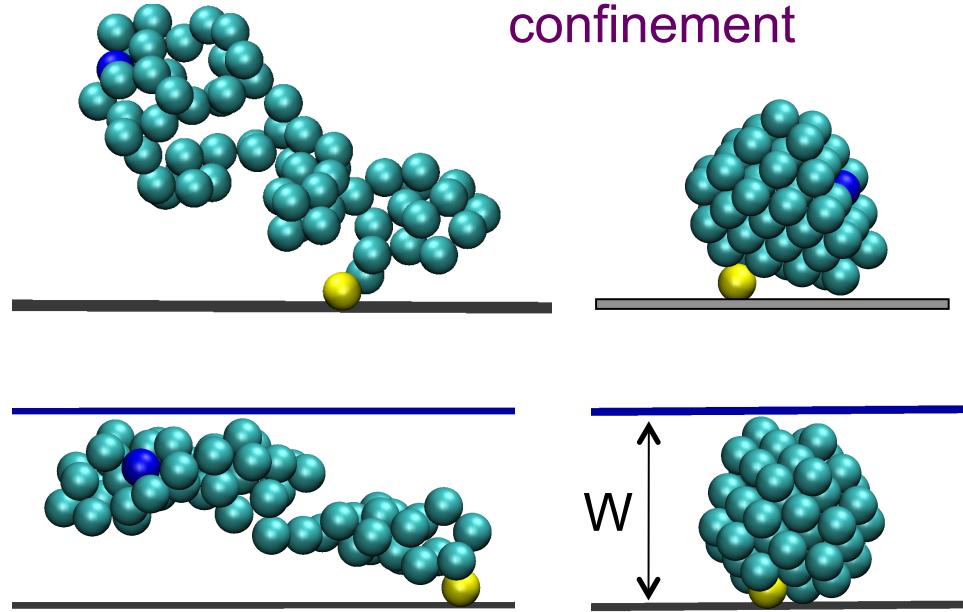


Entropy Reduction & Entropy Driven Folding for Confined Polymers

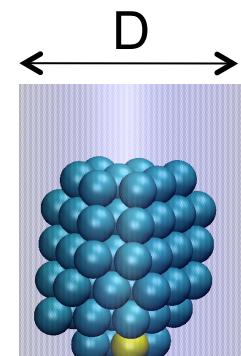


Mark P. Taylor
Dept. of Physics
Hiram College, Hiram OH

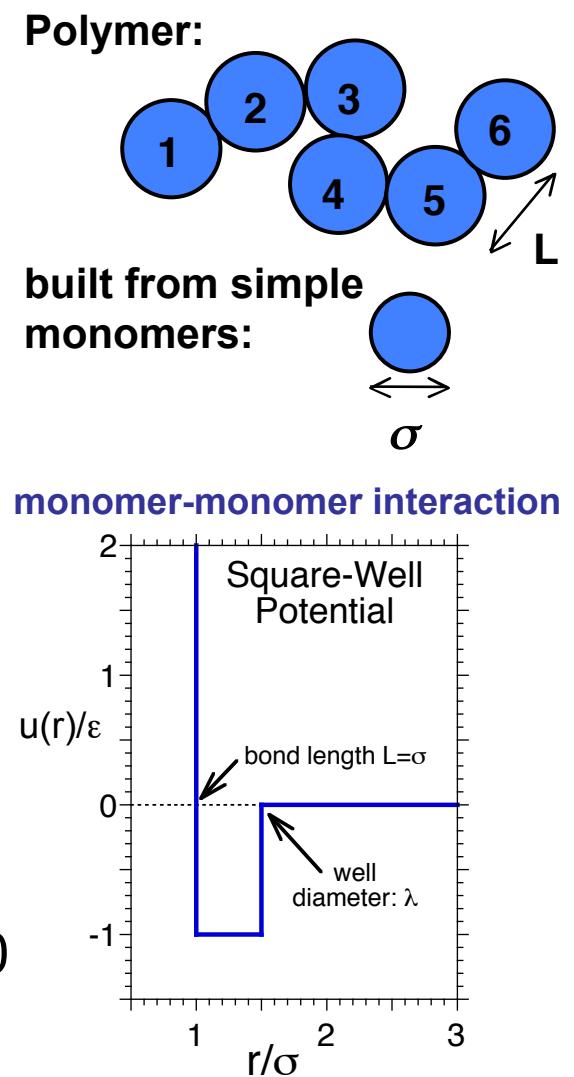
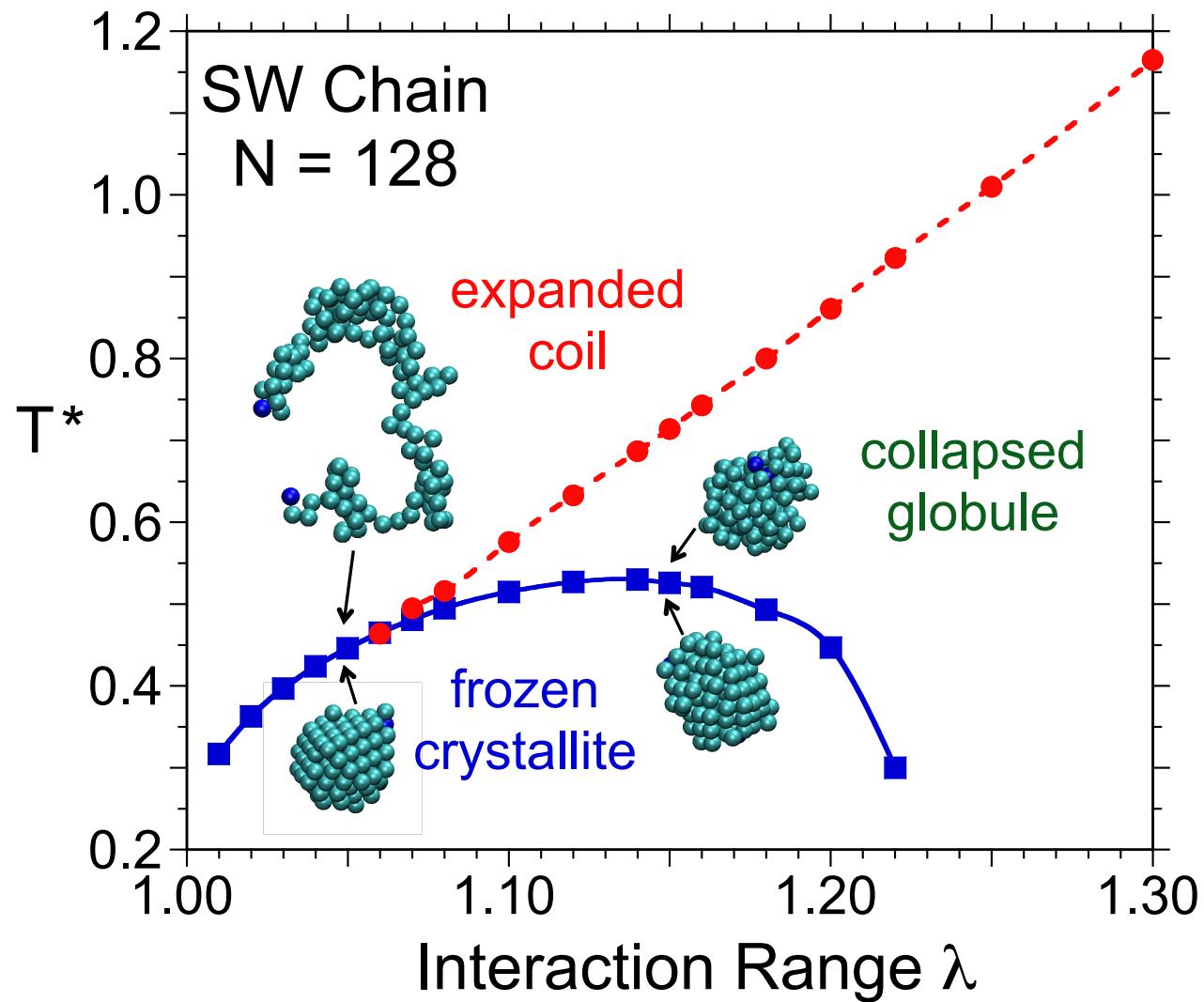
geometric (hard wall)
confinement



$$S = k_B \ln \Omega$$

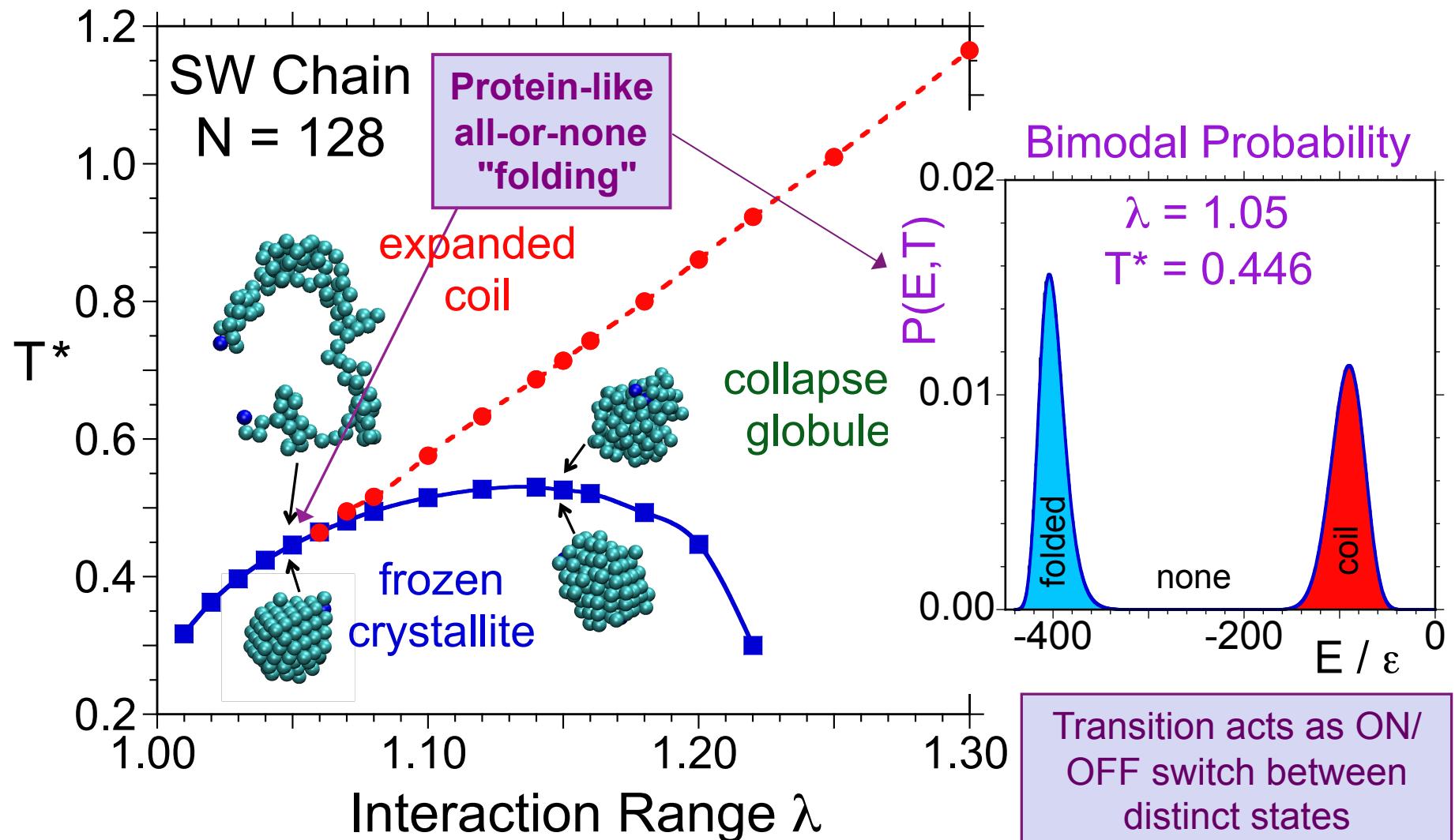


Phase Diagram for a Square-Well-Sphere Polymer



Taylor, Paul, & Binder, J. Chem. Phys. **131**, 114907 (2009);
J. Chem. Phys. **145**, 174903 (2016).

Phase Diagram for a Square-Well-Sphere Polymer



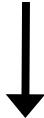
Taylor, Paul, & Binder, J. Chem. Phys. **131**, 114907 (2009);
J. Chem. Phys. **145**, 174903 (2016).

Density of States and Wang-Landau Sampling I

SW Chain Model has a discrete energy spectrum: $E_n = n\epsilon$

Density of States:

$g(E_n)$ = volume of configurational phase space for energy state E_n



Thermodynamics:

microcanonical entropy:
 $S(E) = k_B \ln g(E)$

canonical partition function:
 $Z(T) = \sum g(E) \exp(-E/k_B T)$

See also:

Yan & de Pablo, PRL **90**, 035701 (2003)

Schneider, Mueller, Janke,
Comp. Phys. Comm. **216**, 1 (2017)

iterate
m levels
we use
m=30

Wang-Landau algorithm* ... an iterative simulation method to compute $g(E_n)$:

Starting w/ $g(E_n)=1$, $H(E_n)=0 \forall n$, $f_0 = e$

Generate sequence of chain conformations using acceptance criteria:

$$P_{acc}(a \rightarrow b) = \min\left(1, \frac{g(E_a)}{g(E_b)}\right)$$

Update DOS: $g(E_n) \rightarrow f_m g(E_n)$

Update visitation histogram: $H(E_n) \rightarrow H(E_n)+1$

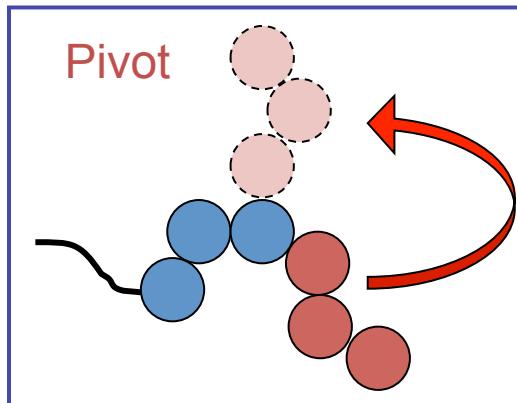
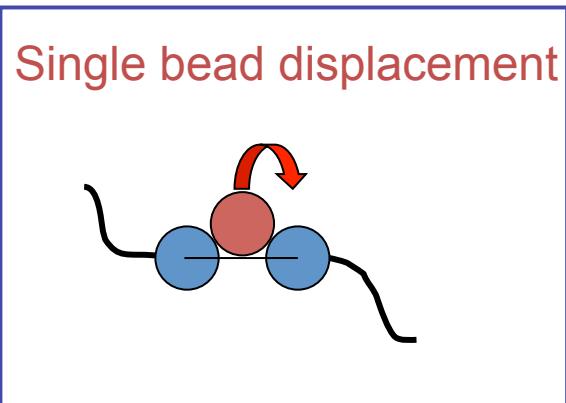
When histogram ~flat ...
reduce modification factor: $f_{m+1} = (f_m)^{1/2}$
reset histogram to zero: $H(E_n) = 0 \forall n$

*Wang & Landau, PRL **86**, 2050 (2001); PRE **64**, 056101 (2001).

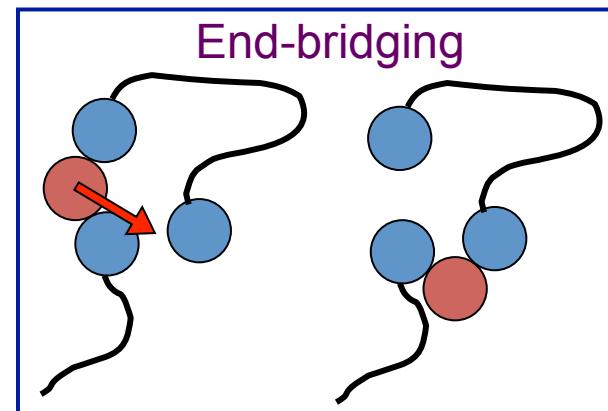
Wang-Landau Sampling II – MC Move Set

Success of the WL methods depends on underlying MC move set

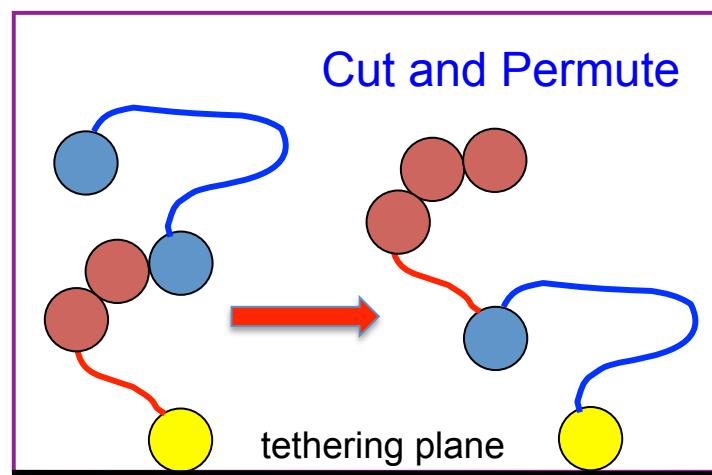
These moves sample most of configuration space:



To achieve highly compact conformations:



For good sampling near the tethering point:



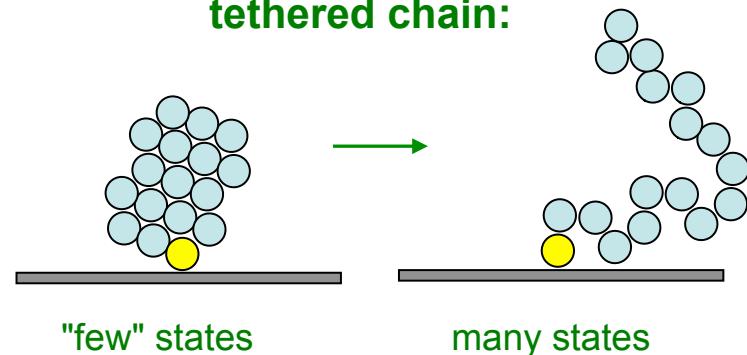
Escobedo & de Pablo, JCP **102**, 2636 (1995)

This move swaps upper/lower segments while preserving chain end-to-end vector.

Causo, J. Stat. Phys. **108**, 247 (2002)

Confined Chain: What do we expect?

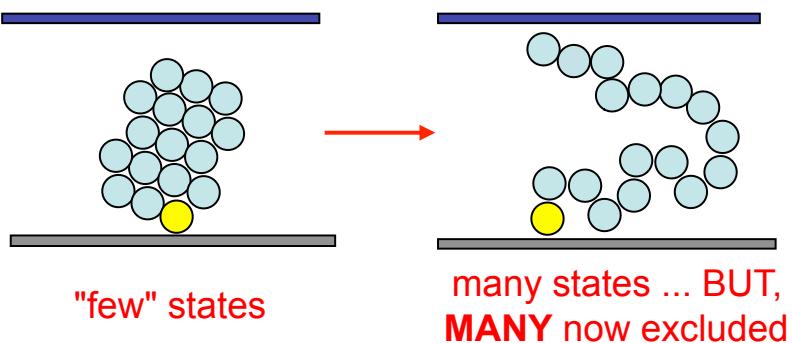
unfolding of a tethered chain:



"few" states

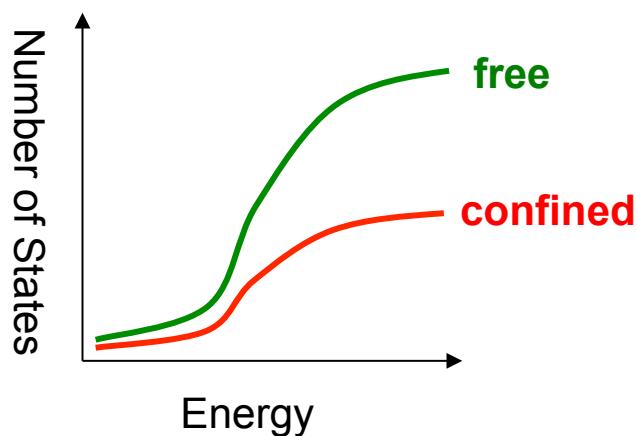
many states

unfolding of a confined chain:



"few" states

many states ... BUT,
MANY now excluded



entropy reduction: $\ln(g_{\text{confined}} / g_{\text{free}})$

Fewer states to "unfold into" shifts equilibrium towards folded state:

**Entropic Stabilization
of Folded State**

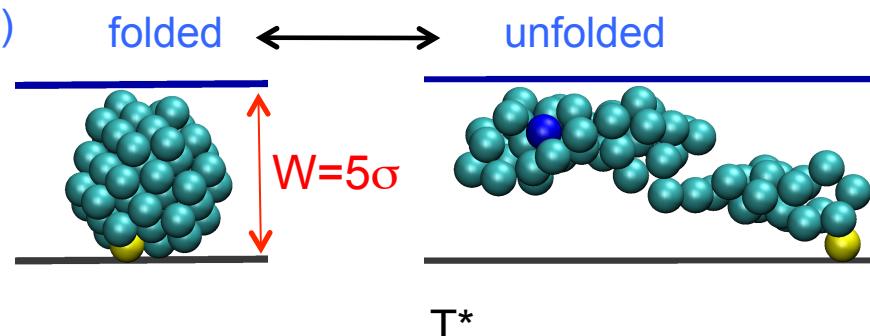
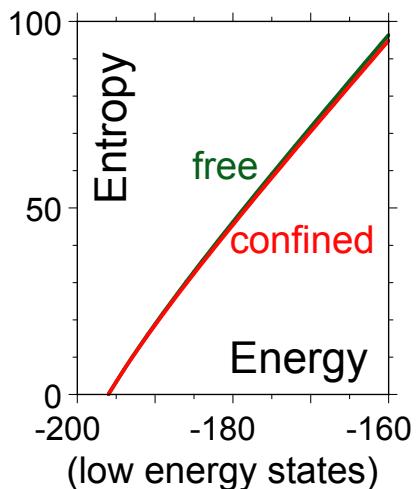
Minton, Biophys. J. **78**, 101 (2000)

Zhou and Dill, Biochem. **40**, 11289 (2001)

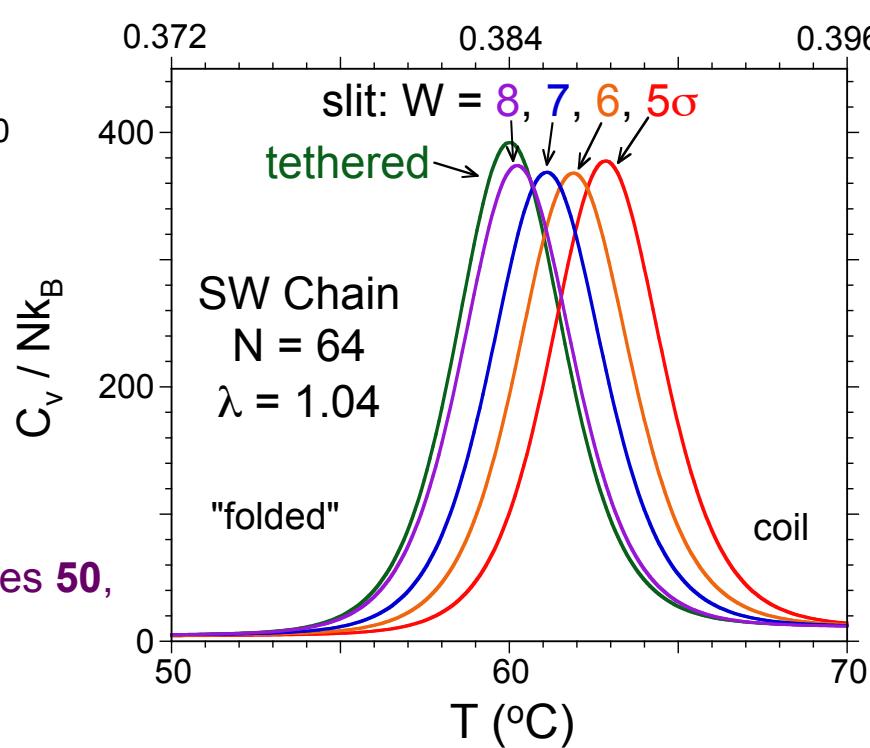
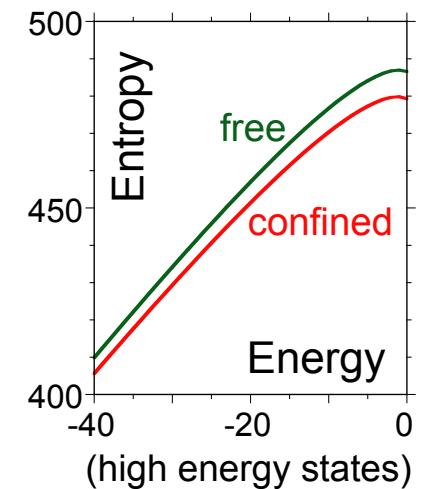
Confinement effects for Square-Well Chain

Unfolding of a tethered chain ($N=64$) in a slit:

Upper wall does not affect
folded ensemble ($W \geq 5\sigma$)

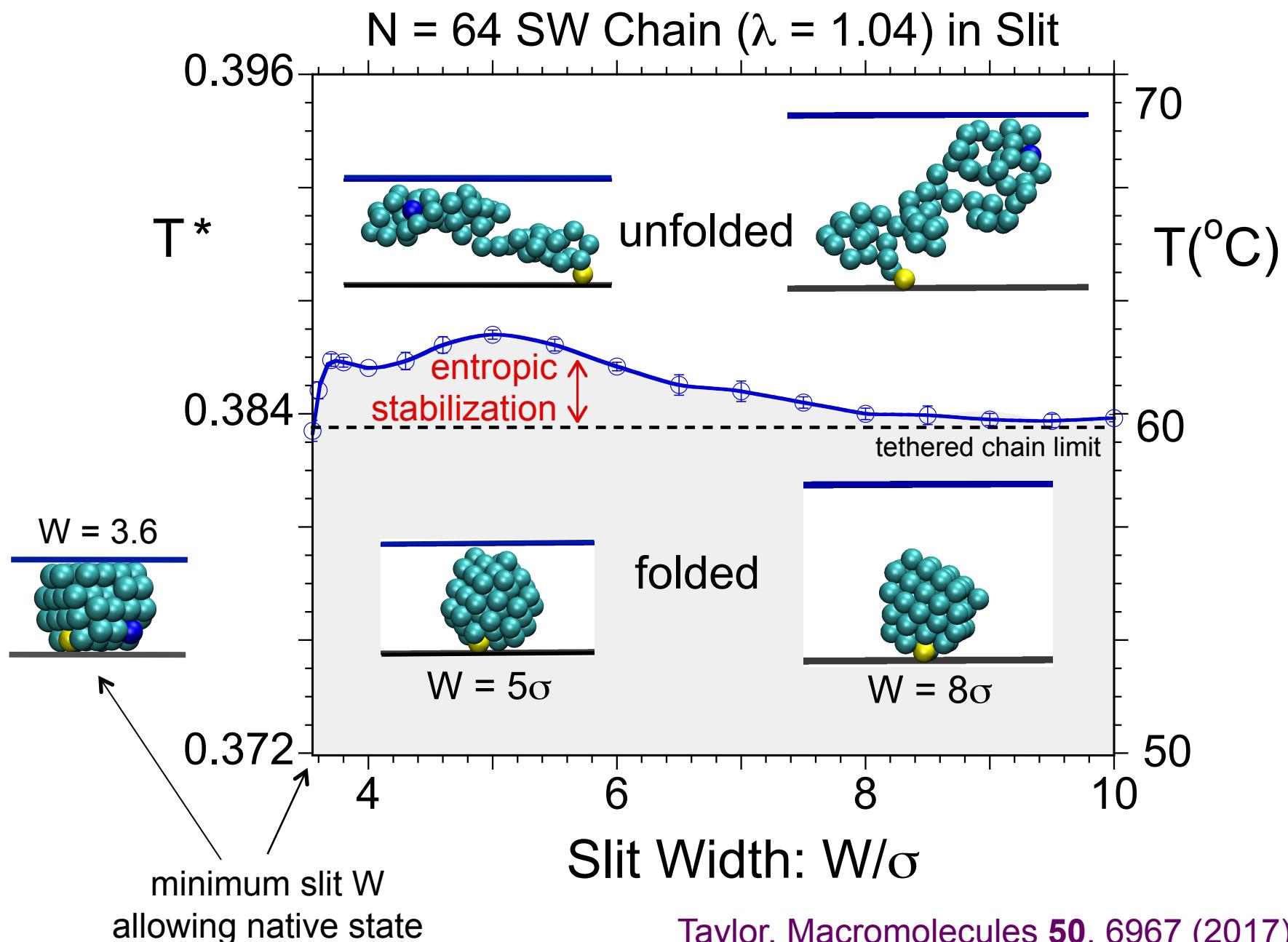


Entropy reduction for
unfolded ensemble



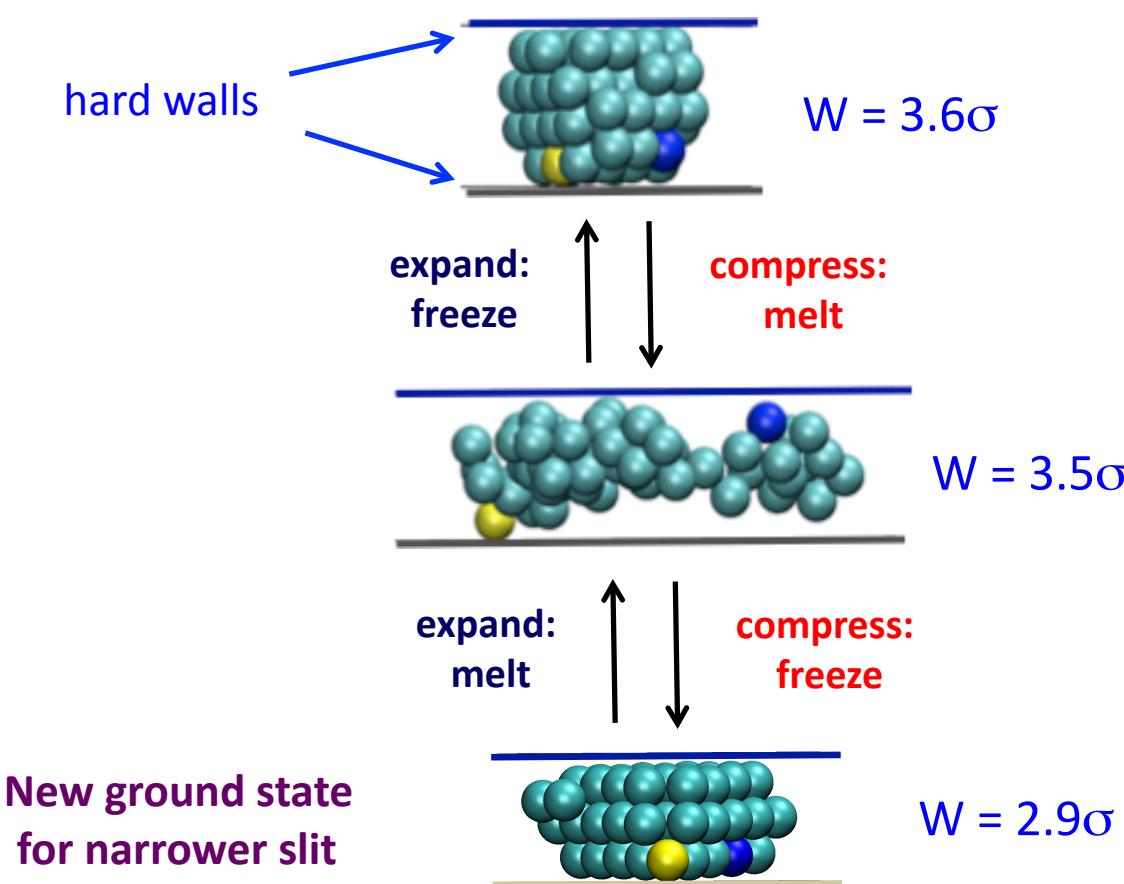
transition shifted
to higher T ...
expected result
for **entropic
stabilization**

Taylor, Macromolecules **50**,
6967 (2017)

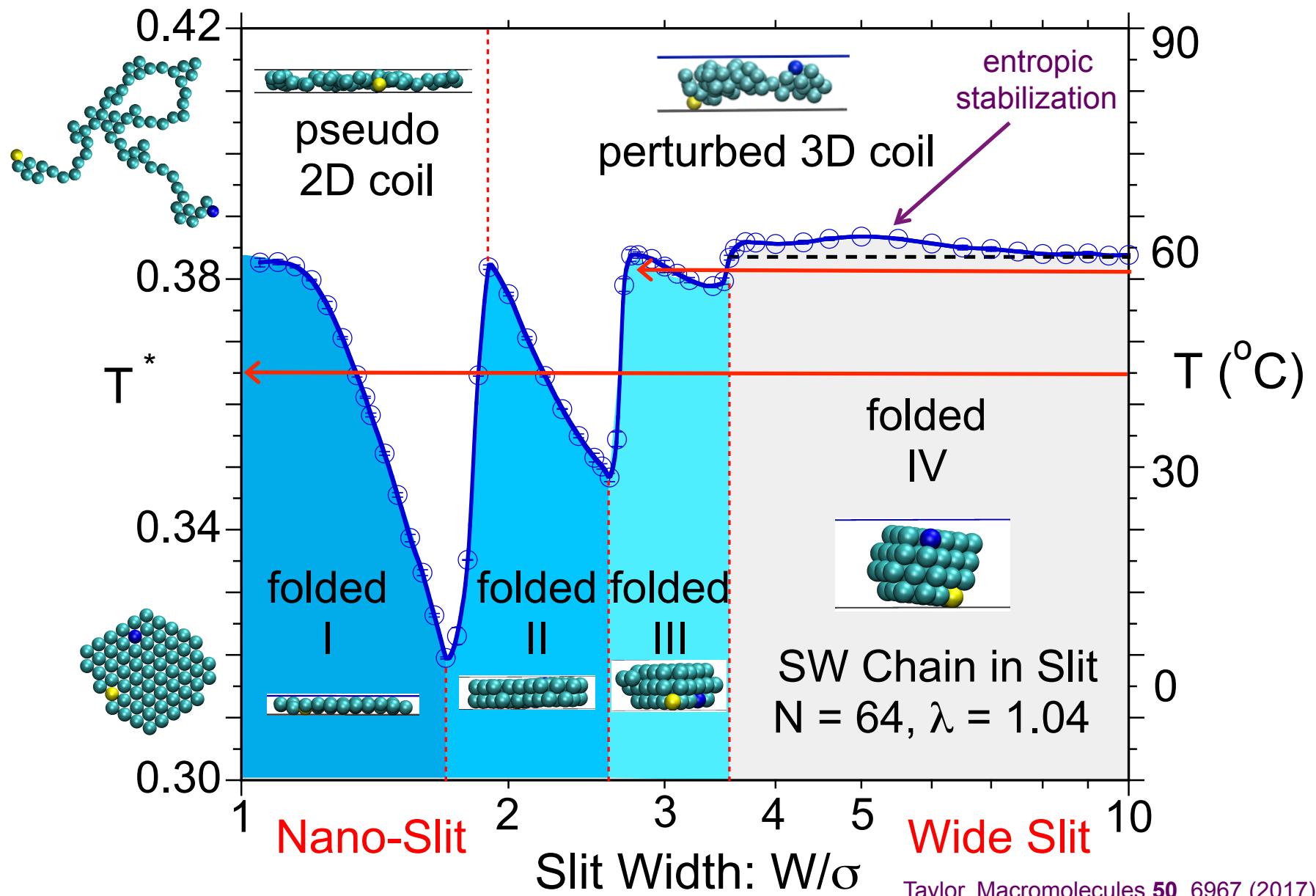


Confinement can drive conformational transitions

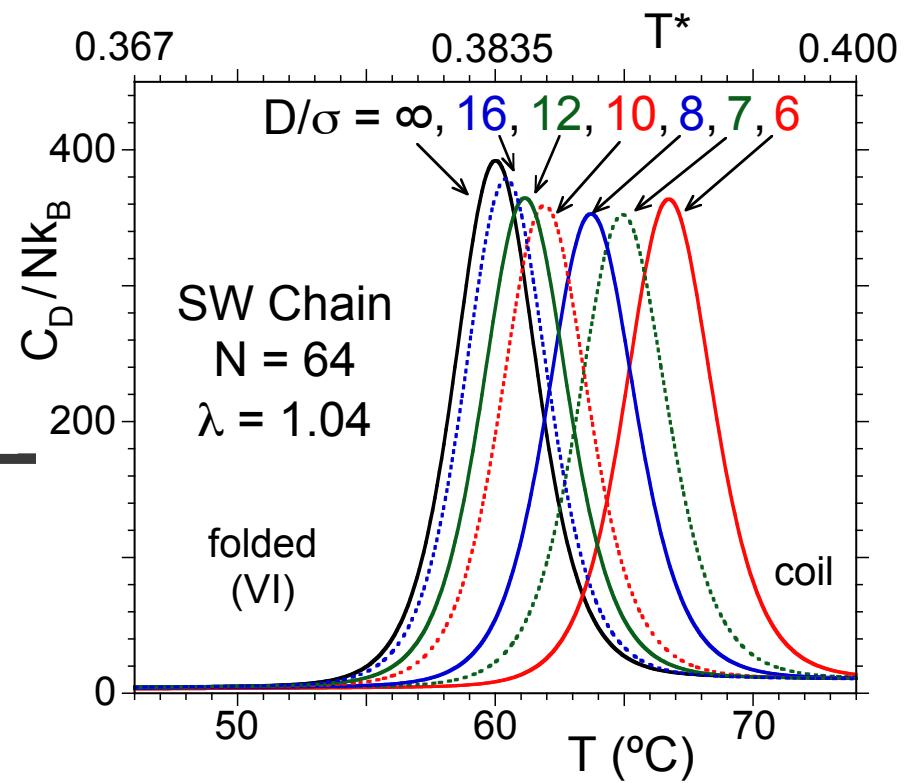
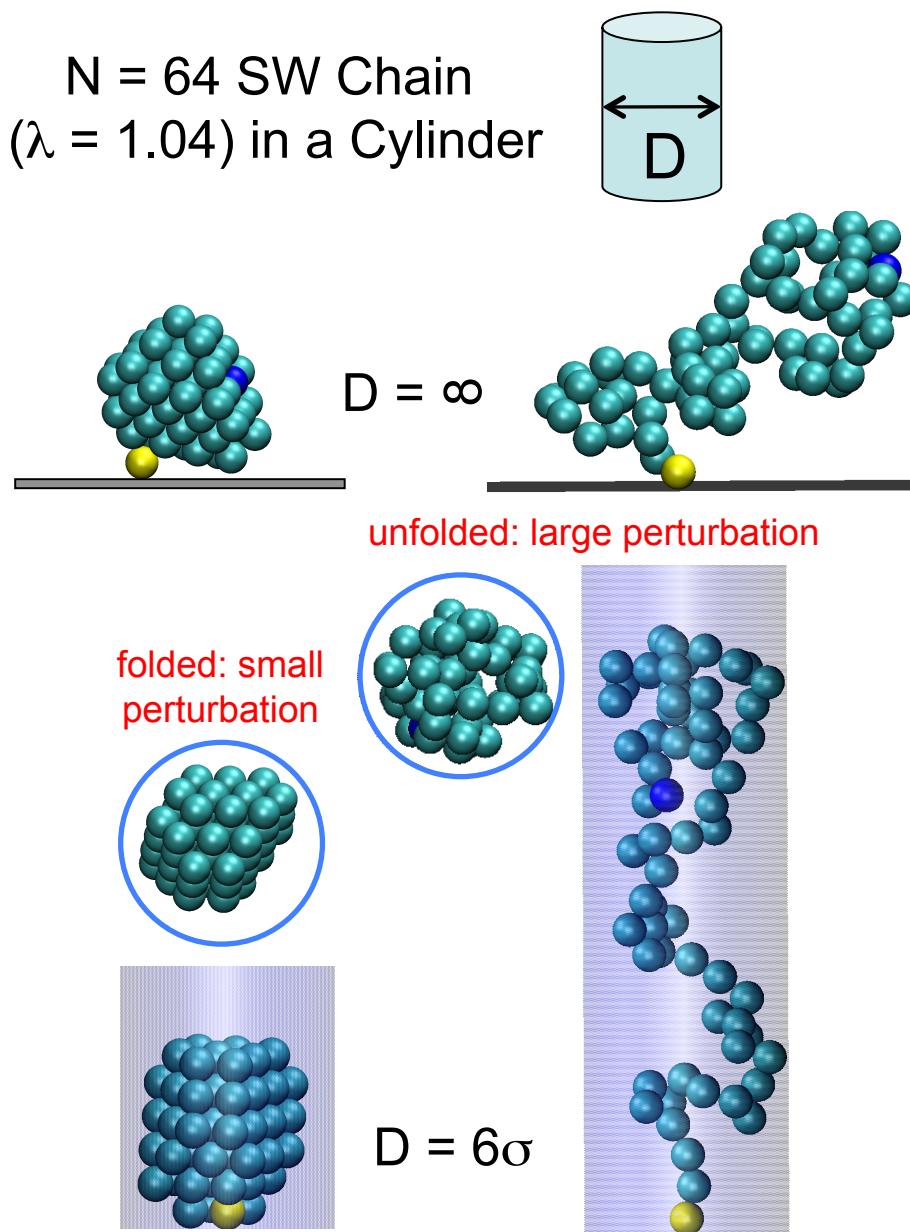
Polymer in a slit:
Isothermal melting/freezing ($T^* = 0.383$)



Chain in a Slit Phase Diagram



Cylindrical Confinement & Entropic Stabilization

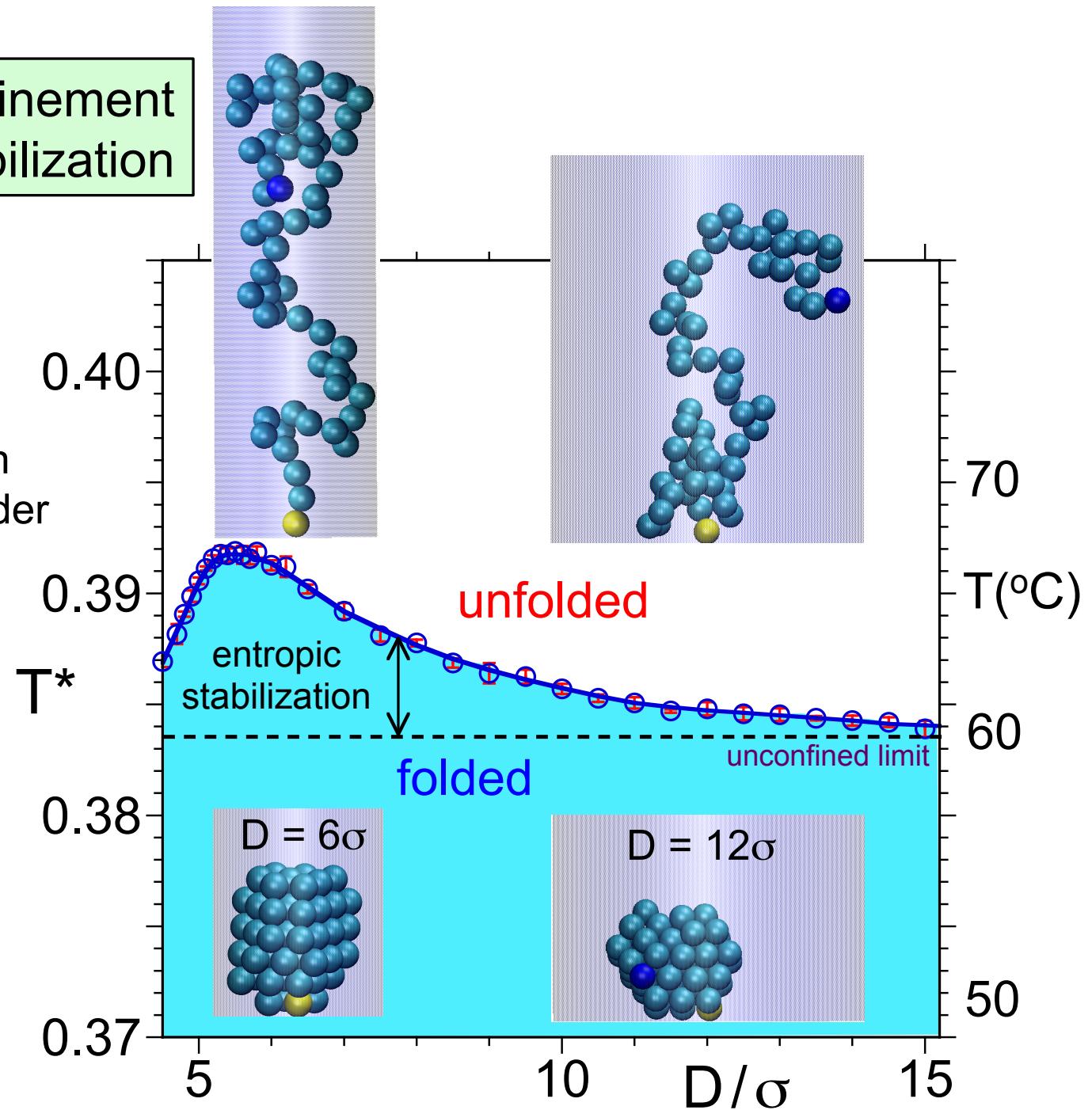
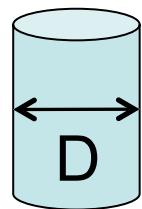


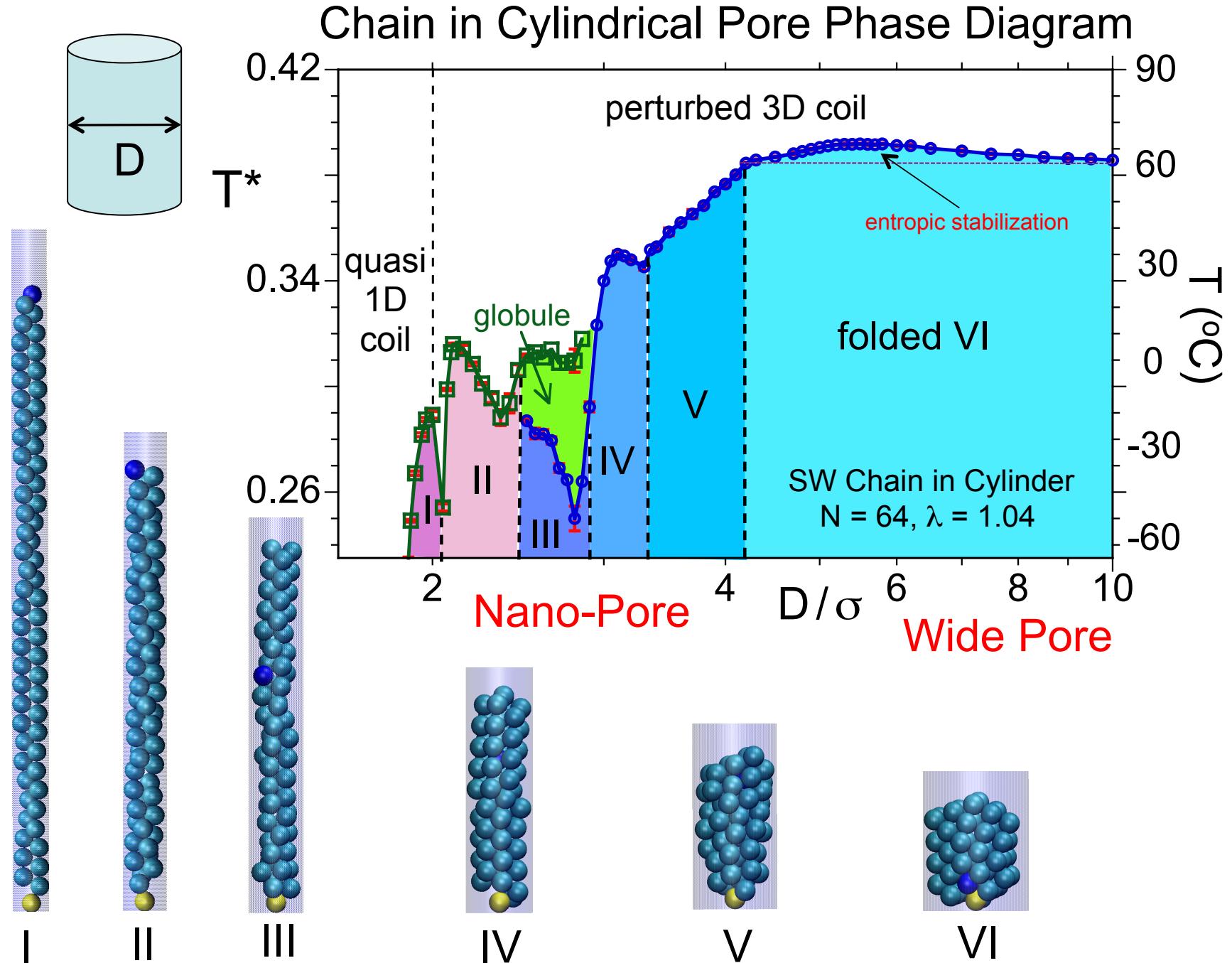
Christian O'Neil

Troy Prunty

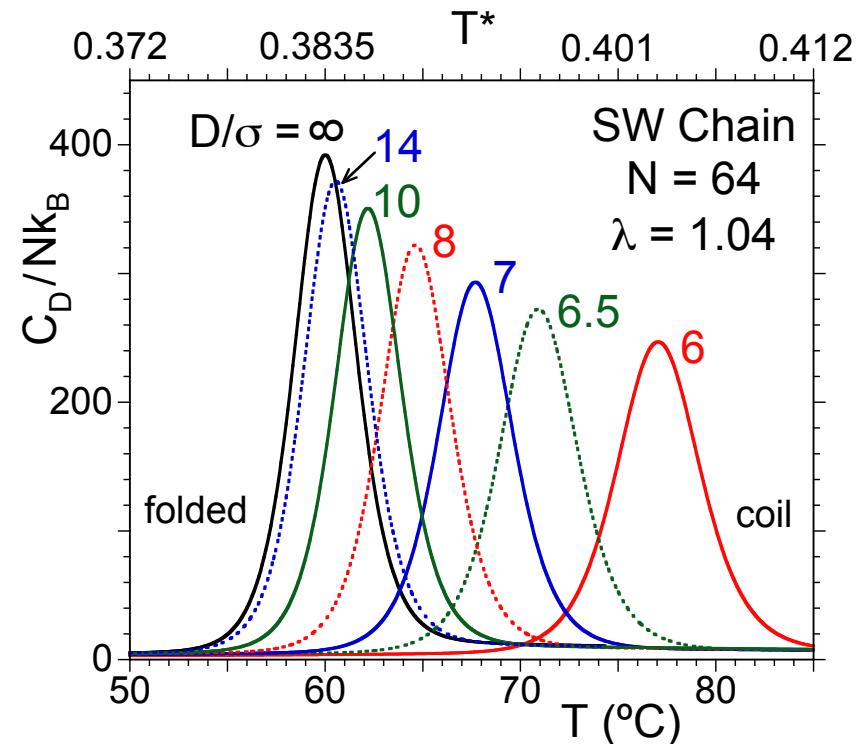
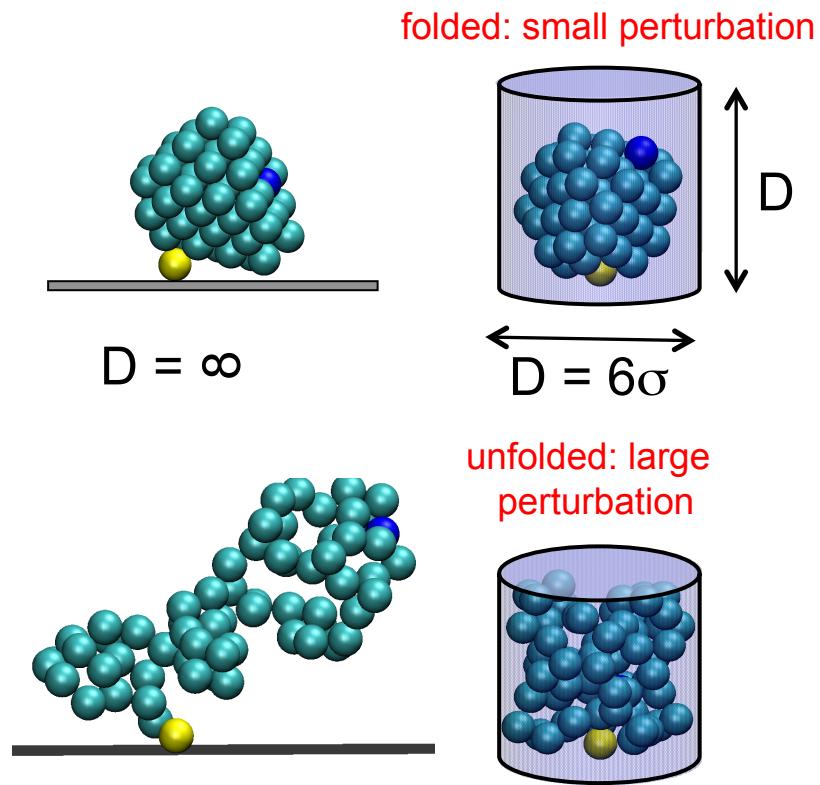
Cylindrical Confinement & Entropic Stabilization

$N = 64$ SW Chain
($\lambda = 1.04$) in Cylinder

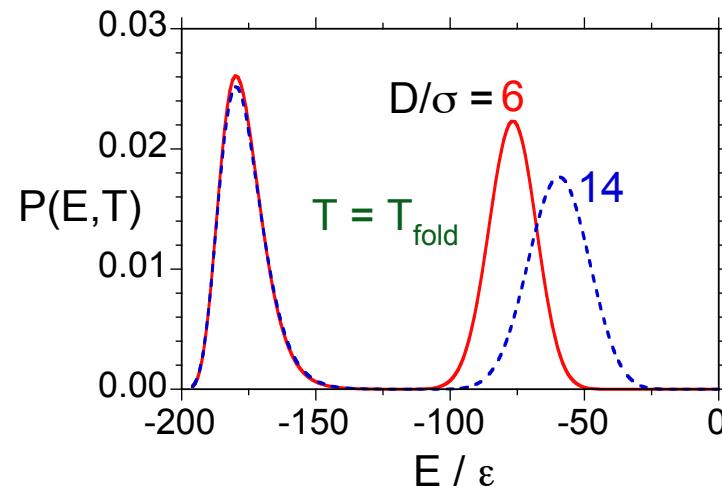




Chain in a Cylindrical Box (3D confinement)



Probability distribution at transition:
Bimodal, even in strong confinement
... 1st order like transition

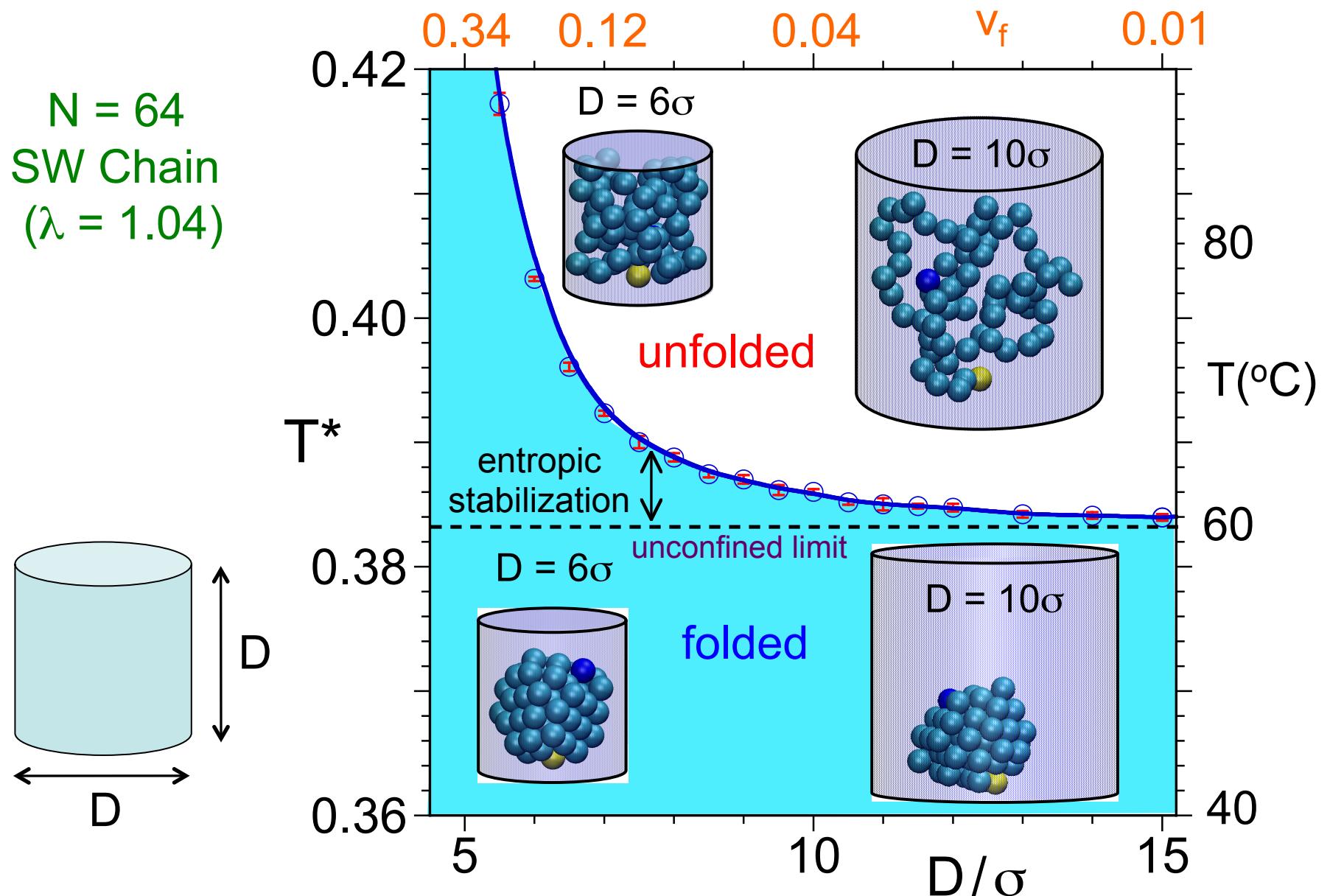


For chain collapse in Spherical Confinement see:

Marnez et al., Condens. Matt. Phys. (2012)

Marnez & Janke, Physics Procedia (2014)

3D Confinement: Folding in a Cylindrical Box ($D=H$)



Entropy reduction for a confined hard-sphere chain

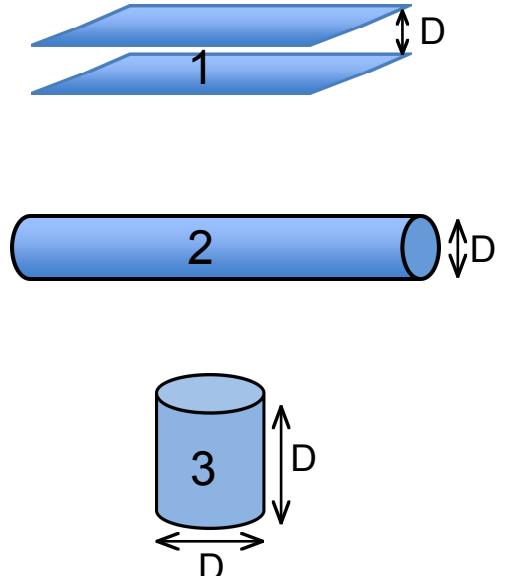
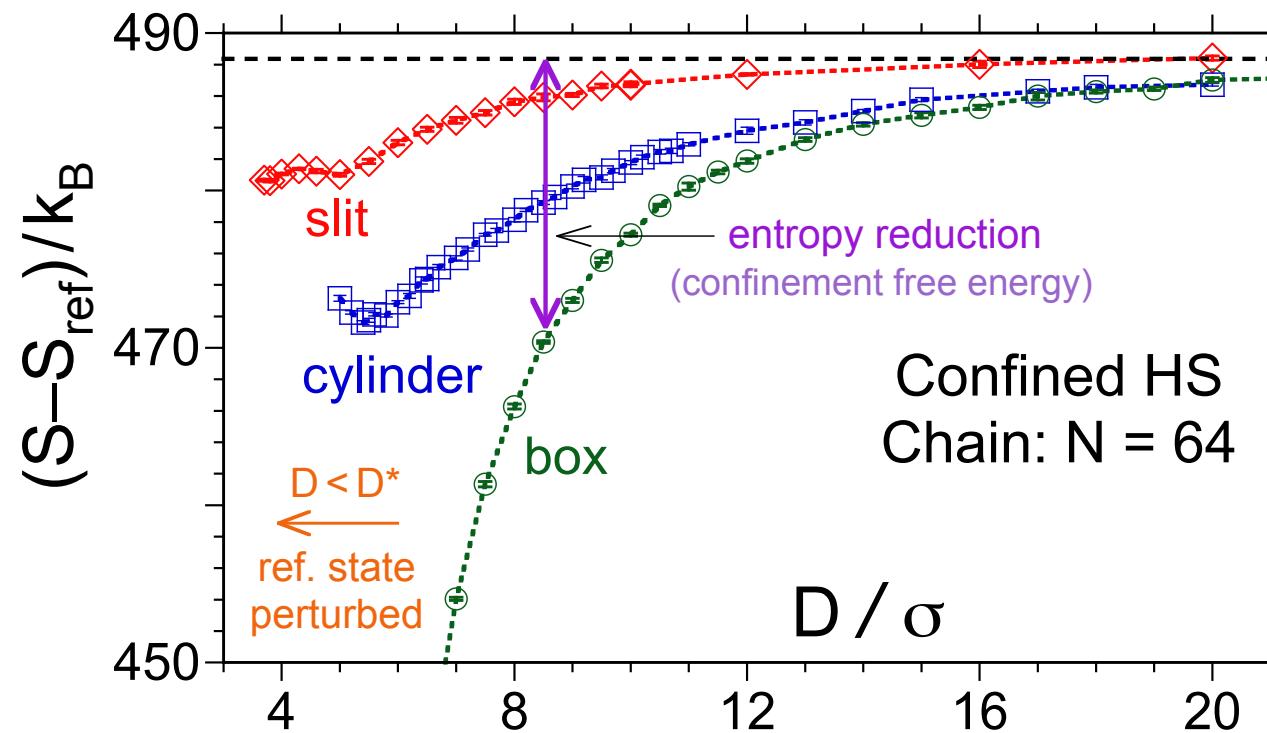
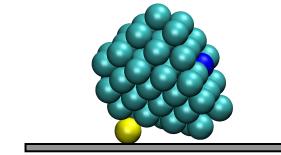
HS-chain free energy (entropy) is given by the SW-Chain density of states:

$$Z_{\text{SW}}(T) = \sum g(E) e^{-E/kT}$$

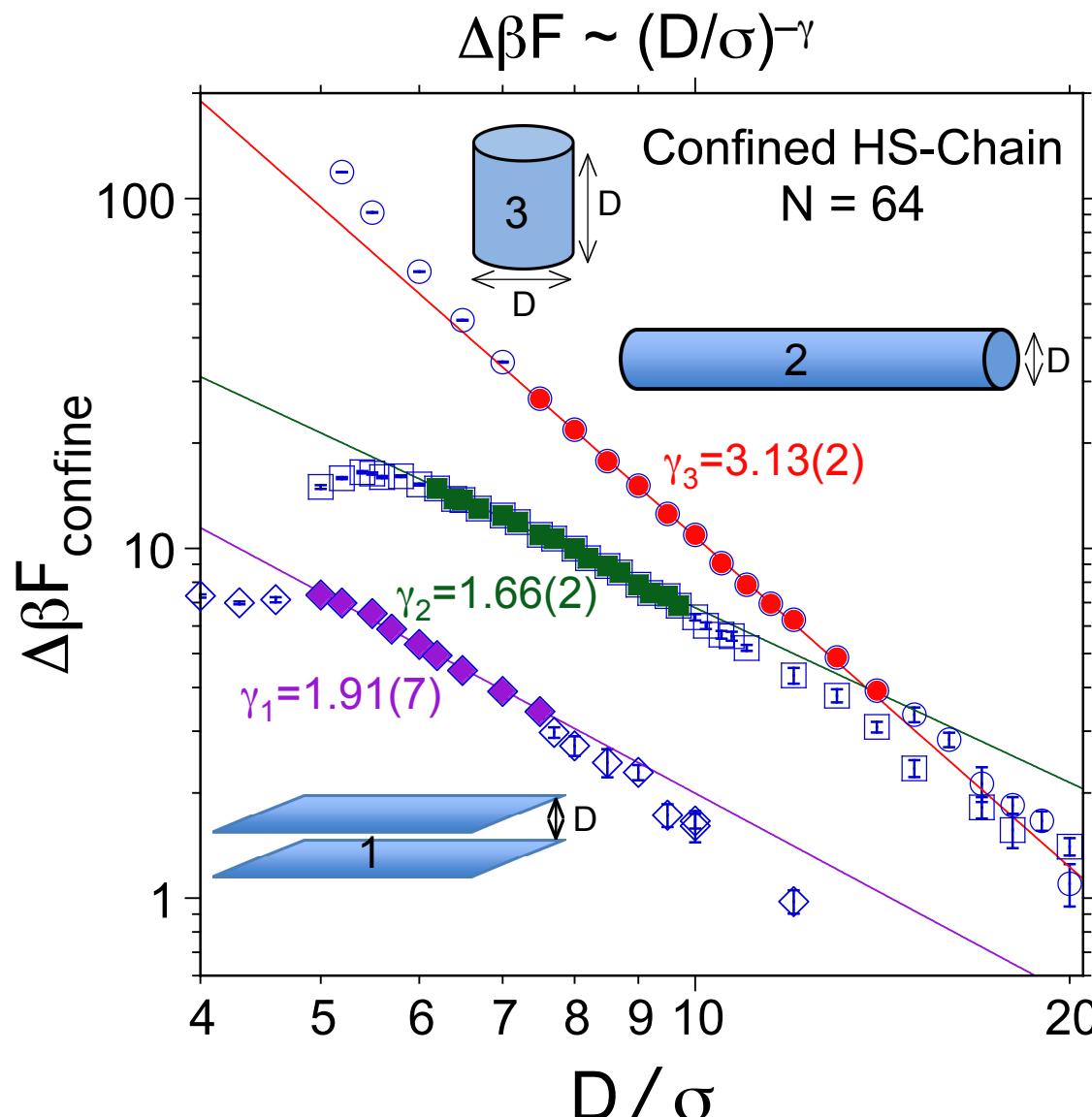
$$T \rightarrow \infty \dots Z_{\text{HS}} = \sum g(E)$$

$$\ln Z_{\text{HS}} = -\beta F_{\text{HS}} = S_{\text{HS}}/k_B$$

In WL we compute $g(E)/g(E_{\text{ref}})$
... reference state used here:



Confinement Free Energy: Scaling Properties*



For 3D confinement $\rho \sim N/D^3$
Density expansion gives:

$$\Delta\beta F \sim N \rho \sigma^3 \sim N^2 (D/\sigma)^{-\gamma}$$

$$\gamma \approx 3$$

Cylinder / Slit set blob size $\approx D$
Scaling analysis gives:

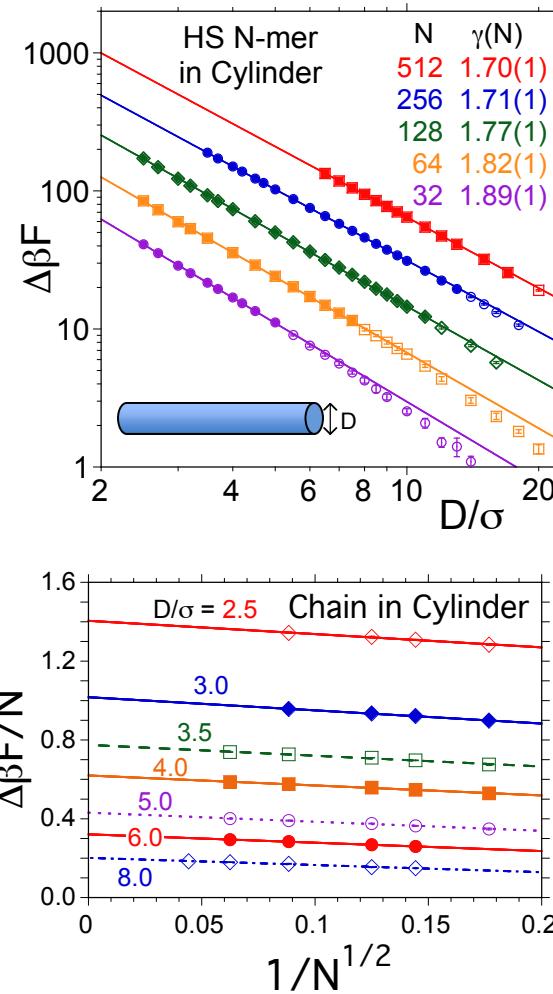
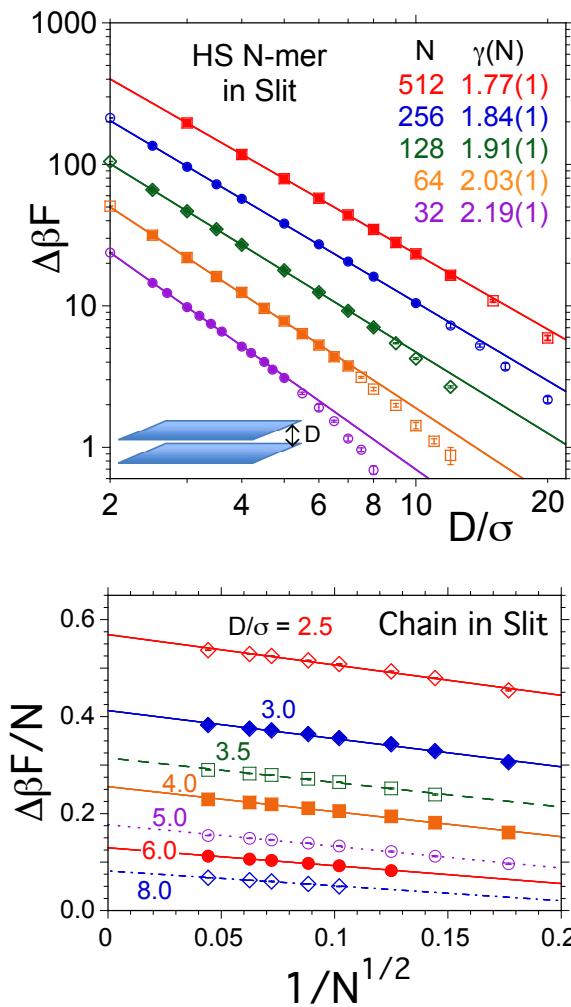
$$\Delta\beta F \sim N (D/\sigma)^{-\gamma}$$

$$\gamma = 1/\nu \approx 1.7$$

*See, for example:
de Gennes, Scaling Concepts in Polymer Physics
Grosberg & Khokhlov, Stat. Phys. of Macromolecules

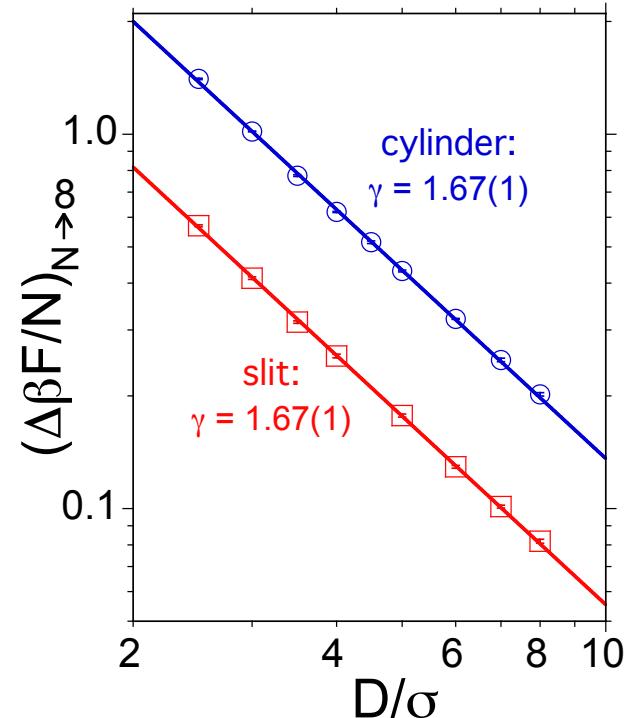
Scaling Properties & Finite Size Effects

Corrections to scaling: $\Delta\beta F/N = A (D/\sigma)^{-\gamma} (1 - b/N^{1/2})$

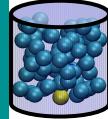
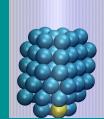
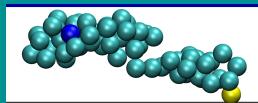


Confined Hard-Sphere Chain
Asymptotic ($N \rightarrow \infty$) scaling

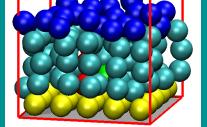
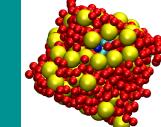
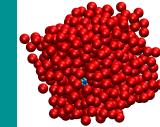
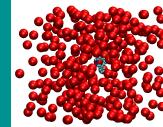
$$\Delta\beta F/N = A(D/\sigma)^{-\gamma}$$



See also: Cacciuto & Luijten,
Nano Letters **6**, 901 (2006)



Summary



Confinement Effects on Polymer Folding/Unfolding

- Hard wall confinement provides entropic stabilization of folded state
- Folding transitions maintain "all-or-none" character in confinement
 - Increasing confinement can lead to isothermal freezing/melting
- Finite size effects important for confinement free energy of HS-chains

Christian O'Neil



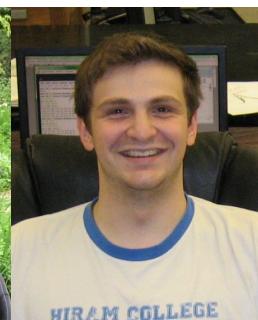
Troy Prunty



Beth Breen



Christopher Vinci



Norah Ali



Jaden Slovensky



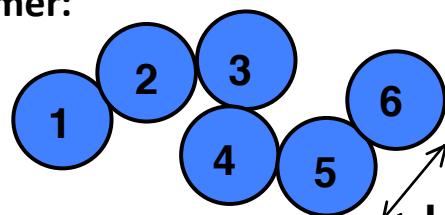
Funding: NSF (DMR-1607143)



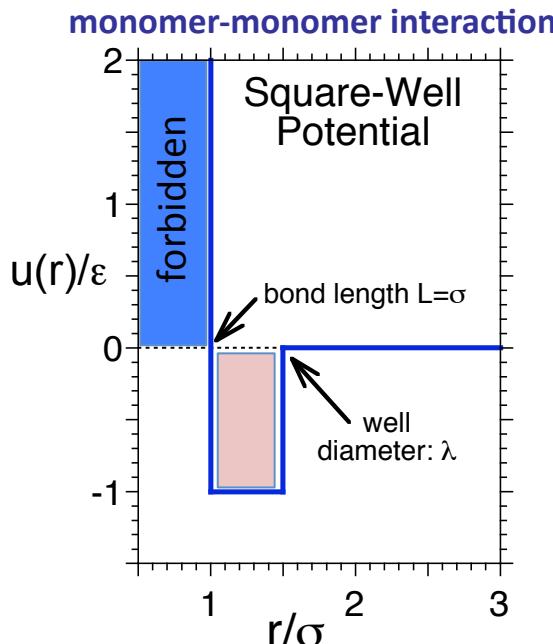
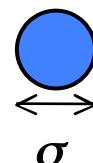
Extra Slides

Square-Well Chain Model

Polymer:



built from simple monomers:



Model Parameters:

ε = well depth (sets energy scale)

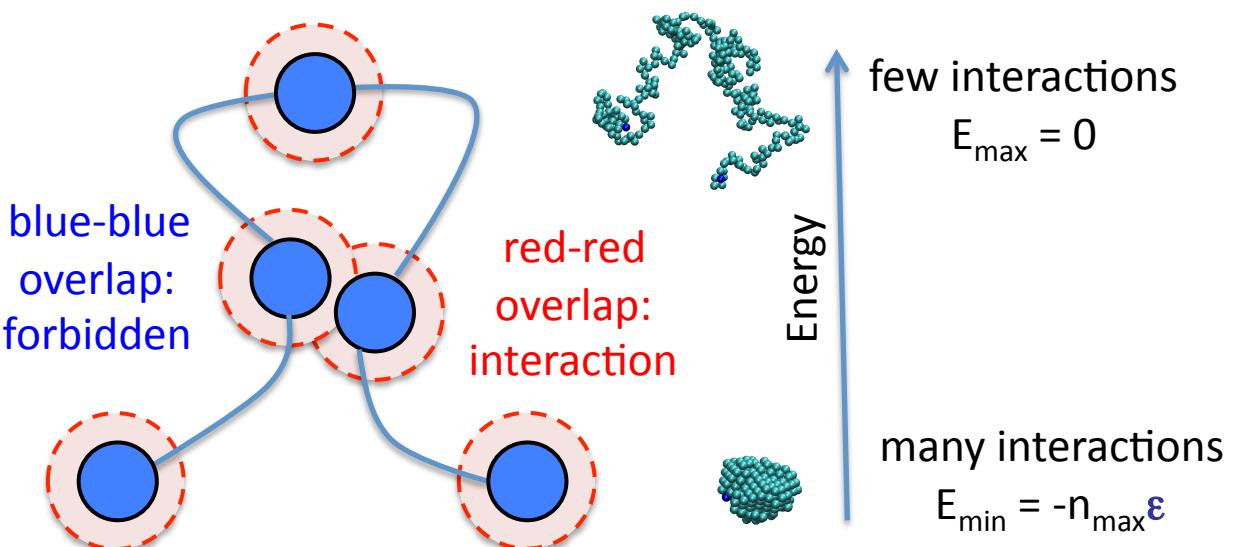
σ = hard-sphere diameter

L = bond length ($L = \sigma$)

$\lambda = \text{interaction range}/\sigma$

$T^* = k_B T/\varepsilon$ = reduced temperature

Model has a discrete energy spectrum: $E_n = -n\varepsilon$
(n = number of monomer-monomer interactions)



The **entropy** of the model is given by

S = k ln [g(E)] where g(E) is :

$$g(E) = \frac{\varepsilon}{V} \int \cdots \int \prod_{k=1}^{N-1} \frac{\delta(r_{k,k-1} - L)}{4\pi L^2} \delta\left(E - \sum_{i < j+1}^N u(r_{ij})\right) d\vec{r}_1 \dots d\vec{r}_N$$

This 3N-dimensional integral (3N=192 for N=64)
presents a daunting numerical challenge!

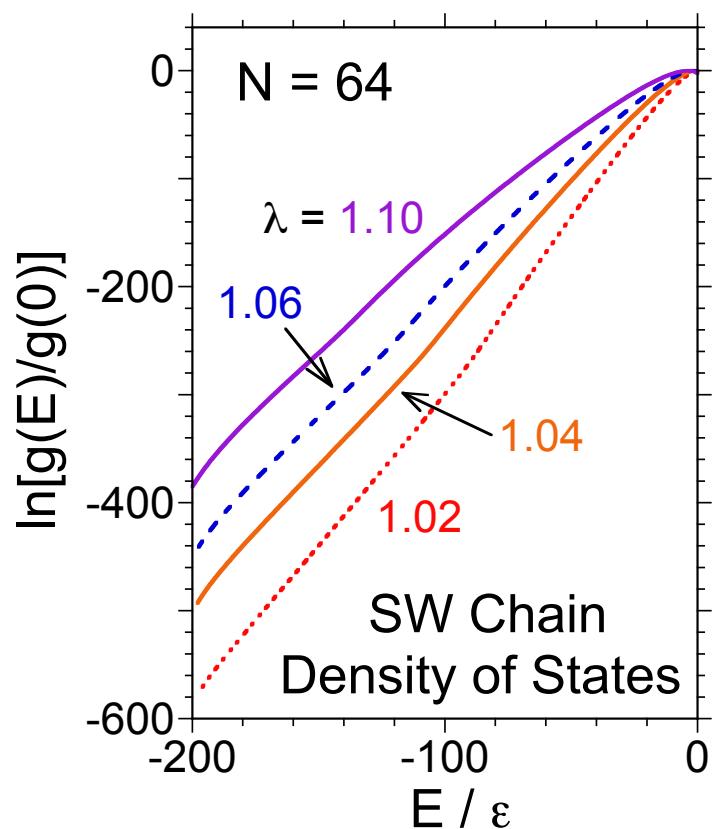
M.P. Taylor, J. Chem. Phys. **118**, 883 (2003)
J. Chem. Phys. **147**, 166101 (2017)

We use **computer simulation methods** (Wang-Landau
algorithm) to compute this integral.

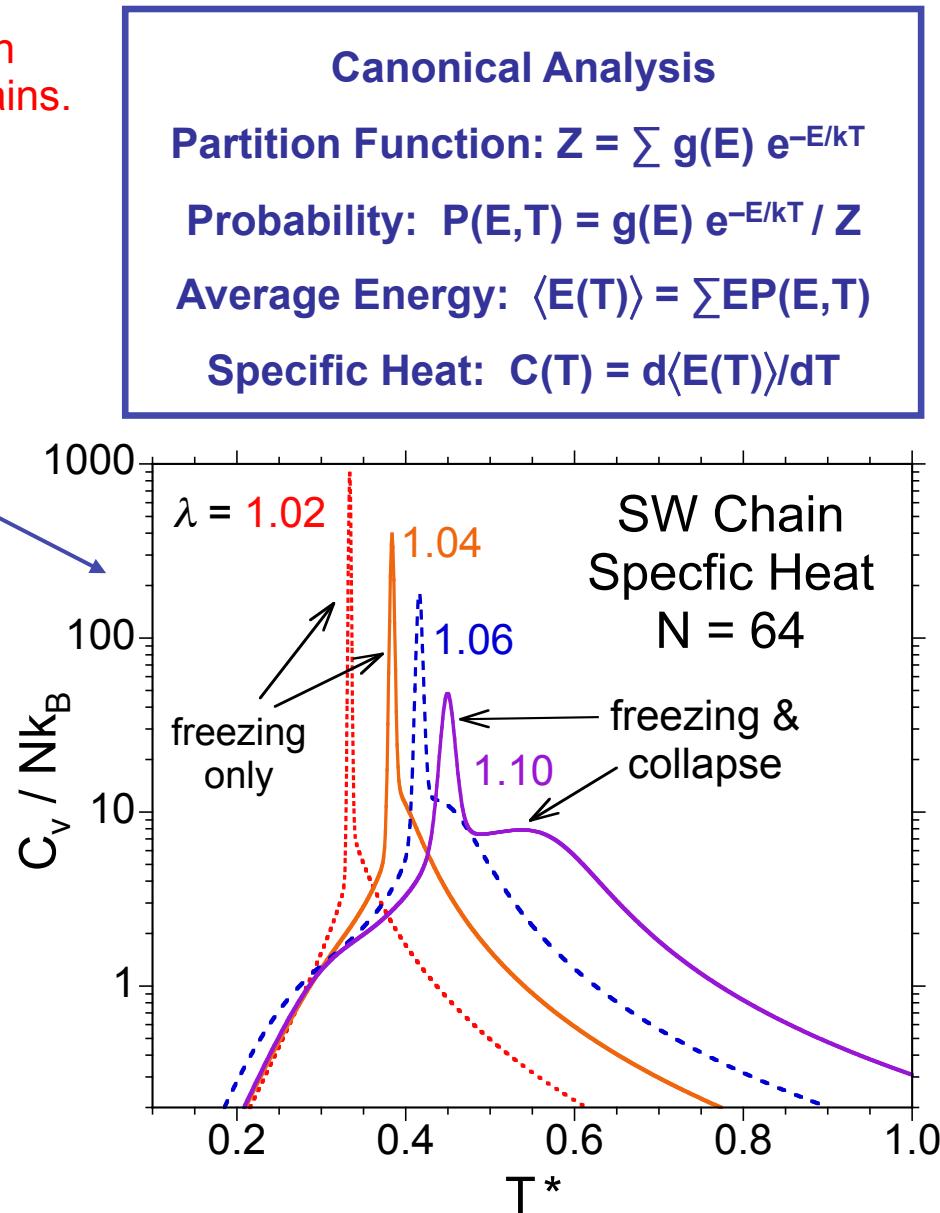
F. Wang and D.P. Landau, Phys. Rev. Lett. **86**, 2050 (2001)

Single Chain DOS and Canonical Analysis

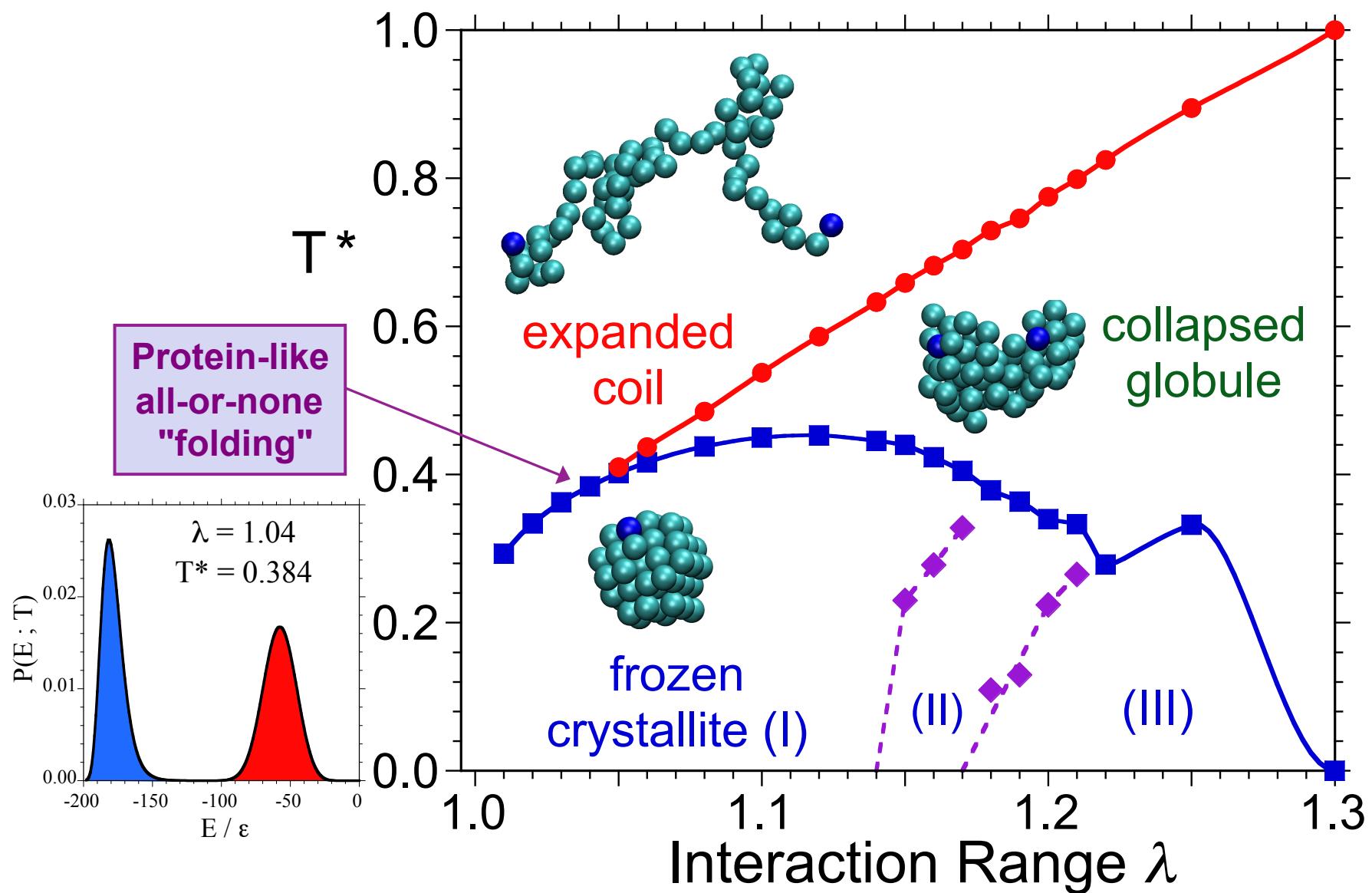
Algorithm validated via comparison with exact DOS results for short ($n \leq 6$) SW chains.



Taylor, Paul, & Binder,
Polymer Sci., Ser. C **55**, 23 (2013)



N=64 Square-Well Chain Phase Diagram



Taylor, Paul, & Binder, Polymer Sci., Ser. C 55, 23 (2013)

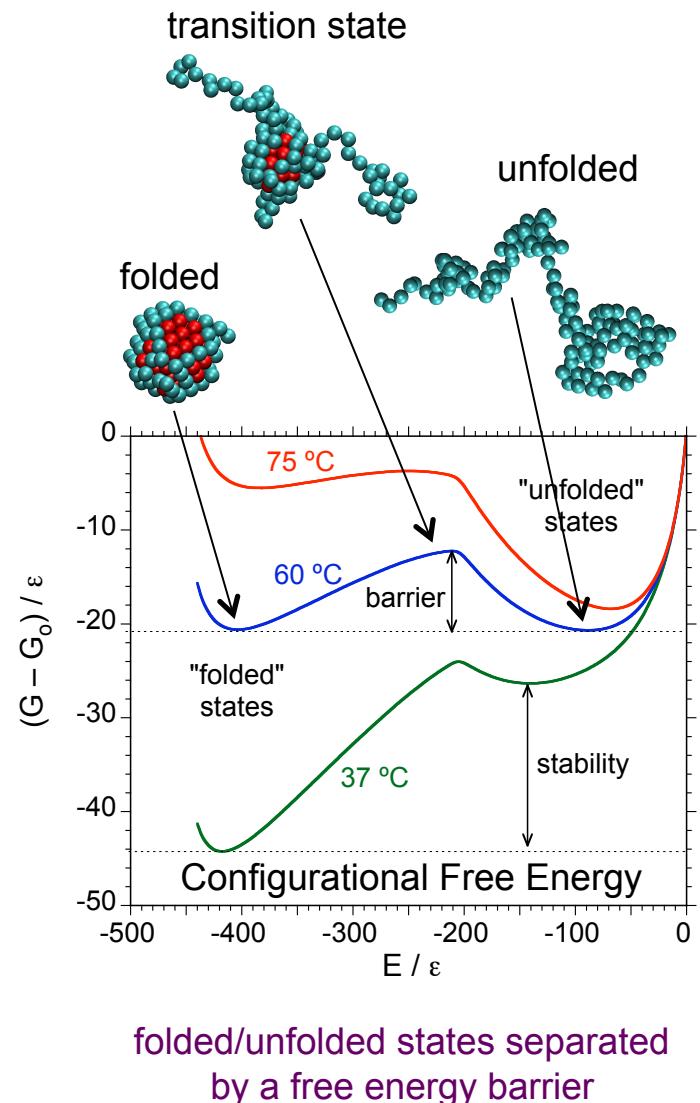
All-or-none polymer "folding" transition

Thermodynamics characteristic of
two-state protein folding ... simple
model for a complex biophysical
process.

Taylor, Paul, Binder, JCP **145**, 174903 (2016)

This type of transition useful for
responsive/smart materials
... provides an "on/off switch"
between two distinct states.

Confinement effects are of interest
for both of these applications.





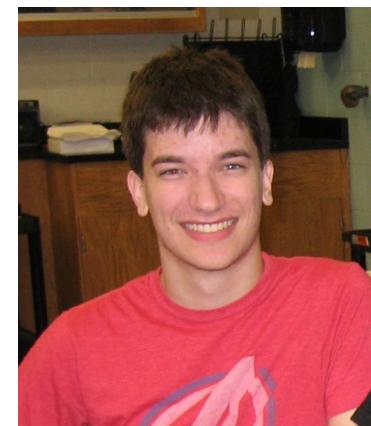
Jutta Luettmer-Strathmann
Akron



Wolfgang Paul
Halle



Jaden Slovensky Beth Breen Norah Ali



Christian O'Neil



Troy Prunty
Hiram