

Free-energy barriers in spin glasses: mean-field vs short-range models

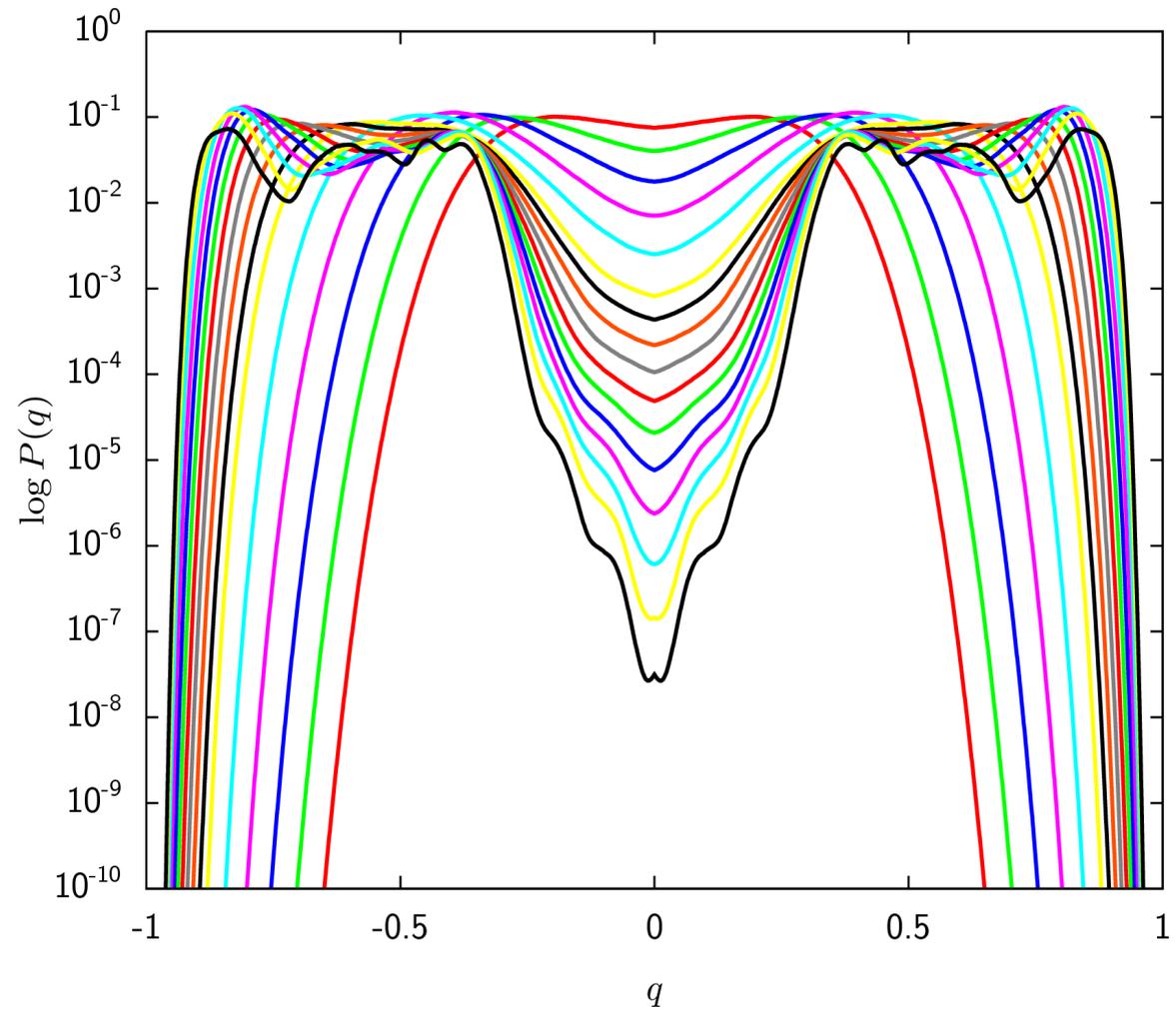
Elmar Bittner
Universität Leipzig

Collaborators: A. Nußbaumer and W. Janke

Spring School on **Monte Carlo Simulations of Disordered Systems**
Leipzig, April 4, 2008.

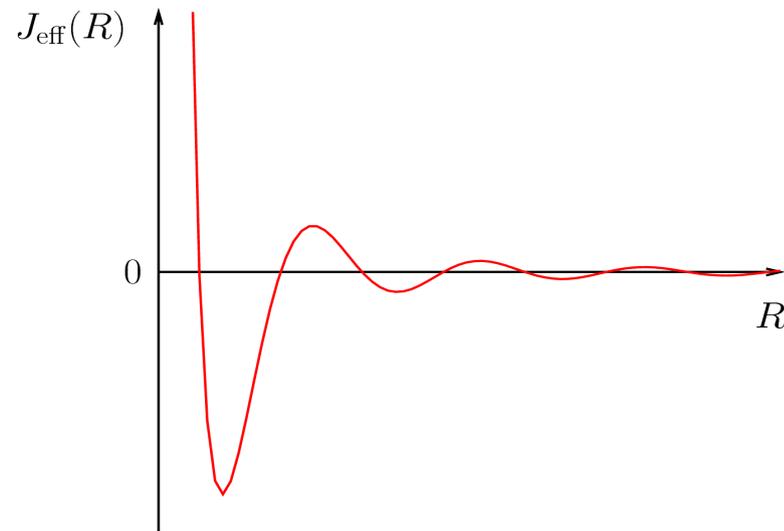
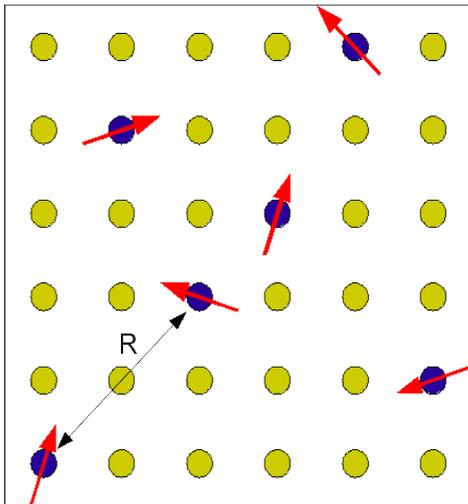
Outlook

- Intro
- Models
- Problems
- Algorithms
- Motivation
- Results



Spin Glass Systems

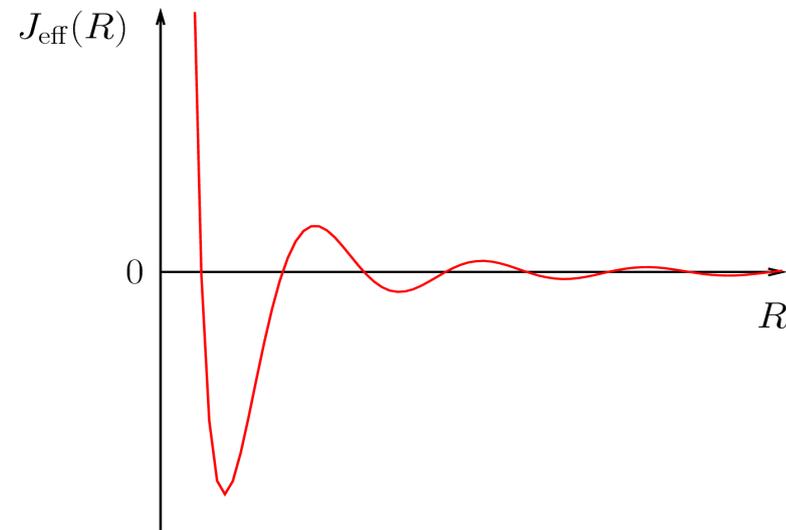
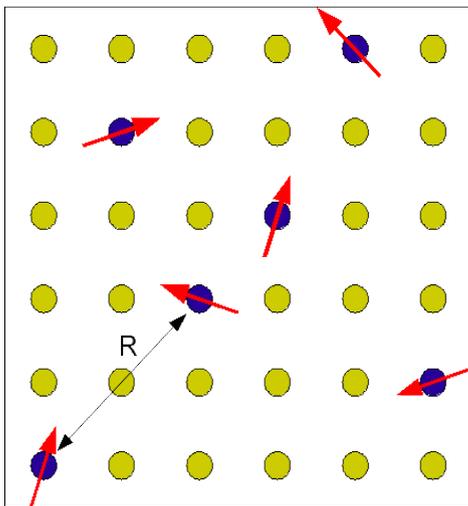
- There are real experimental spin glass systems.
(dilute solutions of magnetic transition metal impurities in noble metal hosts, for instance Au-2.98% Mn or Cu-0.9% Mn)



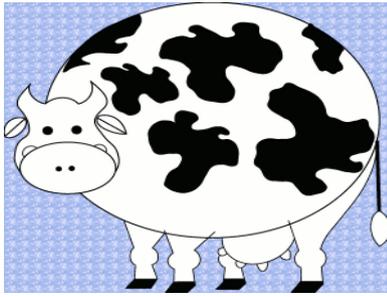
RKKY interaction: $J_{\text{eff}}(R) = J_0 \frac{\cos(2k_F R)}{R^3}$, $k_F R \gg 1$, k_F Fermi wave number

Basic ingredients for spin-glass behaviour

- **randomness** in course of the dilution process the positions of the impurity moments are randomly distributed
- **competing interactions** due to the oscillations in the effective interaction as a function of the distance R



RKKY interaction: $J_{\text{eff}}(R) = J_0 \frac{\cos(2k_F R)}{R^3}$, $k_F R \gg 1$, k_F Fermi wave number



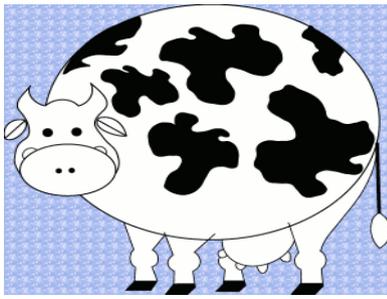
Spherical Cow

Most theory uses the simplest model with these ingredients:

the **Edward-Anderson Model (EA)**

$$H = - \sum_{\langle i, j \rangle} J_{ij} S_i S_j - \sum_i h_i S_i$$

with $S_i = \pm 1$ lie on a regular lattice and the quenched coupling constants J_{ij} .

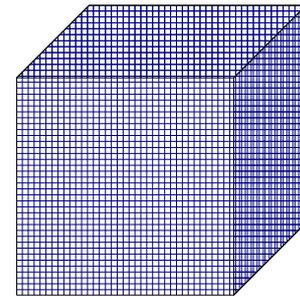


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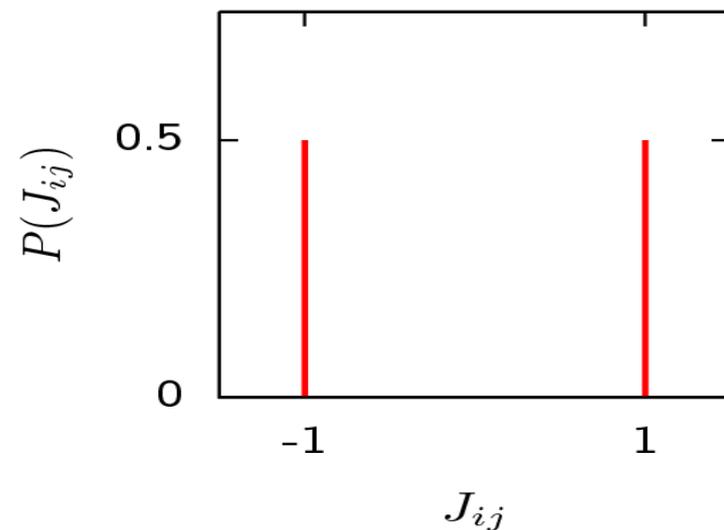


bimodal distribution $J_{ij} = \pm 1$

3D: $T_c \sim 1.16, h_i = 0$

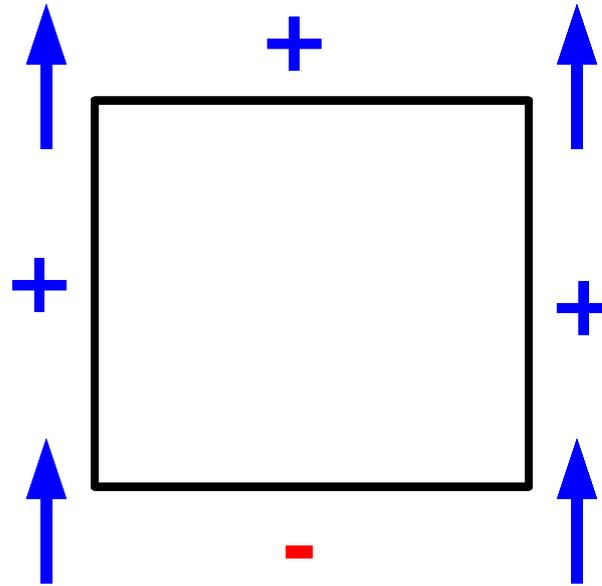
[M. Palassini and S. Caracciolo, Phys. Rev. Lett. 82, 5128 (1999)]

no solution



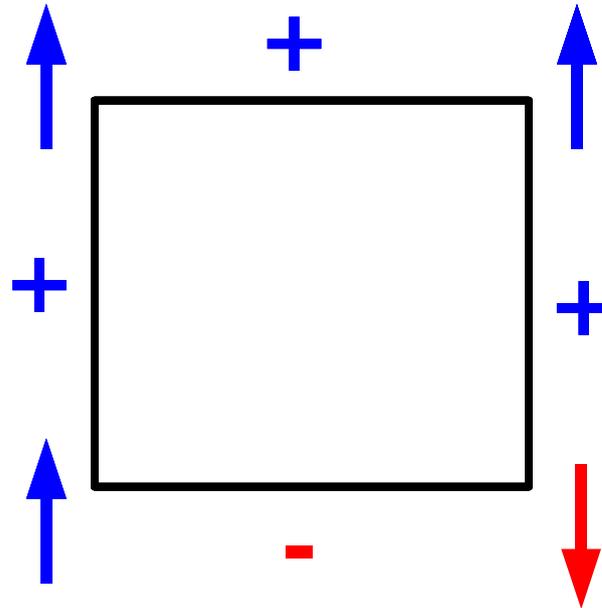
What is a spin glass?

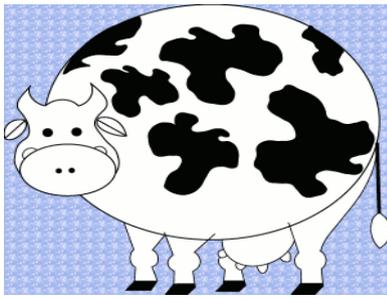
A system with **disorder (randomness)** and **frustration**.



What is a spin glass?

A system with **disorder (randomness)** and **frustration**.

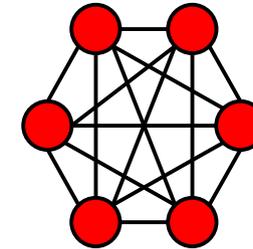




Another Cow

the **Sherrington-Kirkpatrick Model (SK)**

fully connected



$$H = - \sum_{i < j} J_{ij} S_i S_j - \sum_i h_i S_i$$

Gaussian distribution with:

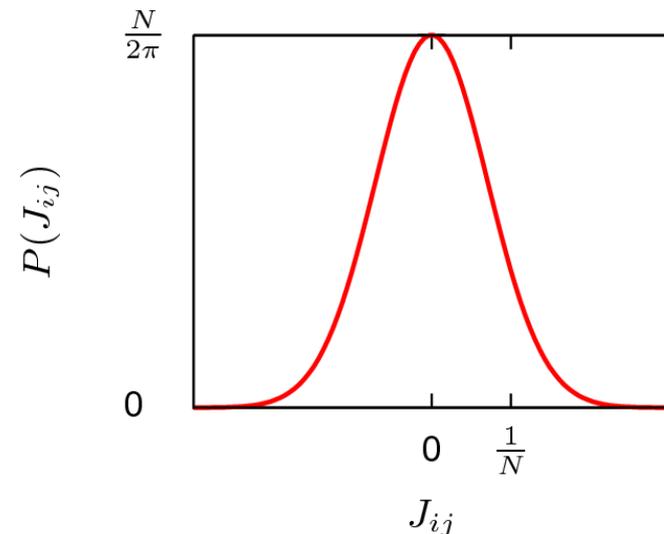
$$\langle J_{ij} \rangle = 0$$

$$\langle J_{ij}^2 \rangle - \langle J_{ij} \rangle^2 = \frac{J}{N}$$

N is the number of spins.

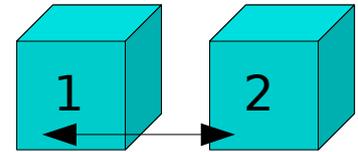
$$T_c = 1, h_i = 0$$

mean field, Parisi's replica solution [PRL 43 (1979) 1754]



Overlap parameter

$$q = \frac{1}{N} \sum_{i=1}^N S_i^1 S_i^2$$



for to (real) replica S_i^1, S_i^2 and given coupling constants $J = \{J_{ij}\}$

$P_J(q)$ probability density of q

$x_J(q) = \int dq' P_J(q')$ cumulative distribution of $P_J(q)$

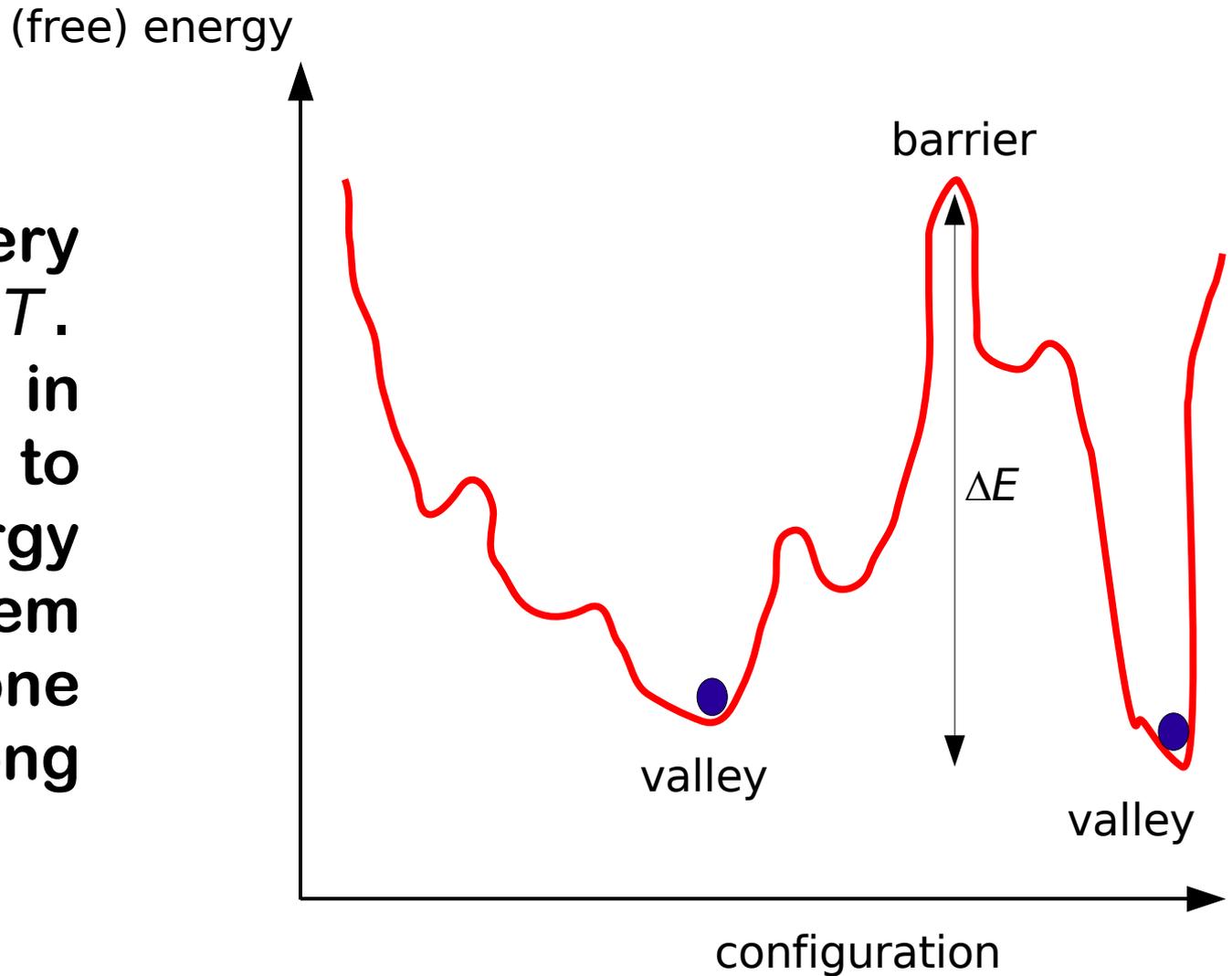
average over the disorder

$$P(q) = [P_J(q)]_{av} = \frac{1}{N_J} \sum_J P_J(q)$$

$$x_J(q) = [x_J(q)]_{av} = \frac{1}{N_J} \sum_J x_J(q)$$

Slow Dynamics

The dynamics is very slow at low T . System is not in equilibrium due to complicated energy landscape: system trapped in one “valley” for long times.



Algorithms

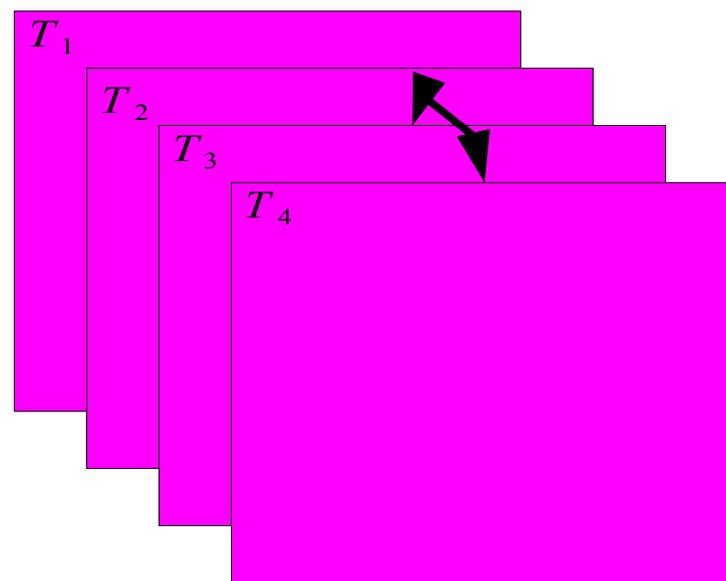
Parallel tempering (PT)

Exchange at regular intervals
system i and $i+1$ with

$$P(i, i+1) = \min[1, \exp(\Delta\beta \Delta E)]$$

expectation values for single system:

$$\langle A \rangle_{T_i} = \langle A_i \rangle$$



System can
decorrelate at high T

[K. Hukushima and K. Nemoto, J. Phys. Soc. Jpn. 65 (1996) 1604]

Talk of H. Katzgraber

Algorithms

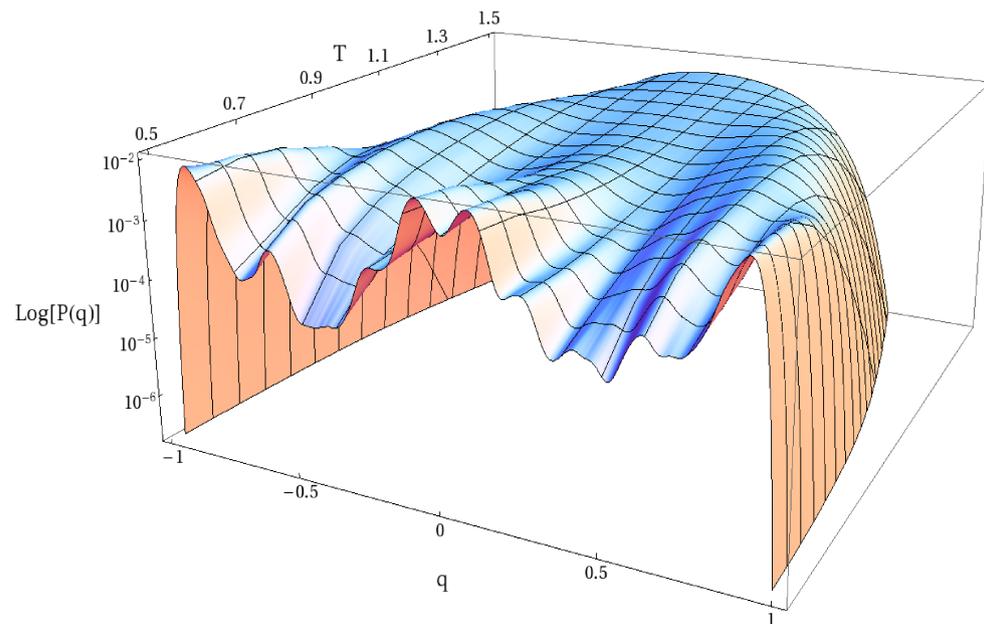
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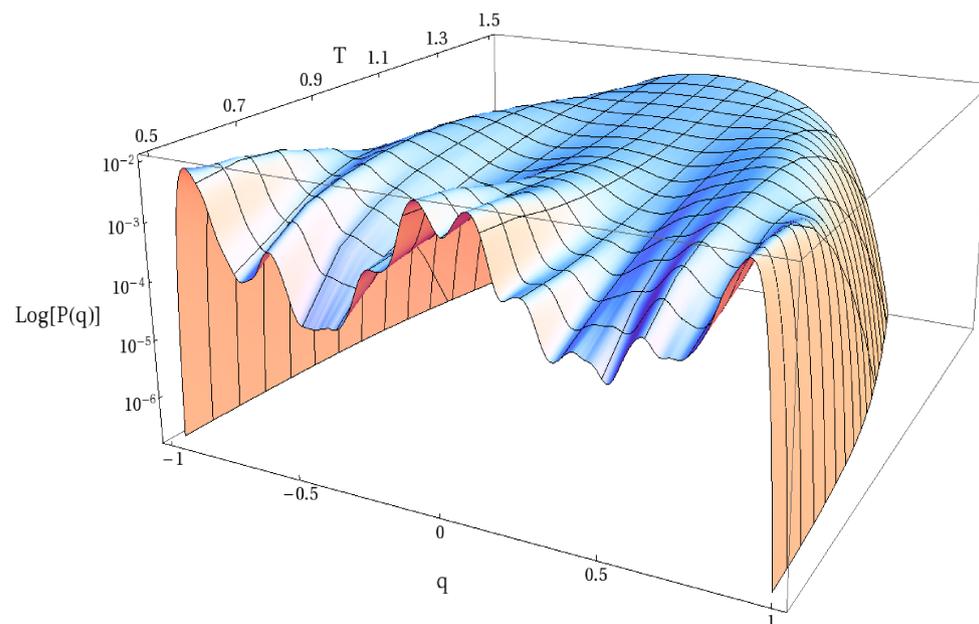
Multioverlap Algorithm (MuQ)

non-Boltzmann sampling with multioverlap weights $W(q)$:

$$\exp[-\beta H] W(q)$$

canonical expectation values:

$$\langle O \rangle^{\text{can}} = \frac{\langle W O \rangle}{\langle W \rangle}$$



System can reach highly suppressed states

[B. Berg, W. Janke, PRL 80 (1998) 4771]

Algorithms

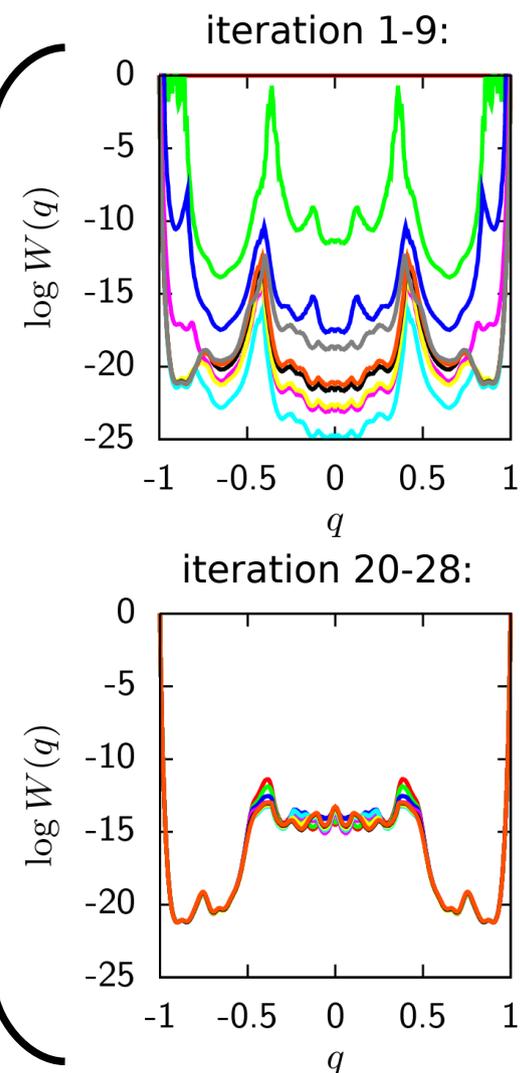
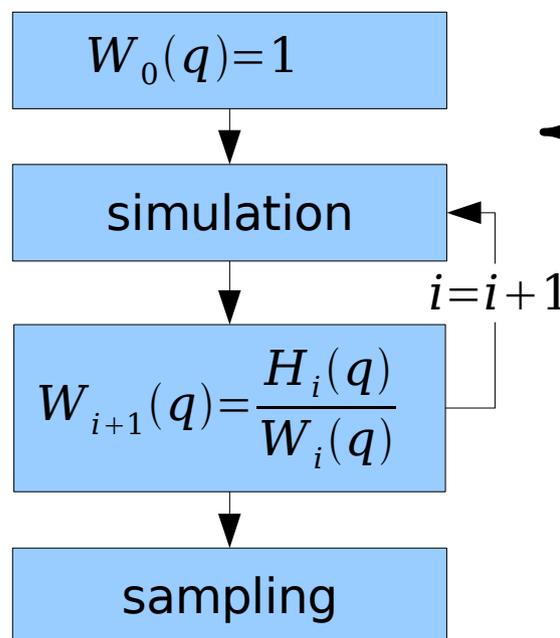
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The 2D Ising Model

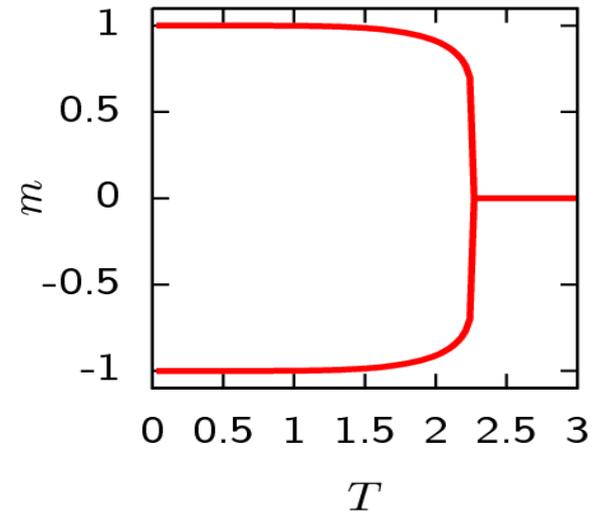
$$H = - \sum_{\langle i, j \rangle} J S_i S_j$$

$$S_i = \pm 1$$

$$T_c = \frac{2}{\log(\sqrt{2}+1)} = 2.269 \dots$$

$$m_0 = \left(1 - \sinh^{-4}\left(\frac{2}{T}\right)\right)^{\frac{1}{8}}$$

Onsager solution, Onsager-Yang solution



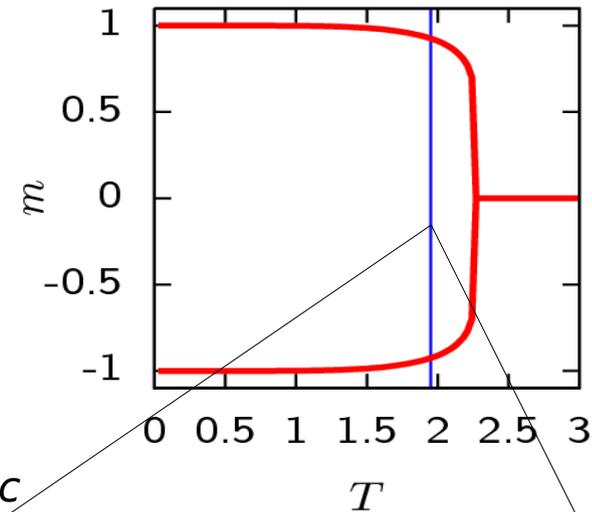
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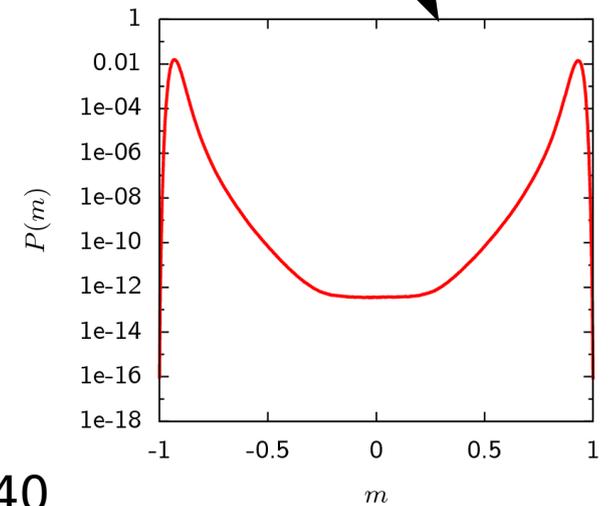
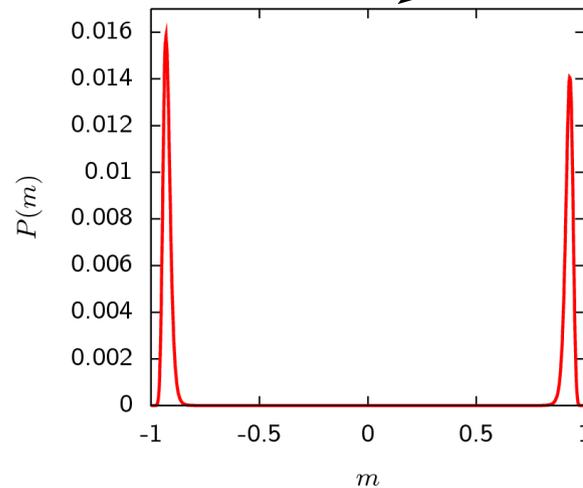
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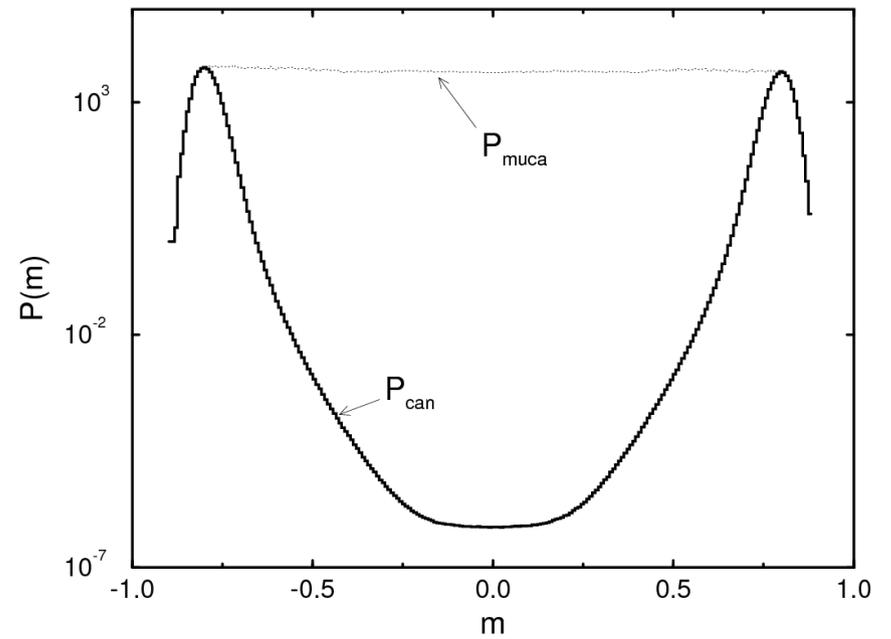
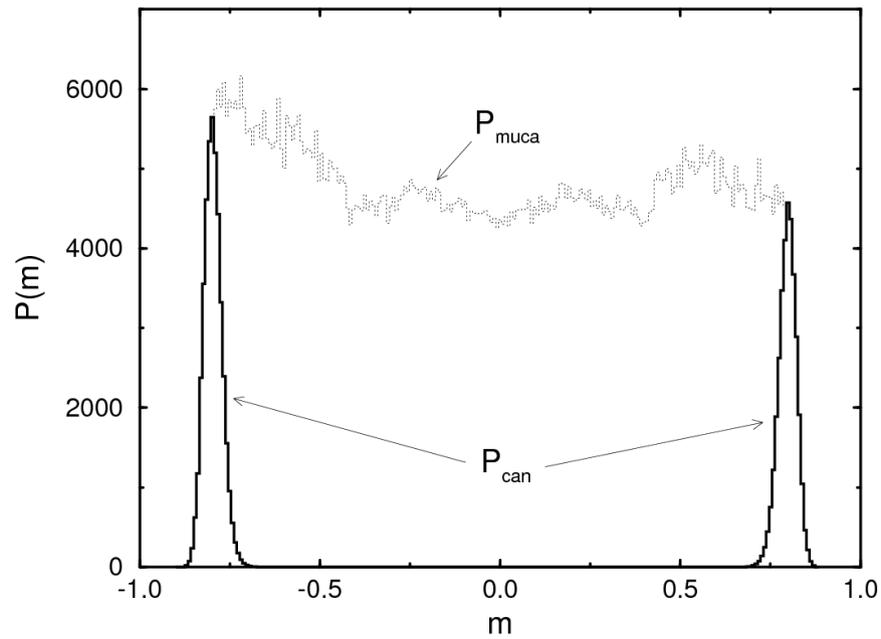


$$T = 1.95 < T_c$$



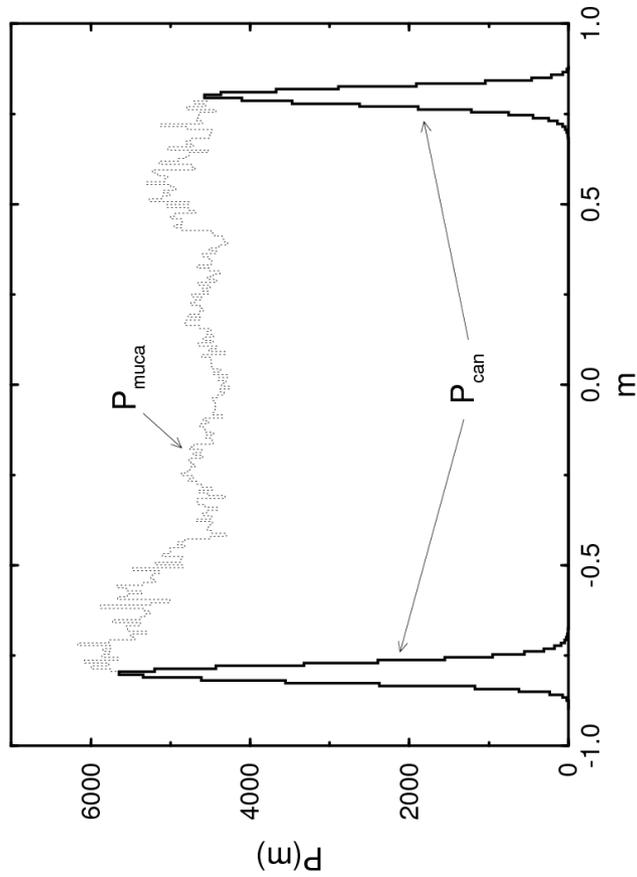
$L=40$

The 2D Ising Model

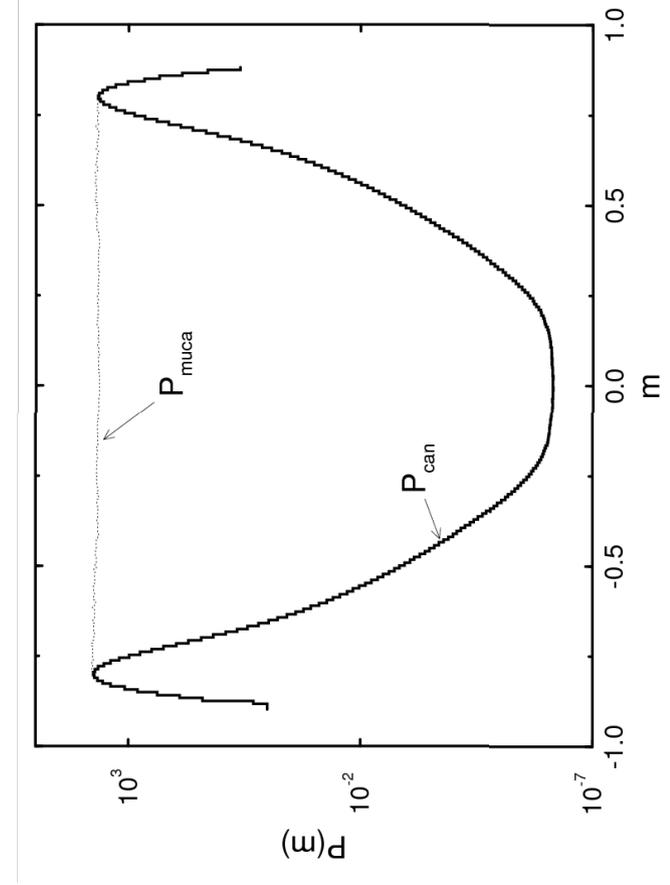


multimagnetical simulation

The 2D Ising Model

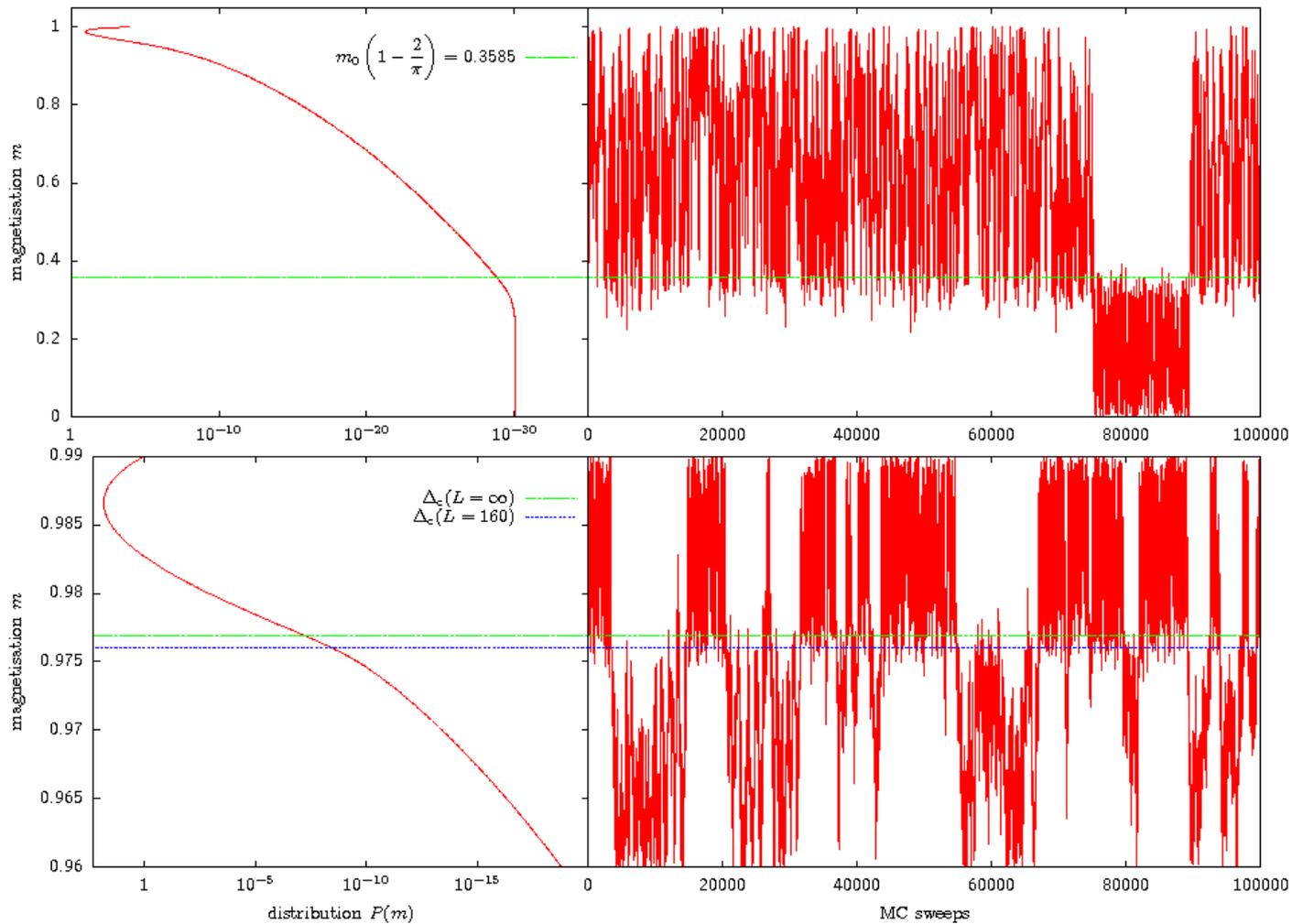


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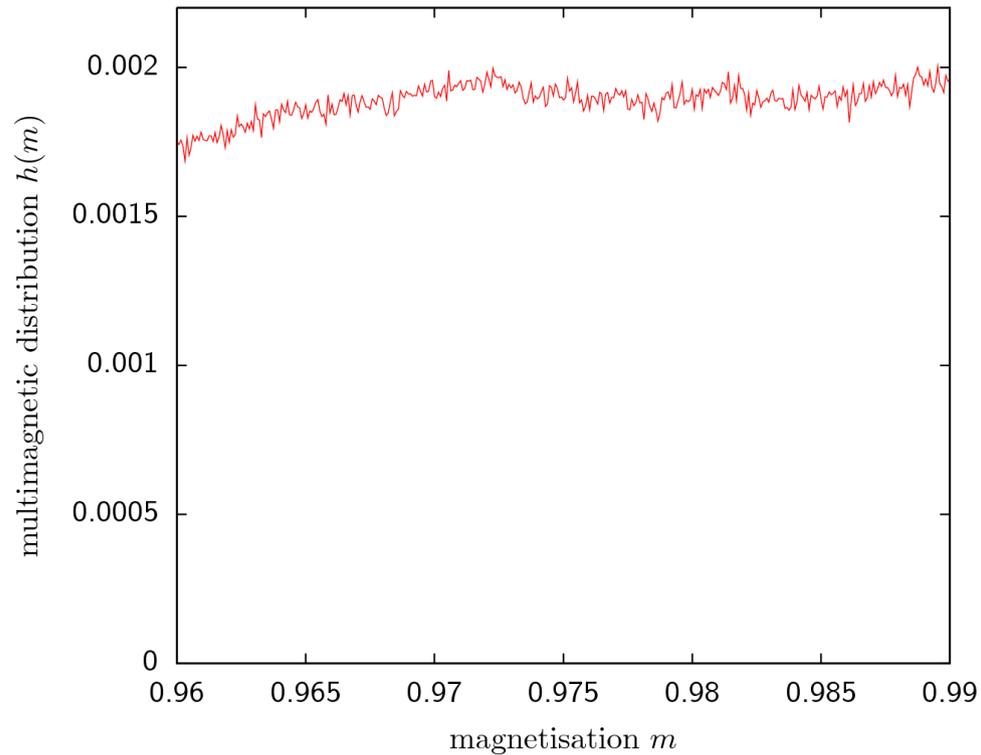
The 2D Ising Model

But there are still problems:



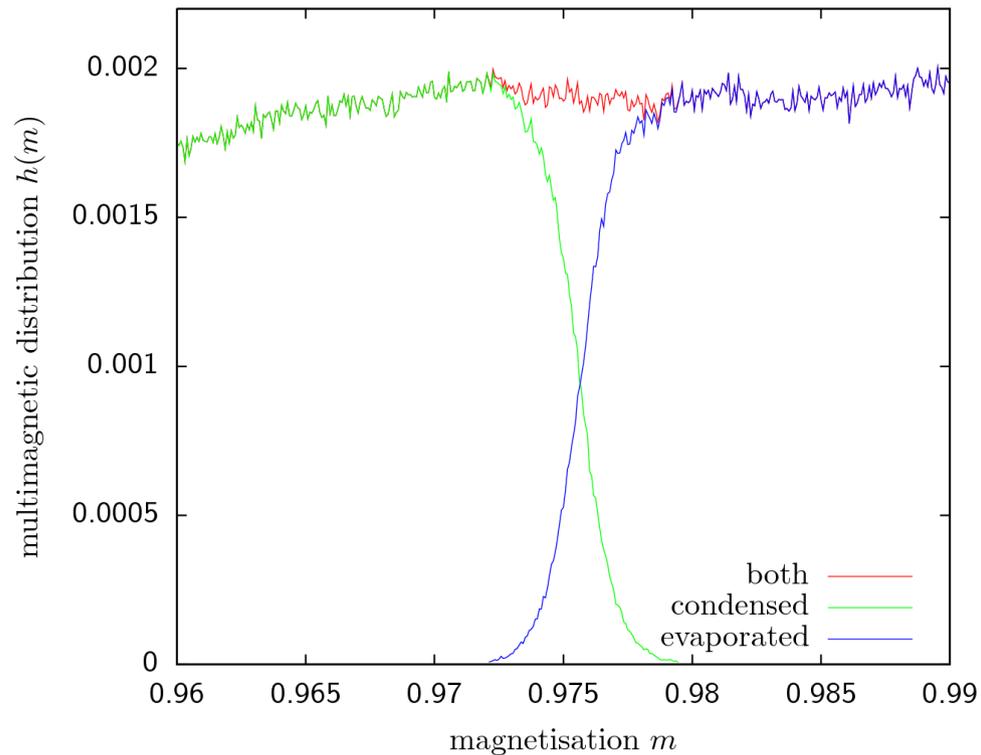
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The 2D Ising Model

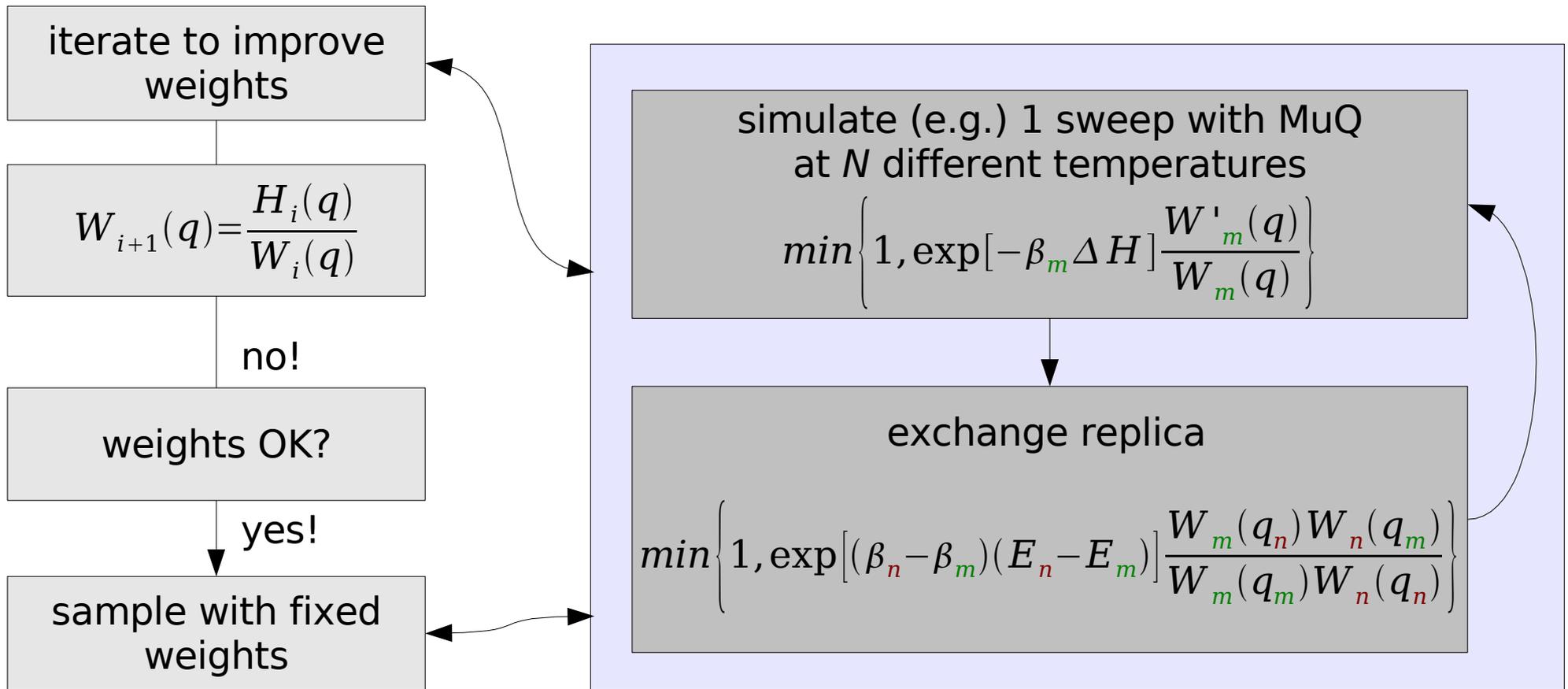
But there are still problems:



T. Neuhaus and J. S. Hager, J. Stat. Phys. 116 (2003) 47,
see also poster from A. Nußbaumer

Algorithms

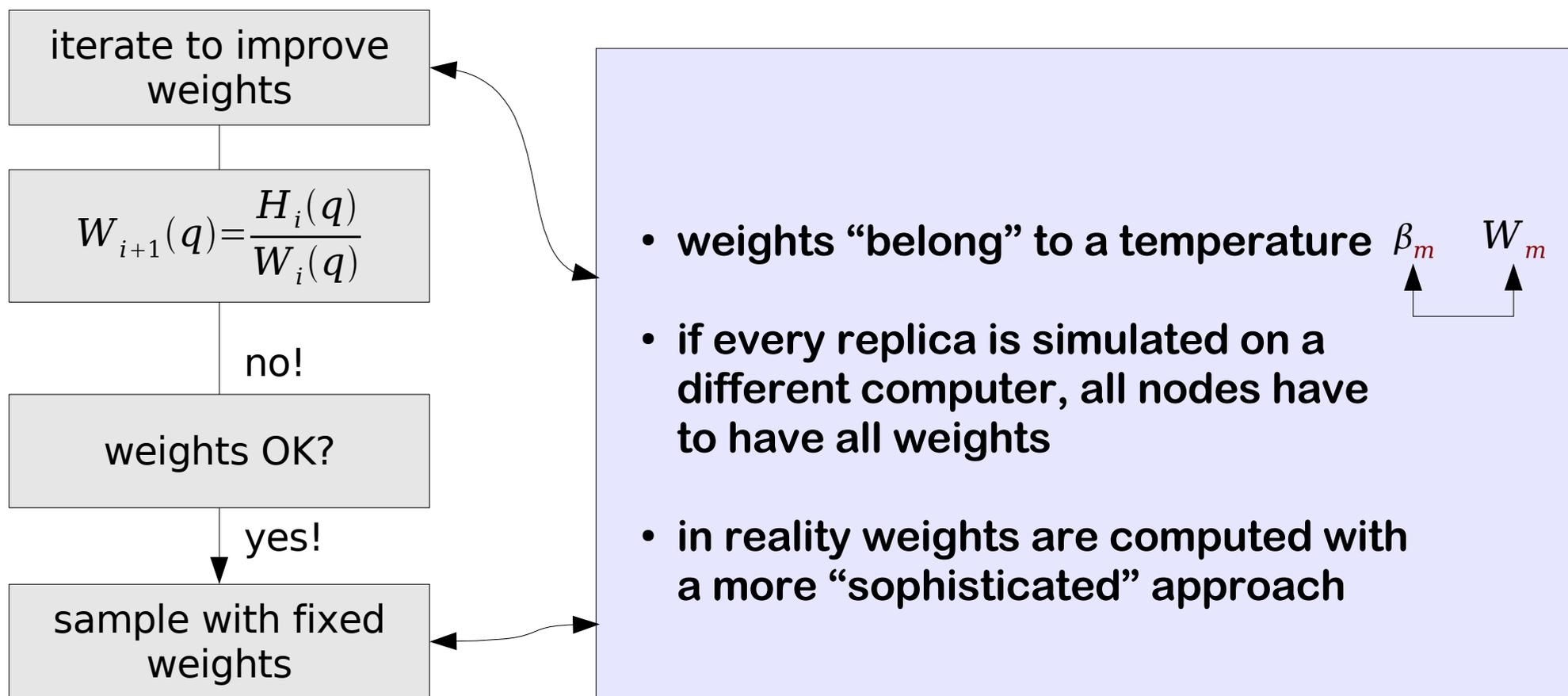
combination of both methods: **PT-MuQ**



[E. Bittner, A. Nußbaumer, W. Janke, in preparation]

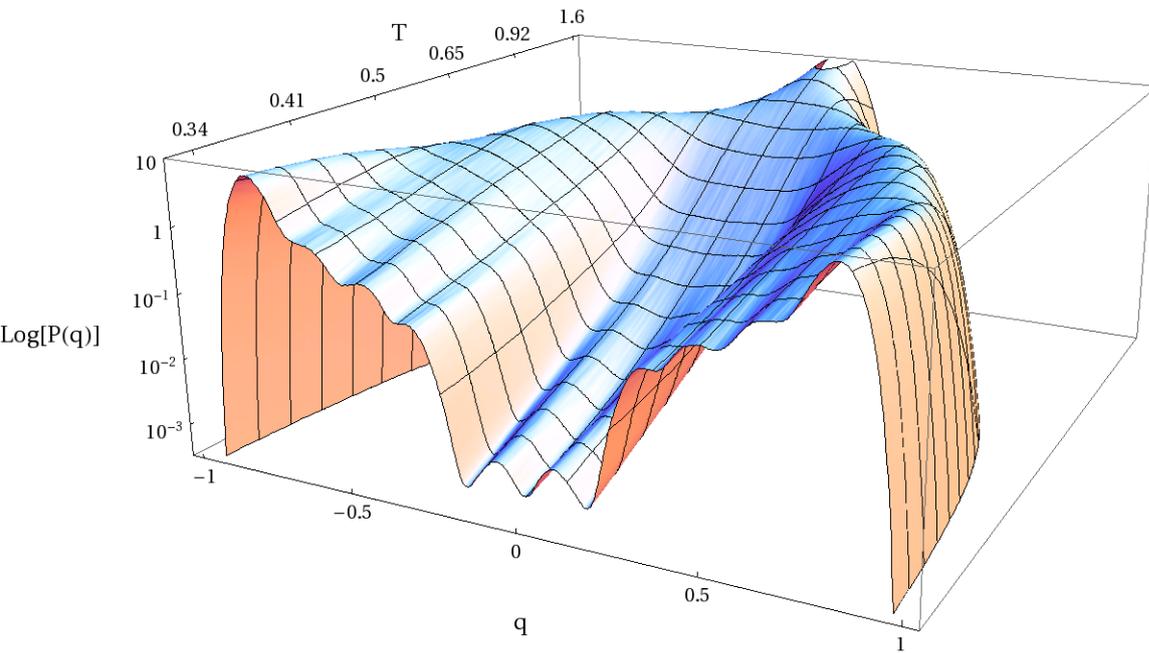
Algorithms

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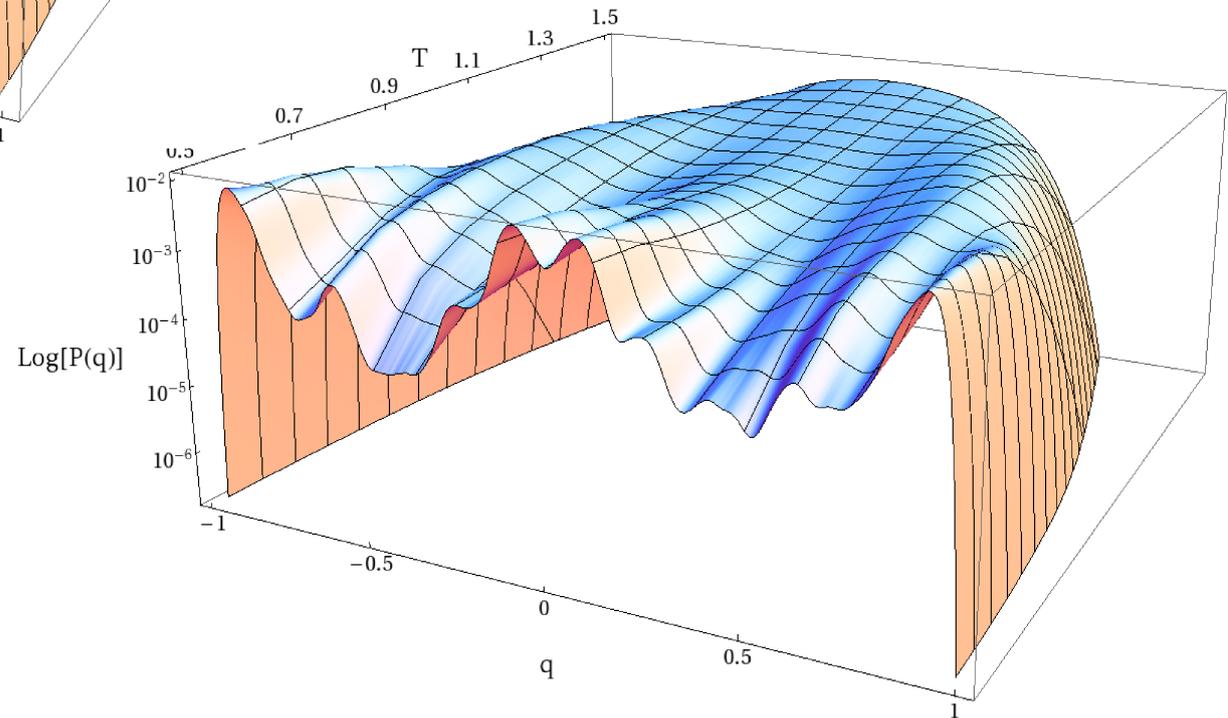
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Algorithms



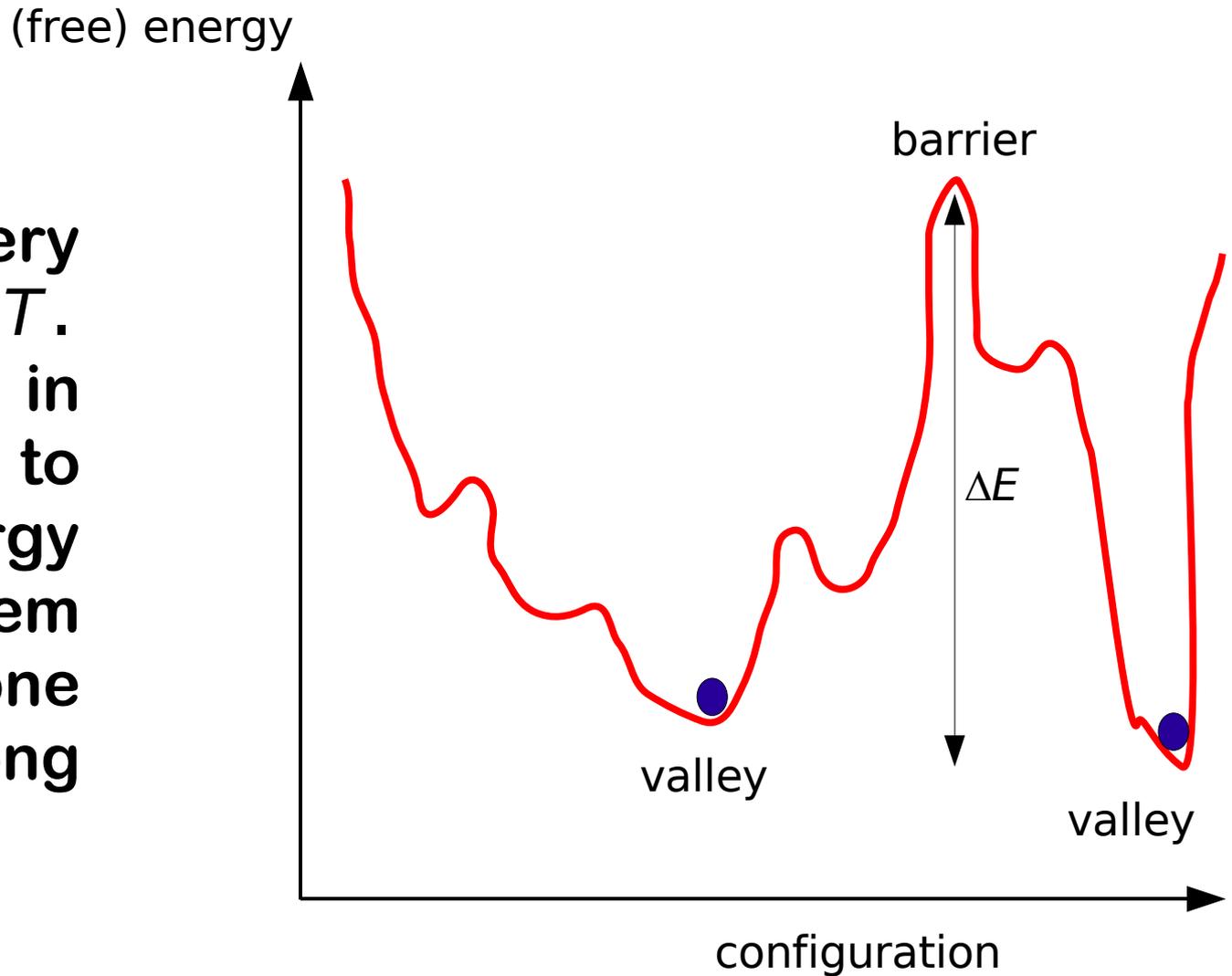
SK model, $N=512$

EA model, $V=8 \times 8 \times 8$



Slow Dynamics

The dynamics is very slow at low T . System is not in equilibrium due to complicated energy landscape: system trapped in one “valley” for long times.



Main objective: barrier heights

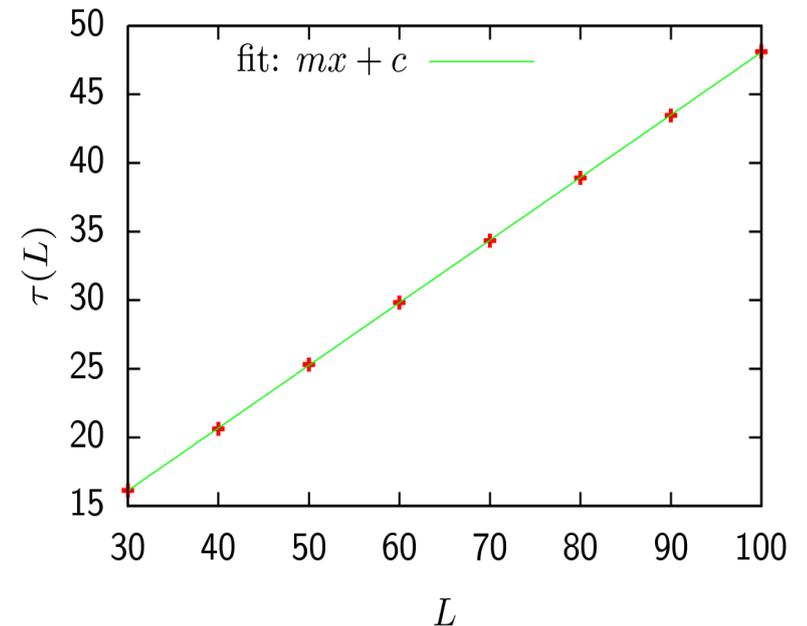
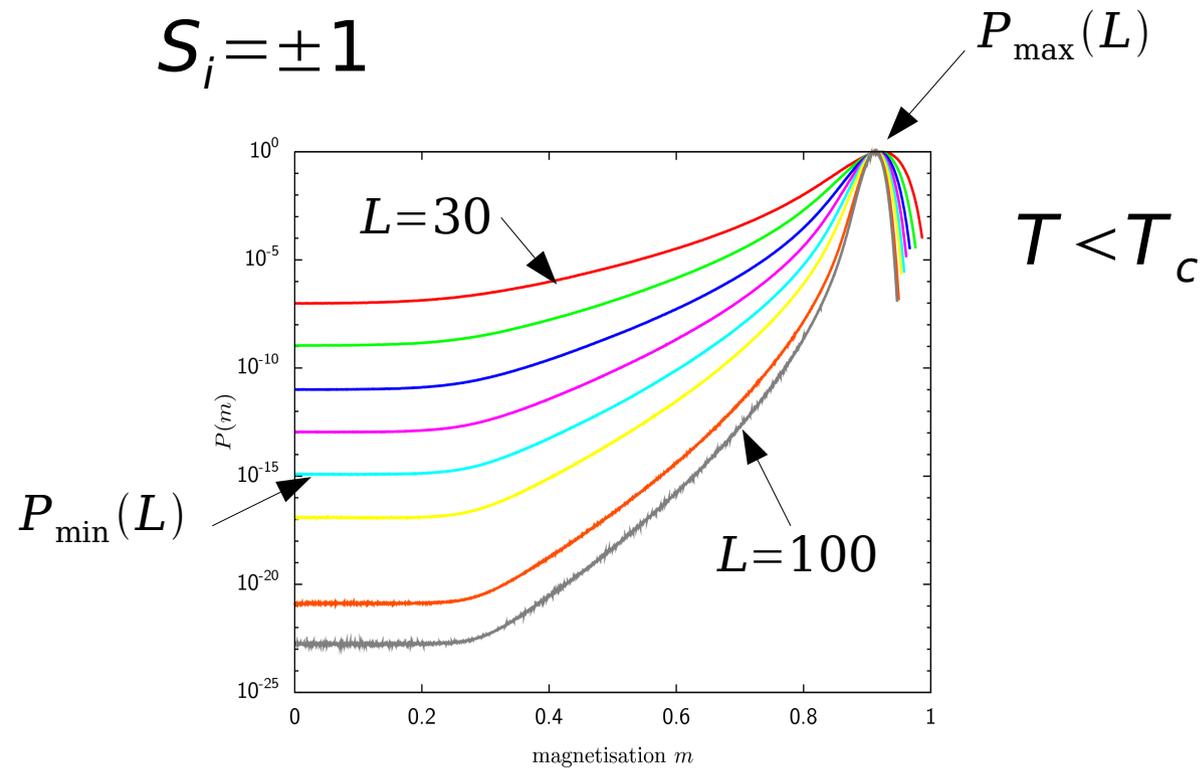
2D Ising Model

$$H = - \sum_{\langle i, j \rangle} J S_i S_j$$

$$\frac{P_{\max}(L)}{P_{\min}(L)} \sim \exp[F_B(L)] \equiv \tau(L)$$

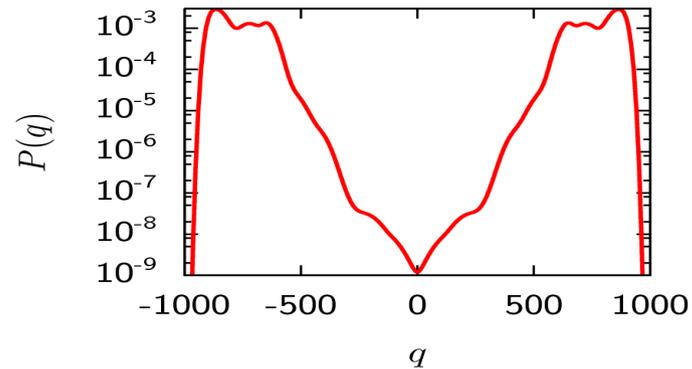
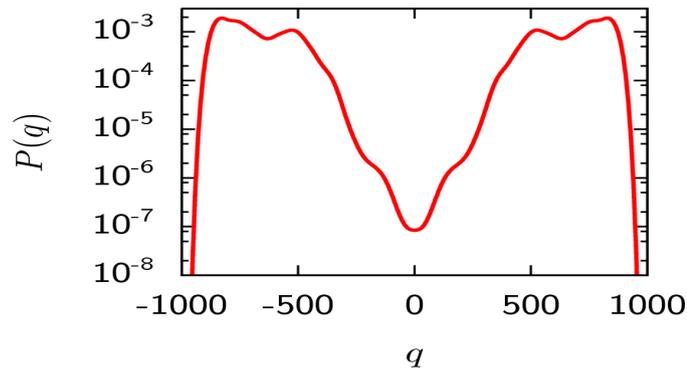
$$S_i = \pm 1$$

$$F_B \sim \sigma(L) \sim 2L$$

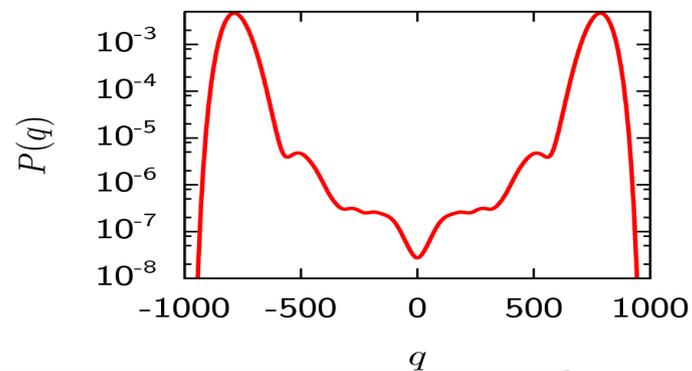
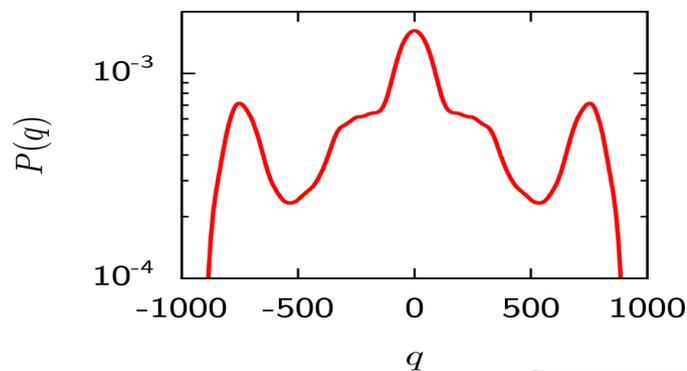


Main objective: barrier heights

Spin glasses:



$T < T_c$



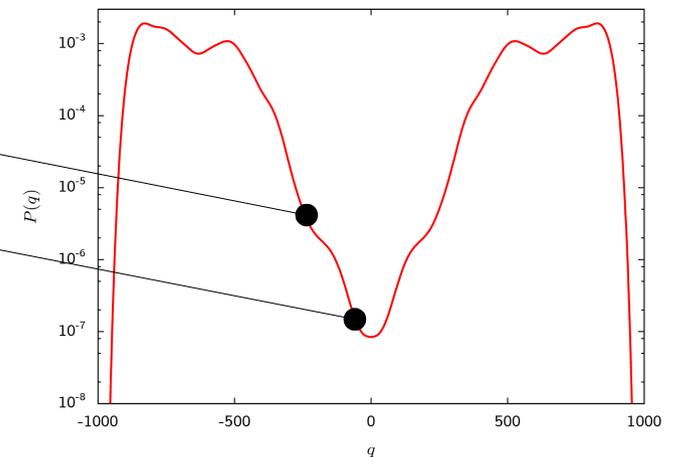
How do we measure
the **size** of the
largest barrier?

1d Markov chain/transition matrix

Definition:

$$T = \begin{pmatrix} 1 - w_{1,2} & w_{1,2} & 0 & \dots \\ w_{2,1} & 1 - w_{2,1} - w_{2,3} & w_{2,3} & \dots \\ 0 & w_{3,2} & 1 - w_{3,2} - w_{3,4} & \dots \\ 0 & 0 & w_{4,3} & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

$$w_{i,j} = \frac{1}{2} \min \left[1, \frac{P(q_j)}{P(q_i)} \right]$$



[B.A. Berg, A. Billoire and W. Janke, PRB 61 (2000), 12143]

1d Markov chain/transition matrix

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T fulfills detailed balance \rightarrow only real eigenvalues

autocorrelation time for q : $\tau_B^q = \frac{1}{N \log(\lambda_1)}$

λ_1 second largest eigenvalue ($\lambda_0 = 1$)

Motivation

theoretical predictions for **mean-field model (SK)**:
barrier between time-reversed states scales with
system size as

$$N^{1/3}$$

(Rodgers and Moore, 1988)

(Kinzelbach and Horner, 1991)

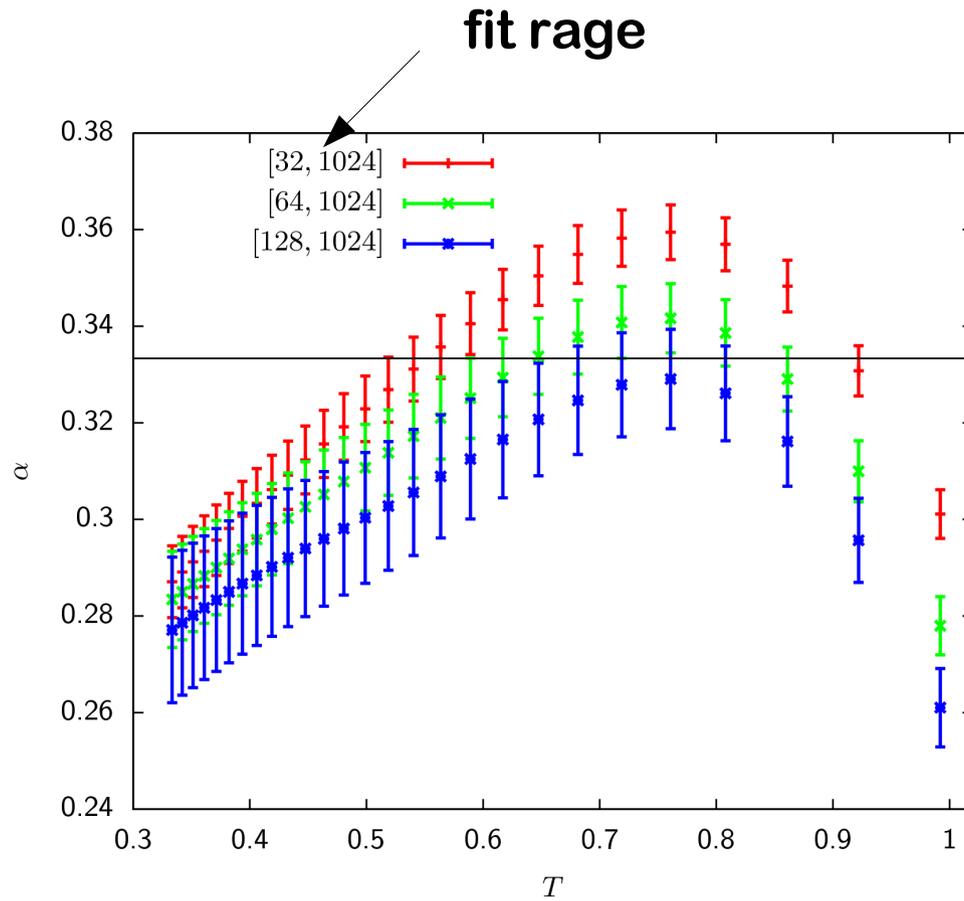
results for **short-ranged models (EA)** are far away
from the mean-field theory limit

$$c_1 + c_2 \ln(N)$$

(Berg, Billoire, and Janke, 2000)

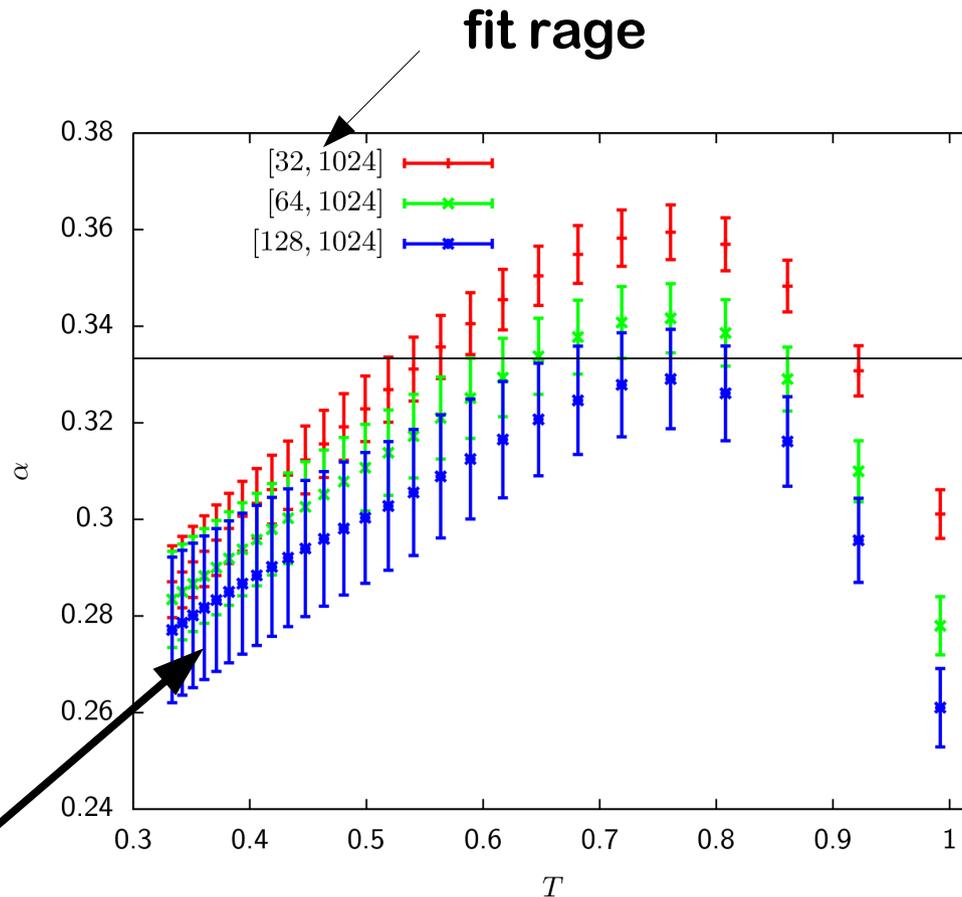
Results for the SK Model

$$\tau_B \propto \exp(cN^\alpha)$$



Results for the SK Model

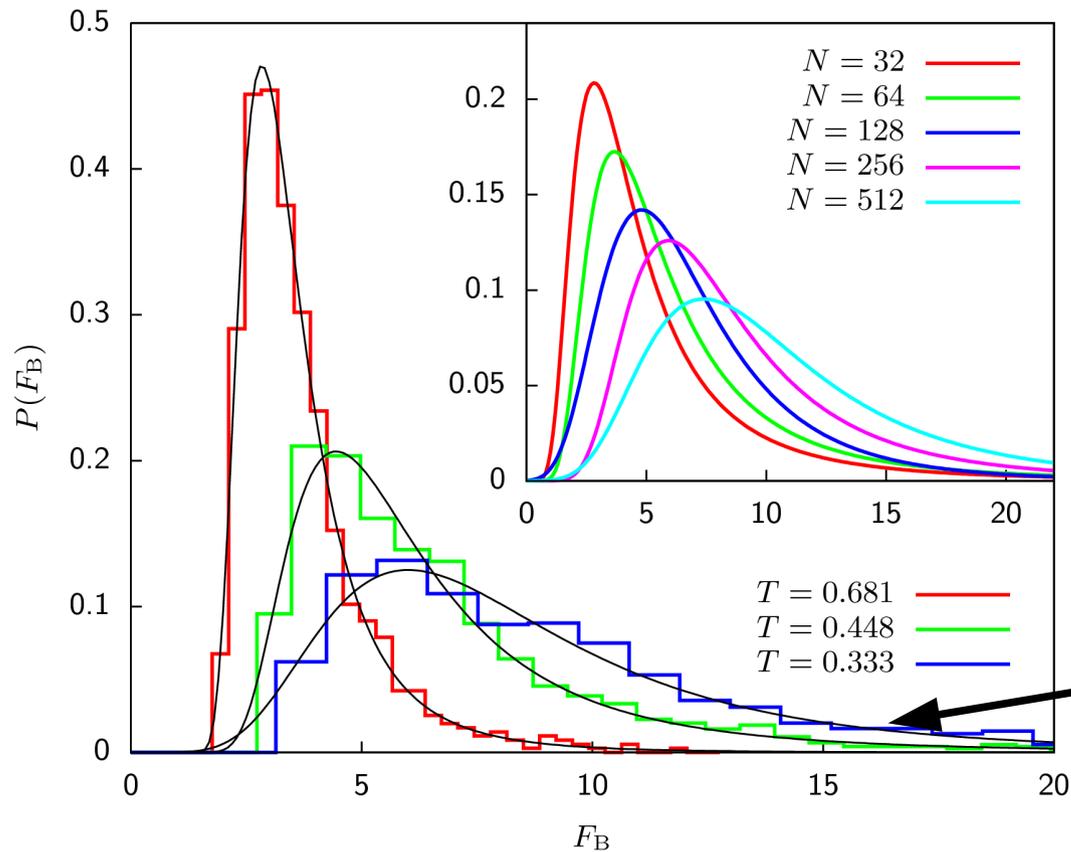
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non-self-averaging (SK)

[L.A. Pastur and M.V. Shcherbina, J. Stat. Phys. 62 (1992) 1]

Results for the SK Model



extreme-value distribution

fat tails

Fit integrated probability density:

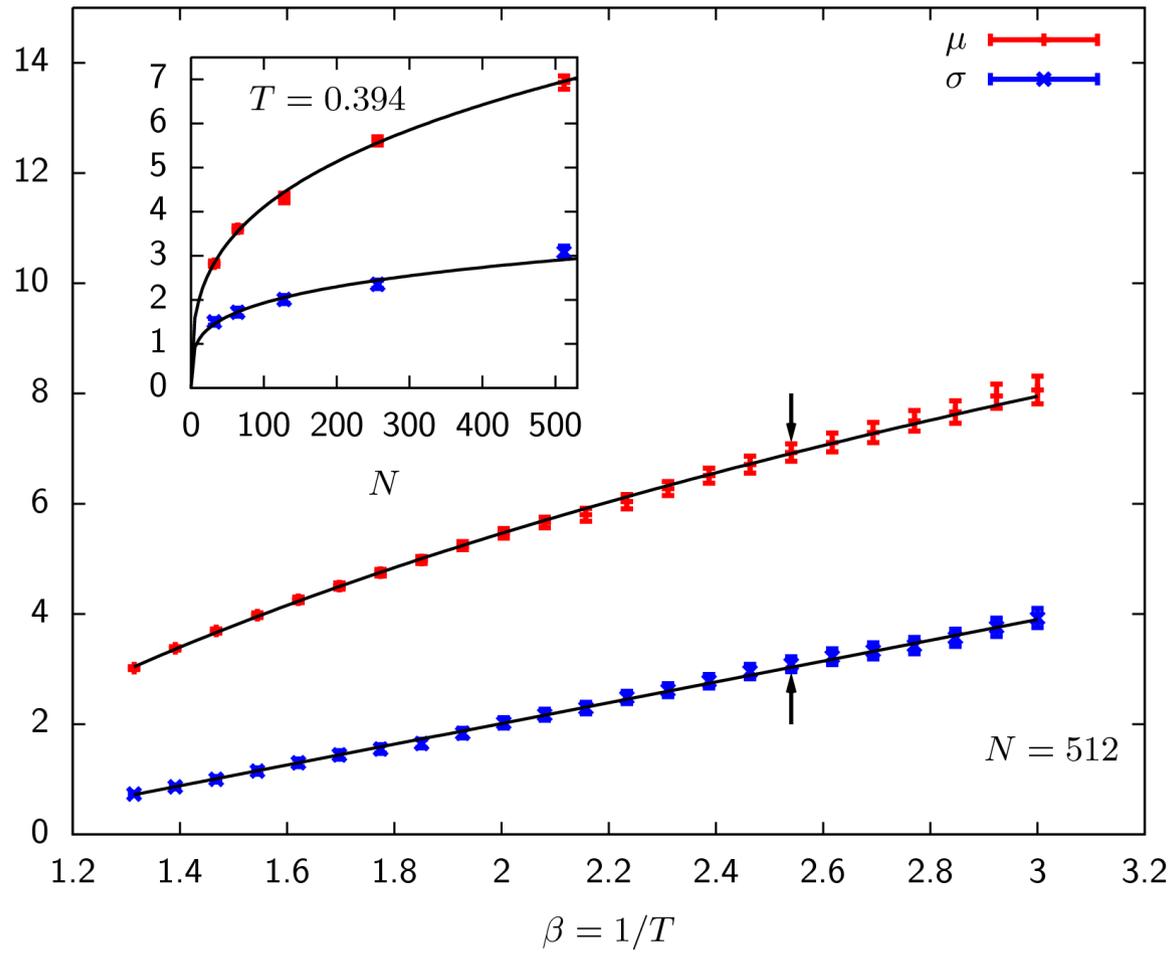
$$F_{\xi; \mu; \sigma}(x) = \exp \left[- \left(1 + \xi \frac{x - \mu}{\sigma} \right)^{-1/\xi} \right]$$

$$T < T_c \longrightarrow$$

$$\xi > 0$$

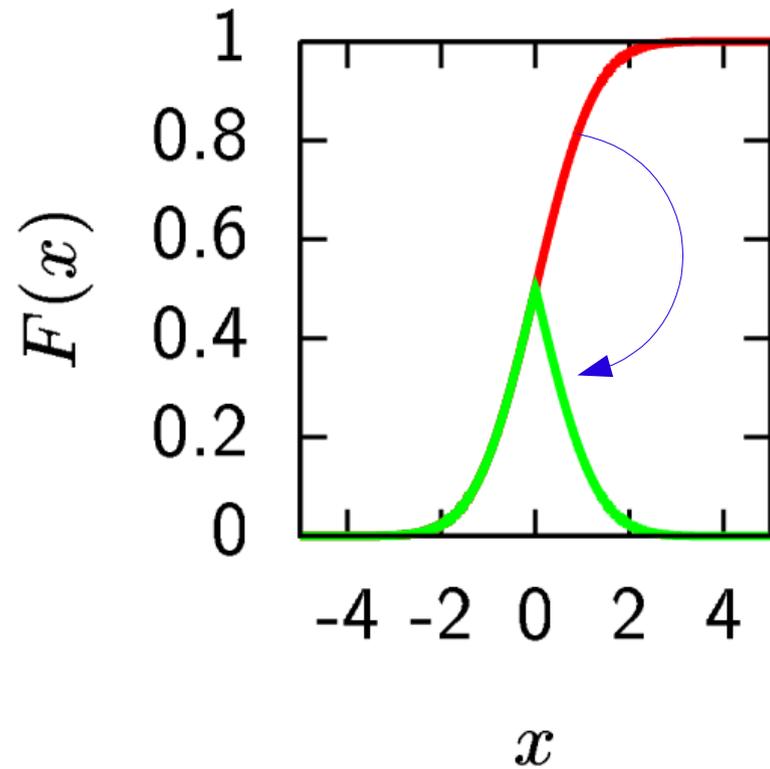
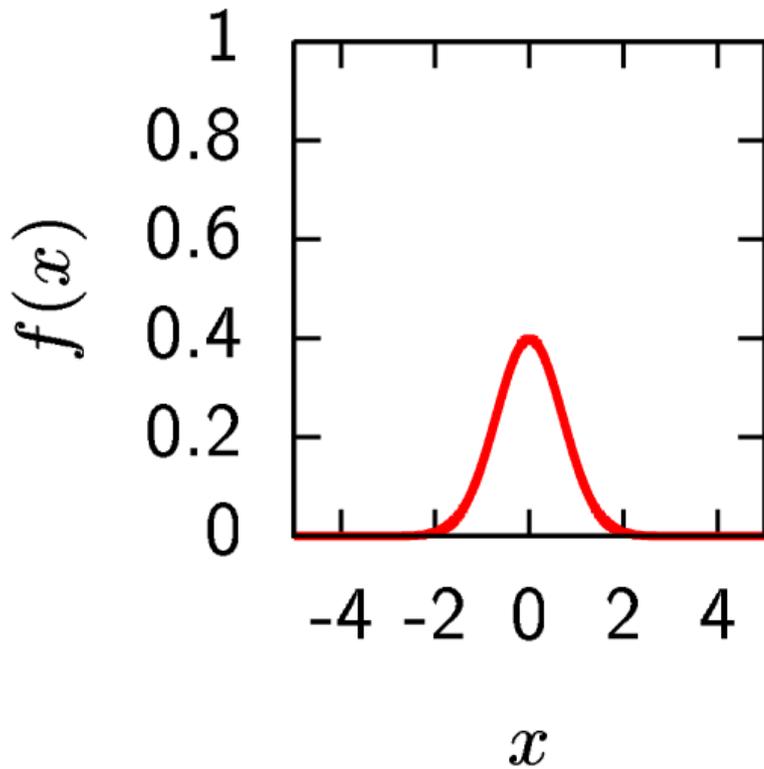
fat tailed (algebraic)
Fréchet distribution

Results for the SK Model



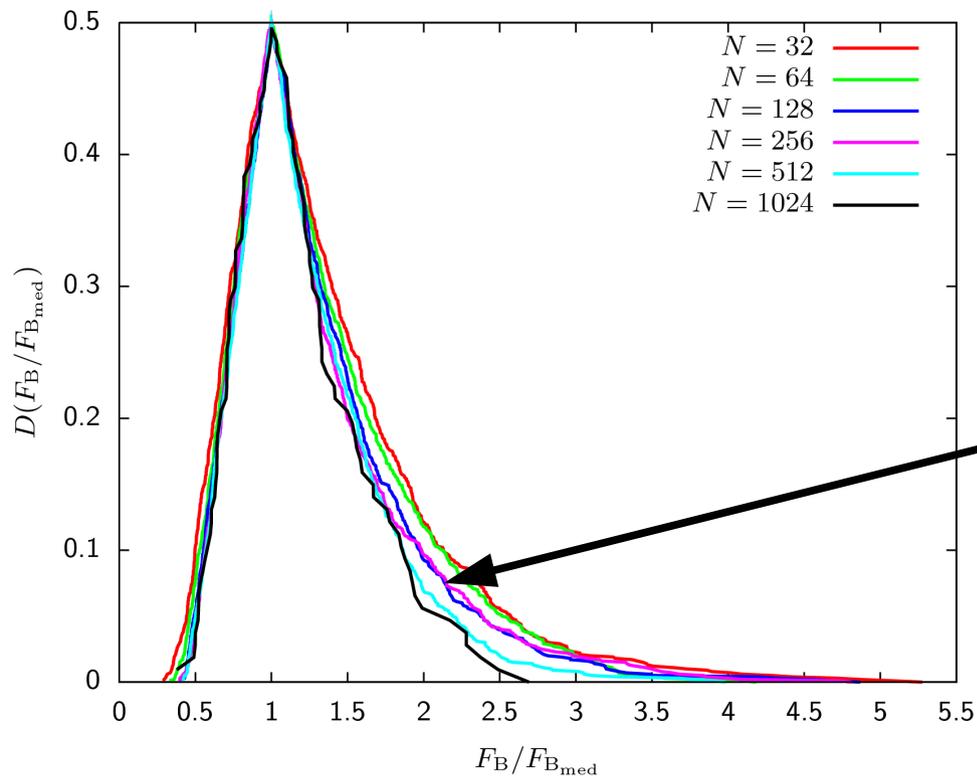
Results for the SK Model

Peaked probability distribution $D(F_B/F_{B_{\text{med}}})$



Results for the SK Model

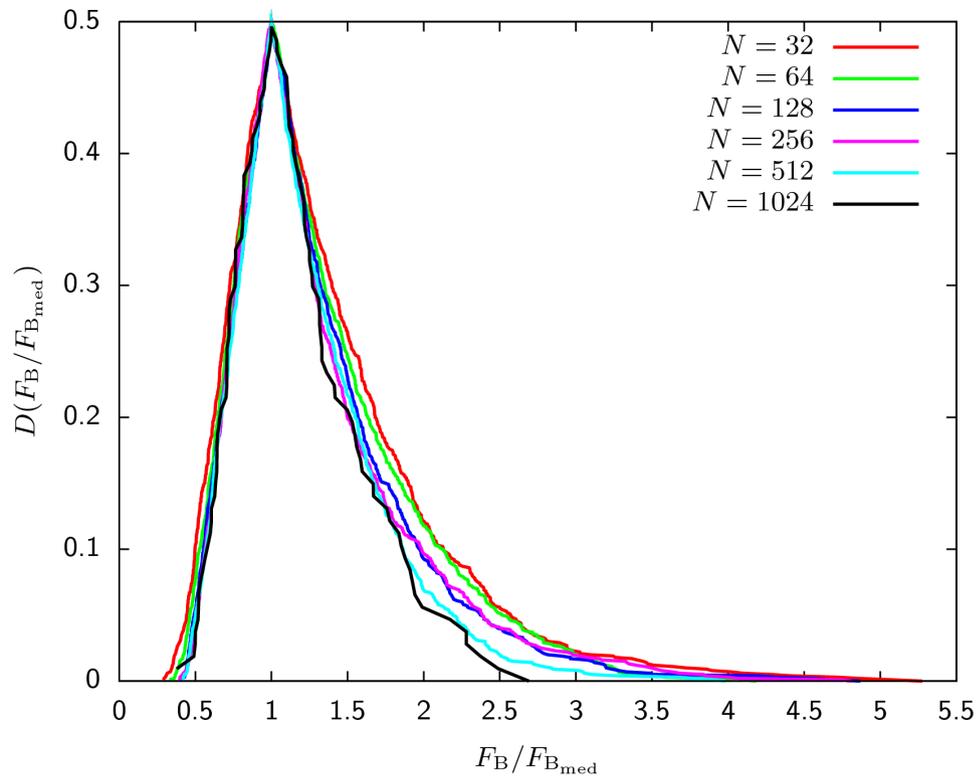
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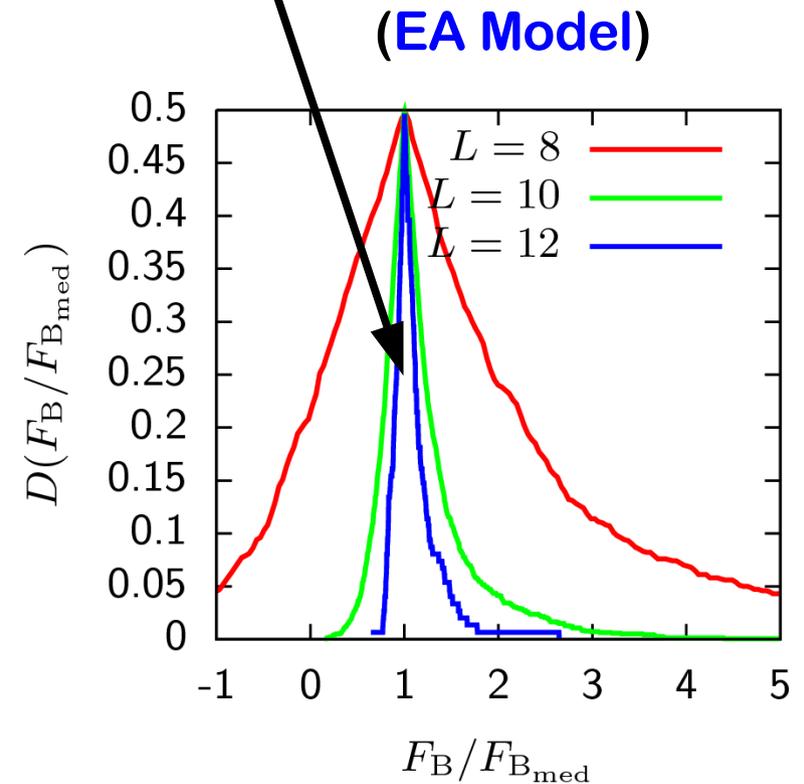
no scaling,
i.e. non-self-averaging

Results for the SK Model

Peaked probability distribution $D(F_B/F_{B_{\text{med}}})$

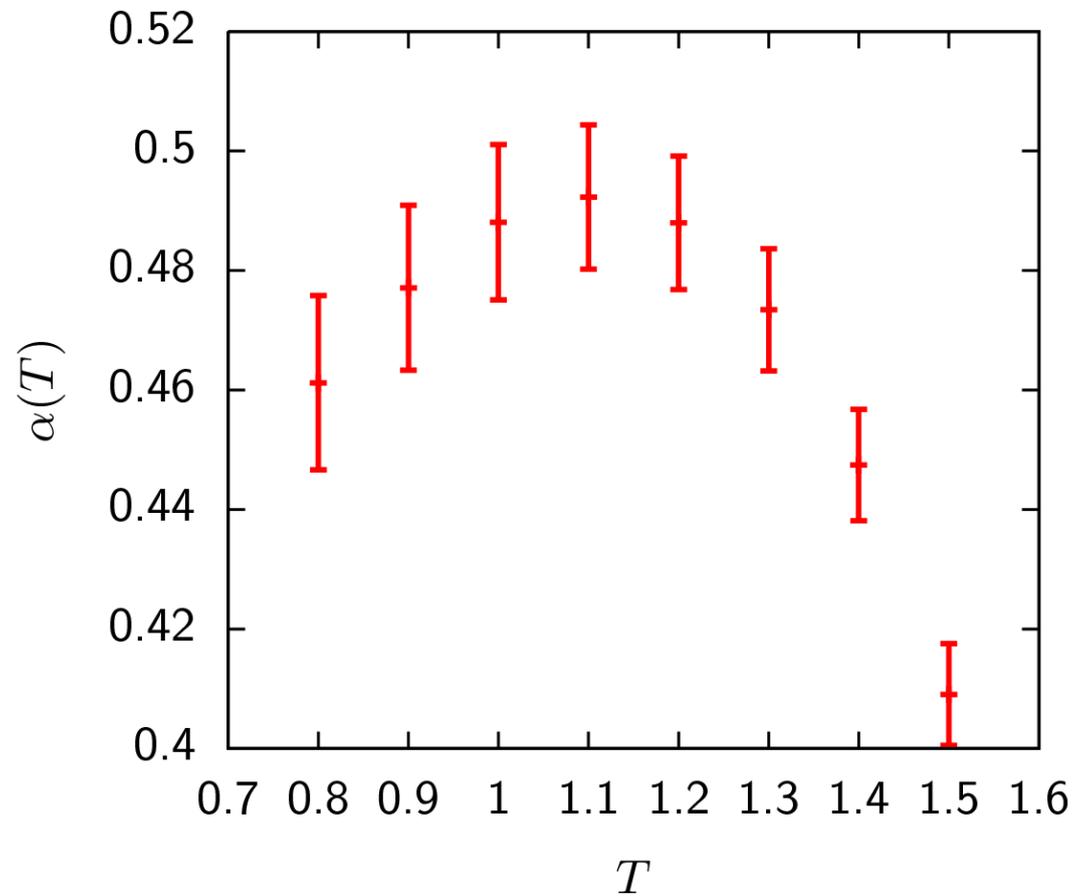


scaling,
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Results for the EA Model

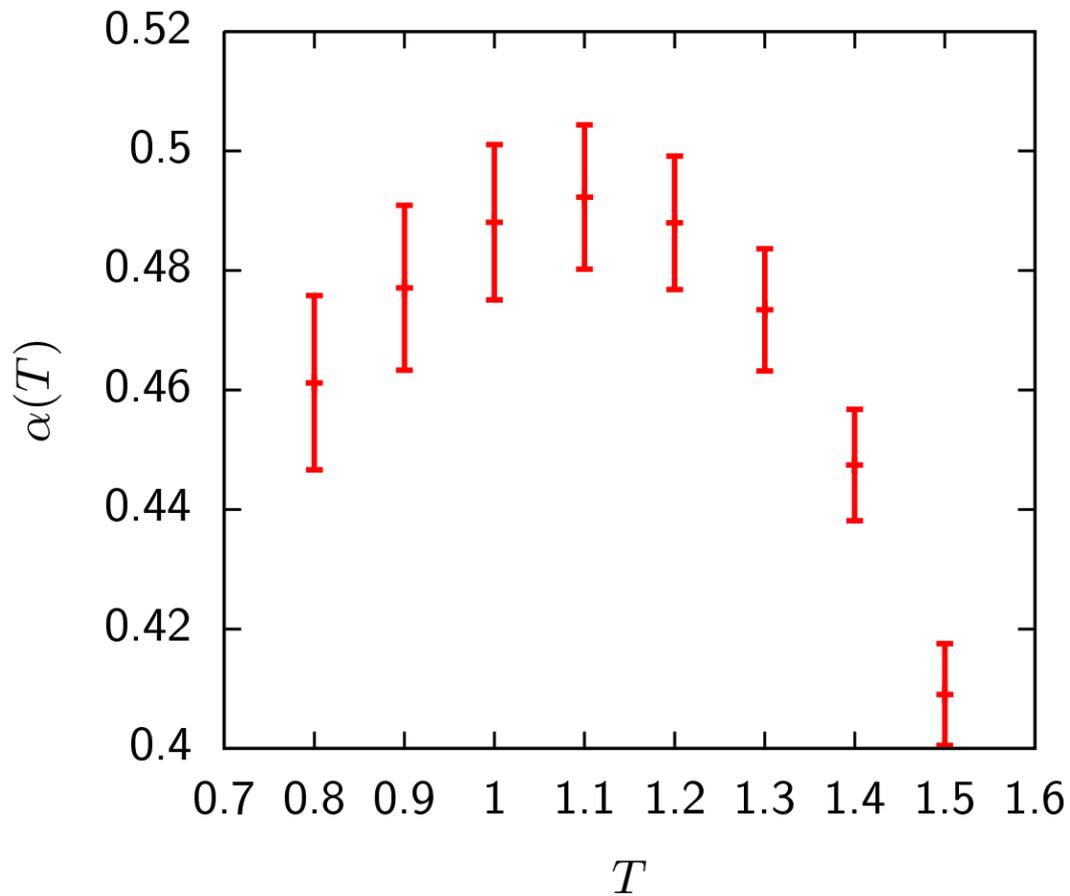
$$\tau_B \propto \exp(cN^\alpha)$$



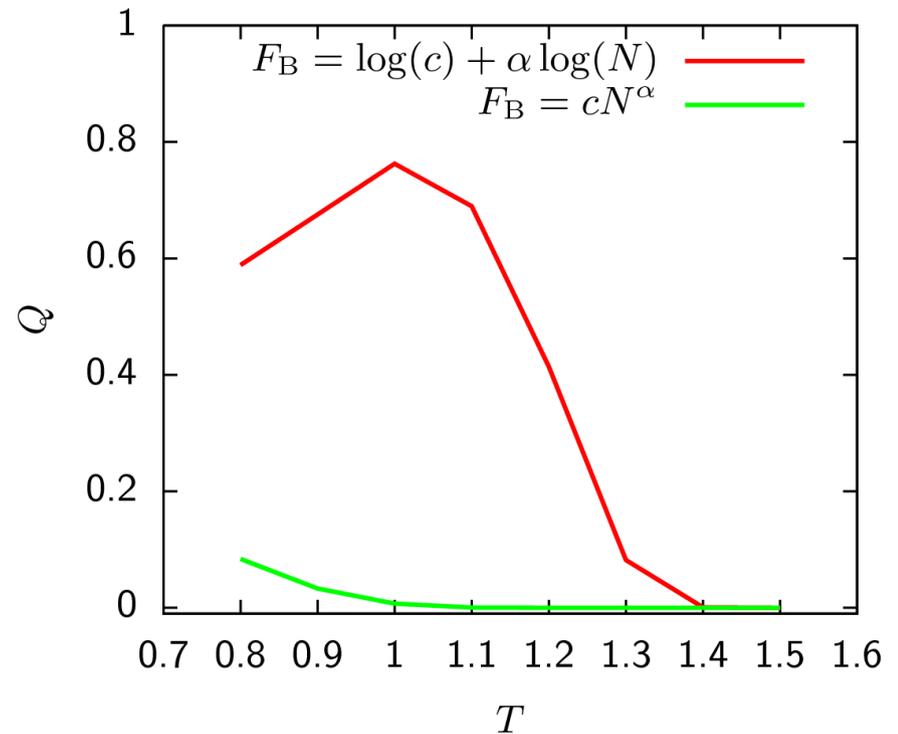
but goodness of fit ...

Results for the EA Model

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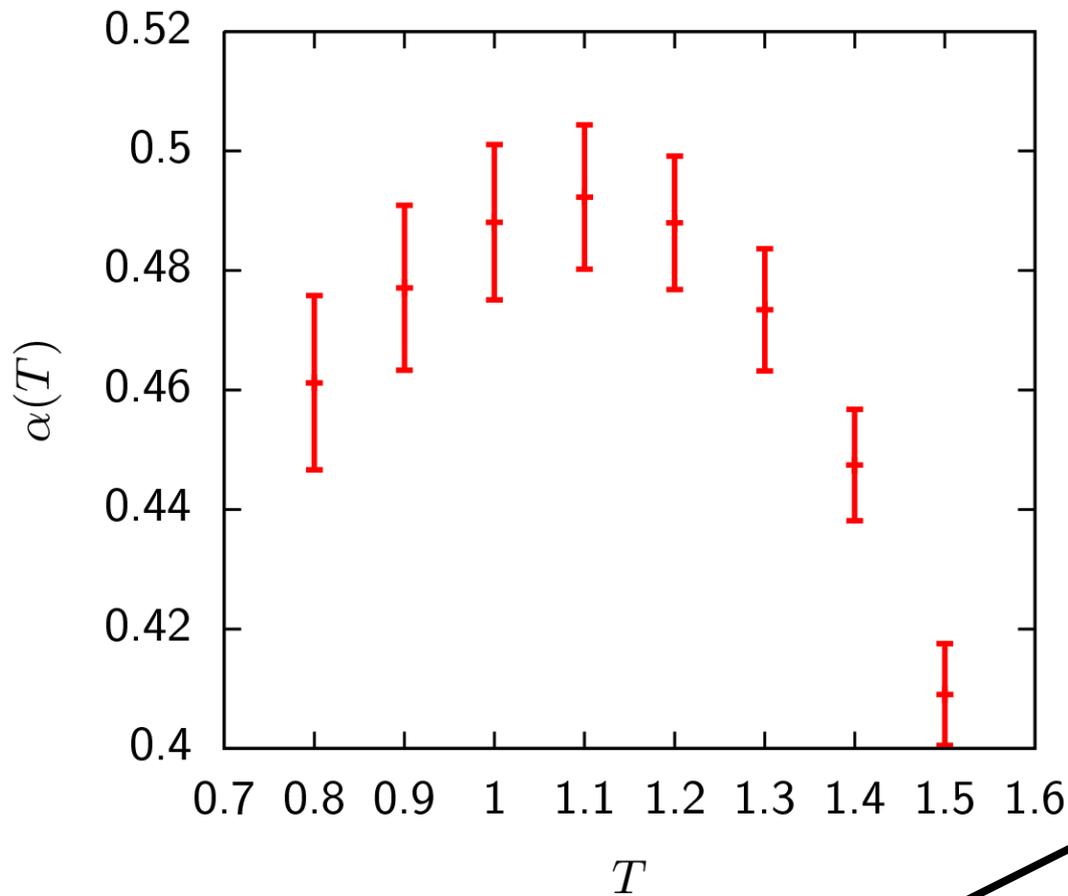


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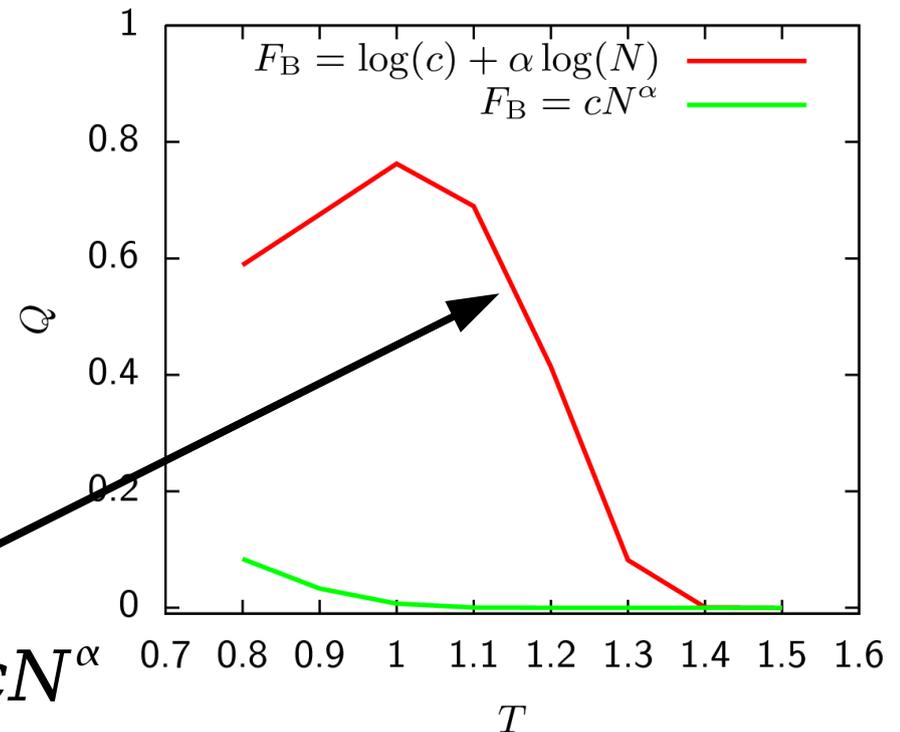


Results for the EA Model

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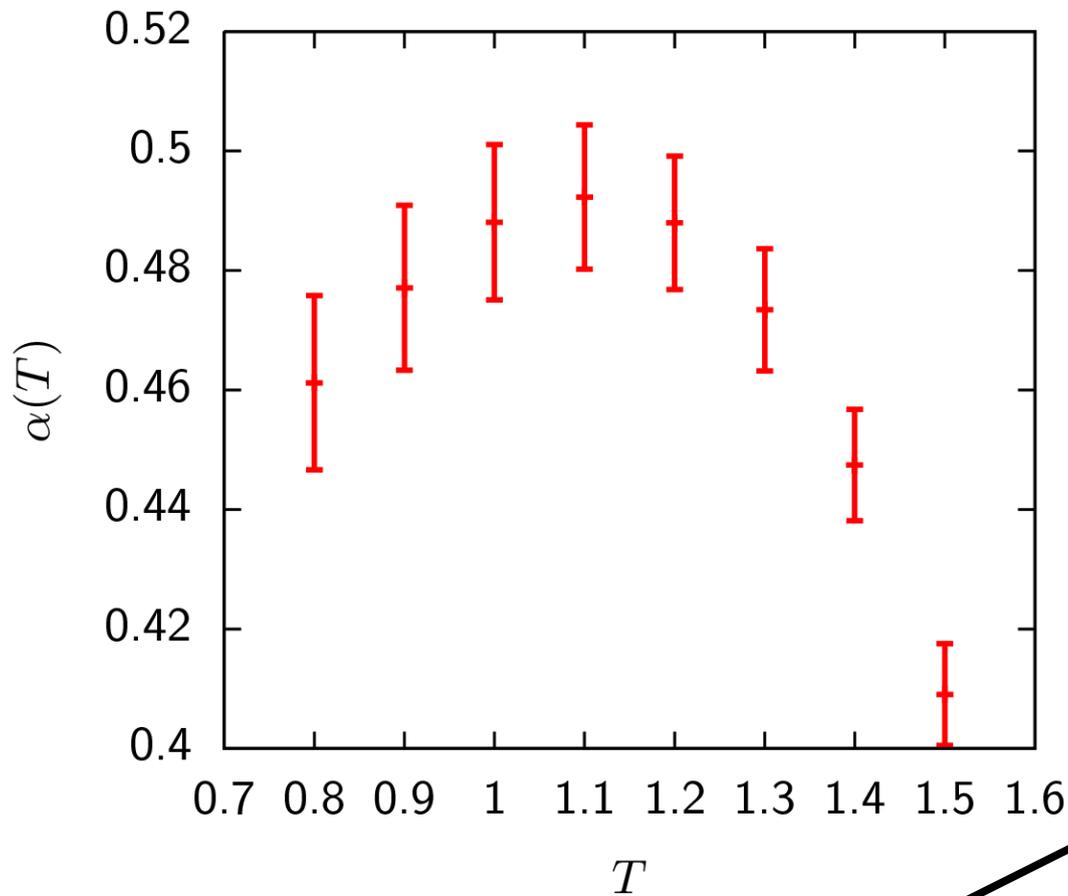
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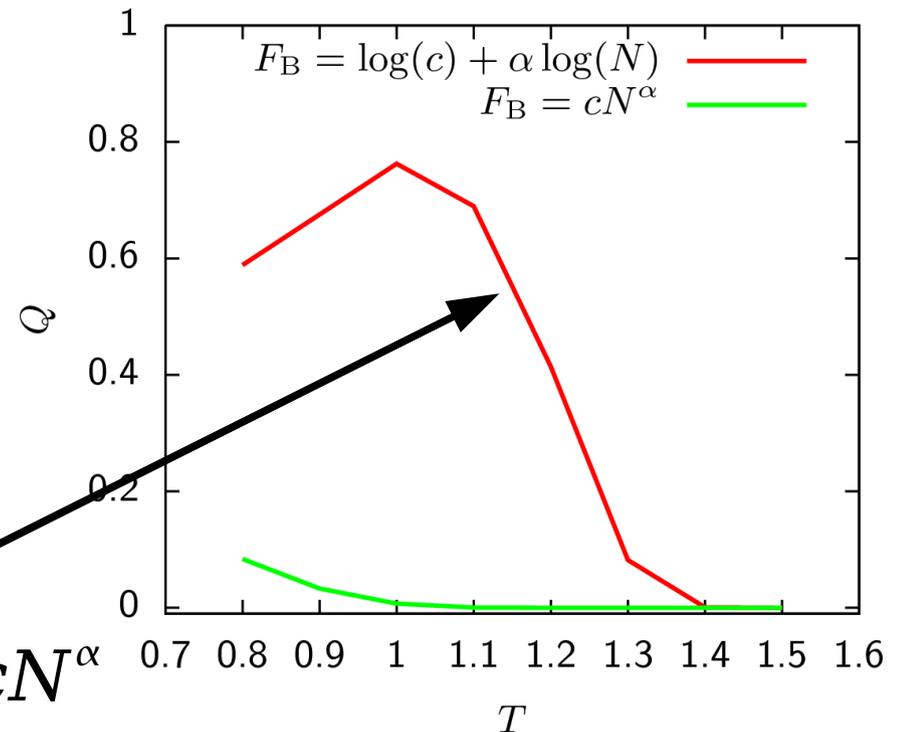
therefore, different fit: $\tau_B \propto cN^\alpha$

Results for the EA Model

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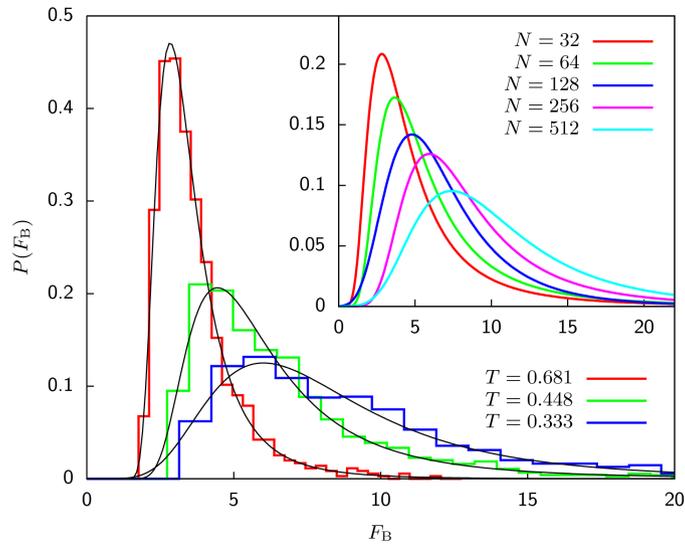


but goodness of fit ...

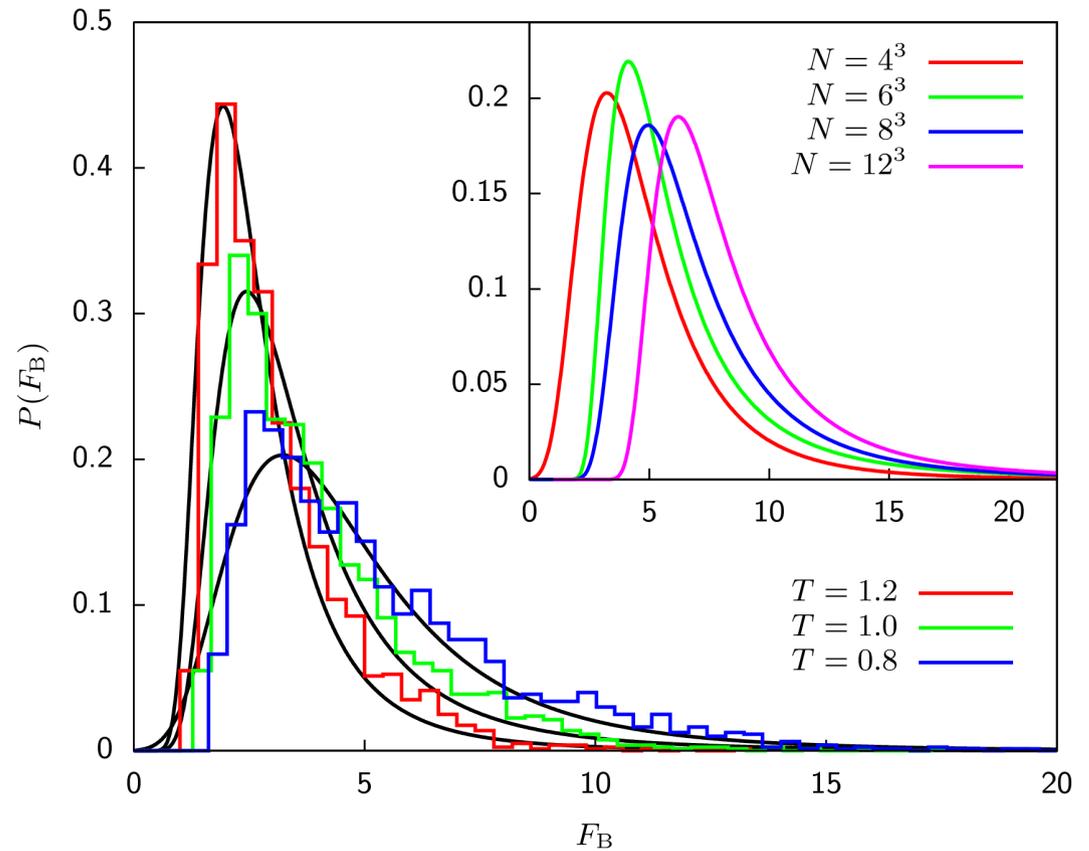


therefore, different fit: $\tau_B \propto cN^\alpha$

Results for the EA Model



SK Model



EA Model

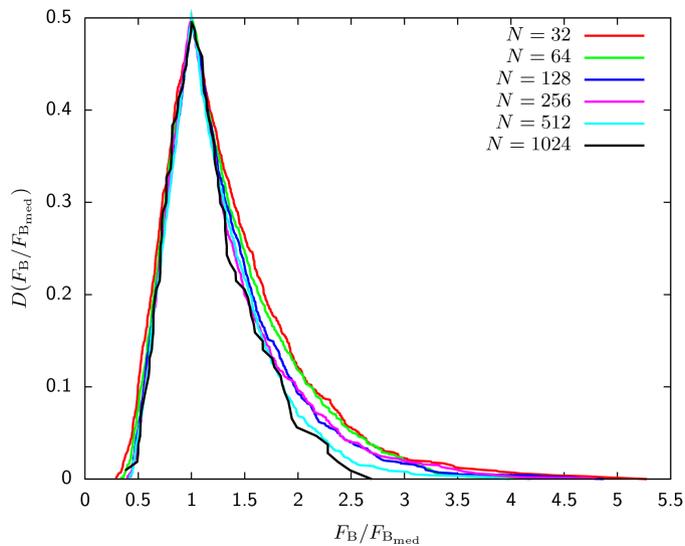
Fit integrated probability density:

$$F_{\xi; \mu; \sigma}(x) = \exp \left[- \left(1 + \xi \frac{x - \mu}{\sigma} \right)^{-1/\xi} \right]$$

$T < T_c \longrightarrow \xi > 0$
 \longrightarrow fat tailed (algebraic)
 Fréchet distribution

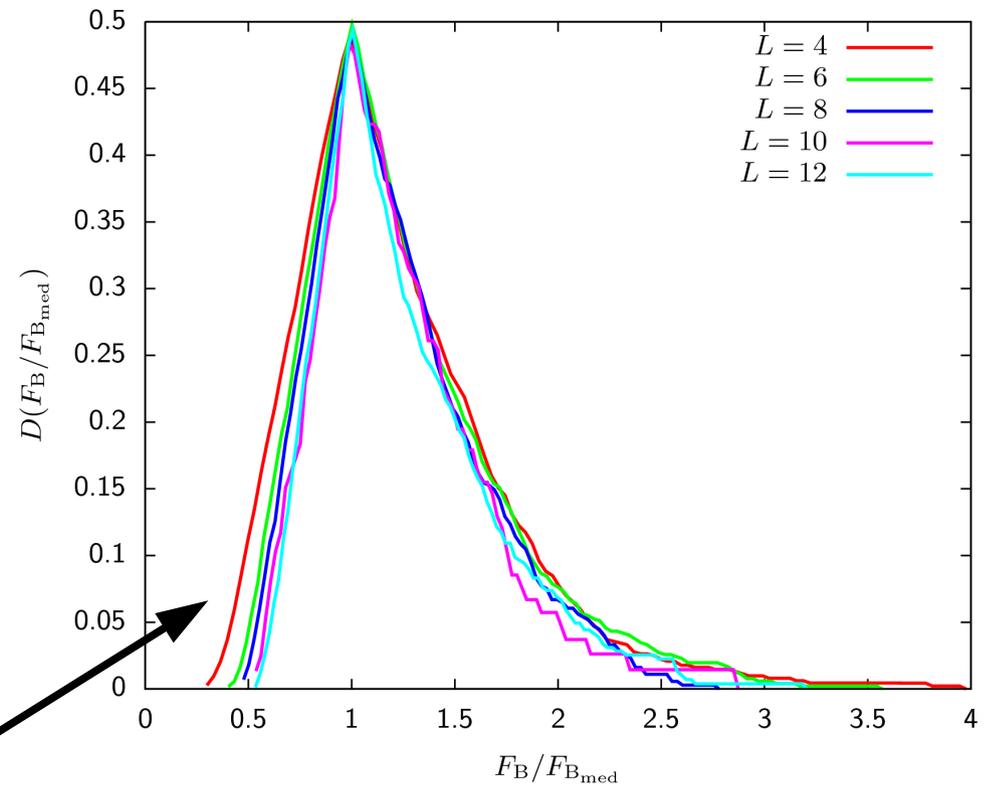
Results for the EA Model

Peaked probability distribution $D(F_B/F_{B_{\text{med}}})$



SK Model

no scaling,
i.e. non-self-averaging



EA Model

Conclusion

- **Algorithmic**
 - **PT** is good to decrease the autocorrelation time
 - **MuQ** gives the full $P(q)$ information
 - the combination of **PT+MuQ** makes it possible to get $P(q)$ down to $T \approx 0.5 T_c$
- **Physical**
 - the free energy barriers of the **SK** and **EA** model are
 - a) non-self-averaging
 - b) follow the Fréchet extreme-value distribution
 - the free energy barriers of the **SK** model diverge with an exponent of $\alpha = 1/3$
 - the last is not true for the **EA** model

Thank you!