

**Finite temperature phase  
transition in the 4d abelian  
LGT?**

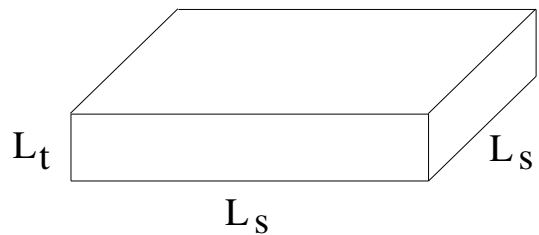
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with **Philippe de Forcrand**

# Definition of the model

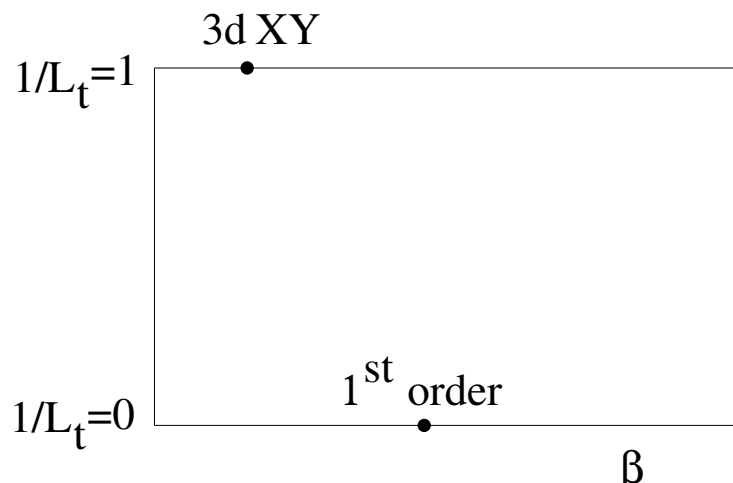
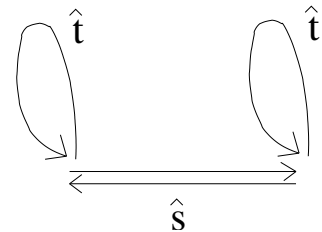
4d Abelian compact U(1) with  
Wilson action (fixed  $L_t \ll L_s, L_s \rightarrow \infty$ )

$$S = \beta \sum_{\square} \cos \theta_{\square}$$



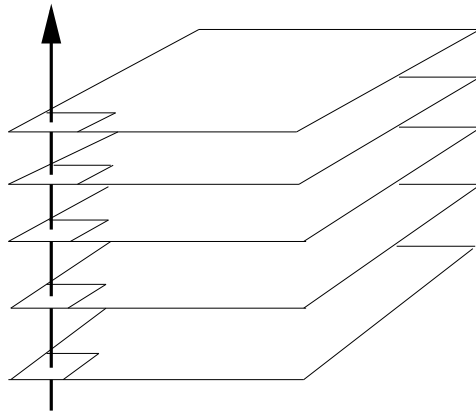
- $L_t = \infty$ : 1<sup>st</sup> order phase transition
- $L_t = 1$ : uncoupled Polyakov loops

$$Z = Z_{3d XY} \cdot Z_{3d U(1) GT}$$



# The order parameter

Response function to  
an external static flux ( $\equiv$  twisted b.c.)



- **Confined** phase: system **insensitive** to an external flux  $\Phi$  ( $\xi$  finite, response  $\sim e^{-L/\xi}$ )
- **Coulomb** phase: system **sensitive** to an external flux ( $\xi = \infty$ )

**Helicity modulus**

$$h(\beta) = \frac{\partial^2 F(\phi; \beta)}{\partial \phi^2} \Big|_{\phi=0}$$

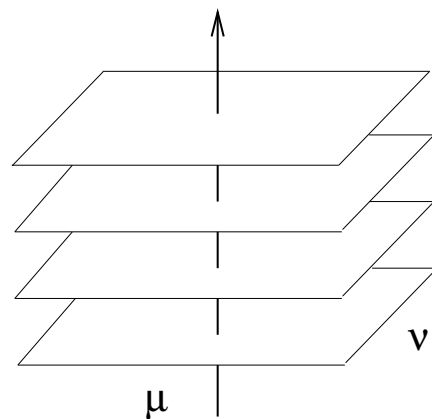
## More about the order parameter

Chose an orientation  $(\mu, \nu)$  to impose  $\Phi$

N.B., it can be

- Spatial orientation
- Temporal orientation

Suppose that  $\Phi$  spreads homogeneously through parallel planes



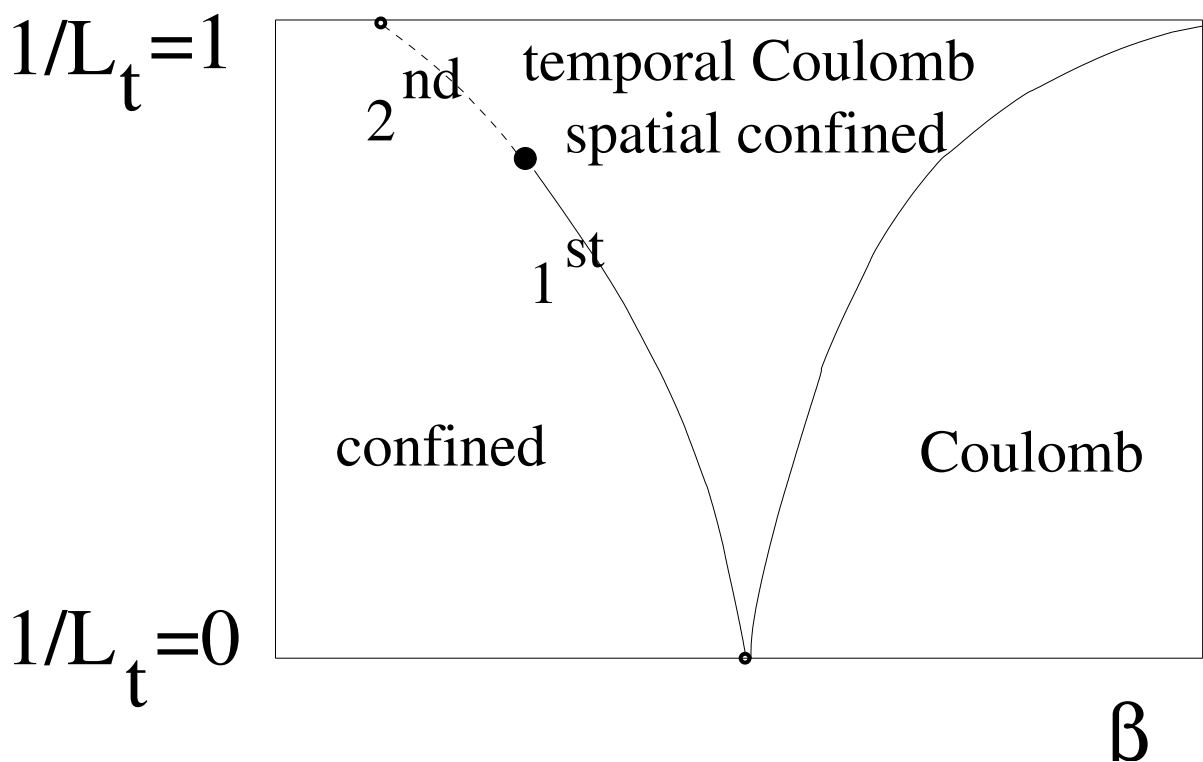
Classical limit  $(\beta \rightarrow \infty)$ ,  $\theta_{\square} = \frac{\Phi}{L_{\mu}L_{\nu}}$

$S = -\beta \sum_{\text{plaq.}} \cos \frac{\Phi}{L_{\mu}L_{\nu}}$ , expand

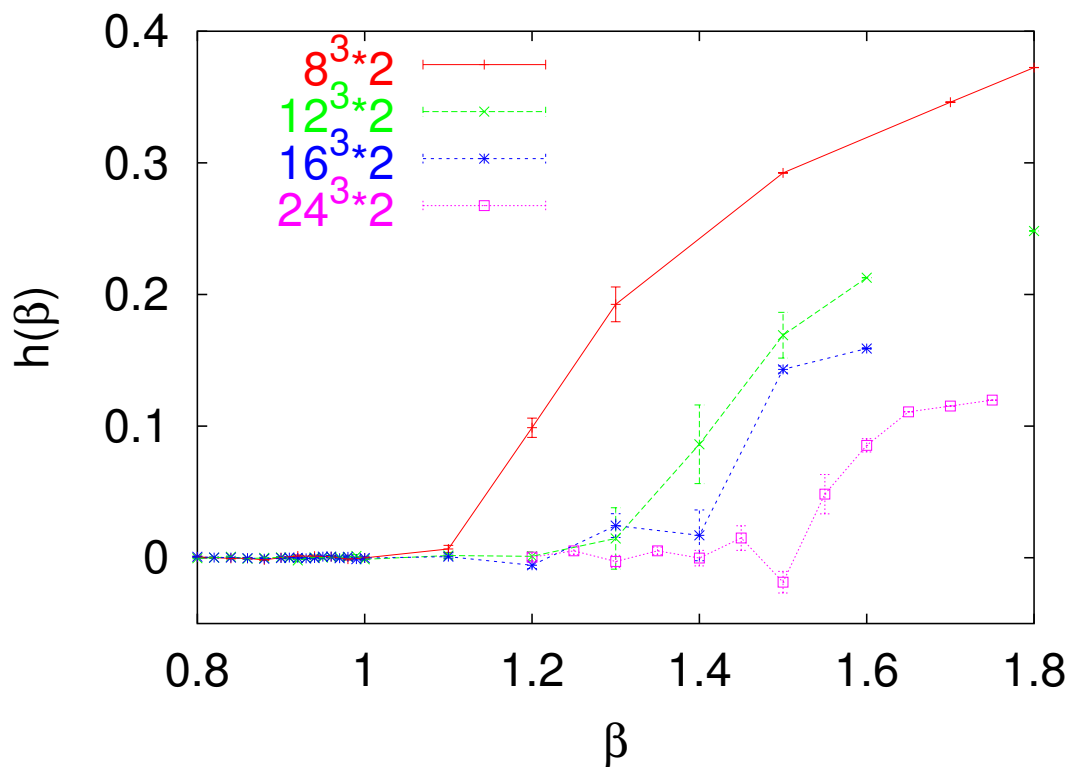
$$F(\Phi) = \frac{\beta}{2} \Phi^2 \frac{L_{\rho}L_{\sigma}}{L_{\mu}L_{\nu}} \rightarrow \frac{\beta_{\mathcal{R}}(\beta)}{2} \Phi^2 \frac{L_{\rho}L_{\sigma}}{L_{\mu}L_{\nu}}, \forall \beta$$

# Conjectured phase diagram

Decoupling of the transition temperature  
for **spatial** and **temporal** loops



# Transition to Coulomb phase



$$\beta_c \propto L_s?$$

If this is the case, this transition disappears in the thermodynamic limit!

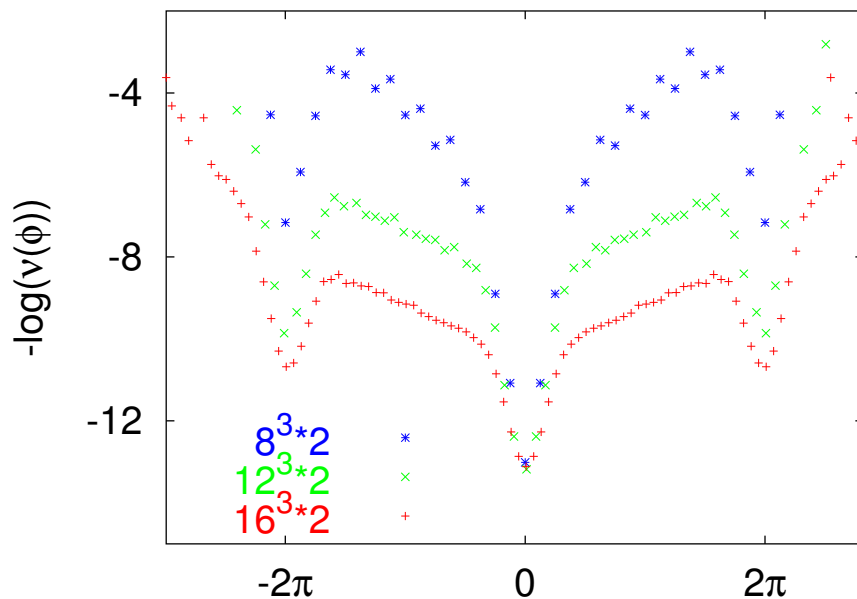
## Flux free energy as $L_s \rightarrow \infty$

$$F(\Phi) = \frac{\beta_R}{2} \Phi \frac{L_s L_t}{L_s L_s} \sim \frac{\beta_R}{L_s}$$

At finite  $V$  competition between  $\beta_R$  and  $L_s$ !

TEST:

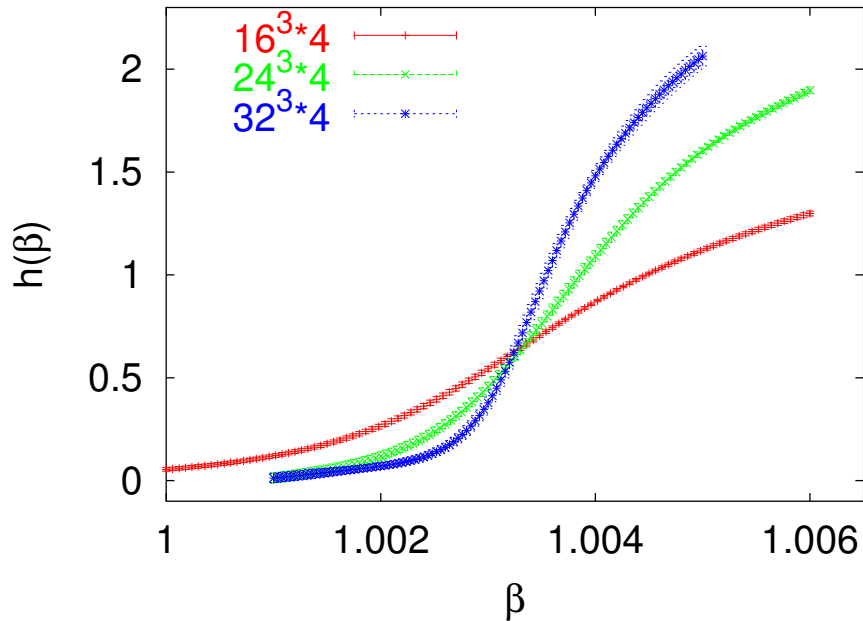
measure flux distribution  $\nu(\Phi)$  ( $\propto e^{-F(\Phi)}$ )  
for different  $L_s$



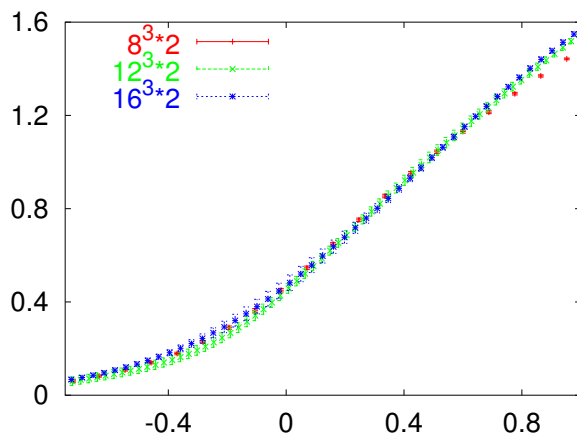
At any finite  $\beta$ , finite density of static monopoles (in time direction) which can  
disorder spatial loops

# Transition to the confined phase

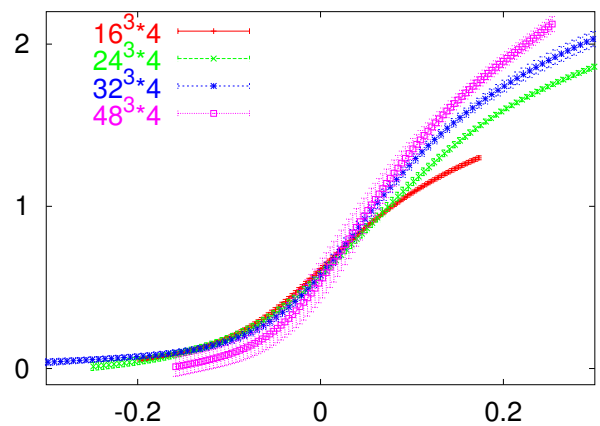
Flux through temporal planes



FSS analysis to determine **position** and **order** of the transition ( $\nu = \nu_{3d XY}$ ).



$L_t = 2$

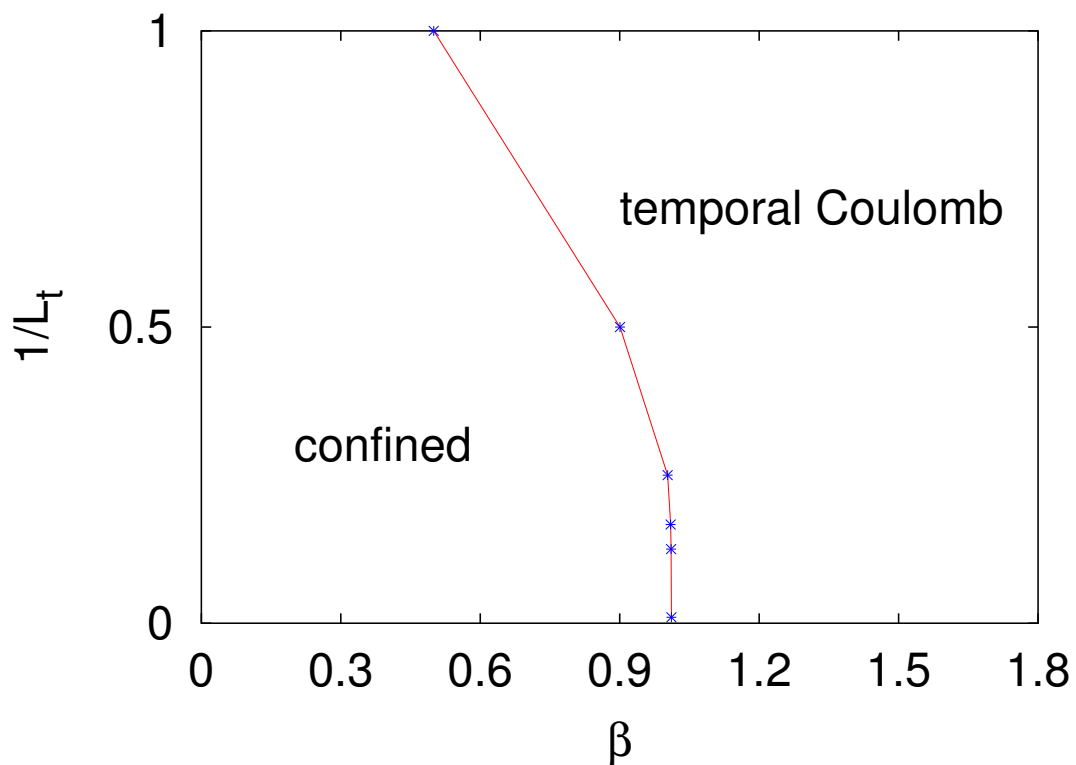


$L_t = 4$



## Location of the transition line

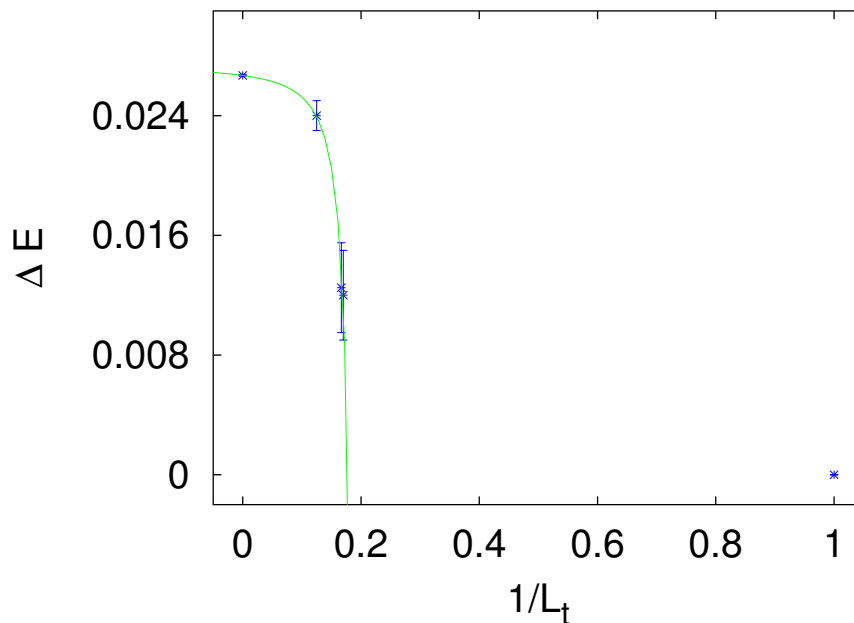
Possible to reconstruct with great precision the position of the phase boundary



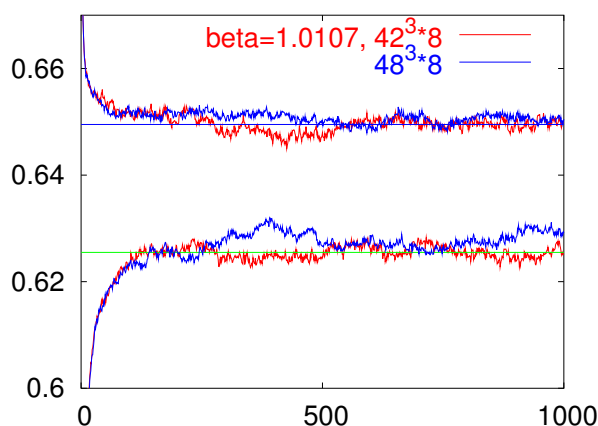
What is the order of the phase transition along this curve?

# Order of the phase transition

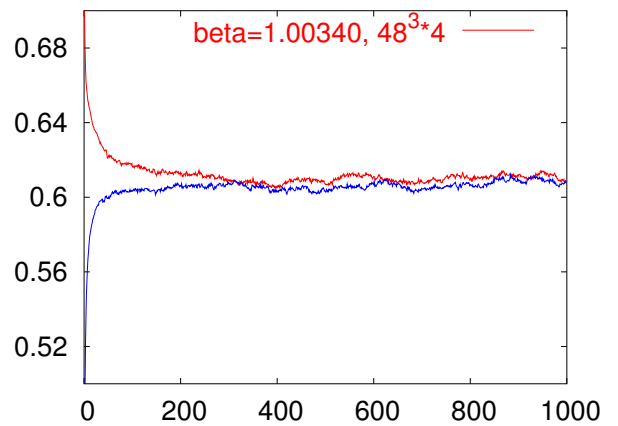
Consider how the latent heat of the (first order) transition changes with  $L_t$



$\bar{L}_t \sim 4$  the latent heat seems to vanish (with anisotropic couplings one could tune  $\bar{L}_t$ )



$L_t = 8$



$L_t = 4$

## Conclusions

Clarified the phase diagram of the system:

- There is only one phase boundary
- Coulomb phase only at  $T = 0$  and  $\beta > \beta_c$
- At finite  $T$  and  $\beta > \beta_c$  spatial Wilson loops obey area law, temporal Wilson loops perimeter law (similar to Yang-Mills)
- At a certain number of temporal slices  $\bar{L}_t$ , the transition seems to turn from first to second order: continuum limit for a system of  $3d$  coupled layers?