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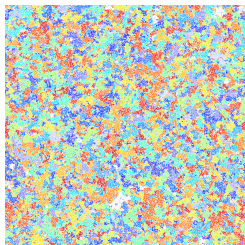


True asymptotics of self-avoiding walks on 3D percolation clusters

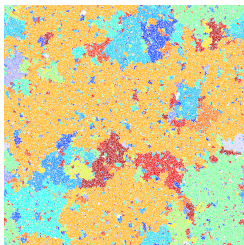
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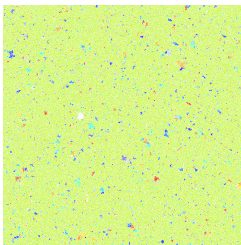
Percolation



$$p = 0.55 < p_c$$



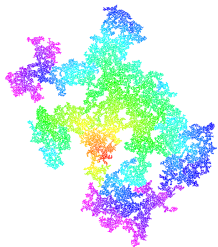
$$p = 0.592746 \approx p_c$$



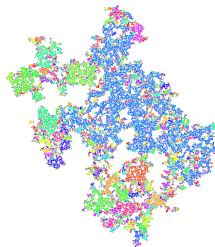
$$p = 0.62 > p_c$$

D	p_c	d_f	d_l	d_{Bf}
2	0.592...	91/48	1.68...	1.64...
3	0.311...	2.52...	1.84...	1.81...
6	0.109...	4	2	2

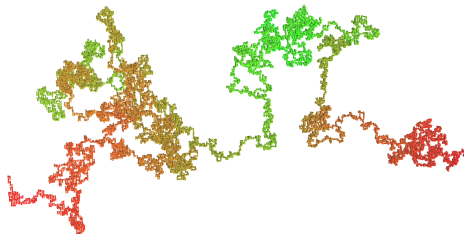
Structure of the critical cluster



10^3 chemical shells around origin

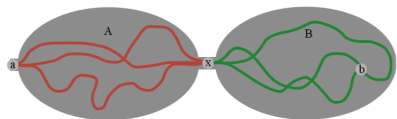


Bi-connected components

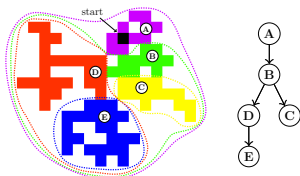


← Backbone of a
3D cluster

Hierarchical factorization

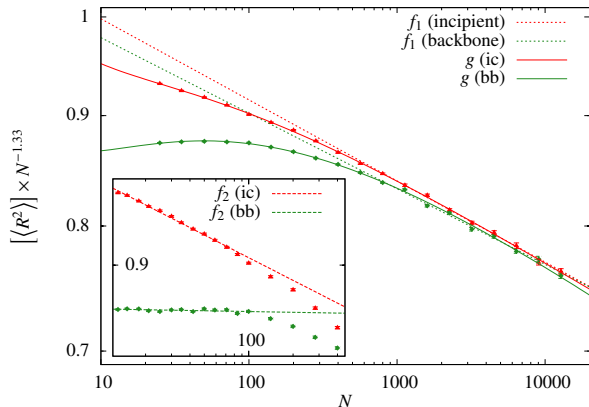


$$Z_{a \rightarrow b}[N] = \sum_{i=0}^N Z_{a \rightarrow x}[i] \cdot Z_{x \rightarrow b}[N - i]$$



- Exact enumeration of SAW configurations in polynomial time.
- 10^4 steps now take ≈ 10 minutes, down from $\approx 10^{960}$ years.

End-to-end distance and Flory exponent ν



$$f(N) \cdot N^{1.33} = a \cdot N^{2\nu}$$

$$g(N) \cdot N^{1.33} = a \cdot N^{2\nu} \left(1 + b/N^\Delta\right)$$

Flory exponent ν – results

Our study¹:

ansatz	range	$\nu_{\text{incipient}}$	ν_{backbone}
$a \cdot N^{2\nu}$	25-84	0.6547(2)	0.6646(2)
$a \cdot N^{2\nu}$	800-12800	0.6433(4)	0.643(1)
$a \cdot N^{2\nu} (1 + b/N^\Delta)$	25-12800	0.644(2)	0.640(3)

Previous numerical results:

study	range	$\nu_{\text{incipient}}$	ν_{backbone}
exact enum. ²	$N \leq 30$	0.660(5)	-
exact enum. ³	$N \leq 40$	-	0.662(6)
MC (PERM). ⁴	$N < 100$	-	0.667(3)

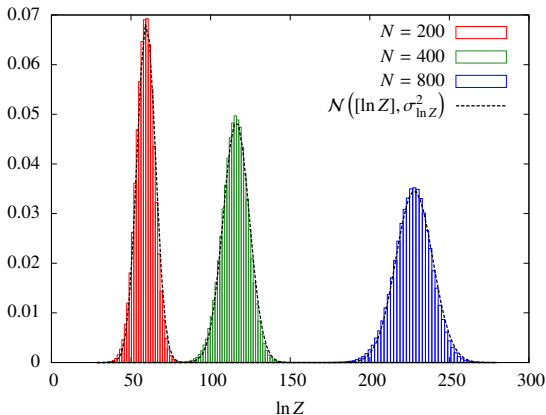
¹N. F. and W. Janke, PRL (in print)

²M. Rintoul et al., PRE 49 2790 (1994)

³A. Ordemann et al., PRE 61 6858 (2000)

⁴V. Blavatska and W. Janke, EPL 82 66006 (2008)

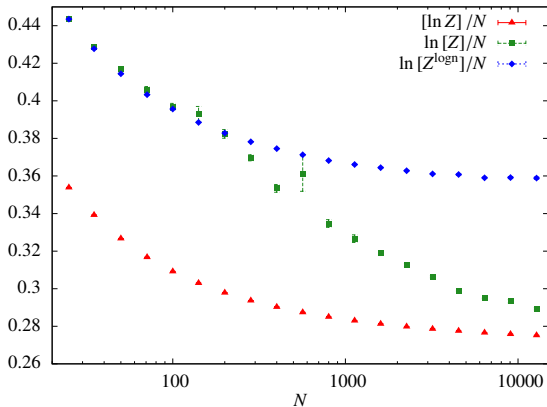
Number of walk conformations Z



$$\sigma_{\ln Z}^2 = A \cdot N^{2\chi}, \quad \chi = 0.500(1)$$

Problem: distribution of Z_N almost log-normal
→ large deviations make estimating $[Z_N]$ difficult.

Number of walk conformations Z



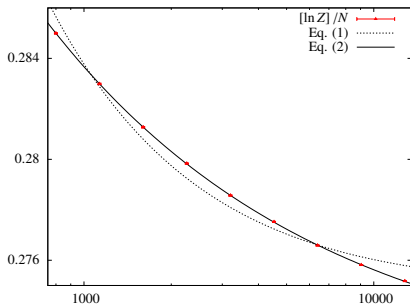
Approximation: $[Z] \approx e^{[\ln Z] + \sigma_{\ln Z}^2/2} =: [Z_{\log n}]$.

Scaling laws for $[Z_N]$ and $[\ln Z_N]$?

Assumption: $e^{[\ln Z_N]} \sim \mu_0^N N^{\gamma_0-1}$

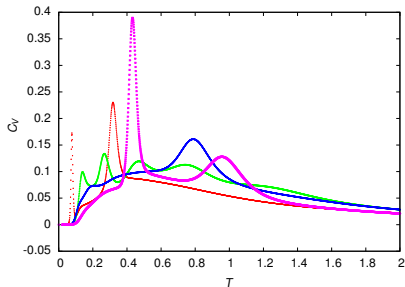
→ $[\ln Z_N] / N = \ln a / N + \ln \mu_0 + (\gamma - 1) \ln N / N$ (1)

Modified: $[\ln Z_N] / N = \ln a / N + \ln \mu_0 (1 + bN^{-\zeta})$ (2)

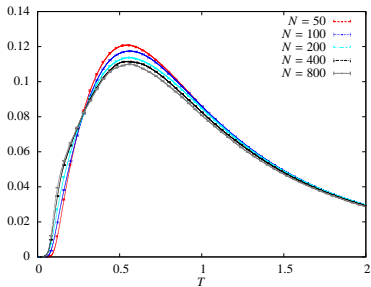


→ $[Z_N] \approx [Z_{\log n}] \sim \cancel{\mu_0^N N^{\gamma_0-1}} \mu^{N(1+b/N^\zeta)}, \quad \mu = e^{\ln \mu_0 + A/2} ?$

Specific heat



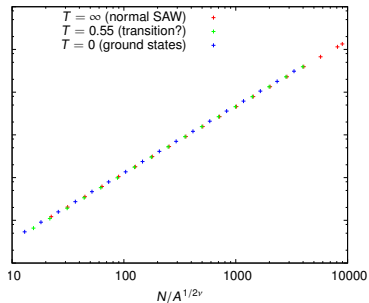
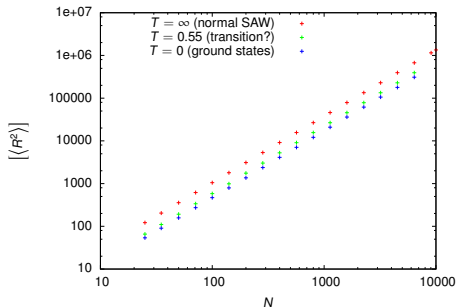
Single clusters



Quenched average

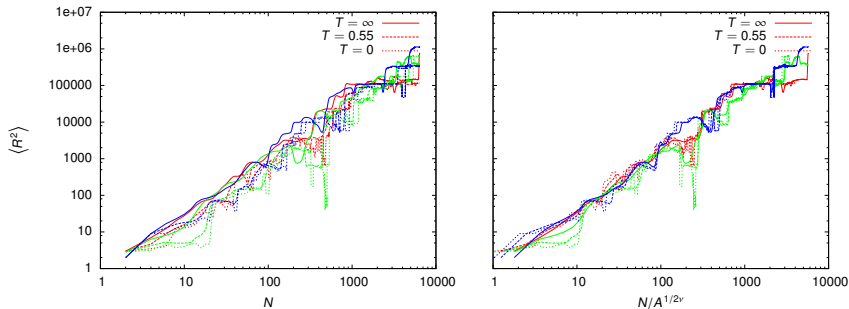
No divergence for average $[C_V] = \frac{1}{NT^2} [\langle E^2 \rangle - \langle E \rangle^2]$.

End-to-end distance



Temperature only changes amplitude: $A_T \cdot \langle R^2 \rangle \sim N^{2\nu}$

End-to-end distance for single clusters



With rescaled N -axis, patterns for different T look very similar!

For once, energy and entropy have shared interests!

Summary

- SAWs on critical percolation clusters: basic model for polymers in disordered media.
- Fractal nature of critical clusters enables SAW enumeration in polynomial time.
- True (?) asymptotic behavior only visible after $\sim 10^3$ steps.
- Flory exponent ν smaller than previously thought, but assumption $\nu_{\text{backbone}} = \nu_{\text{full cluster}}$ holds.
- $[Z_N] \sim \mu^N N^{\gamma-1}$ wrong, instead: $[Z_N] \sim \mu^{N(1+b/N^\zeta)}$?
- Identical scaling behavior for SAWs with short-range attractions at any temperature – no Θ -transition.

Thank you for the attention!

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