

Universality in the $d = 3$ random-field Ising model

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- Some cherished concepts, i.e. the two-exponent scaling scenario ($\bar{\eta} = 2\eta$), have been recently questioned (Tissier and Tarjus, PRL **107**, 041601 (2011)).
- Universality in terms of different field distributions, or even for the same distribution but different values of the disorder strength, has been severely questioned.

Ingredients (I): Hamiltonian and the $T = 0$ scenario

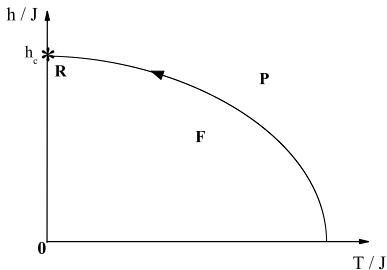
Ingredients (I): Hamiltonian and the $T = 0$ scenario

$$\mathcal{H}^{(\text{RFIM})} = -J \sum_{\langle x,y \rangle} S_x S_y - \sum_x h_x S_x, \quad ; S_x = \pm 1 ; J > 0$$

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- Work at $T = 0$ using efficient optimization algorithms that calculate exact ground states. Avoid statistical errors and equilibration problems.



Ingredients (II): Simulated **continuous** field distributions

- double Gaussian (dG):

$$\mathcal{W}^{(\text{dG})}(h_x; h_R, \sigma) = \frac{1}{2} \frac{1}{\sqrt{2\pi\sigma^2}} \left[e^{-\frac{(h_x - h_R)^2}{2\sigma^2}} + e^{-\frac{(h_x + h_R)^2}{2\sigma^2}} \right]$$

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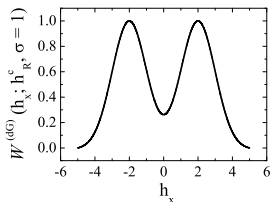
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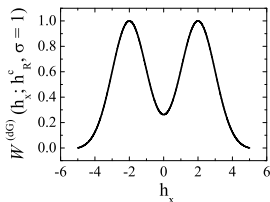
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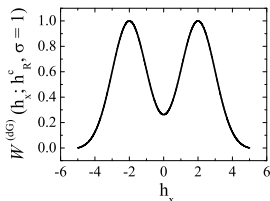


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- $\text{dG}^{(\sigma=2)}$
- Poissonian (P): $\mathcal{W}^{(\text{P})}(h_x; \sigma) = \frac{1}{2|\sigma|} e^{-|h_x|/\sigma}$

Ingredients (III): Computational scheme

- **Fluctuation-dissipation formalism:** Compute *connected* $G_{xy} = \frac{\partial \langle S_x \rangle}{\partial h_y}$ and *disconnected* $G_{xy}^{\text{dis}} = \overline{\langle S_x S_y \rangle}$ correlations functions. For *either* correlator \rightarrow second-moment correlation length.

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- **Re-weighting extrapolation** on h_R and σ : From a single simulation, we extrapolate the mean value of observables to nearby parameters of the disorder distribution. Computing derivatives with respect to σ or h_R is straightforward.

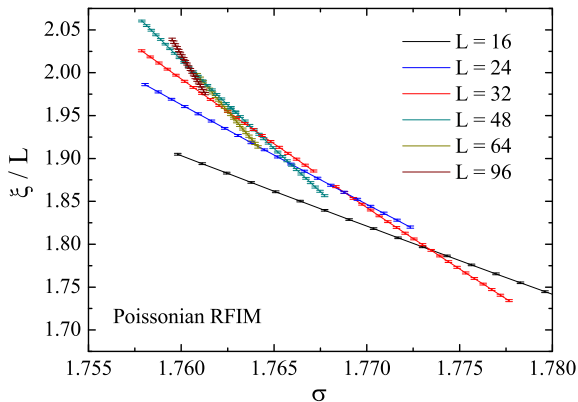
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- For dimensionful quantities O , scaling in the thermodynamic limit as $\xi^{x_O/\nu}$, we consider the quotient $Q_O = O_{2L}/O_L$ at the crossing. For dimensionless magnitudes g , we focus on g_{2L} . In either case, one has:

$$Q_O^{\text{cross}} = 2^{x_O/\nu} + \mathcal{O}(L^{-\omega}), \quad g_{(2L)}^{\text{cross}} = g^* + \mathcal{O}(L^{-\omega}),$$

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- Instances of dimensionful quantities are the derivatives of ξ , $\xi^{(\text{dis})}$ ($x_\xi = 1 + \nu$), the connected and disconnected susceptibilities χ and $\chi^{(\text{dis})}$ [$x_\chi = \nu(2 - \eta)$, $x_{\chi^{(\text{dis})}} = \nu(4 - \bar{\eta})$], and the ratio $U_{22} = \chi^{(\text{dis})}/\chi^2$ [$x_{U_{22}} = \nu(2\eta - \bar{\eta})$].

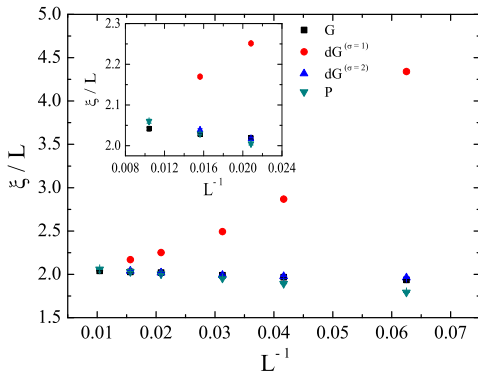
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Huge statistics: $L_{\max} = 192$ and 5×10^7 samples.

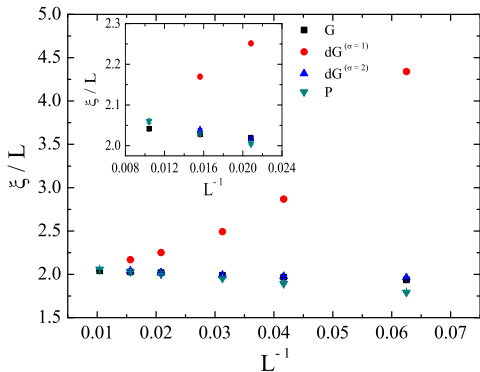
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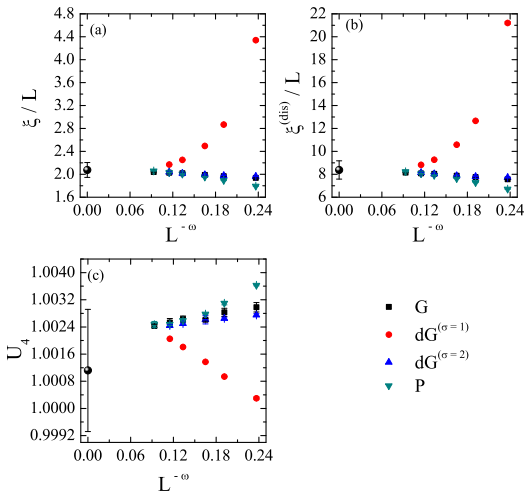
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One needs extrapolation to $L \rightarrow \infty$.

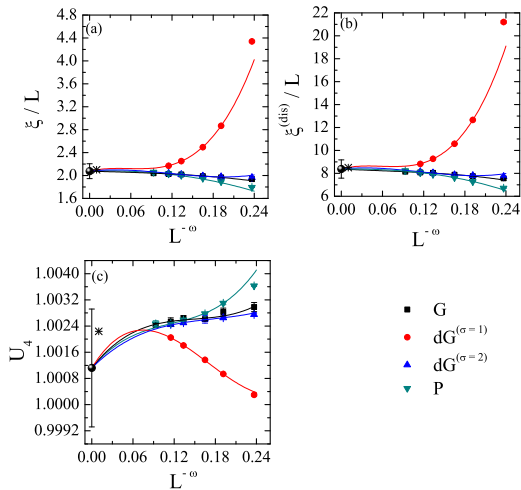
Universality in the $d = 3$ RFIM

$A + BL^{-\omega} + CL^{-2\omega} + DL^{-3\omega}$; $L_{min} = 24$; $\omega = 0.52 \pm 0.11$;
 $\chi^2/dof = 18.83/14$, $Q = 0.17$ (full covariance-matrix!)



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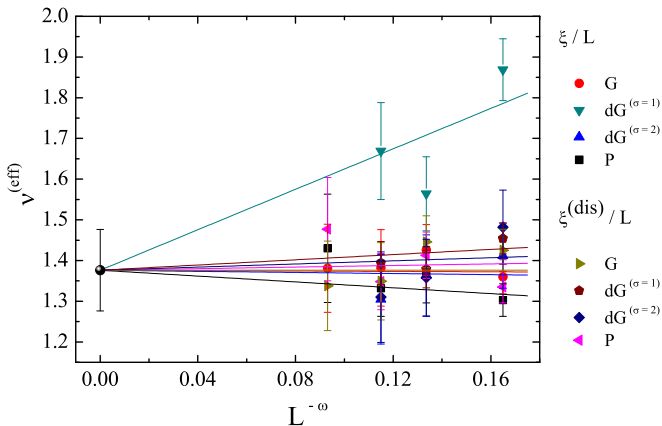


Extrapolation of ν

$$\nu_L = \nu + BL^{-\omega} ; L_{min} = 32 ; \omega = 0.52 ;$$

$$\chi^2/dof = 12.52/10, Q = 0.25$$

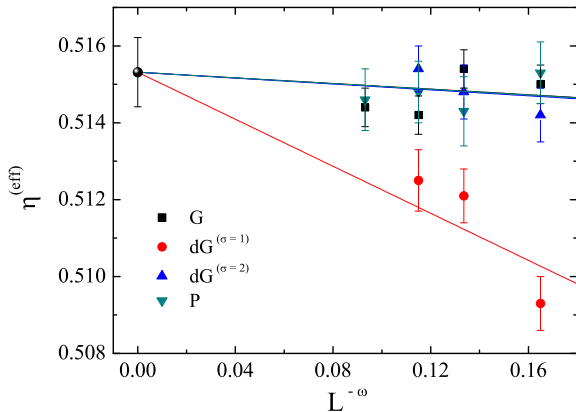
Final estimate: $\nu = 1.38 \pm 0.10$



Extrapolation of η

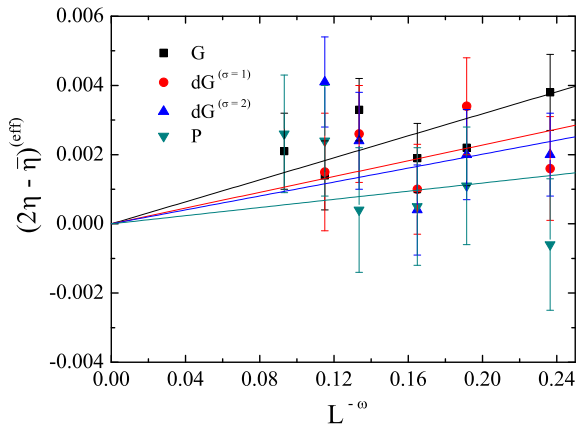
$\eta_L = \eta + BL^{-\omega}$; $L_{min} = 32$; $\omega = 0.52$; $\chi^2/dof = 10/9$, $Q = 0.35$

Final estimate: $\eta = 0.5153 \pm 0.0009$



Extrapolation of $2\eta - \bar{\eta}$

$$(2\eta - \bar{\eta})|_L = BL^{-\omega} ; L_{min} = 16 ; \omega = 0.52 ;$$
$$\chi^2/dof = 18.26/18, Q = 0.44$$



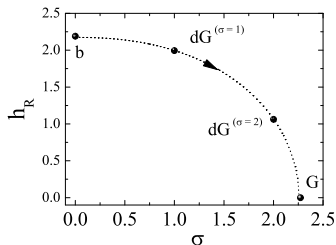
Summary of universal ratios and exponents

| Extrapolation | χ^2/DOF | L_{\min} | Order in $L^{-\omega}$ |
|---|---------------------|------------|------------------------|
| $(\xi/L) _{L=\infty} = 2.08(13)$ | 18.8/14 | 24 | third |
| $(\xi^{(\text{dis})}/L) _{L=\infty} = 8.4(8)$ | | | |
| $U_4 _{L=\infty} = 1.0011(18)$ | | | |
| $\omega = 0.52(11)$ | | | |
| $\nu _{L=\infty} = 1.38(10)(0.03)$ | 12.5/10 | 32 | first |
| $\eta _{L=\infty} = 0.5153(9)(2)$ | 10.0/9 | 32 | first |
| $(2\eta - \bar{\eta}) _{L=\infty} = 0$ (fixed) | 18.3/18 | 16 | first |
| $(2\eta - \bar{\eta}) _{L=\infty} = 0.0026(9)(1)$ | 10.5/17 | 16 | first |
| $\sigma^c[\text{G}] = 2.27205(18)(4)$ | 3.1/3 | 16 | second |
| $h_R^c[\text{dG}^{(\sigma=1)}] = 1.9955(6)(24)$ | 2.5/1 | 24 | second |
| $h_R^c[\text{dG}^{(\sigma=2)}] = 1.0631(7)(10)$ | 0.7/2 | 16 | second |
| $\sigma^c[\text{P}] = 1.7583(2)(2)$ | 3.0/3 | 16 | second |

Conclusions

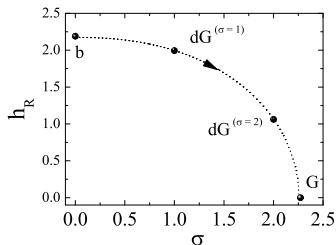
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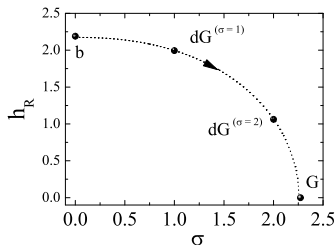
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- The phase diagram of the RFIM is seemingly ruled by a **single fixed point**:



- Existence of **strong scaling corrections** that need to be carefully monitored. Very accurate computation of anomalous dimensions $\eta, \bar{\eta}$.
- The **two-exponent scaling scenario holds** within an accuracy of two parts in a thousand (2/1000).

Acknowledgements

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- Computational time in the clusters *Terminus* and *Memento* (BIFI Institute Zaragoza) and the local cluster of the Department of Physics at the Universidad Complutense de Madrid.

Future work in this direction

- Investigate scaling corrections - the **Achilles' heel** in the study of the RFIM - and universality aspects in higher dimensions. Work in progress with Víctor Martín-Mayor, Marco Picco and Nicolas Sourlas.

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- Royal Society Research Grant in collaboration with Martin Weigel: “Simulating dirty magnets on GPUs”.

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