

\mathbb{Z}_2 Lattice Gerbe Theory

Des Johnston
CompPhys14 Leipzig, Nov 2014

Plan of talk

Gerbes?

Lattice Gauge, Gerbe Theory

\mathbb{Z}_2 Lattice Gerbe Theory

Gerbes?

(Abelian) Gauge theory, one forms:

$$S = \frac{1}{4g^2} \int d^4x F_{\mu\nu} F^{\mu\nu}$$

$$F_{\mu\nu} = \partial_{[\mu} A_{\nu]}$$

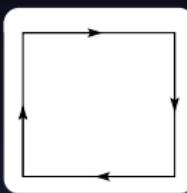
(Abelian) Gerbe theory, two forms:

$$S = \frac{1}{4g^2} \int d^d x H_{\mu\nu\lambda} H^{\mu\nu\lambda}$$

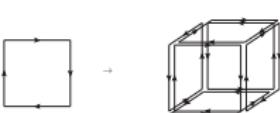
$$H_{\mu\nu\lambda} = \partial_{[\mu} B_{\nu\lambda]}$$

Lattice Gauge/Gerbe Theory

Gauge - Hamiltonian defined on *plaquettes*



Gerbe - Hamiltonian defined on *cubets*



Higher Abelian Gauge theory

General Framework

$$H = - \sum_{C_{n+1}} \left(\prod_{C_n \in \partial C_{n+1}} U(C_n) + c.c. \right)$$

$U(C_n) = \exp(iA(C_n))$ live on the boundaries C_n of cells C_{n+1}

Hamiltonian given by the sum of products of the $U(C_n)$ around the boundary of a C_{n+1}

\mathbb{Z}_2 Gauge theory

Many of the properties are visible already in simplest \mathbb{Z}_2 case

Symmetries, observables (loops)

$$\Gamma(L) = \left\langle \prod_{C_1 \in L} U(C_1) \right\rangle$$

If a confining transition exists

$$\Gamma(L) \sim \begin{cases} \exp(-A(L)) & \beta < \beta_c \\ \exp(-P(L)) & \beta > \beta_c \end{cases}$$

\mathbb{Z}_2 Gerbe theory

Play same game with Gerbe theory

Symmetries, observables (surfaces)

$$\Gamma(S) = \left\langle \prod_{C_2 \in S} U(C_2) \right\rangle$$

If a confining transition exists

$$\Gamma(S) \sim \begin{cases} \exp(-V(S)) & \beta < \beta_c \\ \exp(-A(S)) & \beta > \beta_c \end{cases}$$

So Far, so general

We can use Wegner's results on duality for generalized Ising models (1971) to say more

Lattice N d -dimensional hypercubes

$M_{d,n}$ model, $N_s = \binom{d}{n-1} N$ spins sited at the centres of the $(n-1)$ -dimensional hypercubes

Hamiltonians H_{dn} , product of $2n$ spins on the $(n-1)$ -dimensional faces of the $N_b = \binom{d}{n} N$ n -dimensional hypercubes.

$M_{d,3}$ are lattice Gerbe theories

Schematically

Gauge $M_{d,2}$

$$H = - \sum_{\square} U^4$$

Gerbe $M_{d,3}$

$$H = - \sum_{\boxtimes} U^6$$

(Generalized) Ising duality

$M_{d,n}$ on d -dimensional hypercubic lattice and dual $M_{d,d-n}^*$ on dual lattice

$$Z_{d,n}(\beta) \sim Z_{d,d-n}^*(\beta^*)$$

$$\tanh(\beta) = \exp(-2\beta^*)$$

Duality also in field

$$Z_{d,n}(\beta, h) \sim Z_{d,d-n+1}^*(\beta^*, h^*)$$

$$\tanh(\beta) = \exp(-2h^*)$$

$$\tanh(h) = \exp(-2\beta^*)$$

Consequences of (Generalized) Ising duality $d = 3$

$d = 3$, use in field duality

$$Z_{3,3}(\beta, 0) \sim Z_{3,1}^*(\infty, h^*)$$

Dualize back

$$Z_{3,3}(\beta, 0) \sim 2^{N_s} [\cosh(\beta)^{N_b} + \sinh(\beta)^{N_b}]$$

Analytic, no transition

Consequences of (Generalized) Ising duality $d = 4, 6..$

$d = 4$, \mathbb{Z}_2 Gerbe theory dual to standard Ising model

$$Z_{4,3}(\beta) \sim Z_{4,1}^*(\beta^*)$$

Continuous transition

$d = 6$, self-duality

$$Z_{2n,n}(\beta) \sim Z_{2n,n}^*(\beta^*)$$

6d \mathbb{Z}_2 Gerbe theory is self dual, familiar critical temperature
 $\beta_c = \frac{1}{2} \ln(1 + \sqrt{2})$

Making life more complicated

Gauge Higgs

$$H = - \sum_{\square} U^4 - \lambda \sum_{\langle ij \rangle} U \sigma^2$$

Gerbe Higgs

$$H = - \sum_{\boxtimes} U^6 - \lambda \sum_{\square} U \sigma^4$$

Conclusions

Interest in formulating lattice Gerbe theories

Wegner has very thoughtfully calculated much of what is needed to discuss \mathbb{Z}_2 lattice Gerbe theory

These shed some light on Abelian and non-Abelian Gerbe theories

These models are curiously unsimulated.....

References

- A. E. Lipstein and R. A. Reid-Edwards, *Lattice Gerbe Theory*, JHEP 09 (2014) 034, [arxiv:1404.2634].
- C. Omero, P. A. Marchetti and A. Maritan, J. Phys. A **16**, 1465 (1983).
- F. J. Wegner, J. Math. Phys. **12** 2259 (1971).
- F. J. Wegner, Duality in Generalized Ising Models, [arXiv:1411.5815]
- D. A. Johnston, \mathbb{Z}_2 lattice gerbe theory, Phys. Rev. D 90, 107701 (2014), [arXiv:1405.7890]