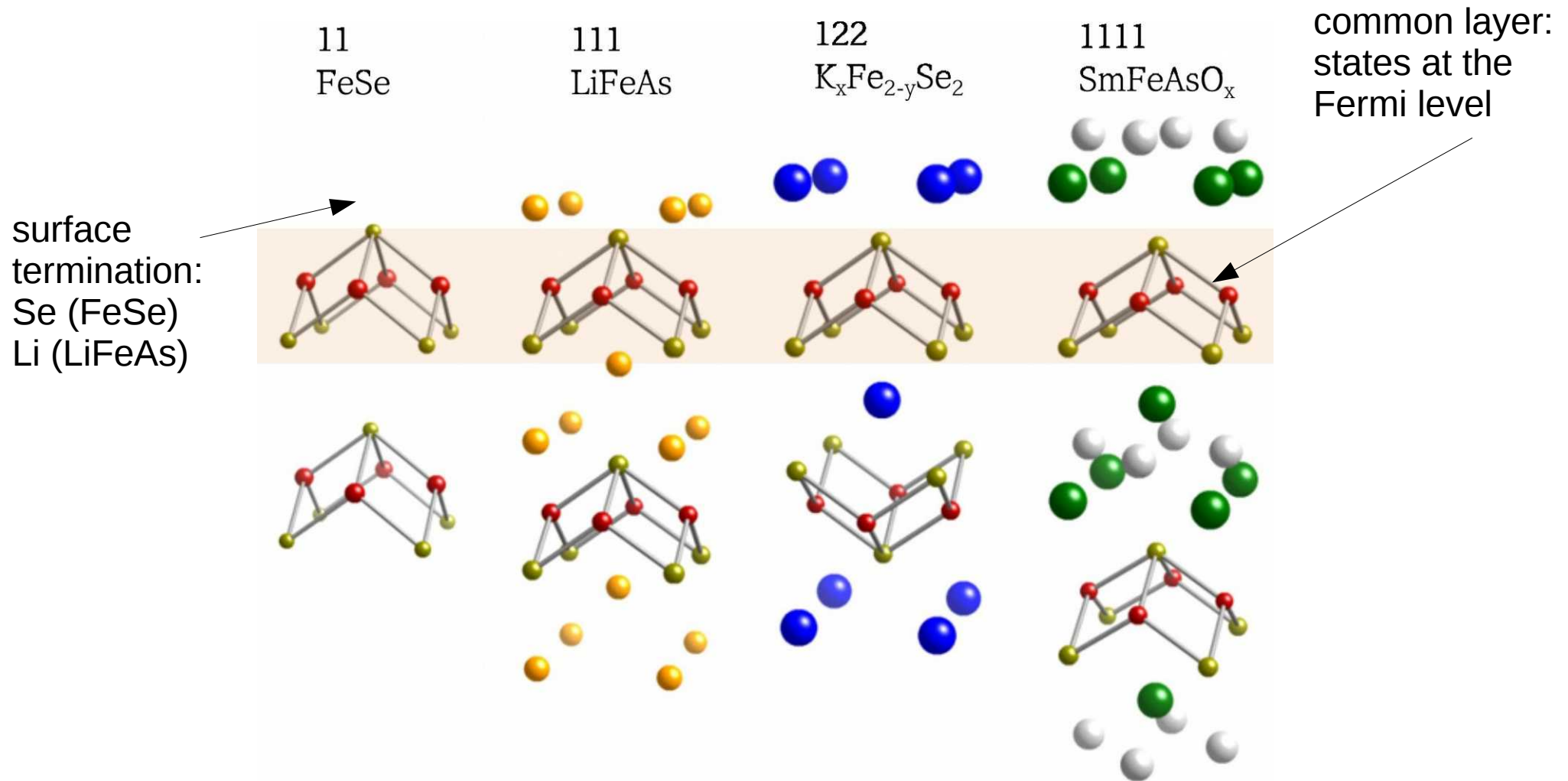


other interesting phases:
charge density wave
C4 symmetric magnetic phases
Nematic phase

Similarities in physical mechanisms?

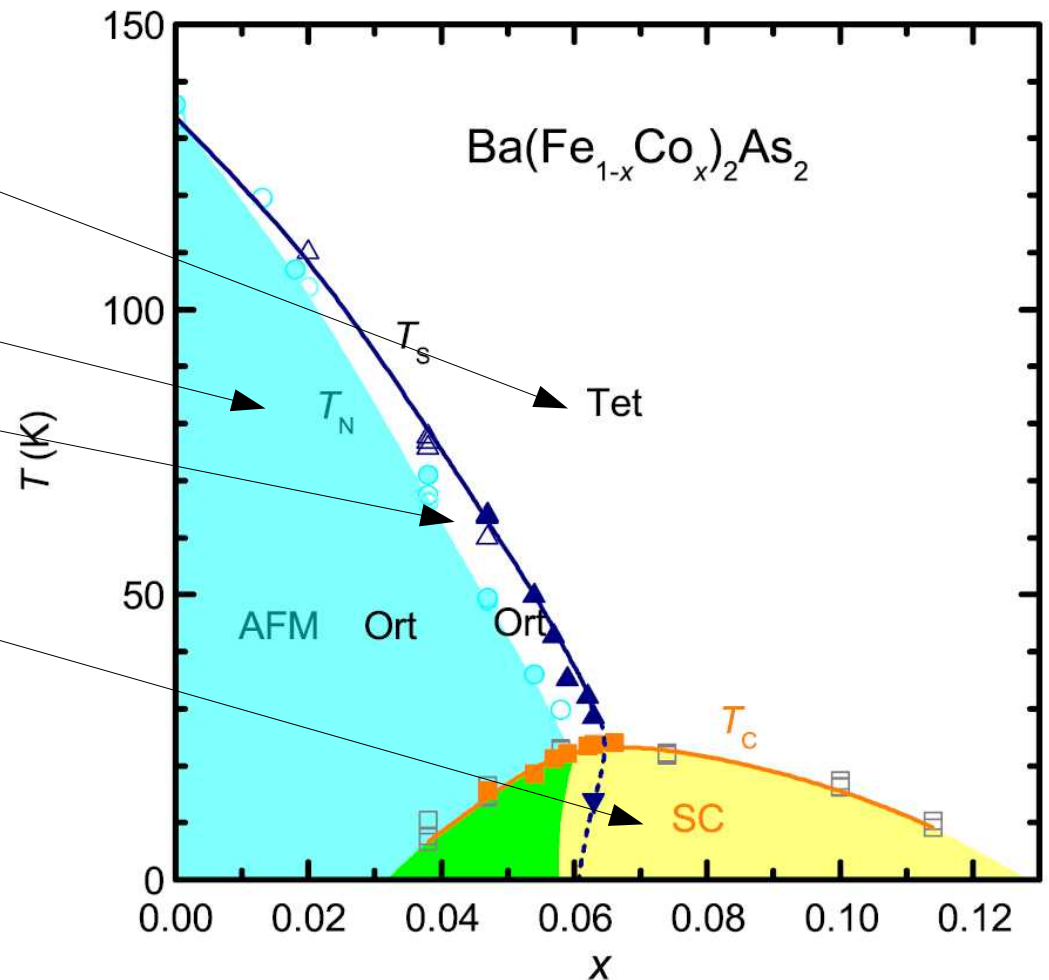
Fe based superconductors

- 2d layered materials: active states Fe(d)



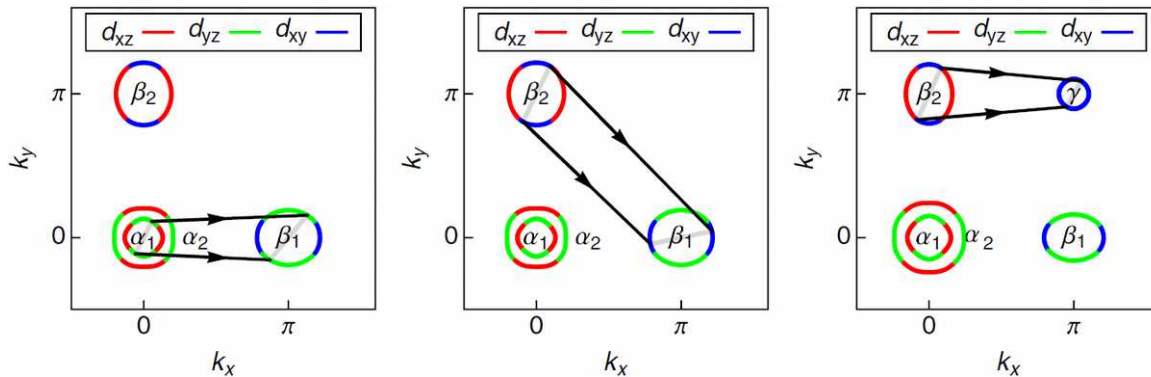
Fe based superconductors

- Band structure
- Magnetism
- Nematicity
- Superconductivity (gap structure)

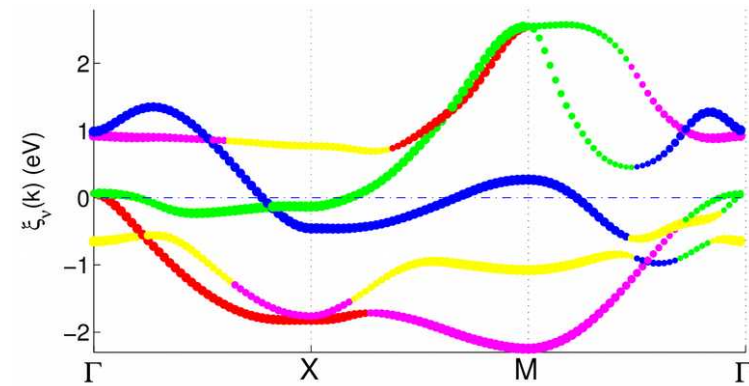


Band structure / Magnetism

- Fe(d) orbitals: quasi 2D electronic structure (5 orbital model)

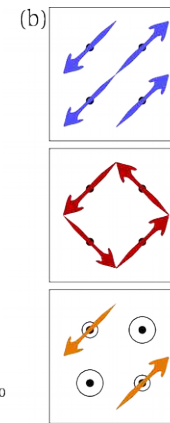
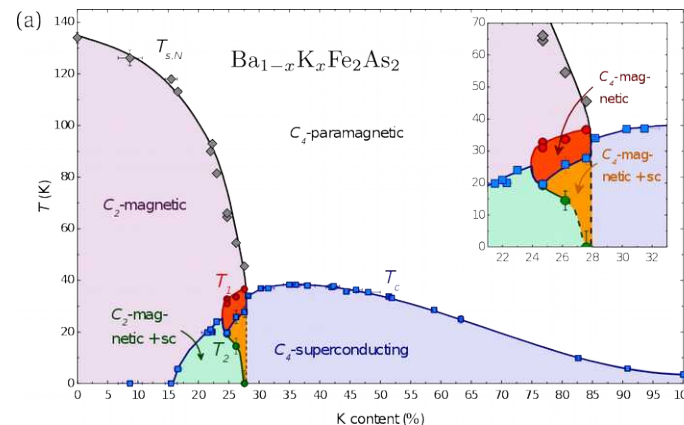
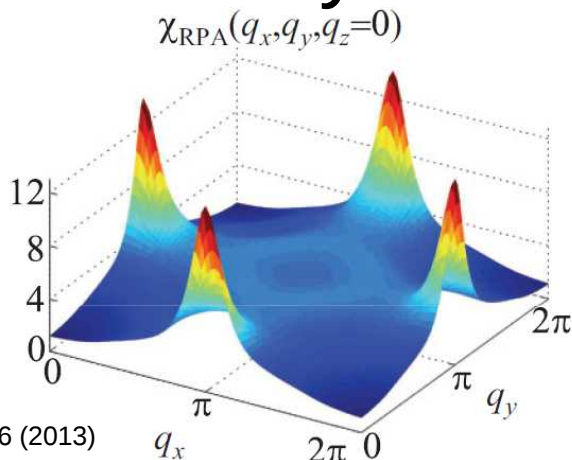


D. J. Scalapino Rev. Mod. Phys. **84**, 1383 (2012)



Chi et al., PRB **94**, 134515 (2016) [LiFeAs]

- Tendency towards $(\pi,0)$, $(0,\pi)$ magnetism

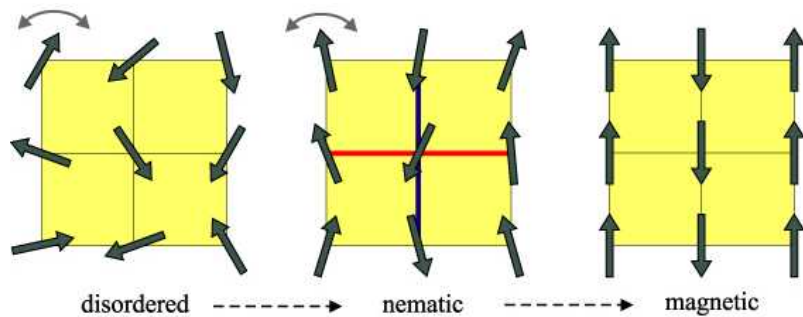


Böhmer, et al., Nat. Commun. **6**, 7911 (2015)
Gastiasoro, Andersen, PRB **92**, 140506 (2015)

Wang, et al., PRB **88**, 174516 (2013)

Nematicity

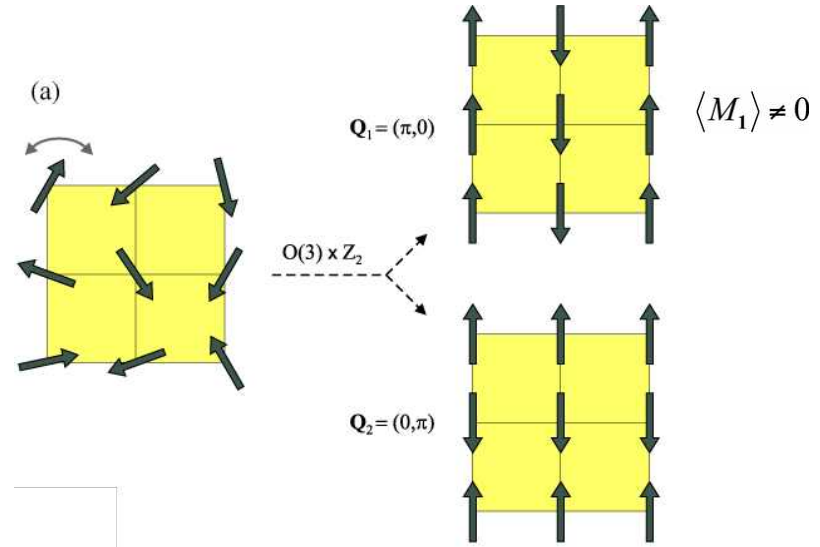
- Magnetic ordering
- Nematic state



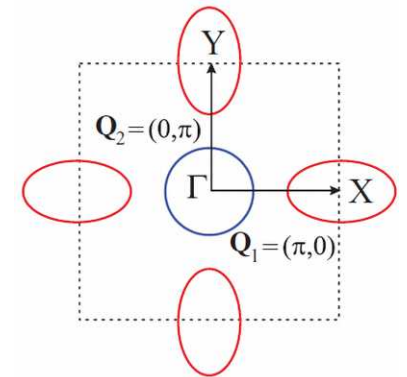
$$\langle M_1^2 \rangle \neq \langle M_2^2 \rangle$$

$$\langle M_1 \rangle = \langle M_2 \rangle = 0$$

$$\langle M_1 \rangle \neq 0$$

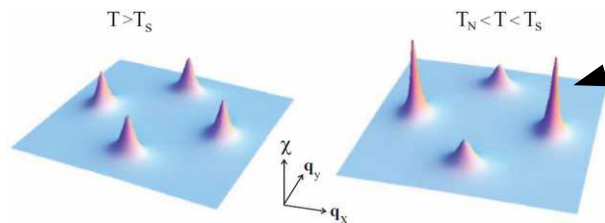


Magnetic fluctuations stronger in x-direction
Tetragonal symmetry breaking



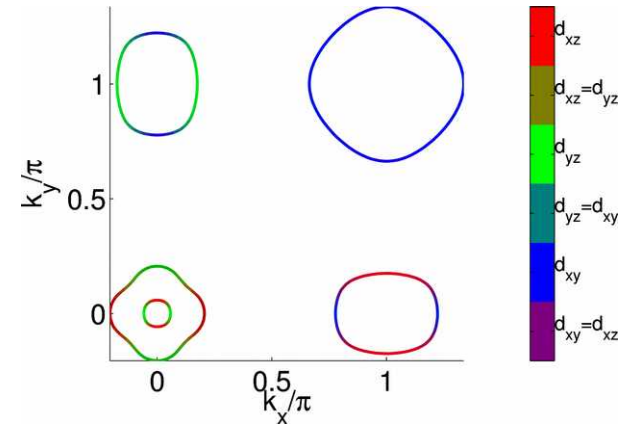
Itinerant approach to nematic state

$$\chi_{\text{nem}} \sim \chi_{\text{mag}}^2$$

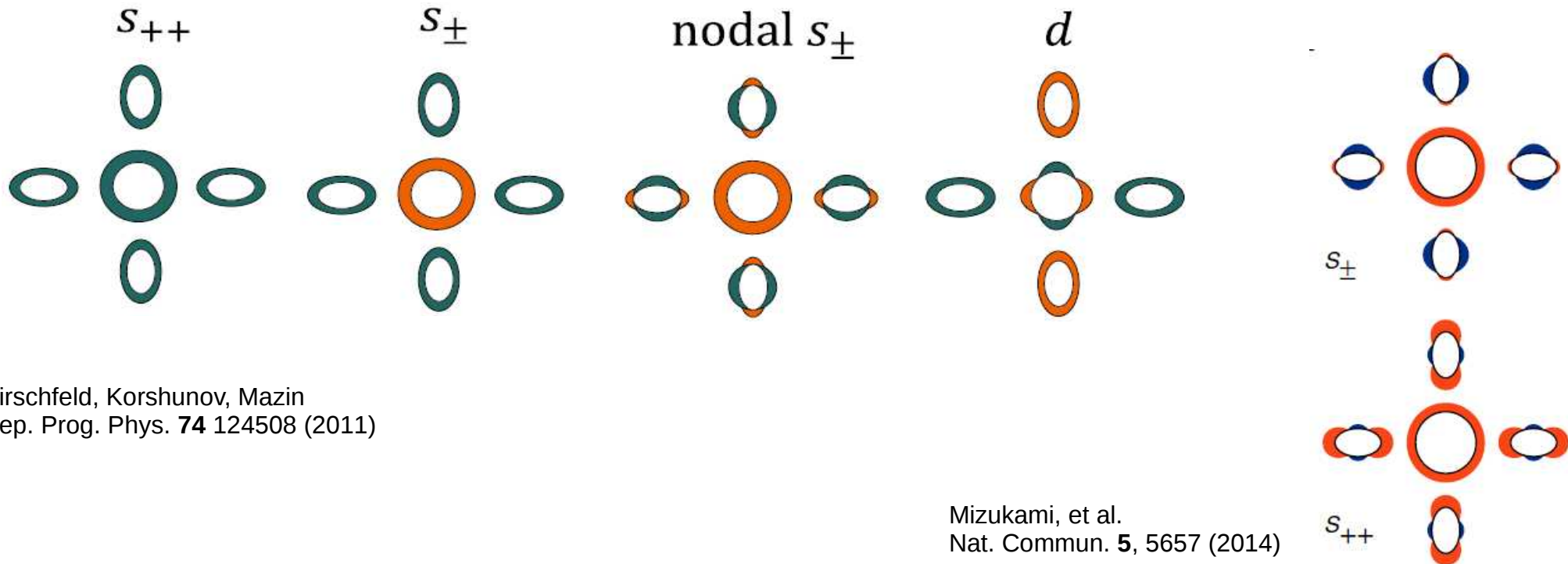


Superconductivity: gap structure

- Fermi surface



- Possible order parameters

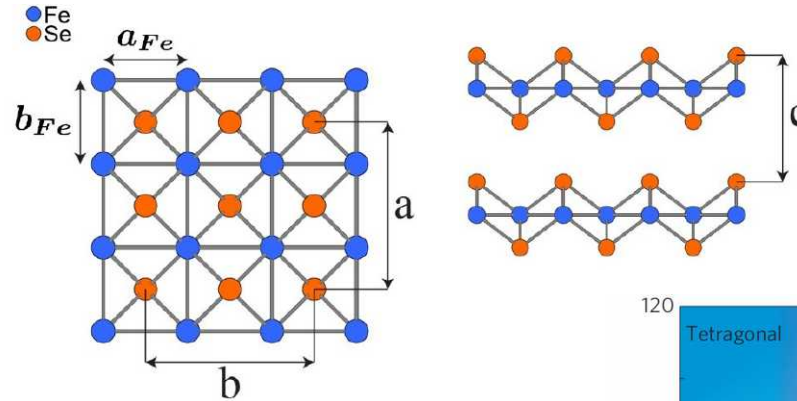


Hirschfeld, Korshunov, Mazin
Rep. Prog. Phys. **74** 124508 (2011)

Mizukami, et al.
Nat. Commun. **5**, 5657 (2014)

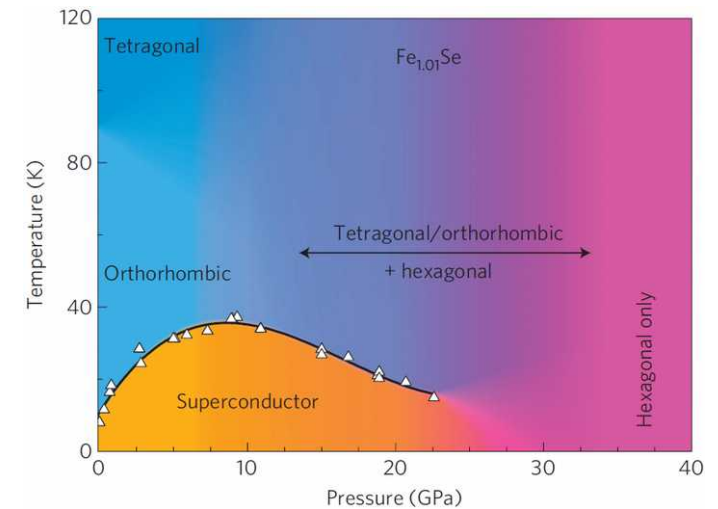
FeSe: Why interesting?

- 11 compound

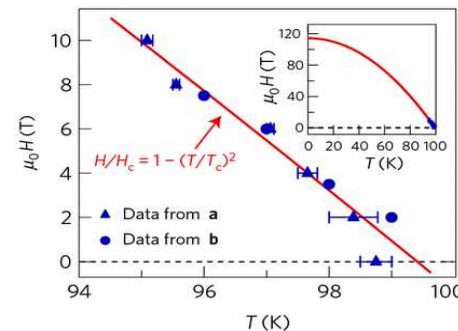
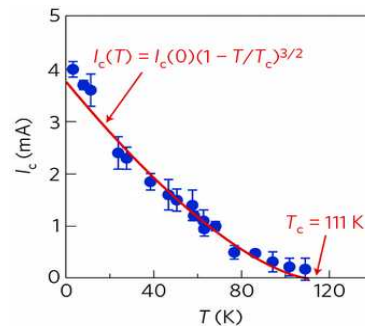
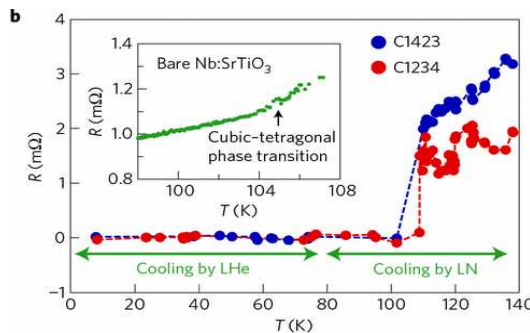


- T_c 8K, under pressure ~ 40 K

Medvedev, et al. Nat. Mater. **8**, 630 (2009)



- T_c 100K (single layer) transport measurement



Ge et al. Nat. Mater. **14**, 285 (2015)

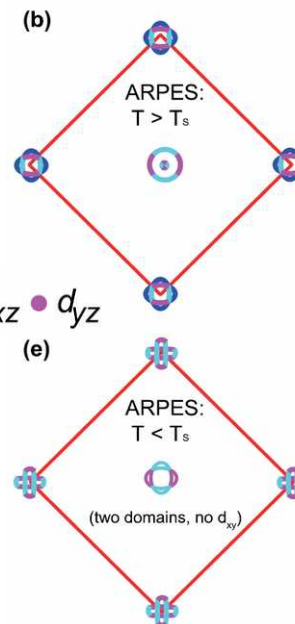
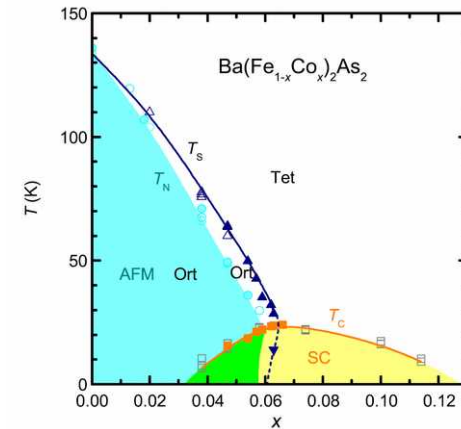
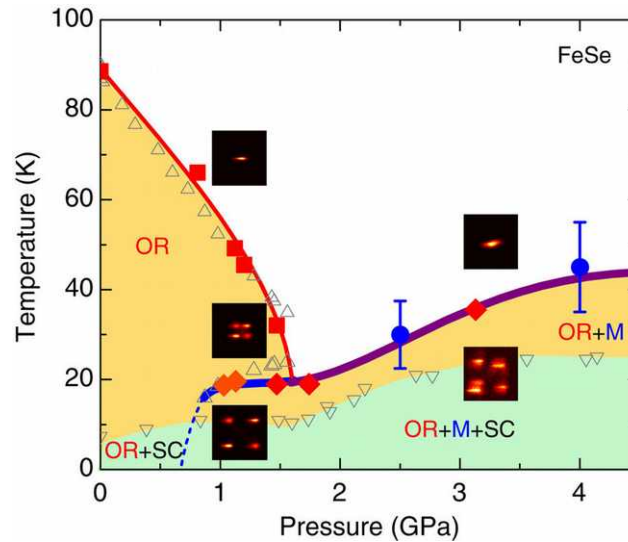
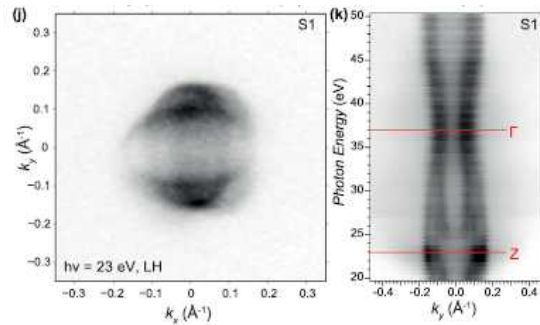
FeSe: Why interesting?

- nematic phase
no magnetism
($p=0$)

K. Kothapalli, et al.,
Nat. Commun. 7, 12728 (2016)

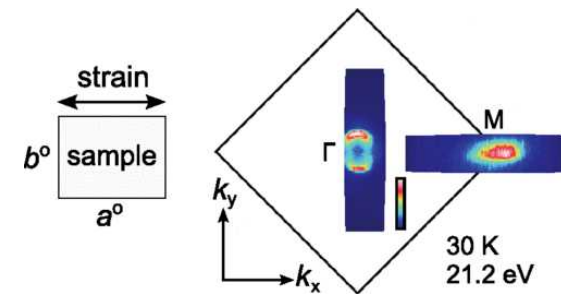
- ARPES

measured band structure
tiny Fermi surface
(far from ab initio results)



Measured
orbital
splitting

detwinned ARPES



Watson, et al., PRB 94, 201107(R) (2016)
Watson, et al., PRB 90, 121111(R) (2014)
Suzuki, et al., PRB 92, 205117 (2015)
Maletz, et al., PRB 89, 220506(R) (2014)
Fedorov, et al., Sci. Rep. 6, 36834 (2016)

FeSe

- Origin of nematic order

- Orbital order

Baek et al., Nat. Mat. **14**, 210 (2015)

Yamakawa, Onari, Kontani, Phys. Rev. X **6**, 021032 (2016)

- Quantum paramagnet

Wang, Kivelson, Lee, Nat. Phys. **11**, 959 (2015)

- Spin ferroquadrupolar / antiferroquadrupolar order

Wang, Hu, Nevidomskyy PRL **116**, 247203 (2016)

Lai, et al., arXiv:1603.03027

- Longer-range Coulomb interactions

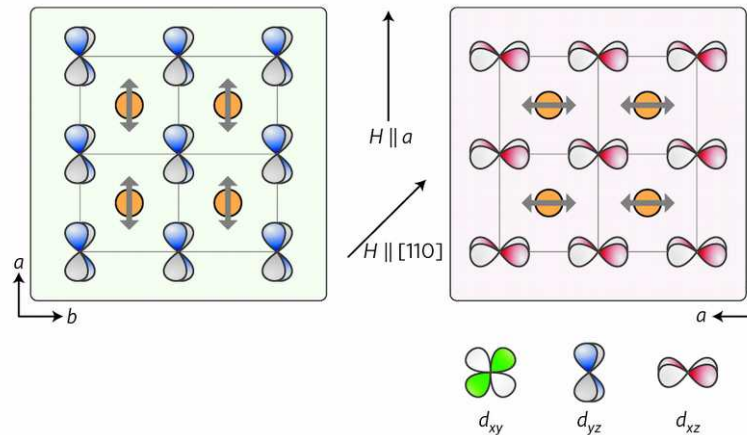
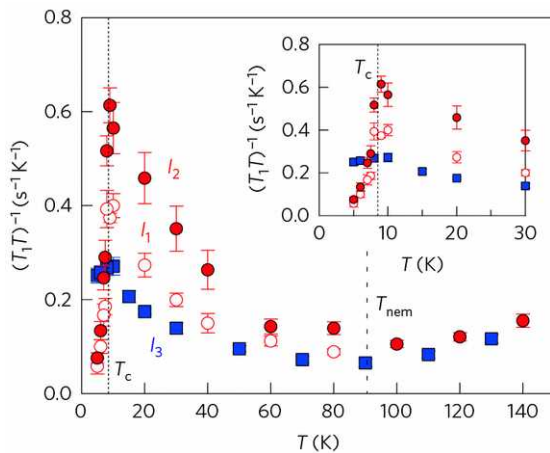
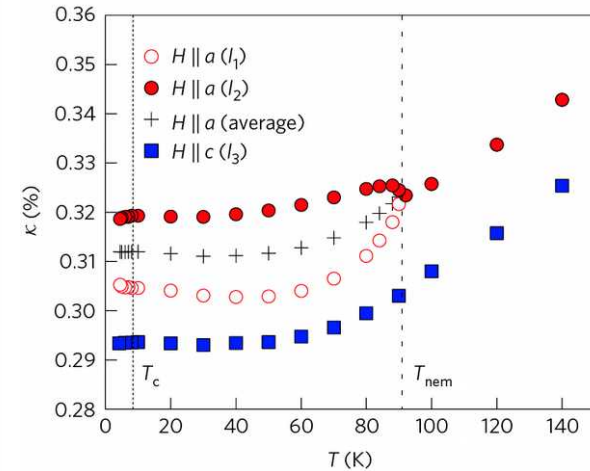
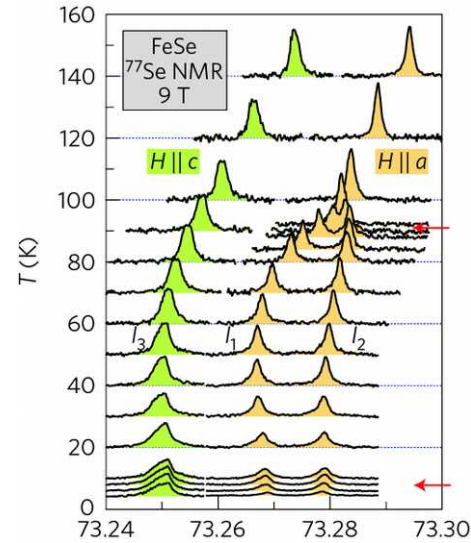
Jiang, et al., PRB **93**, 115138 (2016)

- Competition between magnetism and charge current order

Chubukov, Fernandes, Schmalian, PRB **91**, 201105(R) (2015)

Nematic order

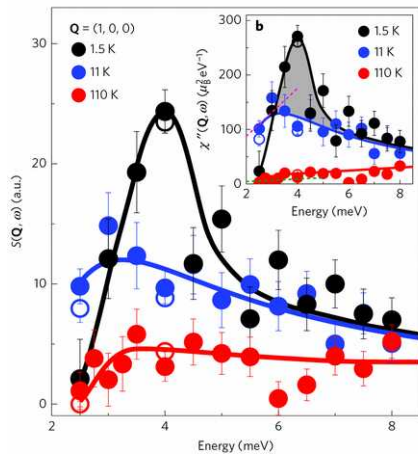
- NMR on FeSe
splitting of the Se line
- No enhanced spin fluctuations below T_s



→ not caused by lattice distortion
→ evidence for orbital order (ferro)

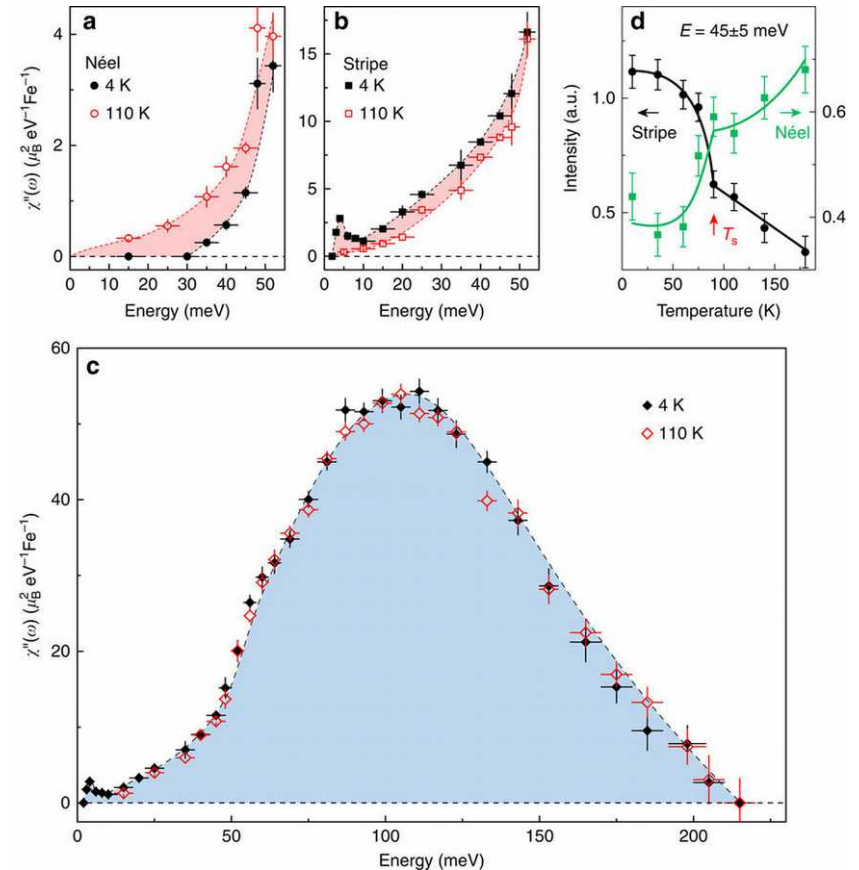
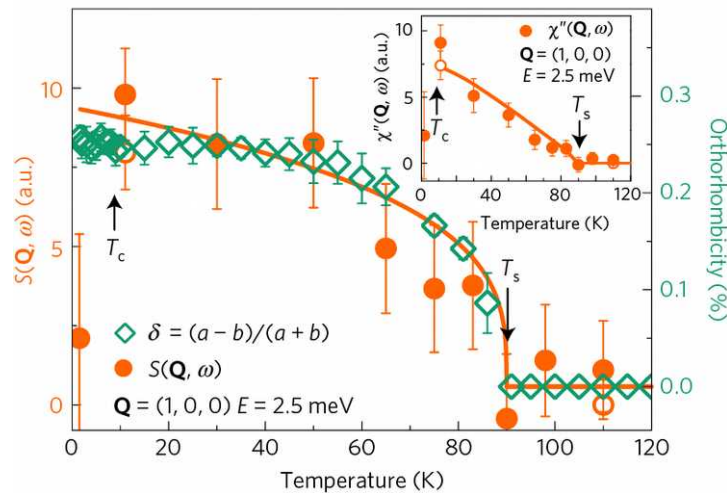
Spin excitations and magnetism

- Competition between stripe and Néel fluctuations



"spin resonance" at $(\pi, 0)$

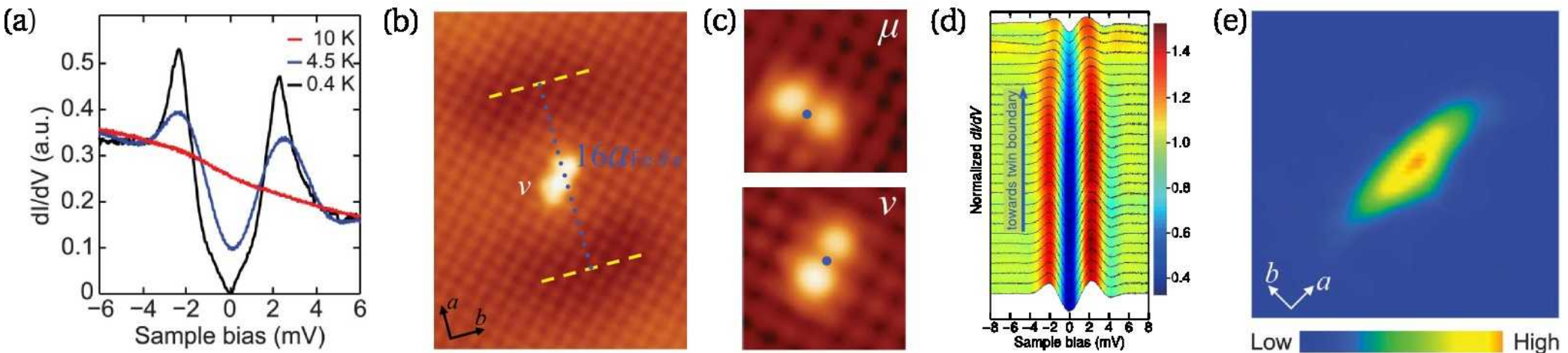
Wang et al., Nat. Mat., 15, 159 (2016)



Wang et al., Nat. Commun. 7, 12182 (2016)

Superconducting gap structure

- consequences: nodal gap structure, anisotropy



Song et al. PRL 109, 137004 (2012)

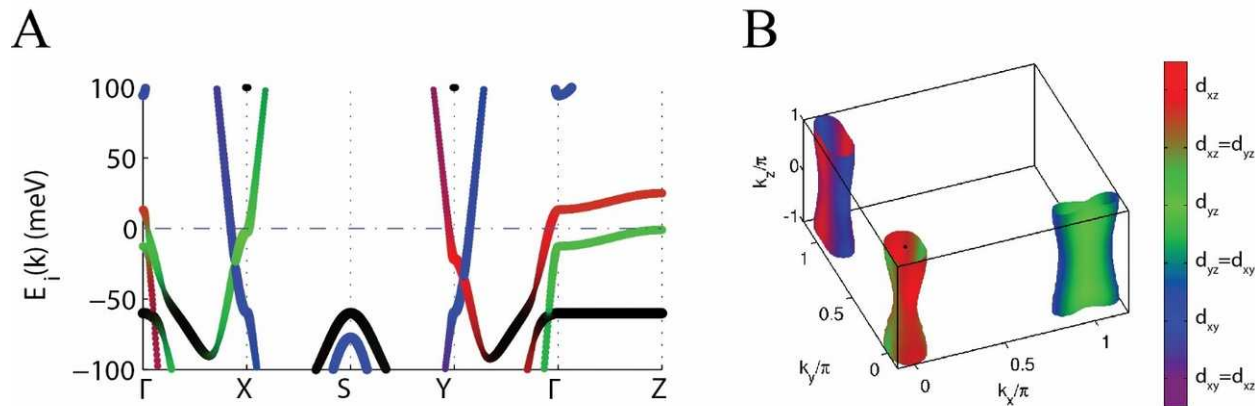
Song et al. Science 332, 1410 (2011)

Modelling

- Band structure:
measured spectral function

$$G(\vec{k}, \omega) = \frac{1}{\omega - E_{\vec{k}} + i0^+}$$

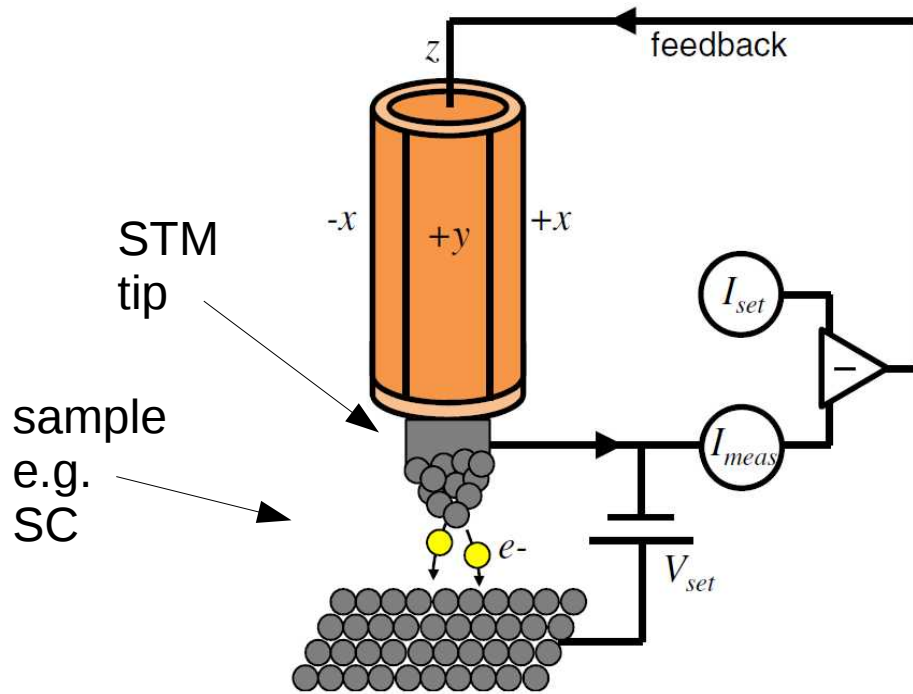
$$A(\vec{k}, \omega) = -\frac{1}{\pi} \text{Im}G(\vec{k}, \omega)$$



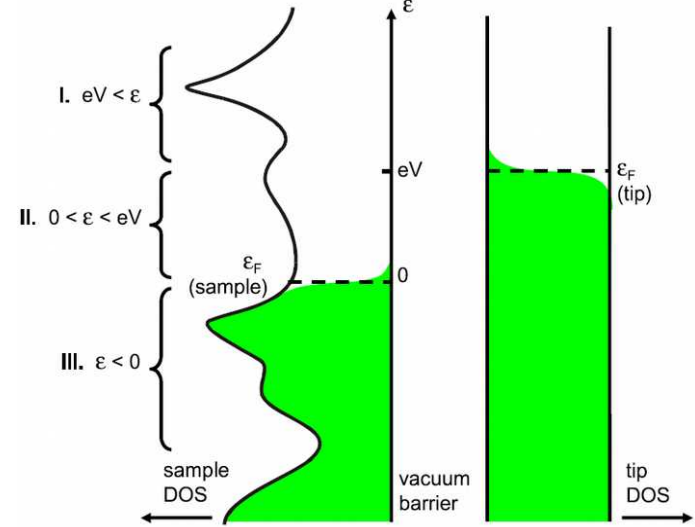
- ARPES
- Quantum oscillations
- Scanning tunnelling microscopy

Watson, et al., PRB 94, 201107(R) (2016)
 Watson, et al., PRB 90, 121111(R) (2014)
 Suzuki, et al., PRB 92, 205117 (2015)
 Maletz, et al., PRB 89, 220506(R) (2014)
 Fedorov, et al., Sci. Rep. 6, 36834 (2016)

Scanning tunnelling microscopy



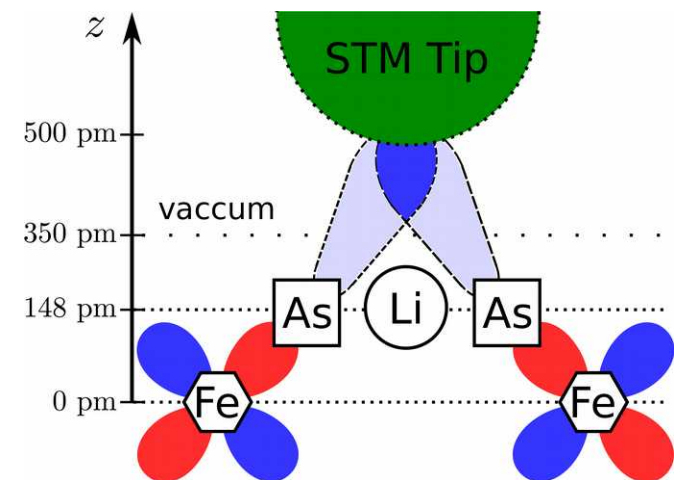
J. Hoffman Rep. Prog. Phys. **74** 124513 (2011)



Tunnelling current:

$$I(V, x, y, z) = -\frac{4\pi e}{\hbar} \rho_t(0) |M|^2 \int_0^{eV} \rho(x, y, z, \epsilon) d\epsilon$$

Local Density Of States (LDOS)
of sample at given energy **at the tip position**

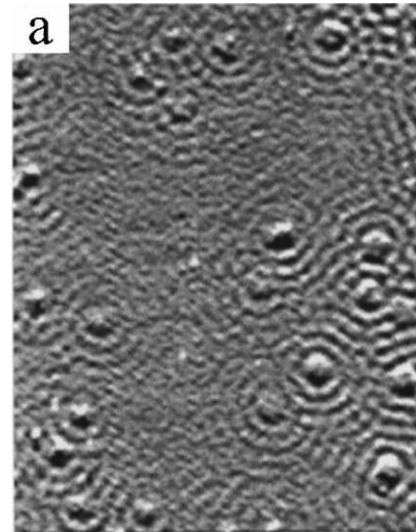


A. Kreisel, et al., Phys. Rev. B **94**, 224518 (2016)

J. Tersoff and D. R. Hamann, PRB **31**, 805 (1985)

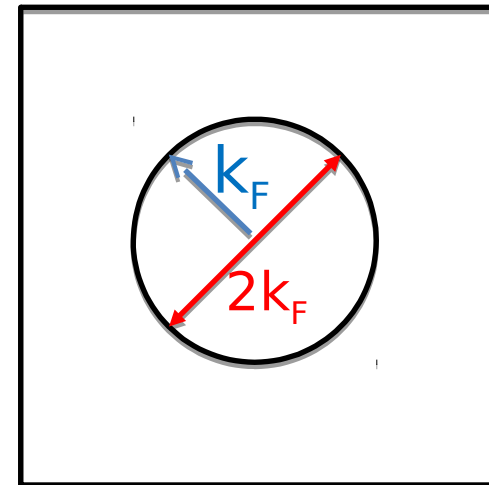
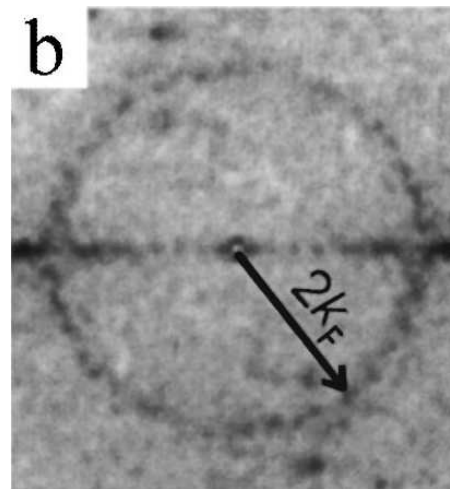
Quasiparticle Interference (QPI)

- STM on normal metal (Cu)
 - impurities
 - Friedel oscillations



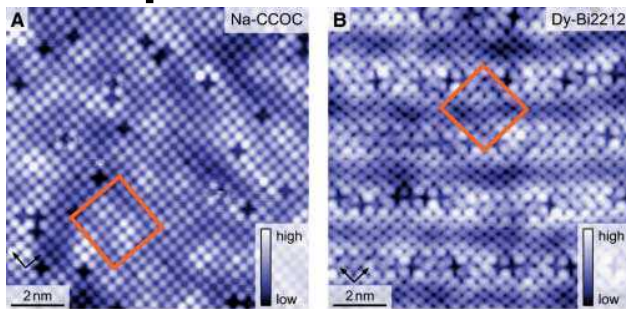
L. Petersen, et al.
PRB **57**, R6858(R)
(1998)

- Fourier transform of conductance map
 - mapping of constant energy contour

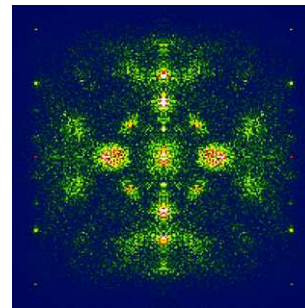


QPI in superconductors

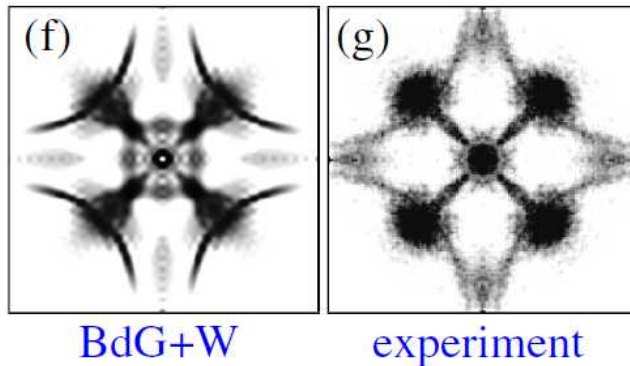
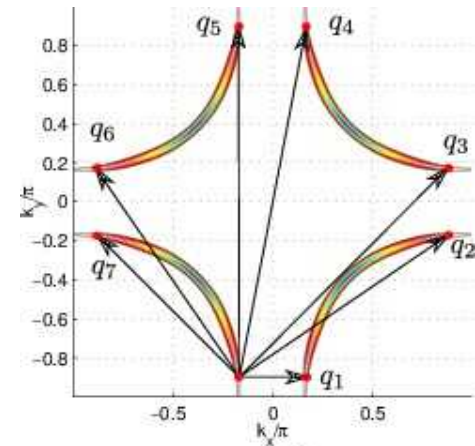
- Fourier transform of differential conductance maps



FT →

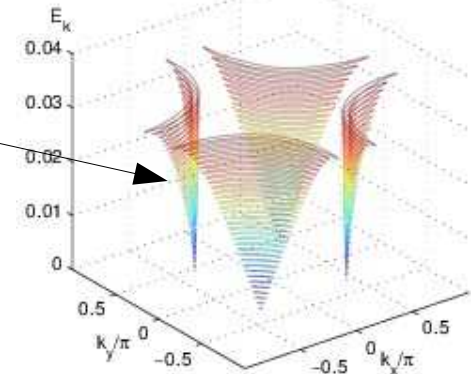


K Fujita et al. Science **344**, 612 (2014)



$$E_k = \pm \sqrt{\epsilon_k^2 + \Delta_k^2}$$

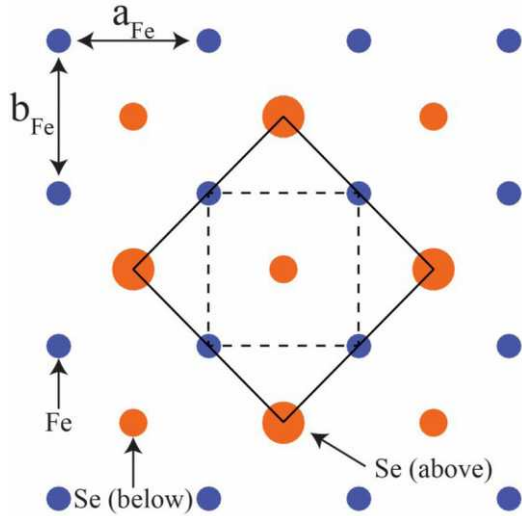
A. Kreisel, et al., PRL **114**, 217002 (2015)



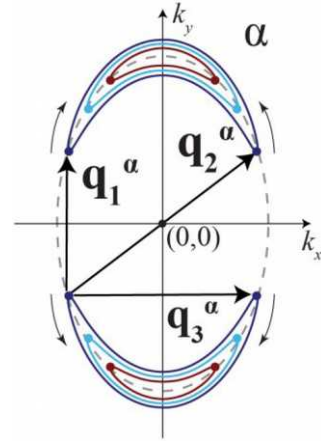
Trace back Fermi surface+measure
superconducting gap function

octet model: 7 scattering
vectors between regions
of high DOS

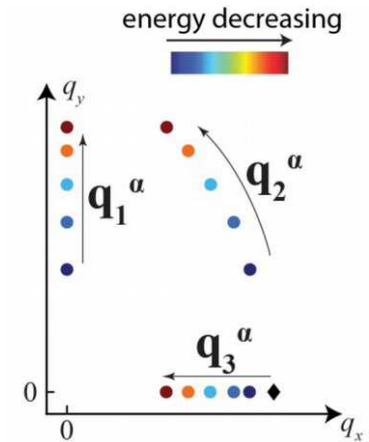
FeSe BQPI



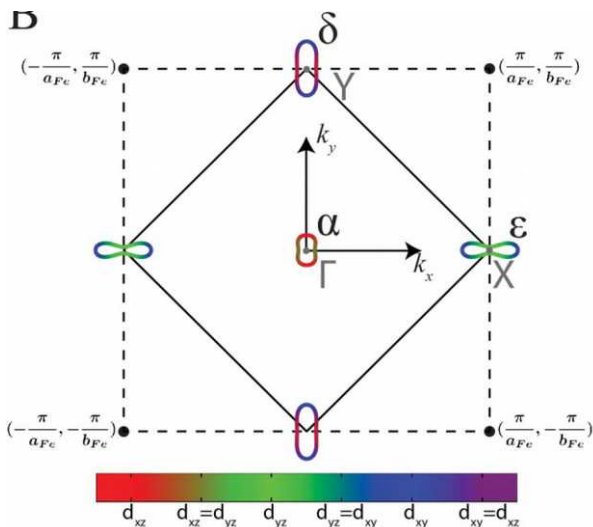
Coordinate system, expected Fermi surface



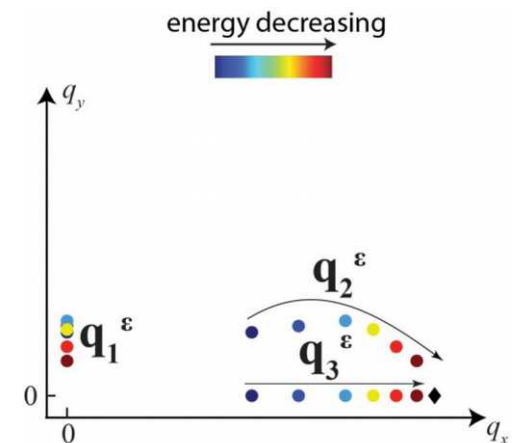
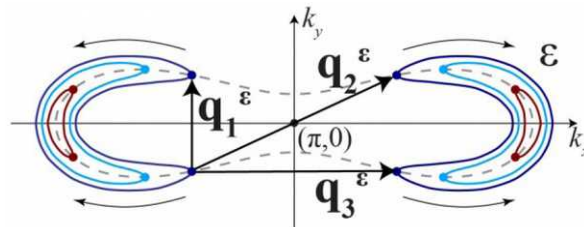
CEC: constant energy contour
Expected scattering vectors



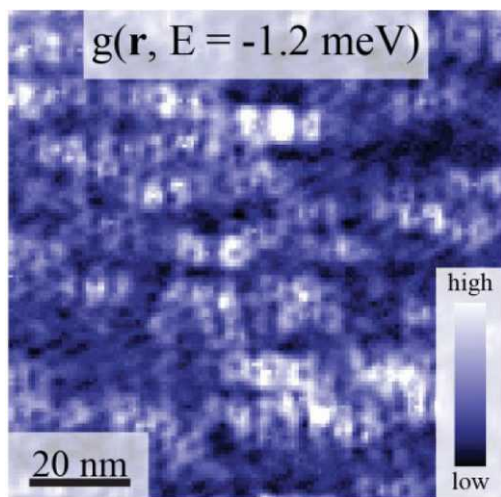
Dispersion of QPI peaks
 $q(E) \rightarrow k(E) \rightarrow E(k)$



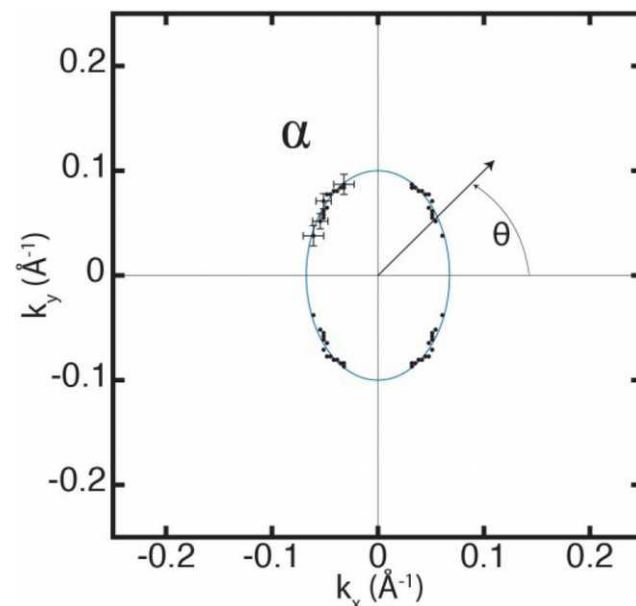
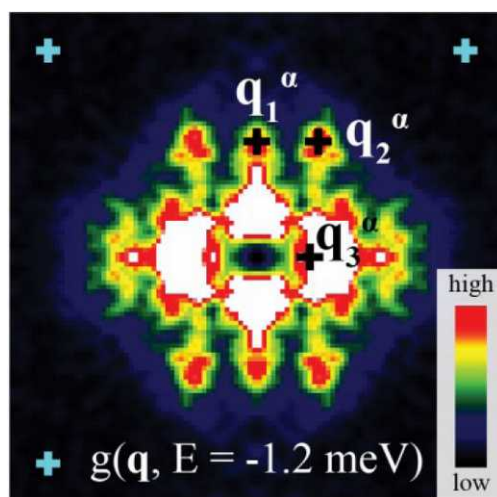
$$E_k = \pm \sqrt{\epsilon_k^2 + \Delta_k^2}$$



FeSe measurement

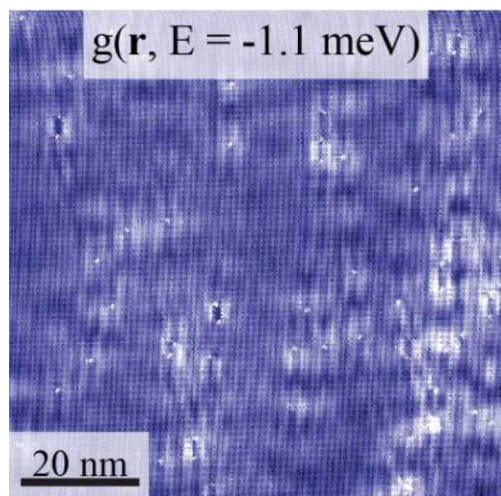


FT

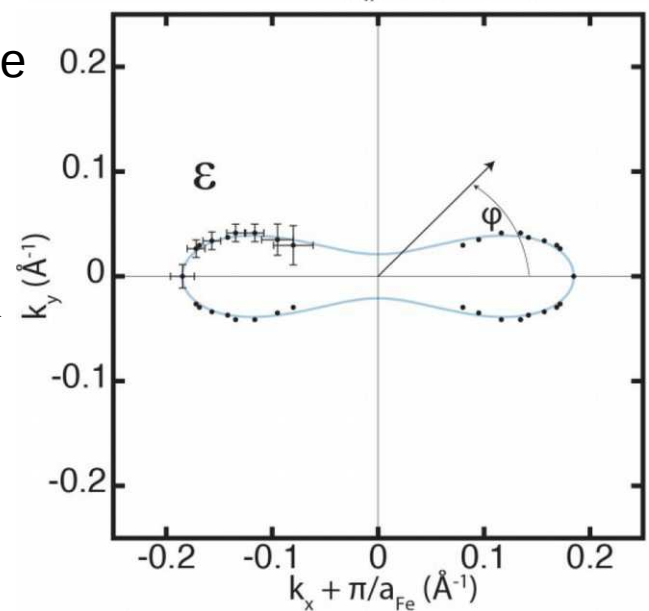
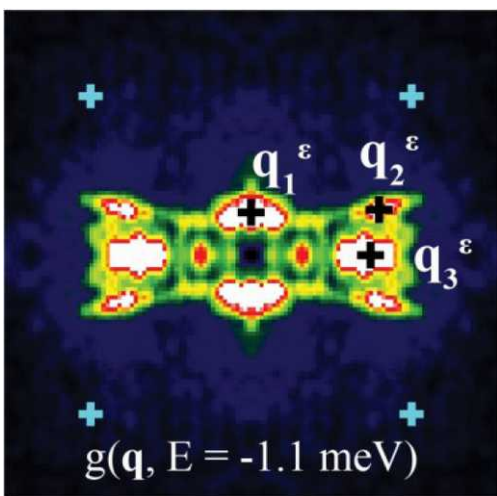


Trace back
Fermi
surface

Conductance maps



FT



Band structure modelling

- Tight binding model

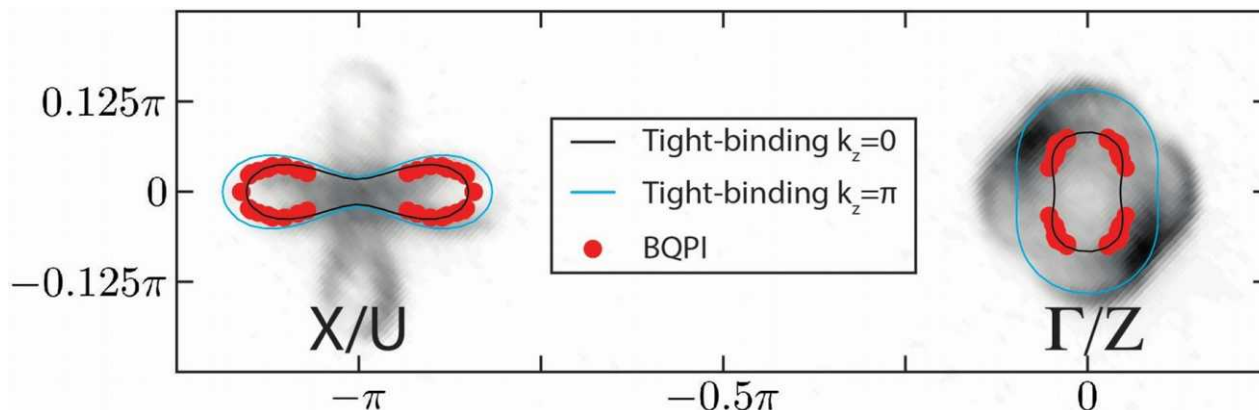
$$H_N = H_0 + H_{OO} + H_{SOC}$$

$$H_0 = \sum_{r,r',a,b} t_{r-r'}^{ab} c_{a,r}^\dagger c_{b,r'}$$

site+bond
centered
orbital order

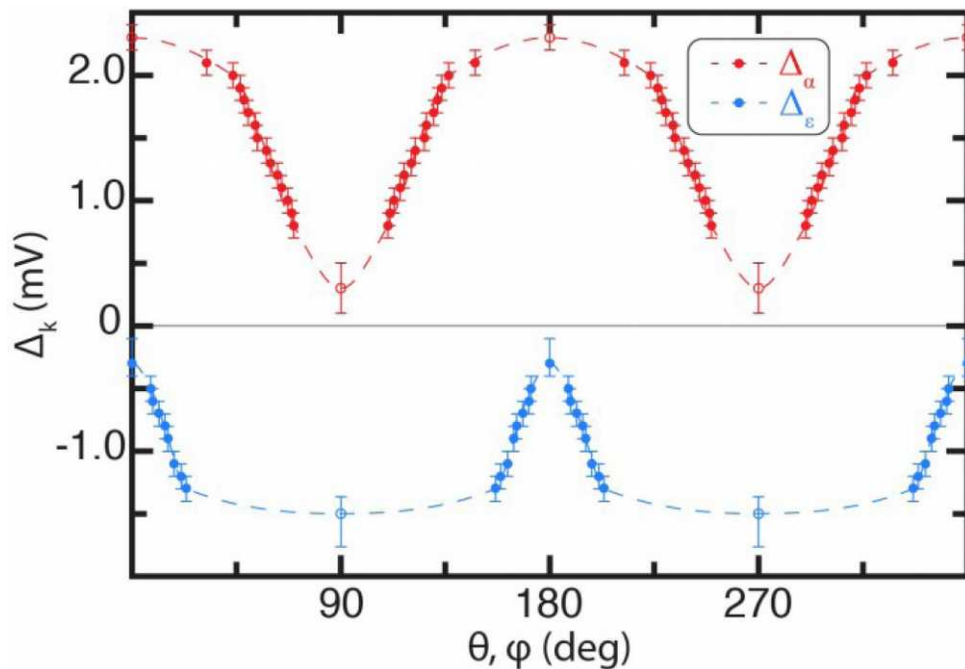
$$H_{SOC} = \lambda \mathbf{L} \cdot \mathbf{S}$$

needed for consistent
splitting at Gamma

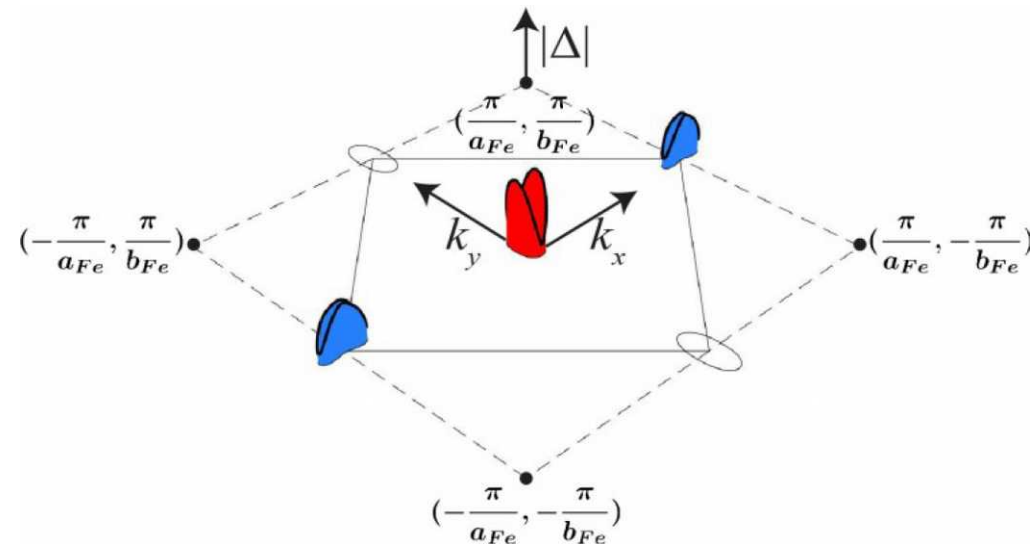
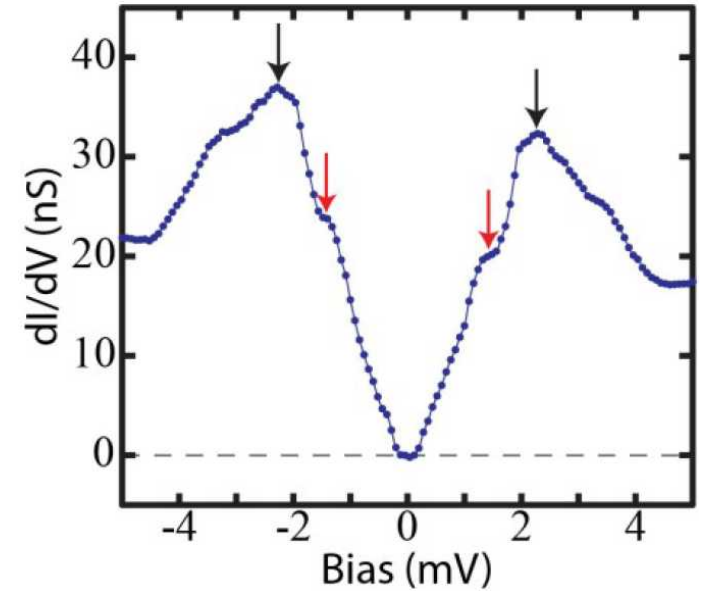


Superconducting gap

- highly anisotropic order parameter, 2 band
- “antiphase” oscillation



Sign change?!

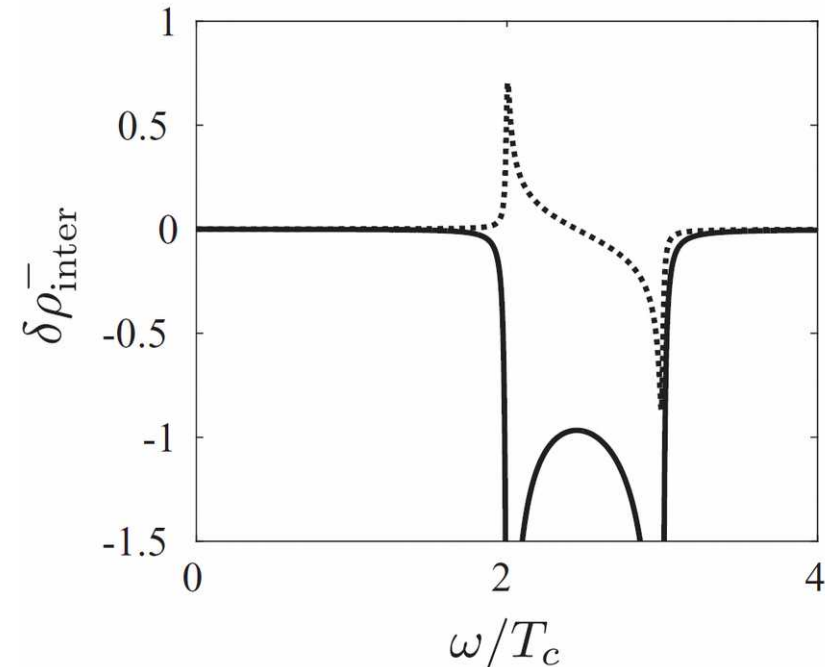


Phase sensitive measurement

- HAEM procedure
- Consider:

$$\rho_-(\vec{q}, \omega) = \text{Re}\{g(\vec{q}, +\omega)\} - \text{Re}\{g(\vec{q}, -\omega)\}$$

- s++: sign change in signal
- s+-: no sign change in signal

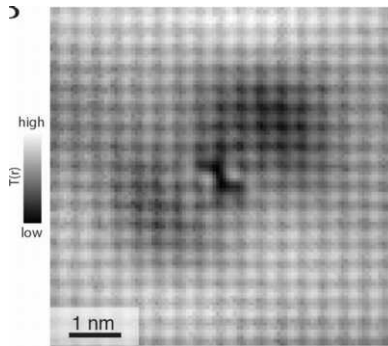


Hirschfeld et al., PRB **92**, 184513 (2015)

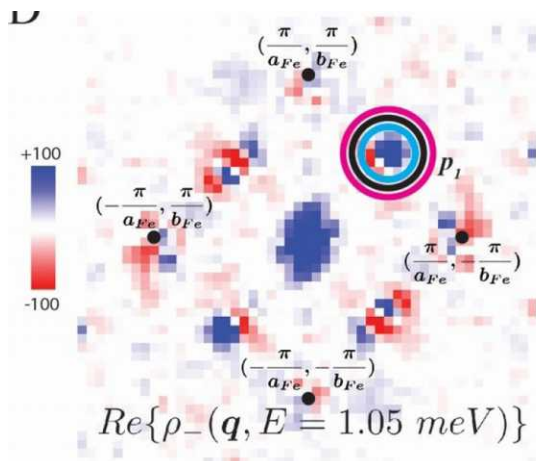
Measurement+modelling

- Problem: shift theorem in FT \rightarrow single impurity (centered)

$$\rho_-(\vec{q}, \omega) = \text{Re}\{g(\vec{q}, +\omega)\} - \text{Re}\{g(\vec{q}, -\omega)\}$$



- separate interband scattering contributions



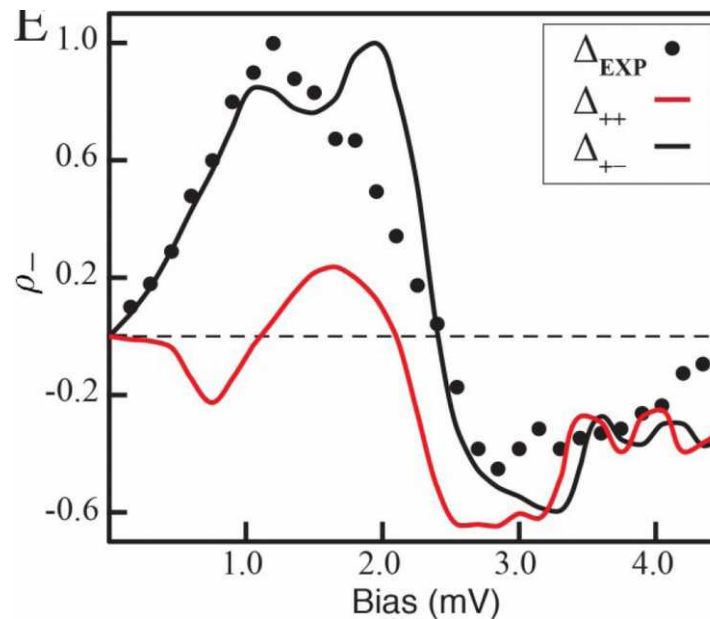
- Theory:
use measured gap
+electronic structure

$$G_{k,k'}(\omega) = G_{k-k'}^0(\omega) + G_k^0(\omega)T(\omega)G_{k'}^0(\omega)$$

$$T(\omega) = [1 - V_{imp}G_0(\omega)]^{-1}V_{imp}$$

$$\delta N(\mathbf{q}, \omega) = \frac{1}{\pi} \text{Tr}\{\text{Im} \sum_k G_k^0(\omega)T(\omega)G_{k+q}^0(\omega)\}$$

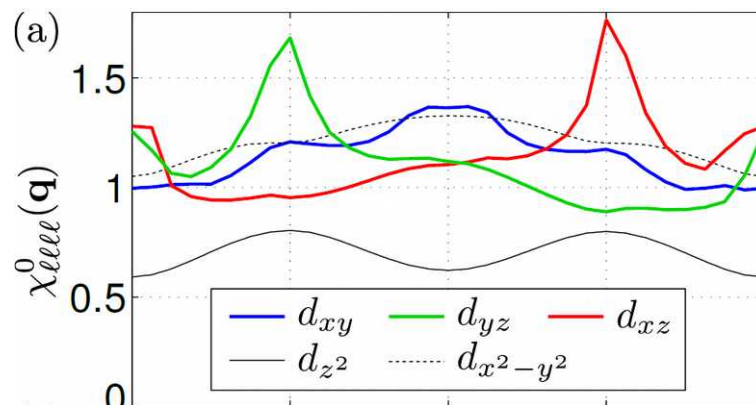
$$\rho(\omega) = \sum'_q \delta N(\mathbf{q}, \omega)$$



\rightarrow no sign change in signal, thus GAP changes sign

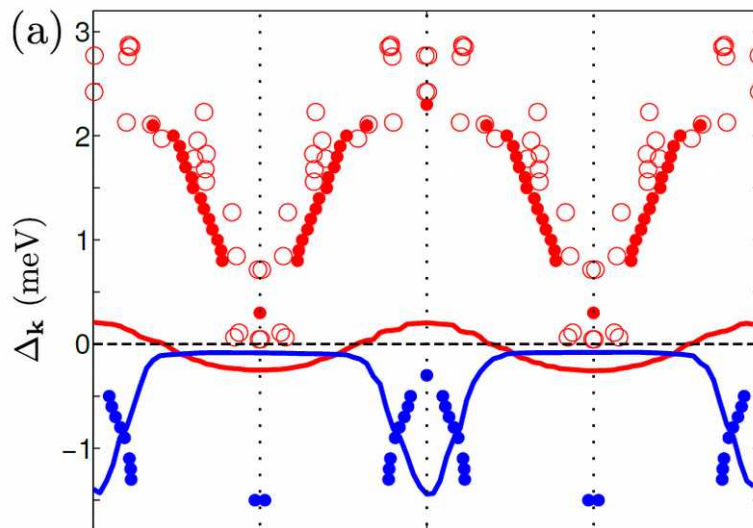
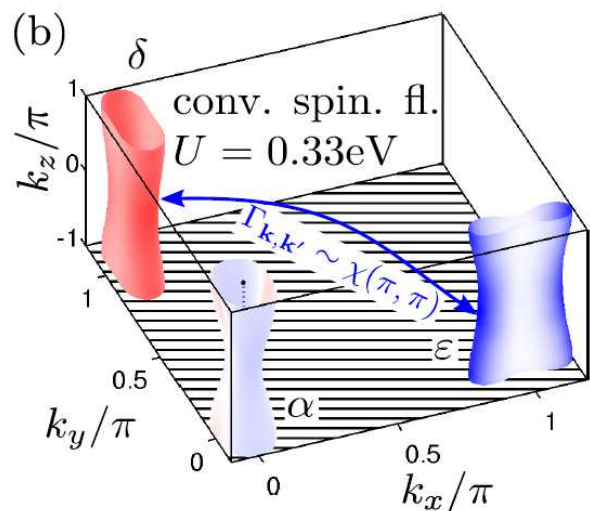
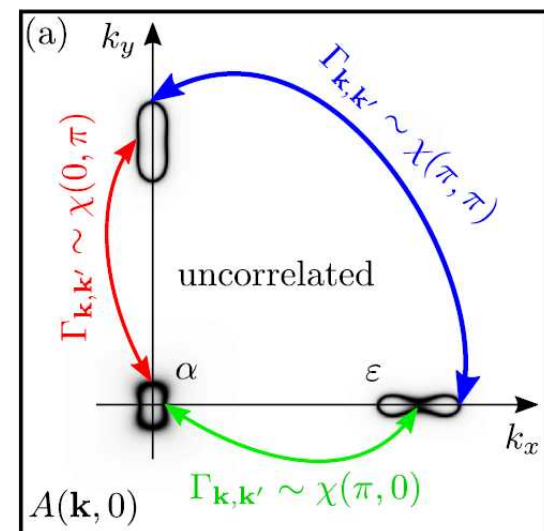
Pairing from spin-fluctuation theory?

- Susceptibility



- Pairing glue

- Solution of BCS equation



Does not work!
 → “in-phase” anisotropy
 → Small anisotropy on Gamma pocket!

What is missing?

- Interactions (standard)
- Electronic structure (measured)
- Pairing mechanism?

$$\begin{aligned}
 H = H_0 &+ U \sum_{i,l} n_{il\uparrow} n_{il\downarrow} + U' \sum_{i,l' < l} n_{il} n_{il'} \\
 &+ J \sum_{i,l' < l} \sum_{\sigma,\sigma'} c_{il\sigma}^\dagger c_{il'\sigma'}^\dagger c_{il\sigma'} c_{il'\sigma} \\
 &+ J' \sum_{i,l' \neq l} c_{il\uparrow}^\dagger c_{il\downarrow}^\dagger c_{il'\downarrow} c_{il'\uparrow},
 \end{aligned}$$

$$\chi_{l_1 l_2 l_3 l_4}^0(q) = - \sum_{k,\mu\nu} M_{l_1 l_2 l_3 l_4}^{\mu\nu}(\mathbf{k}, \mathbf{q}) G^\mu(k+q) G^\nu(k)$$

$$\begin{aligned}
 \Gamma_{l_1 l_2 l_3 l_4}(\mathbf{k}, \mathbf{k}') &= \left[\frac{3}{2} \bar{U}^s \chi_1^{\text{RPA}}(\mathbf{k} - \mathbf{k}') \bar{U}^s \right. \\
 &\quad \left. + \frac{1}{2} \bar{U}^s - \frac{1}{2} \bar{U}^c \chi_0^{\text{RPA}}(\mathbf{k} - \mathbf{k}') \bar{U}^c + \frac{1}{2} \bar{U}^c \right]_{l_1 l_2 l_3 l_4}
 \end{aligned}$$

$$\begin{aligned}
 \Gamma_{\nu\mu}(\mathbf{k}, \mathbf{k}') &= \text{Re} \sum_{l_1 l_2 l_3 l_4} a_\nu^{l_1,*}(\mathbf{k}) a_\nu^{l_4,*}(-\mathbf{k}) \\
 &\quad \times \Gamma_{l_1 l_2 l_3 l_4}(\mathbf{k}, \mathbf{k}') a_\mu^{l_2}(\mathbf{k}') a_\mu^{l_3}(-\mathbf{k}')
 \end{aligned}$$

Fermi liquid description

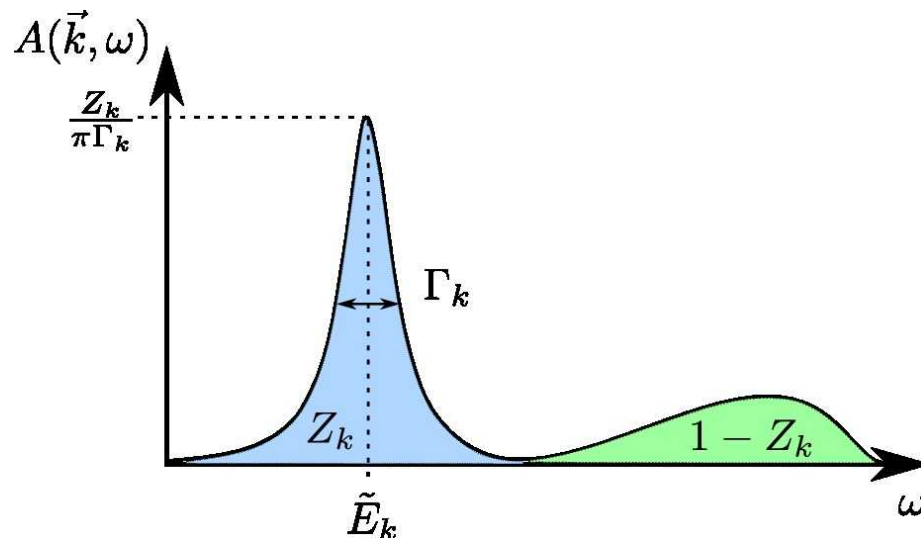
- Coherent electronic states

$$G(\vec{k}, \omega) = \frac{1}{\omega - E_{\vec{k}} + i0^+}$$

$$A(\vec{k}, \omega) = -\frac{1}{\pi} \text{Im}G(\vec{k}, \omega)$$

- Dressed electronic states

$$G(\vec{k}, \omega) = \frac{1}{\omega - E_{\vec{k}} - \Sigma(\vec{k}, \omega) + i0^+}$$



Relevant for Fe based SC:

Yin, Haule, Kotliar, Nat. Mat. 10, 932-935 (2011)
 de' Medici, Giovannetti, Capone. Phys. Rev. Lett. 112, 177001 (2014)
 M. Aichhorn, et al., Phys. Rev. B 82, 064504 (2010)
 Liu et al., Phys. Rev. B 92, 235138 (2015)
 Yi et al., Nat. Comm. 6, 7777 (2015)

...

Orbital selective physics

- States in some orbitals more decoherent than others

Strong renormalizations of the d_{xy} orbital

Yi et al., Nat. Comm. 6, 7777 (2015)

- Spectroscopic probes struggle to detect d_{xy} orbital states

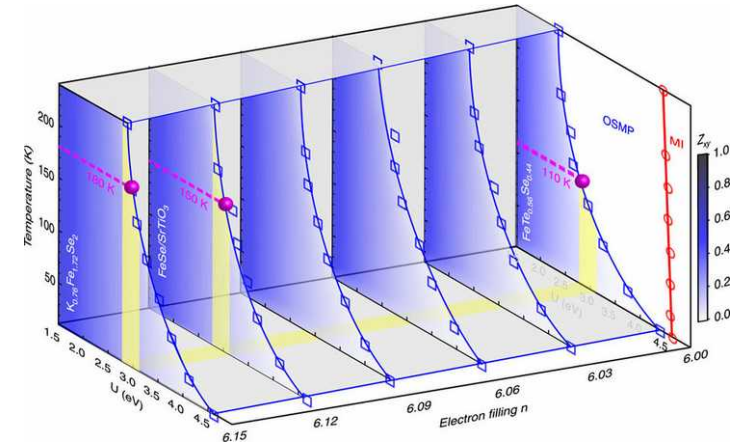
- FeSe: quasiparticle weights for d_{xz} and d_{yz} orbital distinct in nematic phase

$$G_{ab}(\mathbf{k}, \omega) = Z_{ab} G_{ab}^0(\mathbf{k}, \omega)$$

$$c_a \rightarrow \sqrt{Z_a} c_a$$

$$Z_{ab} = \sqrt{Z_a} \sqrt{Z_b}$$

geometric mean of quasiparticle weights



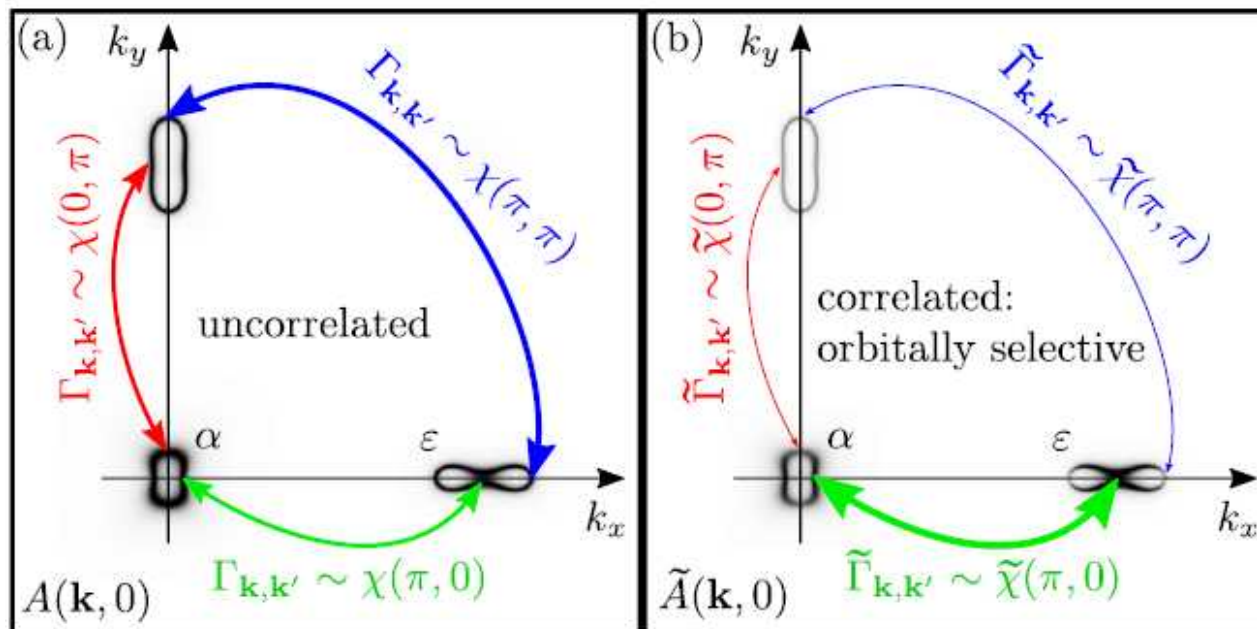
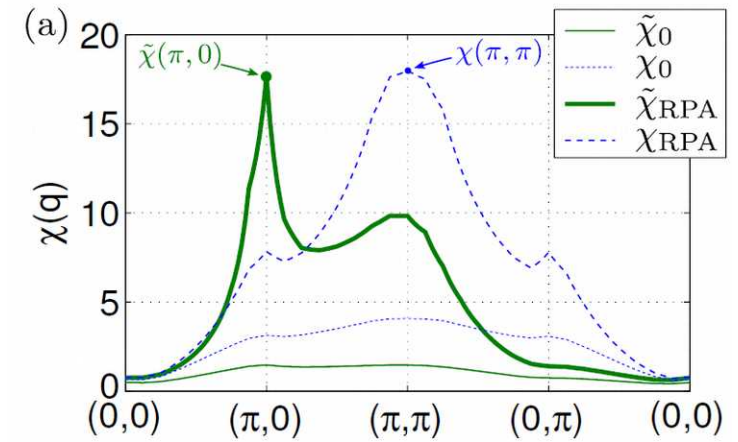
Spin-fluctuation theory

- “dressed susceptibility”

$$\tilde{\chi}_{l_1 l_2 l_3 l_4}^0(\mathbf{q}) = \sqrt{Z_{l_1} Z_{l_2} Z_{l_3} Z_{l_4}} \chi_{l_1 l_2 l_3 l_4}^0(\mathbf{q})$$

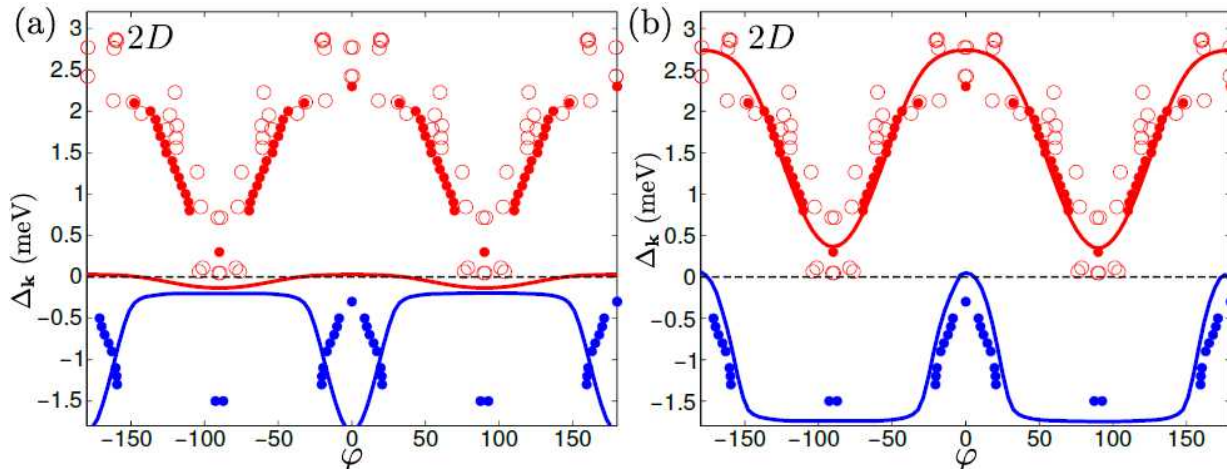
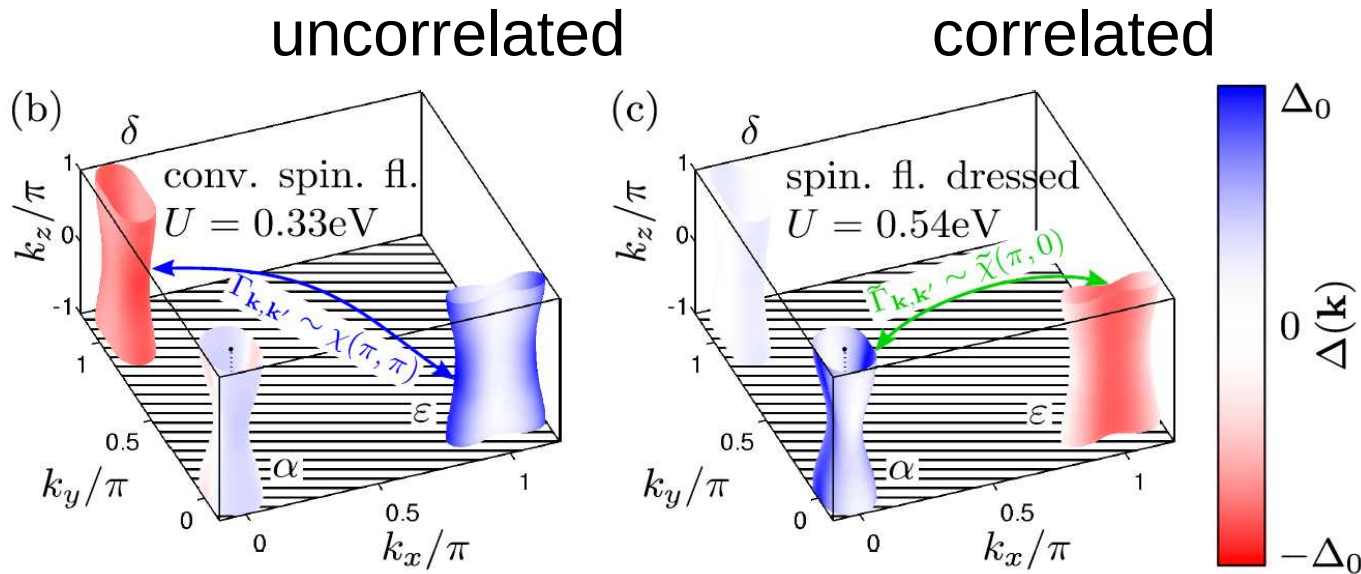
- Dressed pairing interaction

$$\tilde{\Gamma}_{\nu\mu}(\mathbf{k}, \mathbf{k}') = \text{Re} \sum_{l_1 l_2 l_3 l_4} \sqrt{Z_{l_1}} \sqrt{Z_{l_4}} a_{\nu}^{l_1, *}(k) a_{\nu}^{l_4, *}(k) \times \tilde{\Gamma}_{l_1 l_2 l_3 l_4}(\mathbf{k}, \mathbf{k}') \sqrt{Z_{l_2}} \sqrt{Z_{l_3}} a_{\mu}^{l_2}(k') a_{\mu}^{l_3}(-k')$$



Dominant pairing in d_{yz} orbital channel
 → orbital selective pairing

Pairing and gap structure



Fit parameters, so far no microscopic calculation
But: same trends found in microscopic calculations

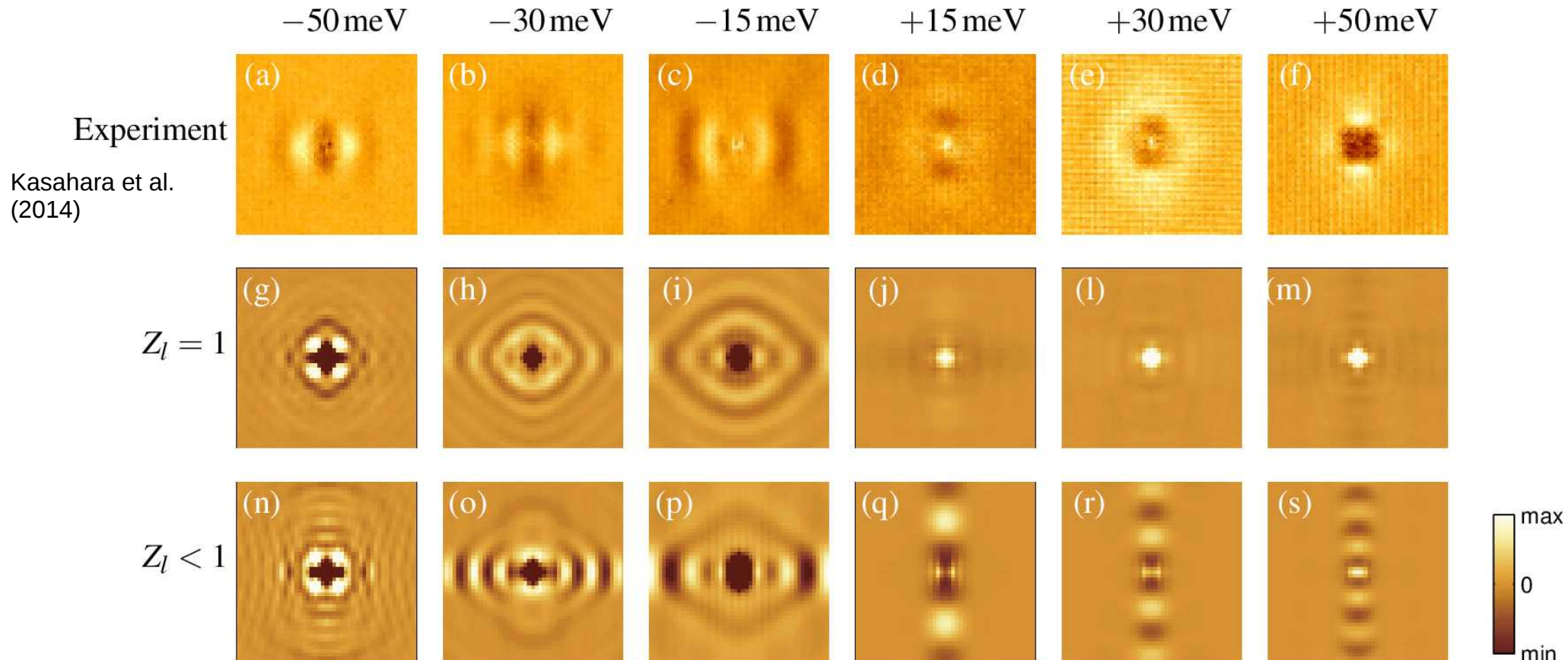
$$(d_{xy}, d_{x^2-y^2}, d_{xz}, d_{yz}, d_{3z^2-r^2})$$

$$\{\sqrt{Z_l}\} = [0.2715, 0.9717, 0.4048, 0.9236, 0.5916]$$

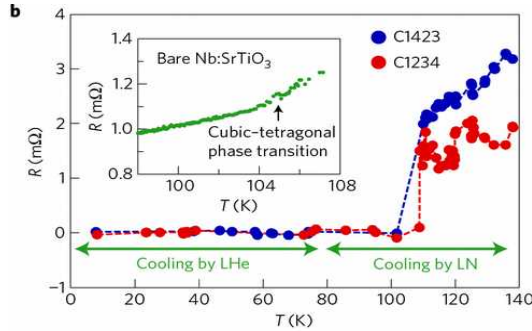
Fermi liquid theory

- Predictions for other experiments: electronic dimer close to impurity

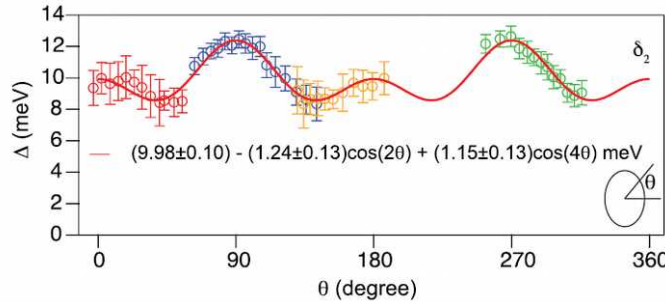
$$\{\sqrt{Z_l}\} = [0.2715, 0.9717, 0.4048, 0.9236, 0.5916]$$



Other systems: FeSe monolayer

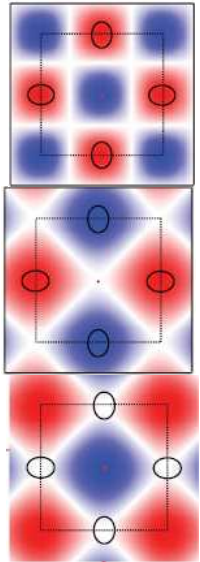


Ge et al. Nat. Mater. **14**, 285 (2015)

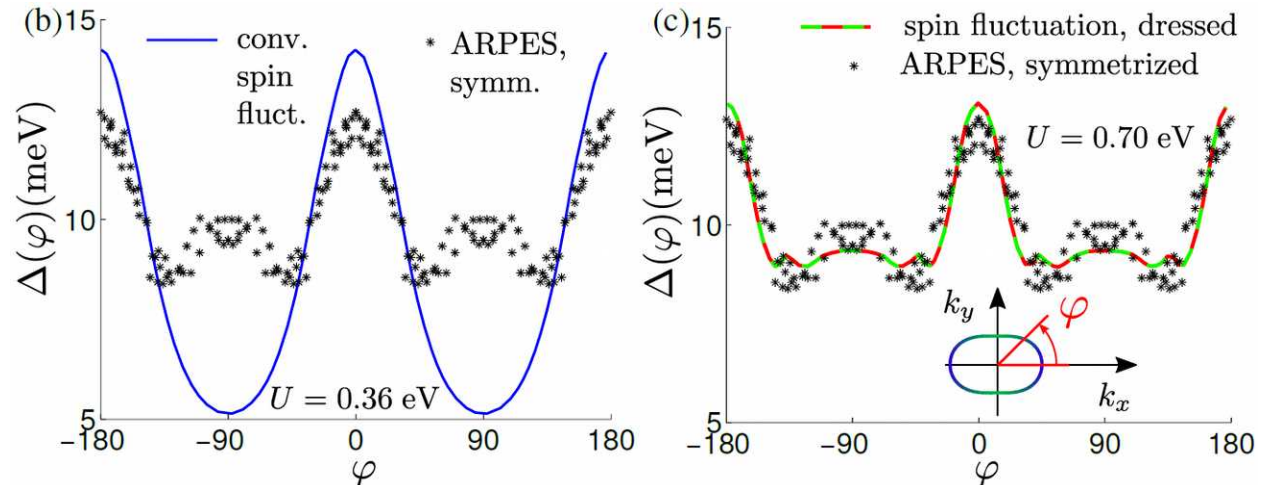
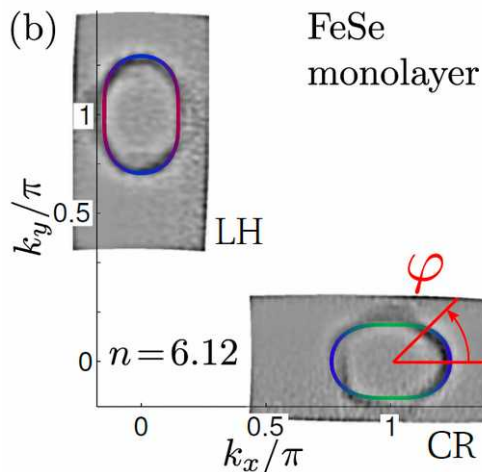


Zhang, et al., Phys. Rev. Lett. **117**, 117001 (2016)

No explanation of the two maxima structure by conventional approaches



- Same model, but: 2D, no orbital order, rigid shift



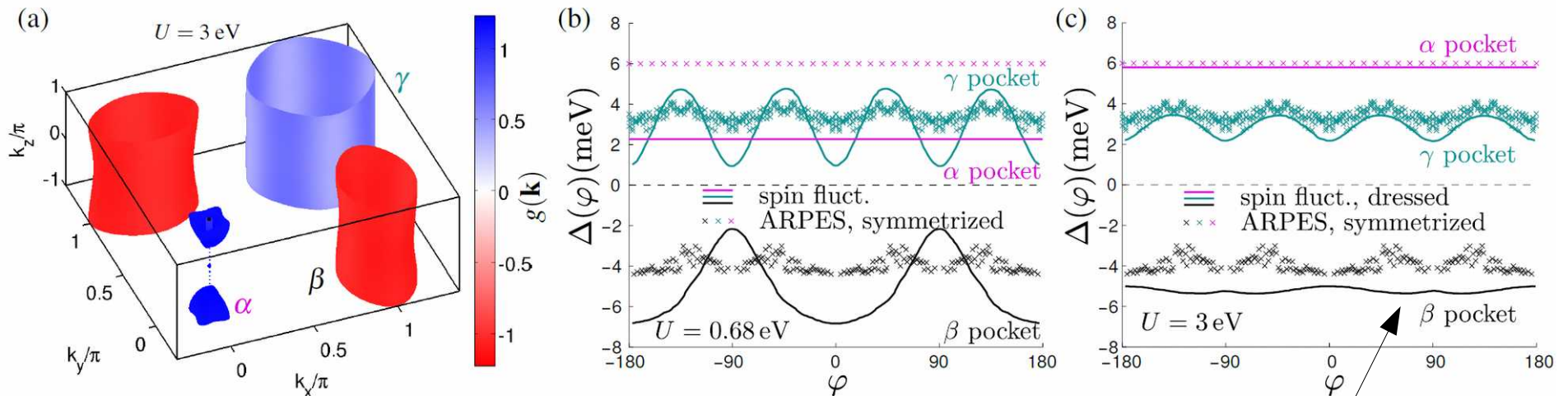
$$\{\sqrt{Z_l}\} = [0.4273, 0.8000, 0.9826, 0.9826, 0.700]$$

$$(d_{xy}, d_{x^2-y^2}, d_{xz}, d_{yz}, d_{3z^2-r^2})$$

LiFeAs

- Large gap on the α pocket

Compare: Y. Wang et al. Phys. Rev. B **88**, 174516 (2013)



$$\{\sqrt{Z_l}\} = [0.5493, 0.969, 0.5952, 0.5952, 0.9267]$$

$$(d_{xy}, d_{x^2-y^2}, d_{xz}, d_{yz}, d_{3z^2-r^2})$$

“antiphase variation” of gap does not come out spot on

Borisenko et al. Symmetry 4, 251 (2012)

Conclusions

- Many interesting and open questions remain in the field of high-Tc superconductivity
- FeSe is extremely interesting due to nematicity and orbital selective pairing due to strong correlations
- A modified spin-fluctuation approach allows for quantitative description of gap functions and other observables

$$G_{ab}(\mathbf{k}, \omega) = Z_{ab} G_{ab}^0(\mathbf{k}, \omega)$$

Acknowledgements



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Davis



Valentin
Taufour



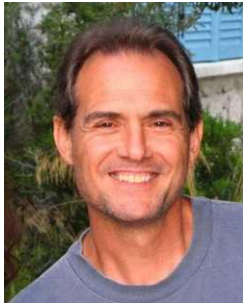
Anna E.
Böhmer



Paul C.
Canfield



Shantanu
Mukherjee



Peter J.
Hirschfeld



Brian M.
Andersen

Peter O. Sprau, Andrey Kostin, Andreas Kreisel, Anna E. Böhmer, Valentin Taufour, Paul C. Canfield, Shantanu Mukherjee, Peter J. Hirschfeld, Brian M. Andersen, J.C. Séamus Davis

arXiv:1611.02134

Discovery of Orbital-Selective Cooper Pairing in FeSe

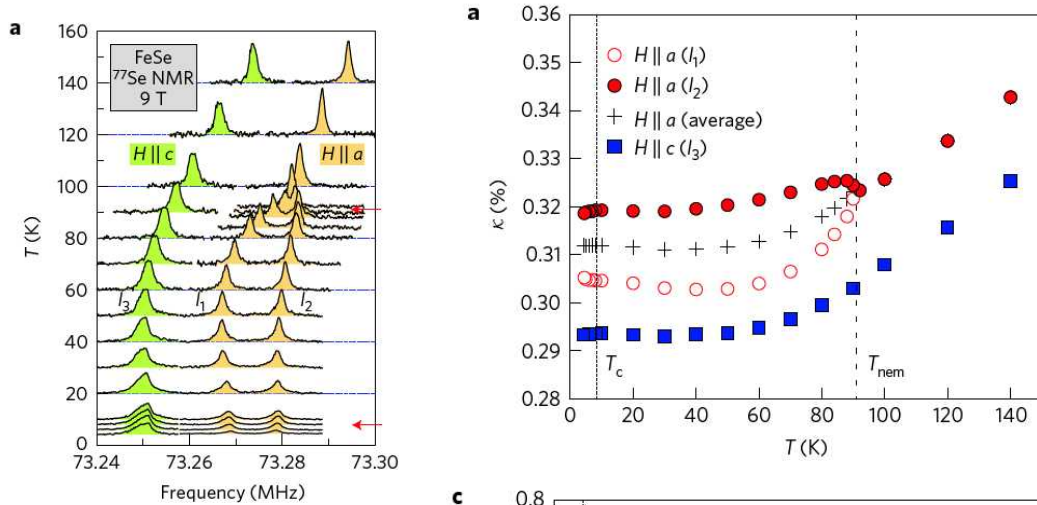
A. Kreisel, Brian M. Andersen, Peter O. Sprau, Andrey Kostin, J.C. Séamus Davis, P. J. Hirschfeld

arXiv:1611.02643

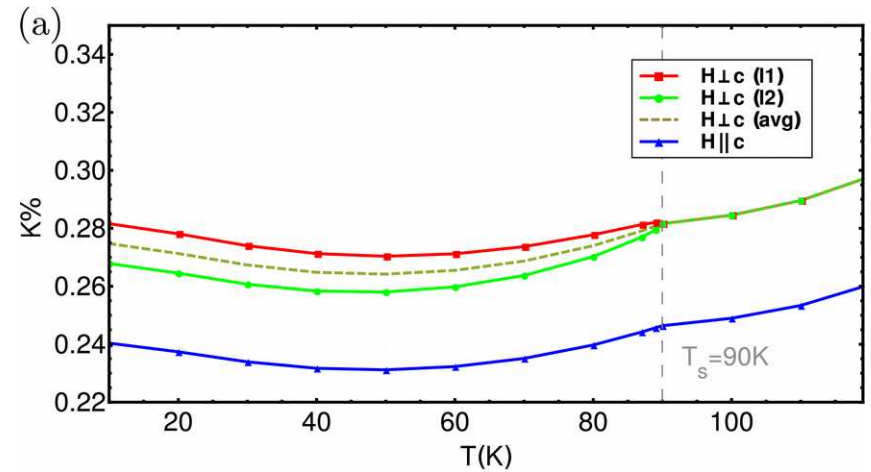
Orbital selective pairing and gap structures of iron-based superconductors

NMR: Knight shift, $1/T_1T$

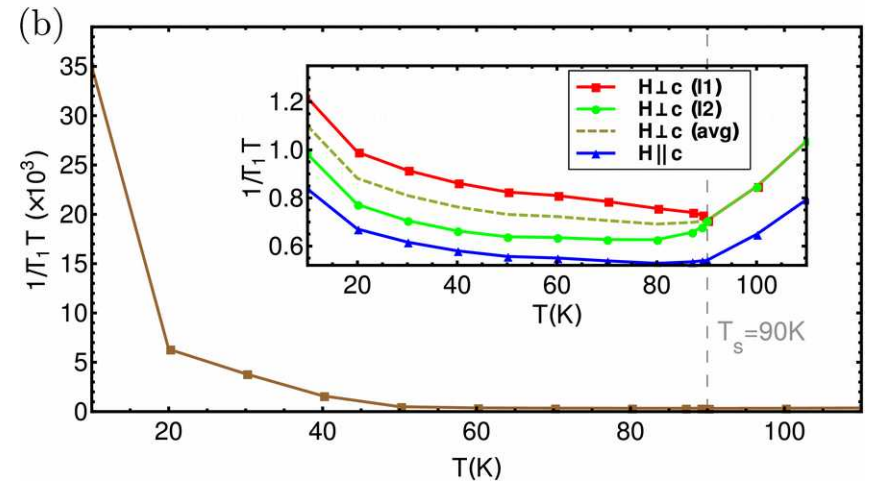
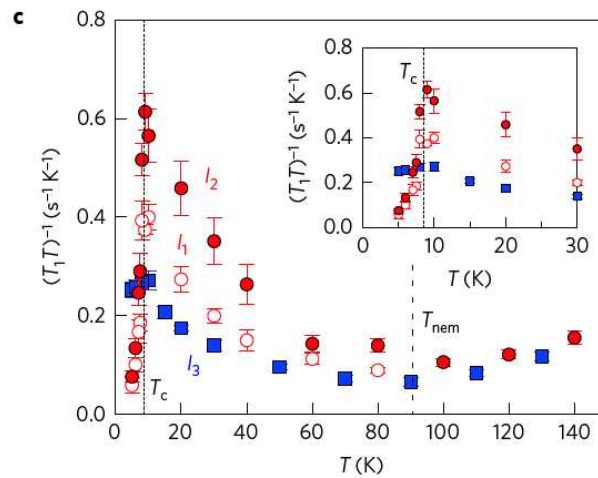
orbital order visible in Knight shift



no enhanced low-energy spin fluctuations visible in NMR



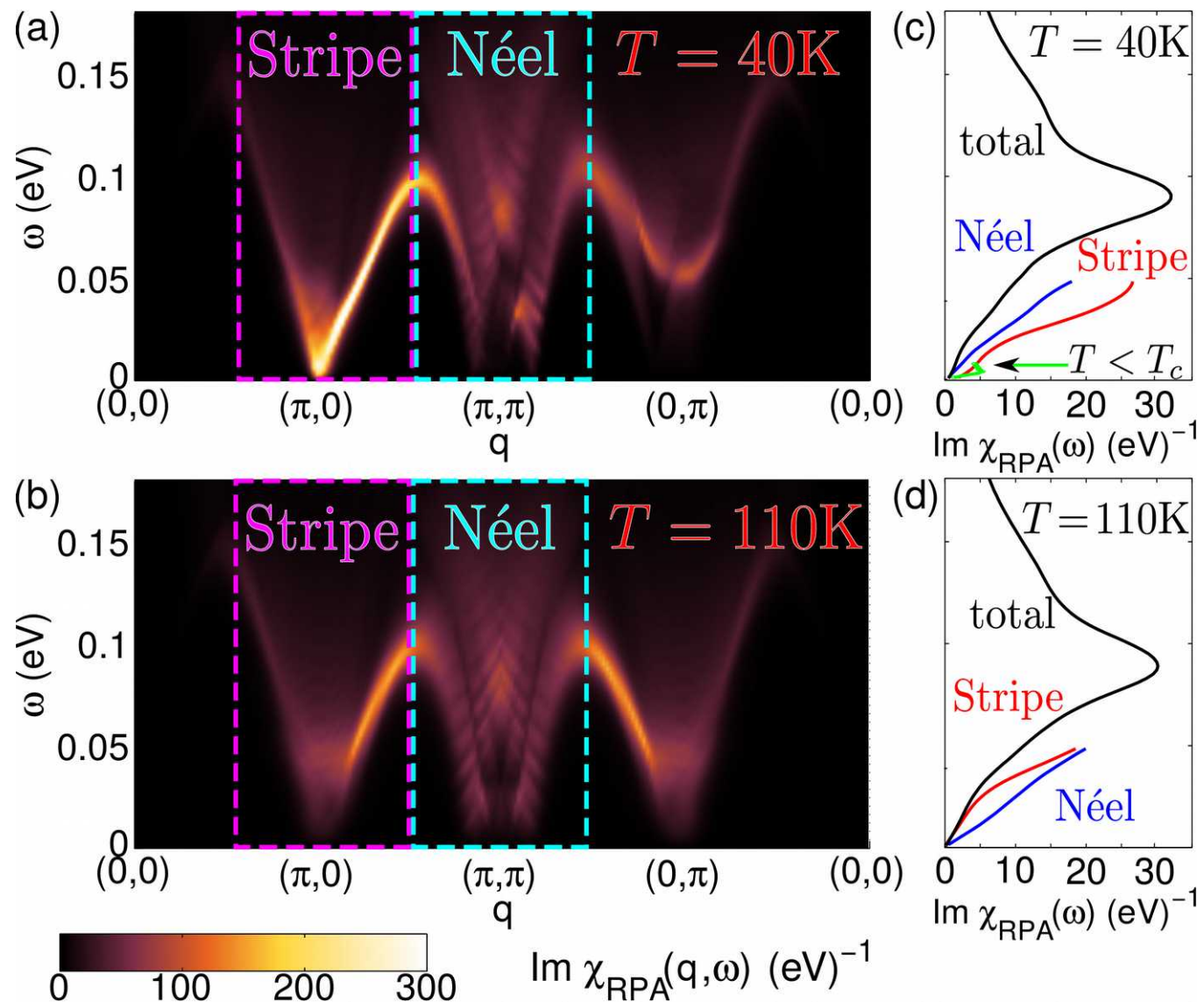
Baek, *et al.* Nature Materials **14**, 210 (2015)



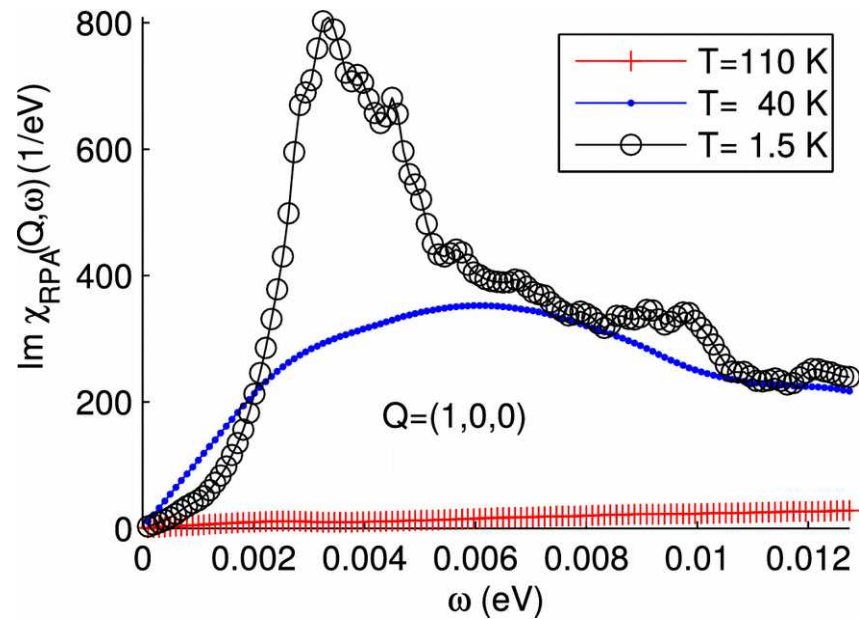
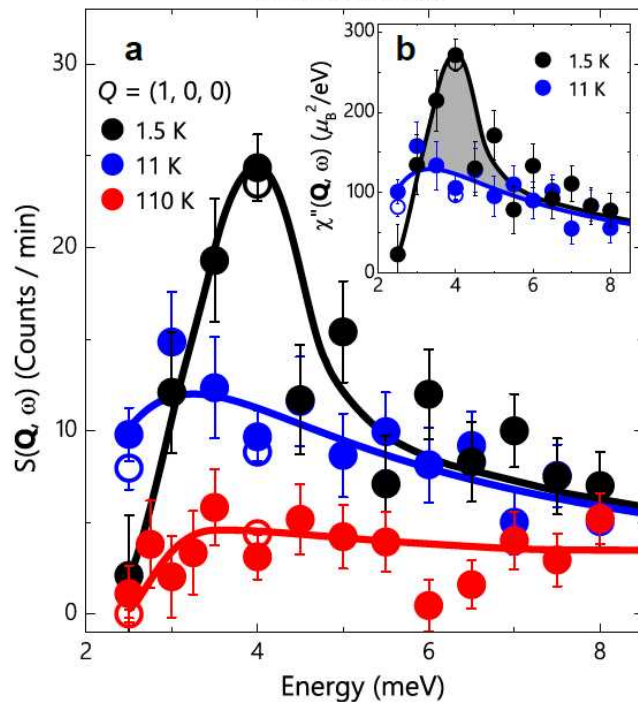
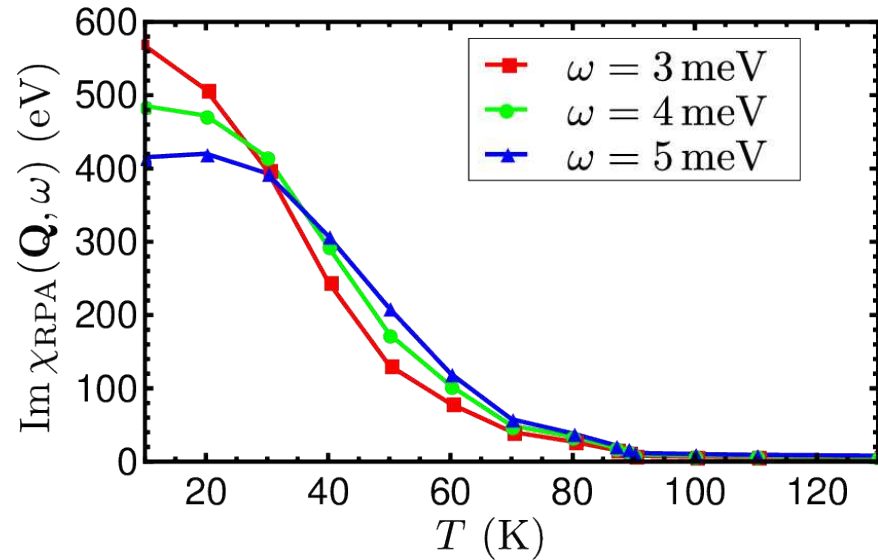
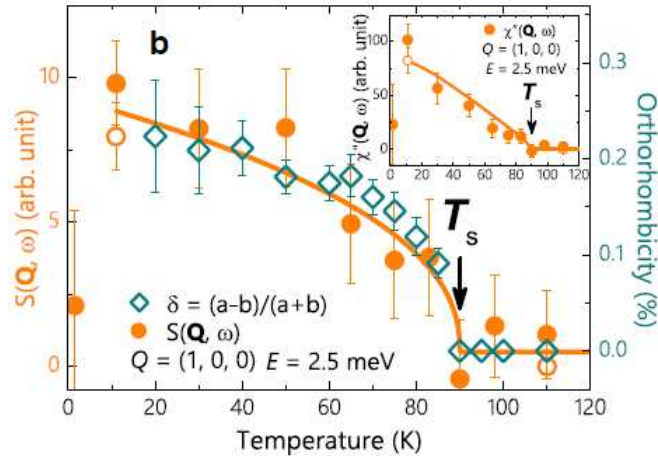
$$\frac{1}{T_1T} = \lim_{\omega_0 \rightarrow 0} \frac{\gamma_N^2}{2N} k_B \sum_{\mathbf{q}\alpha\beta} |A_{hf}^{\alpha\beta}(\mathbf{q})|^2 \frac{\text{Im}\{\chi_{\text{RPA}}^{\alpha\beta}(\mathbf{q}, \omega_0)\}}{\hbar\omega_0}$$

Spin fluctuations at higher energies

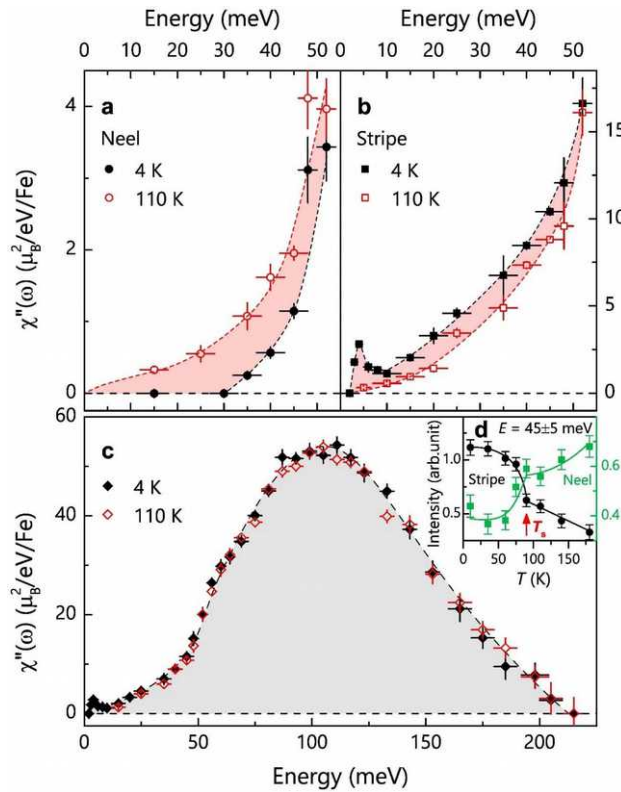
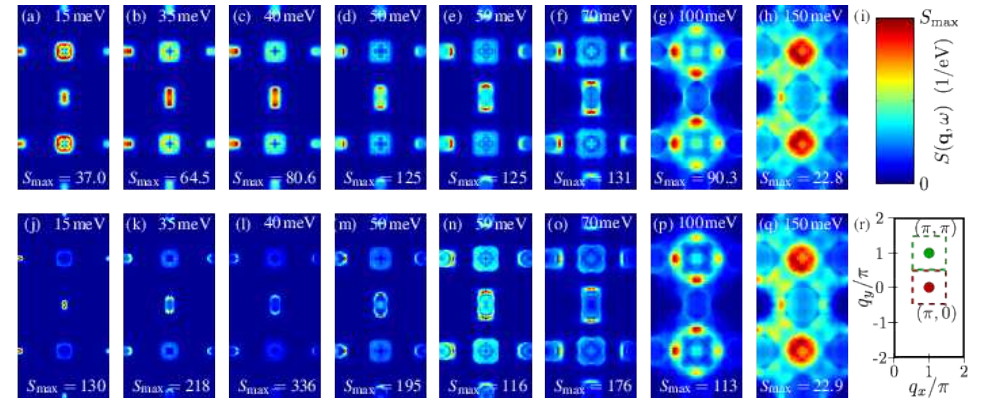
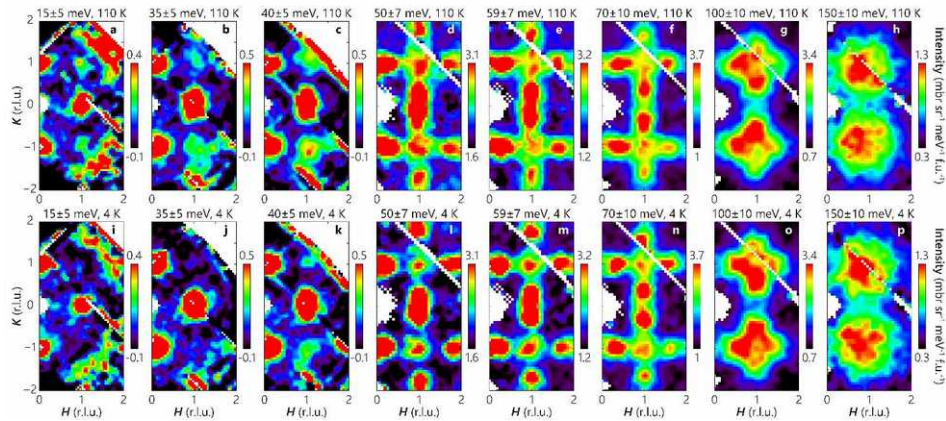
- FeSe: close to magnetic instability (tune interactions accordingly)
- transfer from Néel fluctuations to Stripe fluctuations on lowering temperature
- spin resonance at low energies from transfer of spectral weight in the superconducting state



Inelastic neutron scattering



Inelastic neutron scattering



Q. Wang, *et al*,
arXiv:1511.02485
(2015)

