

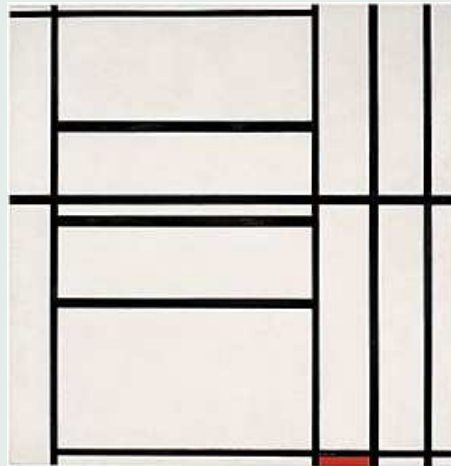
# Fractal Structure of Field Theories

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w/

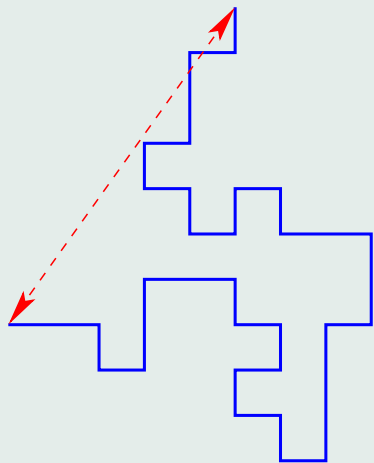
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PIET MONDRIAN, 1938-39

[DE GENNES, 1972]:



- SAW:  $\lim_{N \rightarrow 0} O(N)$
- # of walks of length  $l$  (configurational entropy):

$$c_l \sim l^{\gamma-1}$$

- End-to-end distance:

$$\langle \mathbf{R}^2 \rangle_l \sim l^{2\nu},$$

implying **fractal, or Hausdorff** dimension:

$$D_H = 1/\nu = 4/3$$

Correlation function **SAW**:

$$G(\mathbf{x} \rightarrow \mathbf{x}') = \sum_{\text{Paths: } \mathbf{x} \rightarrow \mathbf{x}'} \beta^l = \sum_l c_l(\mathbf{x} \rightarrow \mathbf{x}') \beta^l$$

Susceptibility:

$$\chi = \sum_{\mathbf{x}'} G(\mathbf{x} \rightarrow \mathbf{x}') = \sum_l \underbrace{\sum_{\mathbf{x}'} c_l(\mathbf{x} \rightarrow \mathbf{x}') \beta^l}_{= c_l}$$

$$\chi = \sum_l l^{\gamma-1} \left( \frac{\beta}{\beta_c} \right)^l \sim \sum_l l^{\gamma-1} e^{-\theta l} \sim |\beta - \beta_c|^{-\gamma}$$

w/ **line tension**  $\theta \sim |\beta - \beta_c|$

Can this be generalized to arbitrary  $-2 \leq N \leq 2$ ?

**Critical exponents**  $O(N)$  models in 2D:

For general  $N$ :

| Model    | $N$ | $c$           | $\gamma$        | $\eta$         | $\nu$         | $D_H$          |
|----------|-----|---------------|-----------------|----------------|---------------|----------------|
| Gaussian | -2  | -2            | 1               | 0              | $\frac{1}{2}$ | $\frac{5}{4}$  |
| SAW      | 0   | 0             | $\frac{43}{32}$ | $\frac{5}{24}$ | $\frac{3}{4}$ | $\frac{4}{3}$  |
| Ising    | 1   | $\frac{1}{2}$ | $\frac{7}{4}$   | $\frac{1}{4}$  | 1             | $\frac{11}{8}$ |
| XY       | 2   | 1             | $\infty$        | $\frac{1}{4}$  | $\infty$      | $\frac{3}{2}$  |

- Fractal dimension:

$$\nu \neq \frac{1}{D_H}$$

- Configurational entropy:

$$c_l \propto l^{\gamma-1}$$

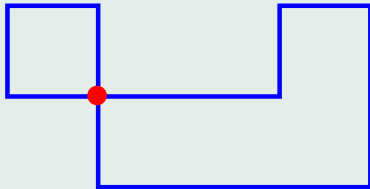
Partition function of  $O(N)$  models in high-temperature rep. w/  $\phi = \begin{pmatrix} \phi_1 \\ \vdots \\ \phi_N \end{pmatrix}$ :

$$Z = \int \prod_{\mathbf{x}} d\phi(\mathbf{x}) \exp \left( -a^d \sum_{\mathbf{x}} \left\{ \frac{1}{2a^2} \sum_{\mathbf{i}} [\phi(\mathbf{x} + a\mathbf{i}) - \phi(\mathbf{x})]^2 + \frac{m^2}{2} \phi^2(\mathbf{x}) + \frac{\lambda}{4} \phi^4(\mathbf{x}) \right\} \right)$$

$$(\lambda \equiv 0) \int \prod_{\mathbf{x}} d\phi(\mathbf{x}) \exp \left[ \beta a^{d-2} \sum_{\mathbf{x}, \mathbf{i}} \phi(\mathbf{x} + a\mathbf{i}) \cdot \phi(\mathbf{x}) - \frac{1}{2} a^{d-2} \sum_{\mathbf{x}} \phi^2(\mathbf{x}) \right], \quad \beta = \frac{1}{2d + m^2 a^2}$$

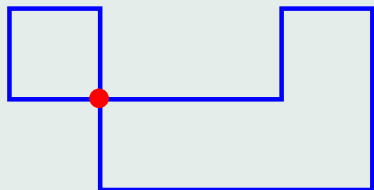
$$= \int \prod_{\mathbf{x}} d\phi(\mathbf{x}) \left[ 1 + \beta a^{d-2} \sum_{\mathbf{x}, \mathbf{i}} \phi(\mathbf{x} + a\mathbf{i}) \cdot \phi(\mathbf{x}) + \dots \right] \exp \left[ -\frac{1}{2} a^{d-2} \sum_{\mathbf{x}} \phi^2(\mathbf{x}) \right]$$

Loop gas:



- Each link:  $\beta$
- $\lambda = 0$ : phantom loops
- $\lambda \neq 0$ : extra Boltzmann weight for crossings

## Loop gas:



- $D_H$ : fractal dimension of loops
- Phase transition:
  - line tension  $\theta$  vanishes
  - proliferation of loops
- $N = 0$ : no loops,  $Z = 1$

In general:

- Line tension vanishes as:

$$\theta \sim |\beta - \beta_c|^{1/\sigma}, \quad \text{w/ } \sigma \neq 1$$

- Susceptibility:

$$\chi \sim \sum_l l^{\sigma\gamma-1} e^{-\theta l}$$

As in percolation theory:

$$\nu = \frac{1}{\sigma D_H}$$

Scaling relation:

$$\gamma = \nu(2 - \eta) \quad \Rightarrow \quad \sigma\gamma = \frac{2 - \eta}{D_H}$$

Partition function of Ising model ( $N = 1$ ) in high-temperature rep. ( $J = 1$ ):

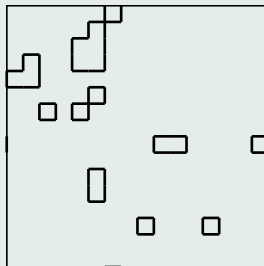
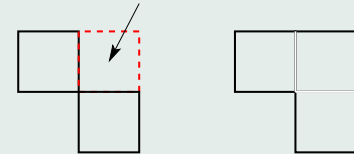
$$Z = (\cosh \beta)^{2N} 2^N \sum_{\text{Closed Graphs}} v^l, \quad e^{\beta S_i S_j} = \cosh \beta (1 + v S_i S_j)$$

w/  $v = \tanh \beta$ ,  $l$ : # links in graph

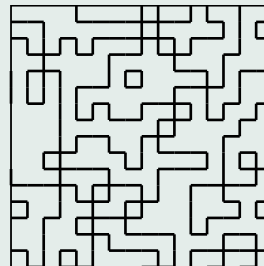
Monte Carlo [JANKE & A.S., 2003 (NPB to appear)]:

- Single plaquette  update
- Acceptance rate:  $p_{\text{HT}} = \min(1, v^{l'-l})$

selected plaquette



$\beta < \beta_c$

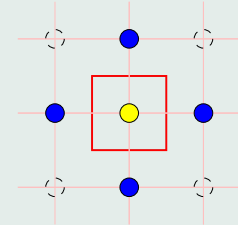


$\beta > \beta_c$

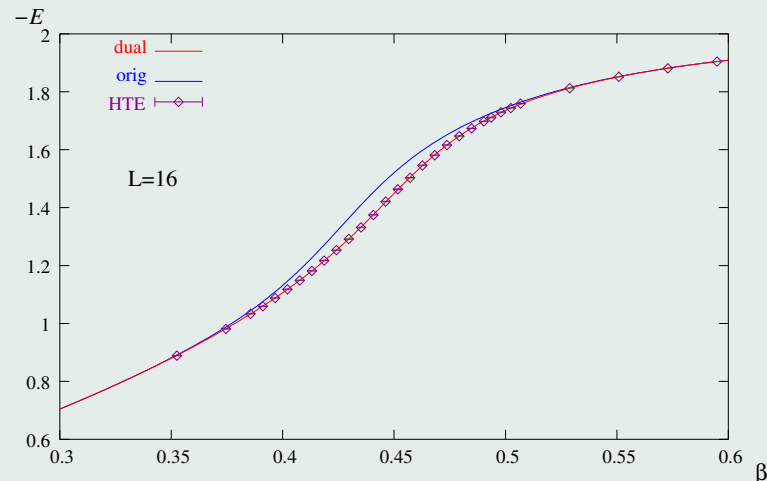
(Periodic  
Boundary  
Conditions)

Phase transition: Proliferation of HT graphs

Single **plaquette**  $\square$  update  $\sim$  Single **spin** update on dual lattice:

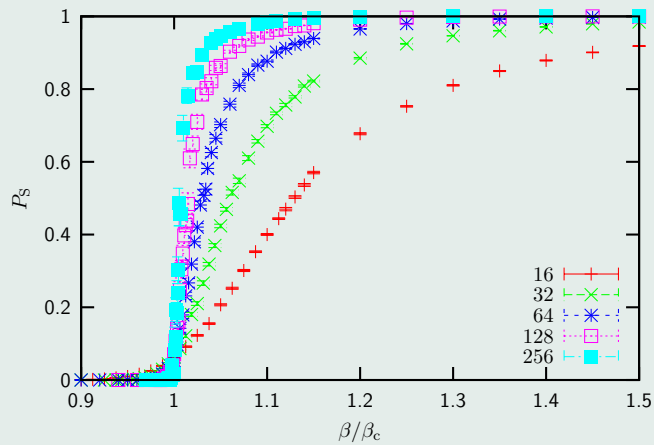


**PBC**: mismatch w/  $\tilde{\beta} = \frac{1}{2} \ln \coth \beta$  [KRAMERS & WANNIER, 1941]:

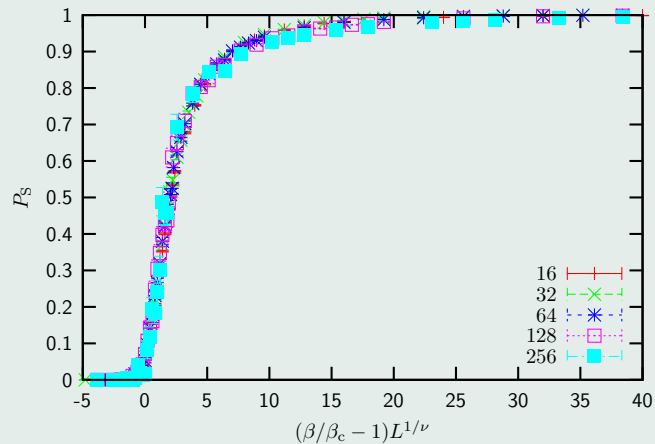


[PEIERLS, 1936]:

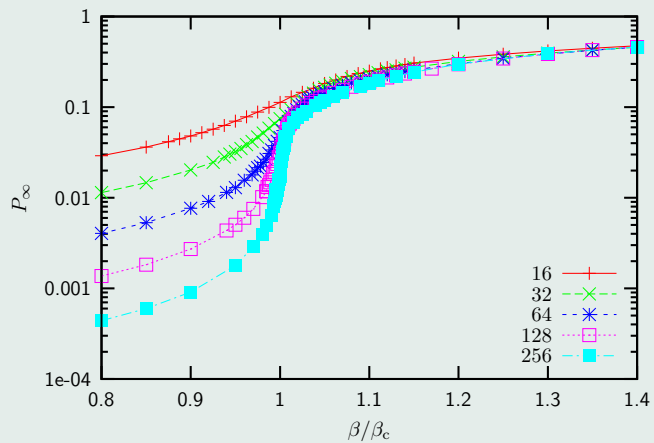
**HT** graphs  $\sim$  Boundaries of geometrical clusters (**Domain Walls**)



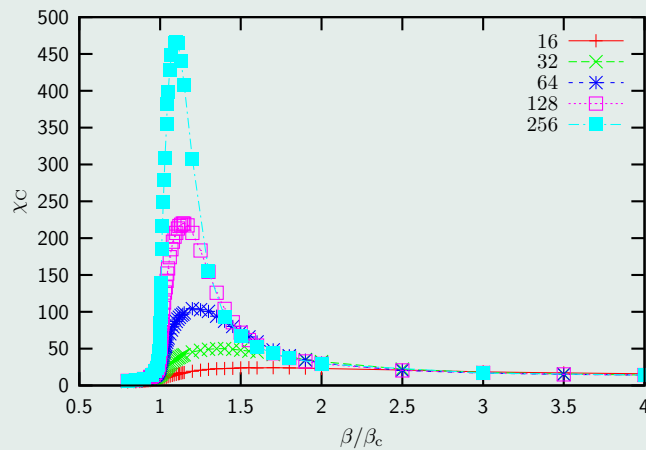
$\therefore$  Graph proliferation @  $T_c$



$\therefore \nu = 1$



$\therefore \beta_G = 0.627(8)$



$\therefore \gamma_G = 0.748(6)$

MC results consistent w/:

$$\sigma = \frac{8}{11}, D_H = \frac{11}{8} \Rightarrow \nu = \frac{1}{\sigma D_H} = 1$$

2D  $O(N)$  models:

| Model    | $N$ | $c$           | $\gamma$        | $\eta$         | $\nu$         | $D_H$          | $\sigma$       |
|----------|-----|---------------|-----------------|----------------|---------------|----------------|----------------|
| Gaussian | -2  | -2            | 1               | 0              | $\frac{1}{2}$ | $\frac{5}{4}$  | $\frac{8}{5}$  |
| SAW      | 0   | 0             | $\frac{43}{32}$ | $\frac{5}{24}$ | $\frac{3}{4}$ | $\frac{4}{3}$  | 1              |
| Ising    | 1   | $\frac{1}{2}$ | $\frac{7}{4}$   | $\frac{1}{4}$  | 1             | $\frac{11}{8}$ | $\frac{8}{11}$ |
| XY       | 2   | 1             | $\infty$        | $\frac{1}{4}$  | $\infty$      | $\frac{3}{2}$  | 0              |

XY-model ( $N = 2$ ):

- $\sigma = 0$ : algebraic behavior
- Configurational entropy:

$$c_l \sim l^{\sigma\gamma-1}$$

$$\text{w/ } \sigma\gamma - 1 = \frac{1}{6}$$