

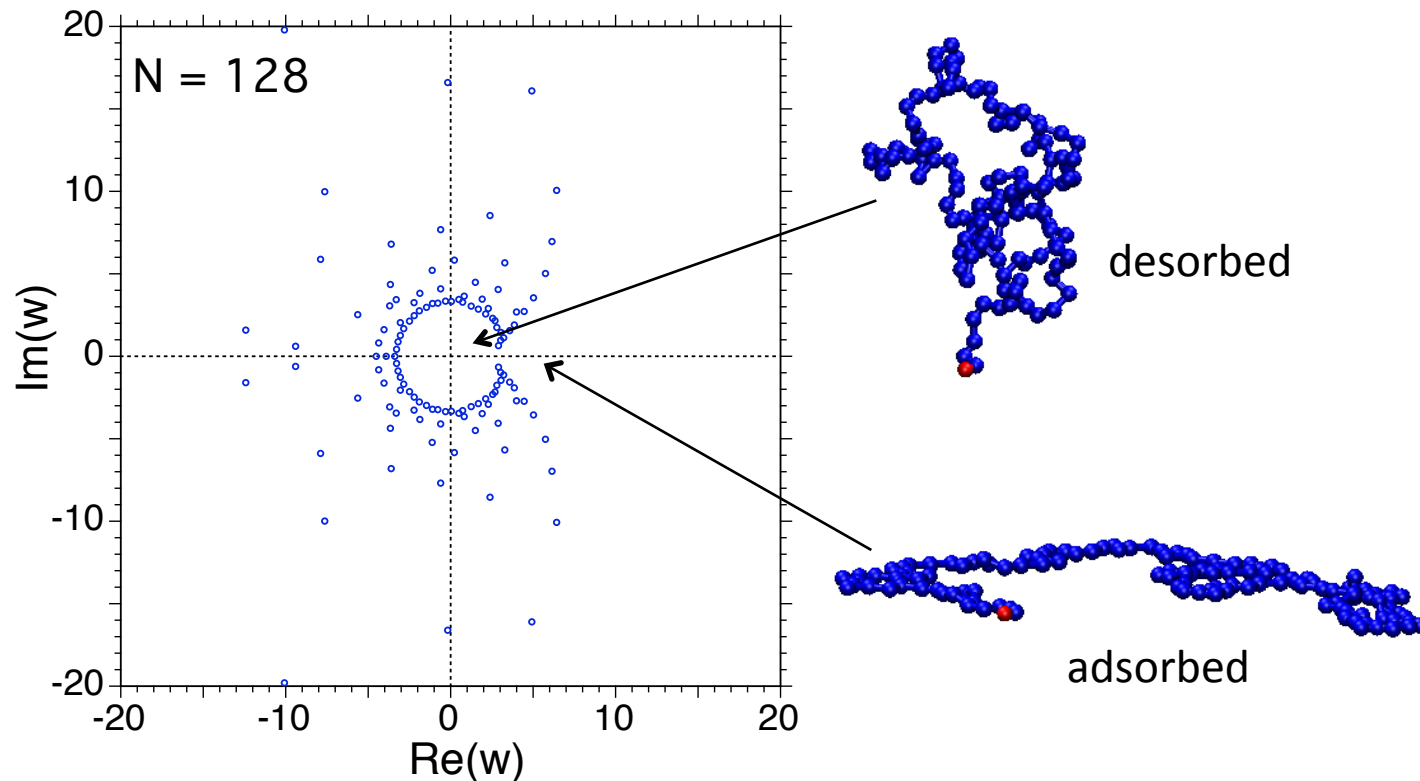
Partition function zeros and finite size scaling for polymer adsorption

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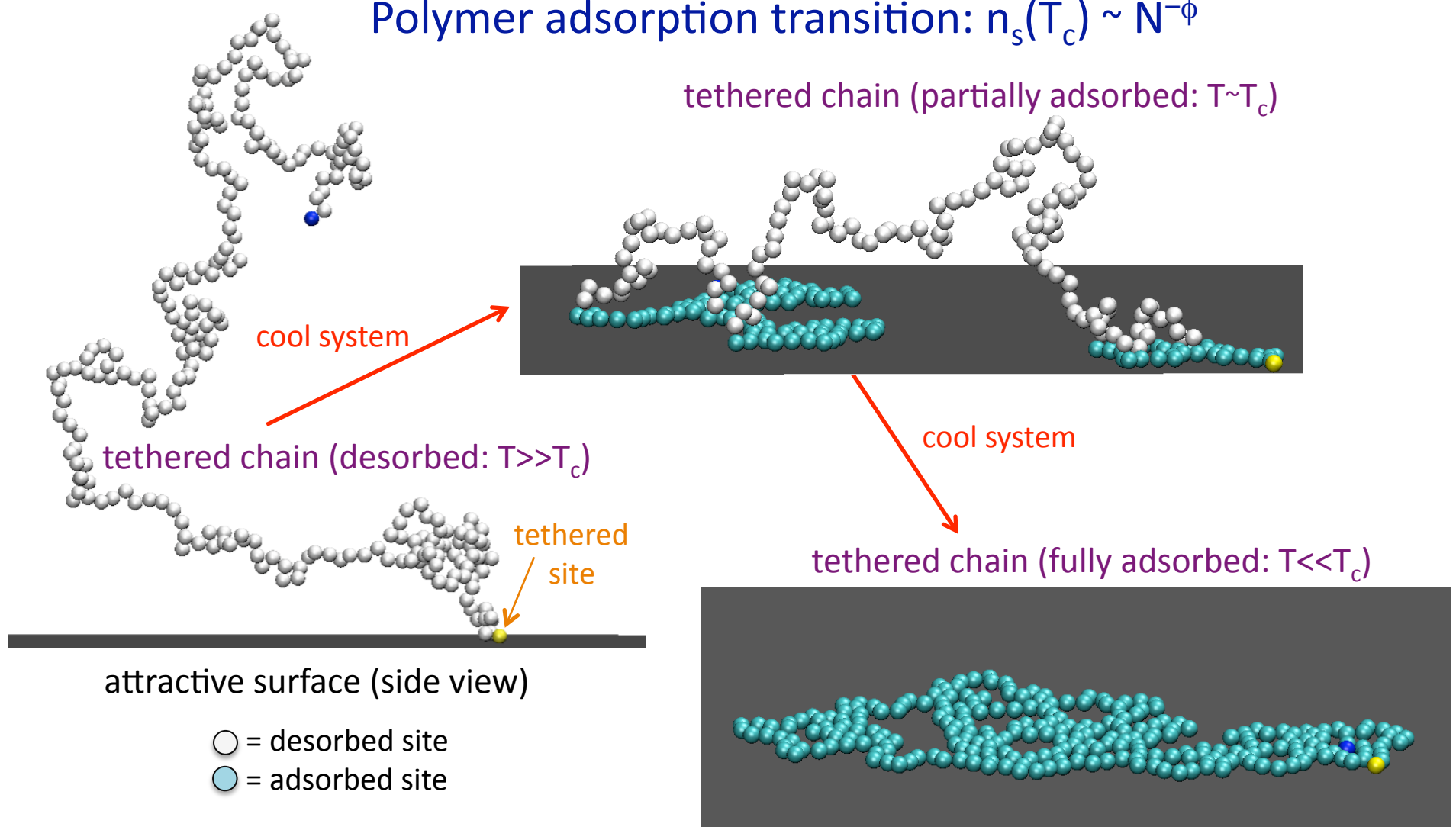


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Polymer adsorption transition: $n_s(T_c) \sim N^{-\phi}$



Eisenriegler, Kremer, Binder, JCP **77**, 6296 (1982) [$\phi = 0.58$]

Metzger, Muller, Binder, Bashnagel, Macromol Theory Sim **11**, 985 (2002) [lit. review: $\phi = 0.40-0.67$; $\phi=0.50$]

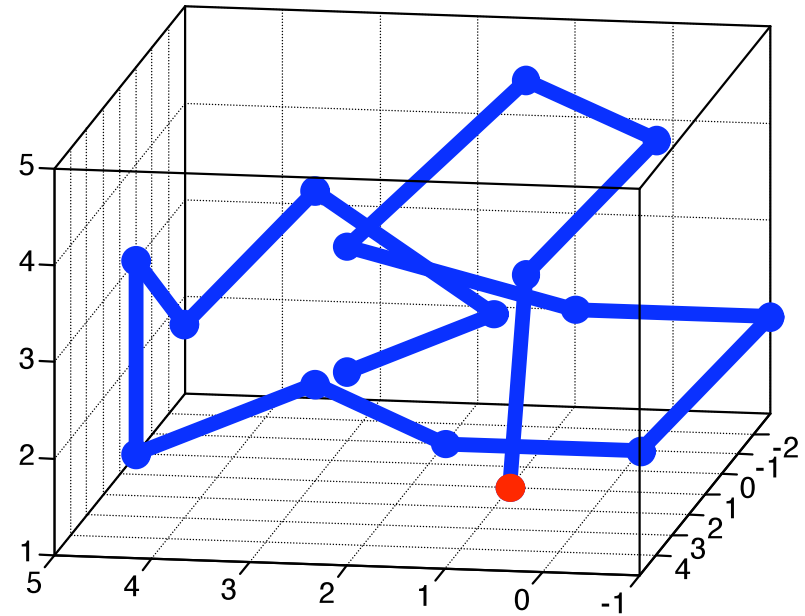
Grassberger, J Phys A: Math Gen **38**, 323 (2005) [$\phi = 0.48$]

Decas, Sommer, Blumen, Macromol Theory Sim **17**, 429 (2008) [$\phi = 0.50-0.59$]

Klushin, Polotsky, Hsu, Markelov, Binder, Skvortsov, PRE **87** 022604 (2013) [$\phi = 0.48$]

Bond fluctuation model for tethered chains

- Simple cubic lattice ($a = 1$)
- Bond lengths: $\sqrt{4}$ to $\sqrt{10}$
- Bead 1 fixed at $(1,1,1)$
- 108 bond vectors
- Interactions between beads i and j with distance r_{ij} :
 - hard core repulsion for $r_{ij}^2 < 4$
- Interactions with the surface:
 - contribution to the internal energy of $-\varepsilon_s$ for $z_i = 1$



Total energy:

$$E = -n_s \varepsilon$$

n_s = number of surface contacts

$T = k_B T / \varepsilon$ = reduced temperature

Density of States and Wang-Landau Sampling I

Density of States:

$g(E_n)$ = volume of configurational phase space for energy state E_n



Thermodynamics:

microcanonical entropy:

$$S(E) = k_B \ln g(E)$$

canonical partition function:

$$Z(T) = \sum g(E) \exp(-E/k_B T)$$

Wang-Landau algorithm* ... an iterative simulation method to compute $g(E_n)$:

Starting w/ $g(E_n)=1$, $H(E_n)=0 \forall n$, $f_0 = e$

Generate sequence of chain conformations using acceptance criteria:

$$P_{acc}(a \rightarrow b) = \min\left(1, \frac{g(E_a)}{g(E_b)}\right)$$

Update DOS: $g(E_n) \rightarrow f_m g(E_n)$

Update visitation

histogram: $H(E_n) \rightarrow H(E_n)+1$

When histogram \sim flat ...

reduce modification factor: $f_{m+1} = (f_m)^{1/2}$

reset histogram to zero: $H(E_n) = 0 \forall n$

iterate
m levels

m=20 is
standard

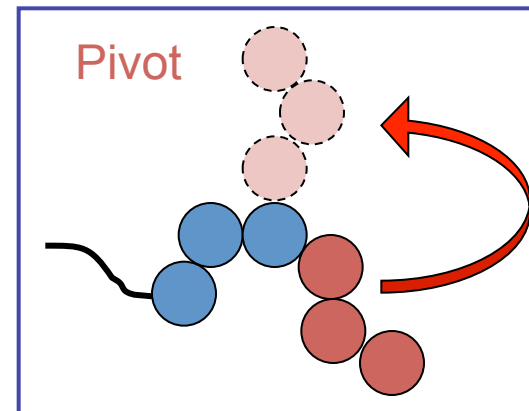
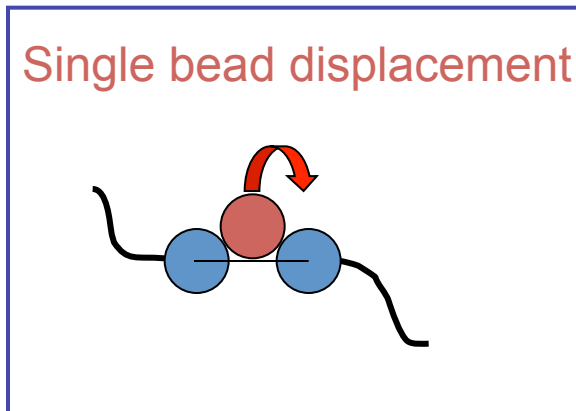
we need
m>25

*Wang & Landau, PRL 86, 2050 (2001); PRE 64, 056101 (2001).

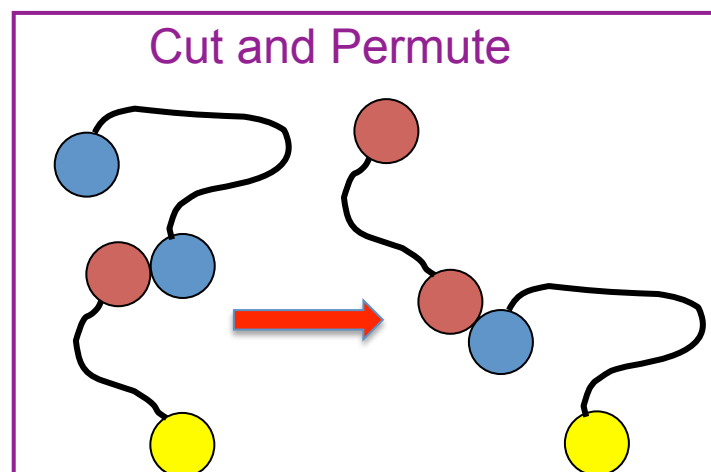
Wang-Landau Sampling II

Success of the WL methods depends on underlying MC move set

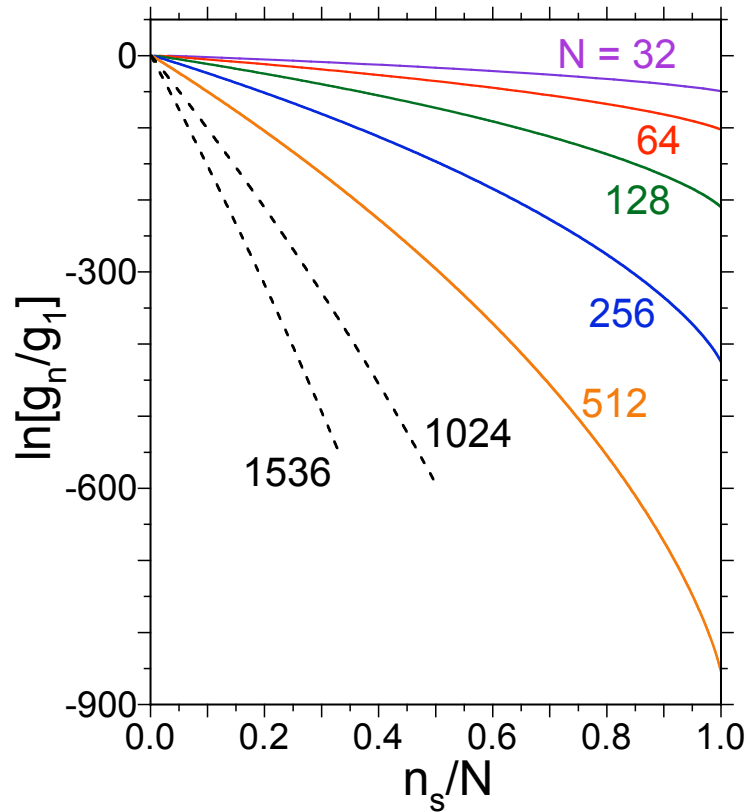
These "standard" moves easily sample most of configuration space:



But for good sampling of chain configurations near the grafting point we also use:



Polymer Adsorption: Single Chain DOS and Specific Heat



n_s = number of surface contacts

expected behavior:

$$C_{\text{inf}}/Nk_B \sim |T - T_c|^{-\alpha}$$

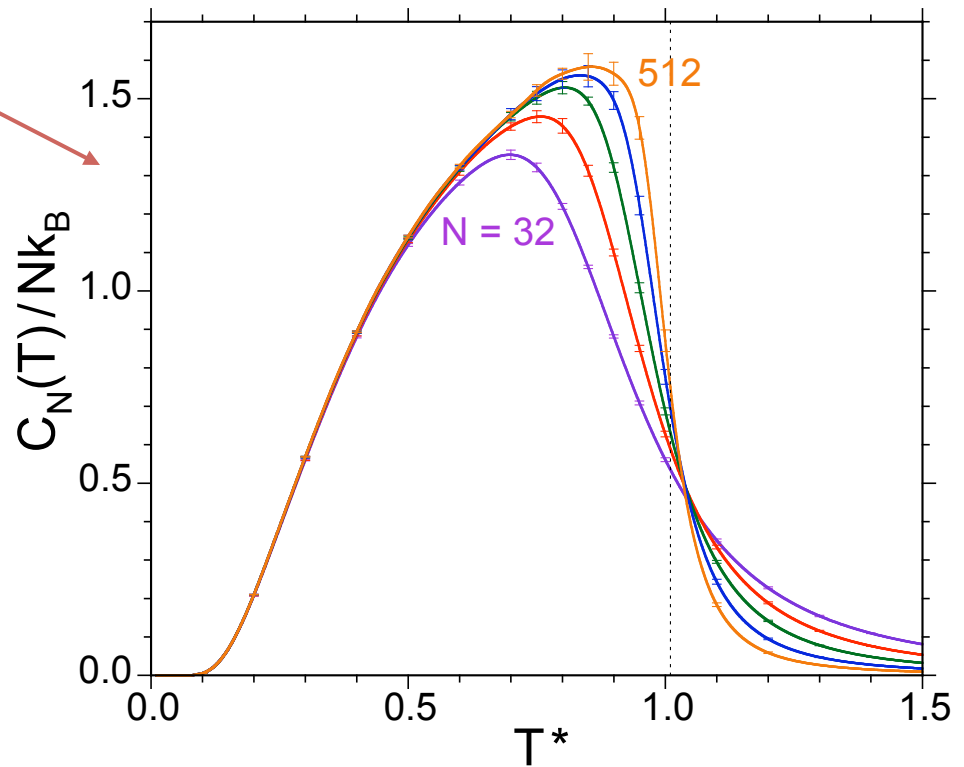
Canonical Analysis

Partition Function: $Z = \sum g(E) e^{-E/kT}$

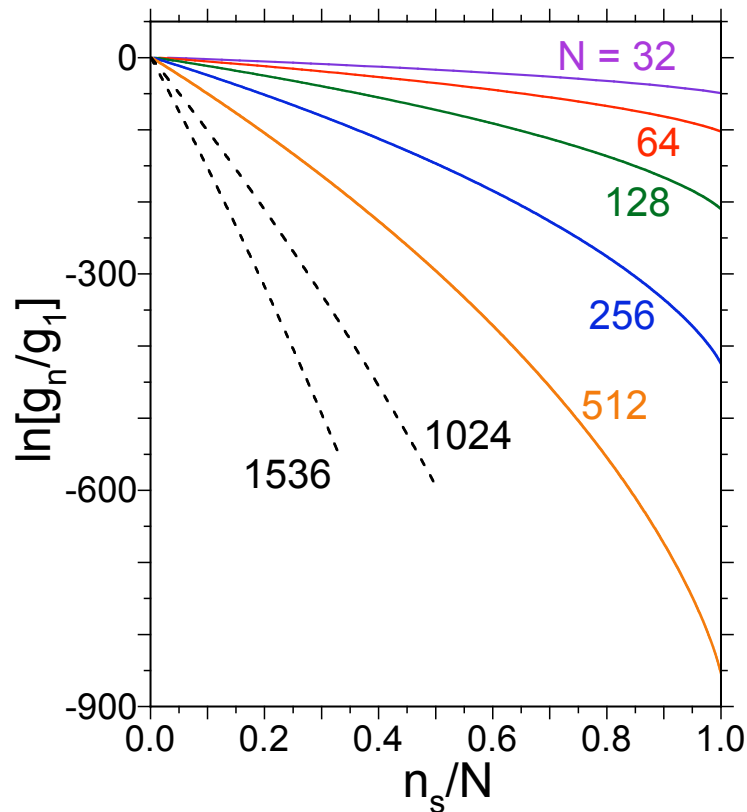
Probability: $P(E,T) = g(E) e^{-E/kT} / Z$

Average Energy: $\langle E(T) \rangle = \sum EP(E,T)$

Specific Heat: $C(T) = d\langle E(T) \rangle/dT$



Polymer Adsorption: Single Chain DOS and Fraction Adsorbed



n_s = number of surface contacts

expected behavior:

$$M_{\text{inf}}(T) \sim |T - T_c|^\beta$$

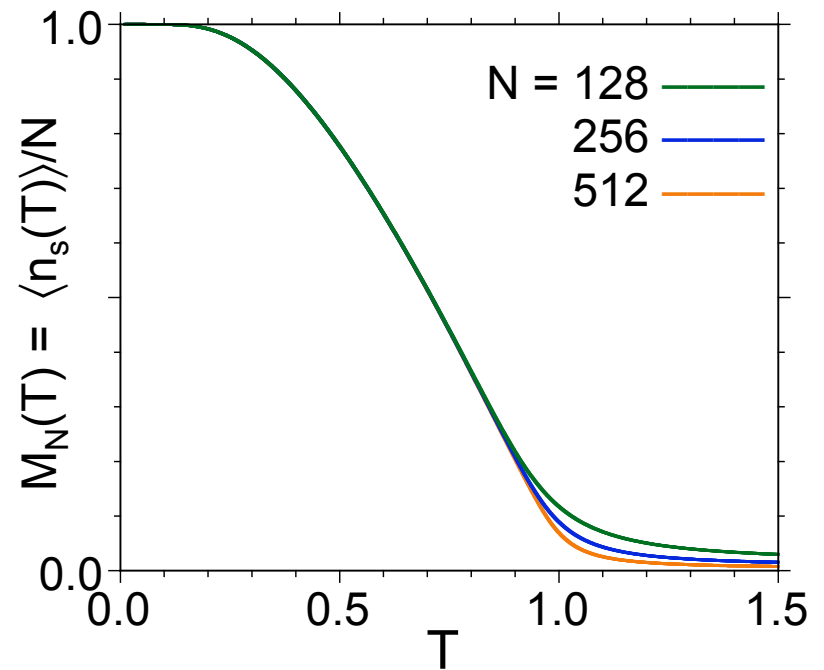
Canonical Analysis

Partition Function: $Z = \sum g(E) e^{-E/kT}$

Probability: $P(E,T) = g(E) e^{-E/kT} / Z$

Average Energy: $\langle E(T) \rangle = \sum E P(E,T)$

Order Parameter: $M(T) = \langle n_s(T) \rangle / N = \langle E(T) \rangle / N\varepsilon$



Polymer Adsorption: Partition Function Zeros

Model has discrete energy spectrum: $-\varepsilon, -2\varepsilon, \dots, -N\varepsilon$

Partition function is a polynomial in $y=e^{\varepsilon/kT}$:

$$Z(T) = \sum g(E) e^{-E/kT} = \sum_n g_n y^n$$

or

$$Z(T) = \prod_k (y-w_k) \quad \text{where } w_k = a_k + ib_k \text{ are the complex zeros of } Z(T)$$

Properties: complex zeros come in pairs $a \pm ib$
any real zeros must be negative
sum of $\text{Re}(w_k)$ is negative, i.e., $\sum_k a_k < 0$

All thermodynamics can be expressed in terms of the zeros $\{w_k\}$

Example: Heat Capacity (physical temp. range: $y > 1$)

$$\frac{C(y)}{k_B} = \beta^2 \frac{\partial^2 \ln Z}{\partial \beta^2} = (\ln y)^2 \sum_{k=0}^{k_{\max}} \frac{-y w_k}{(y - w_k)^2}$$

Zeros near real axis contribute most

Partition Function Zeros: Yang-Lee theory

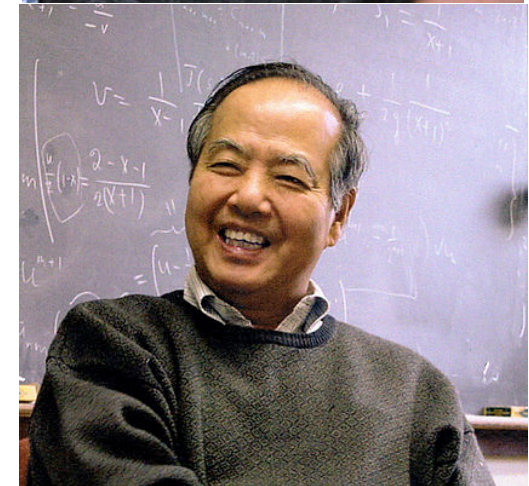
In 1952 C. N. Yang and T. D. Lee proposed a very general theory for phase transitions based on the distribution of the zeros of the grand canonical partition function in the complex plane.

Non-analytic behavior of thermodynamic functions arises when some zeros move onto the positive real axis in the thermodynamic limit.



In 1965 M. Fisher extended the approach to the canonical partition function and zeros in the complex temperature plane.

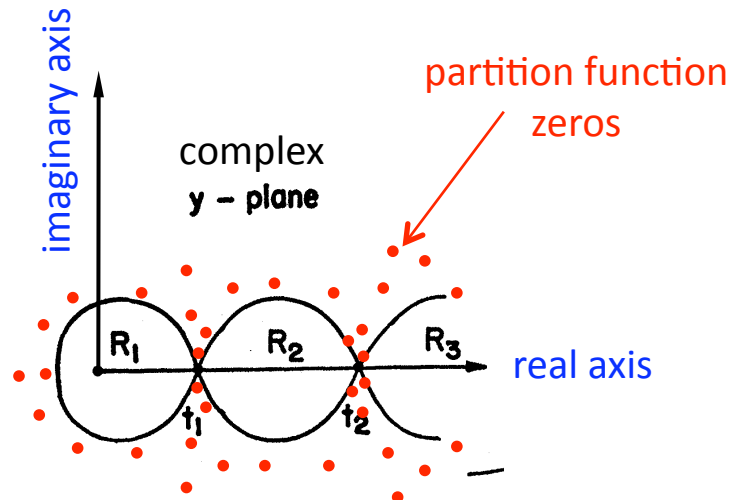
Fisher, in *Lecture Notes in Theoretical Physics* (U. of Colorado Press, 1965).



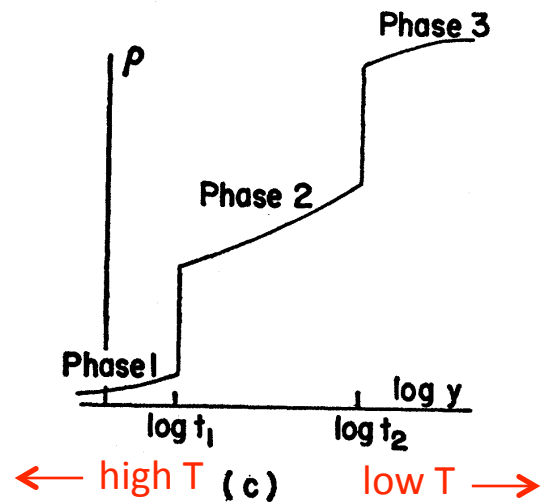
Yang & Lee, *Phys. Rev.* 87, 404, 410 (1952).

Yang-Lee theory of phase transitions

System with two phase transitions

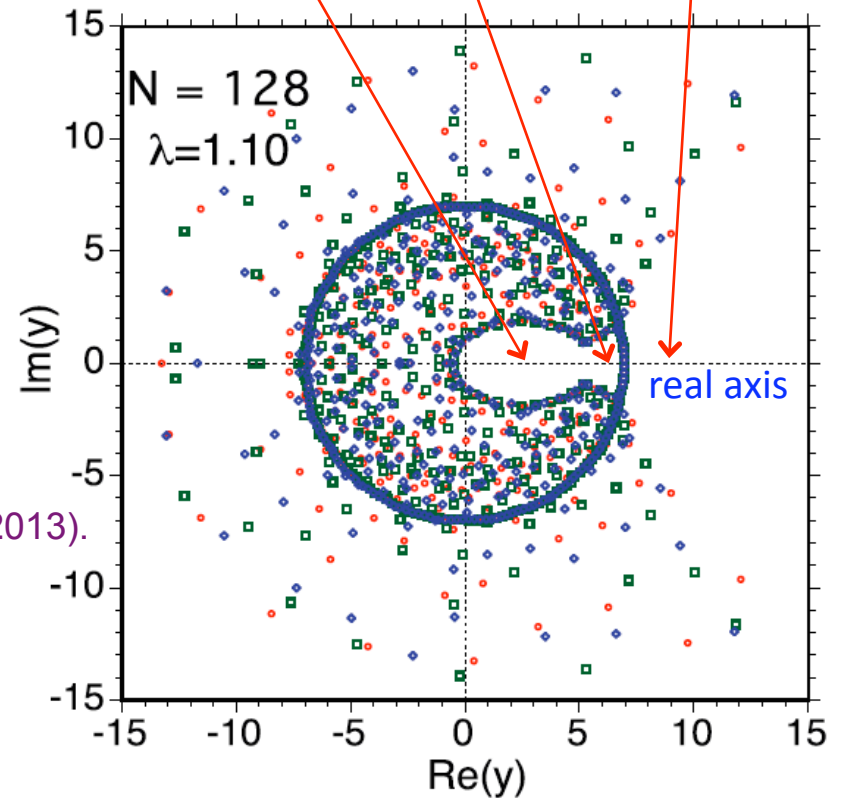
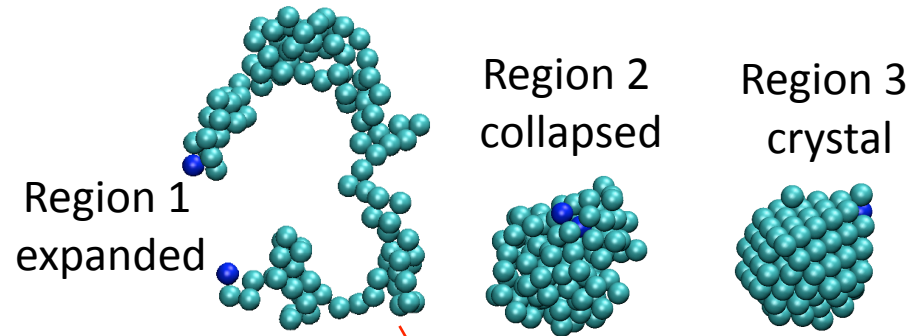


Yang & Lee, Phys. Rev. **87**, 404, 410 (1952).

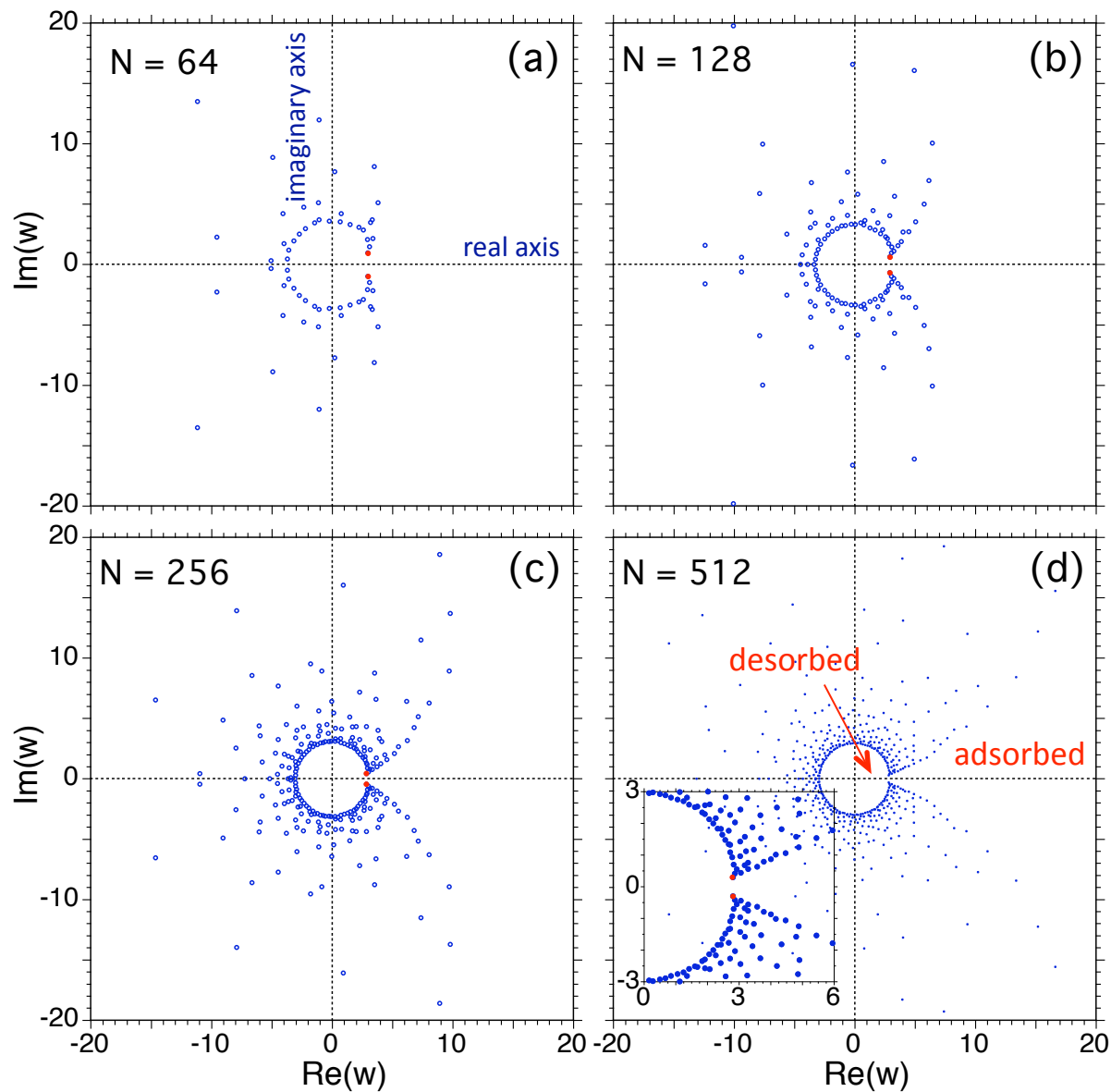


Taylor, Aung, Paul, PRE **88**, 012604 (2013).

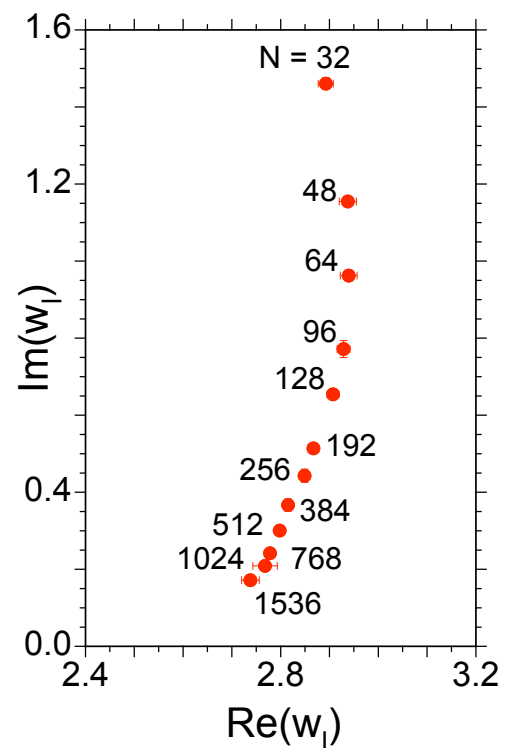
Collapse and freezing of a SW-chain



Roots maps for polymer adsorption follow Yang-Lee behavior ...
leading zeros pinch down onto positive real axis:



Approach of the leading zeros towards the real axis:



Finite size scaling of the leading zeros

Expected scaling form:

$$w_1 - y_c \sim DN^{-\phi}$$

critical point (real: $y=e^{\beta}=e^{1/\tau}$)

crossover exponent

leading root (complex)

complex constant ($D=d_a+id_b$)

Itzykson, Pearson, Zuber, Nucl. Phys. B **220**, 415 (1983)

Scaling of $\text{Im}[w_1]$ gives ϕ :

$$\text{Im}[w_1] \sim d_b N^{-\phi}$$

Linear correlation between

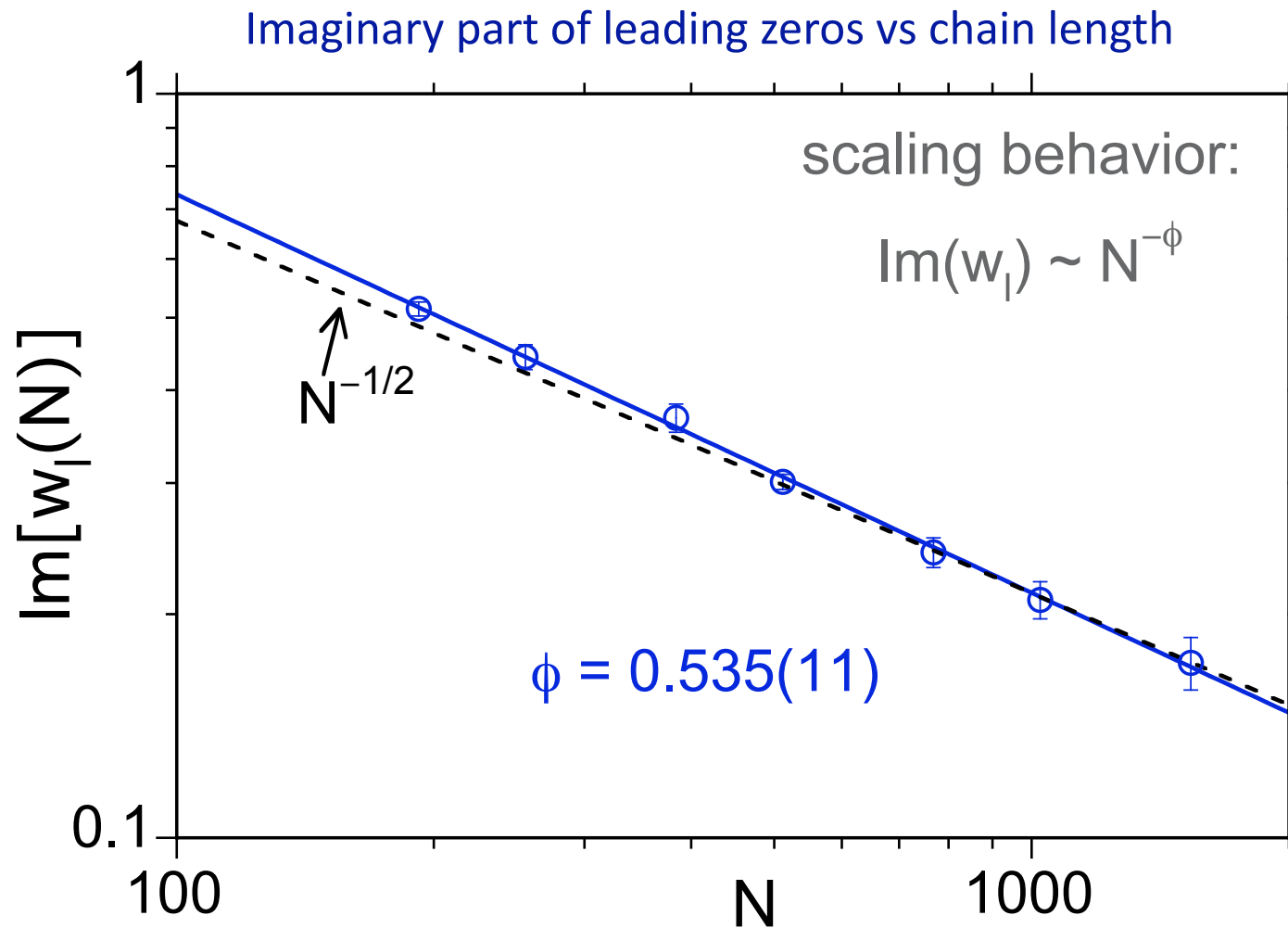
$\text{Im}[w_1]$ and $\text{Re}[w_1]$ gives y_c : $\text{Im}[w_1] \sim (\text{Re}[w_1] - y_c)d_b / d_a$

Alternately, linear correlation

between $\text{Im}[w_1]$ and β_1 gives β_c : $\text{Im}[w_1] \sim (\beta_1 - \beta_c)y_c d_b / d_a$

Complex inverse temperature: $\ln(w_1) = \beta_1 + i\tau_1$

Adsorption transition: The crossover exponent ϕ

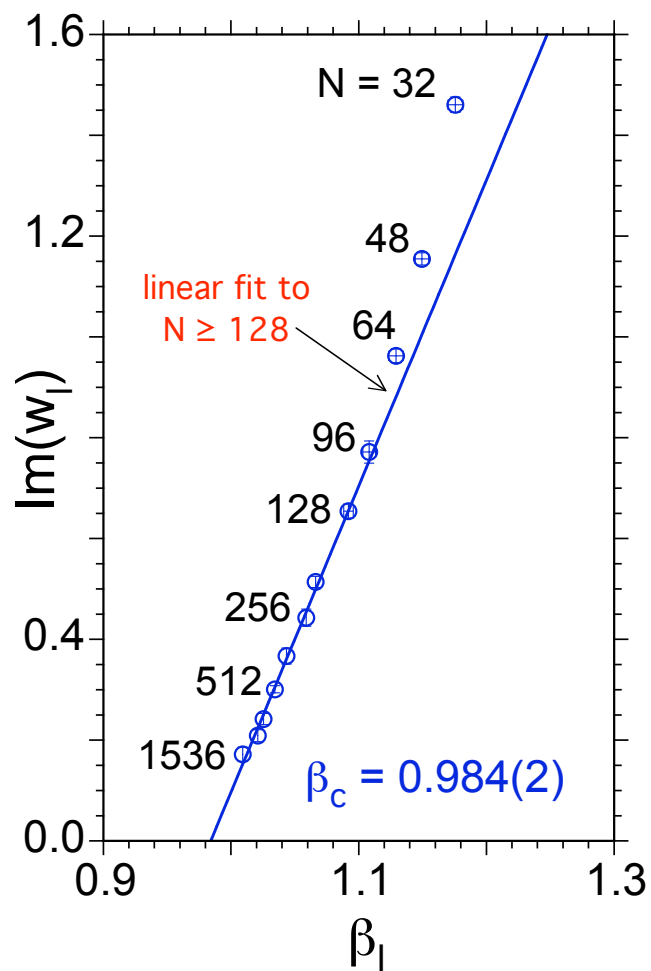
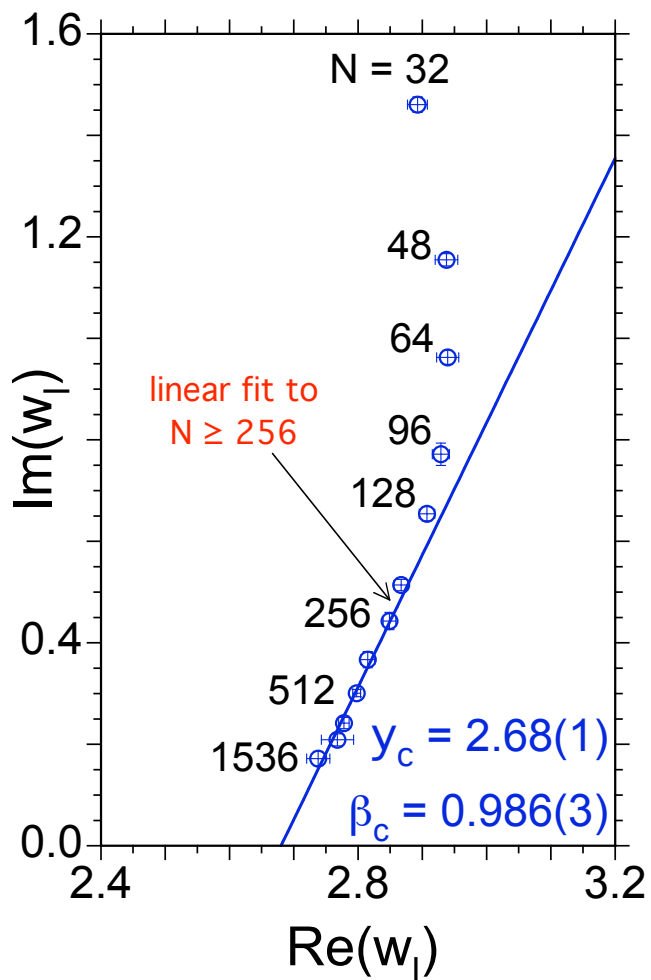


Adsorption transition: The transition temperature T_c

For long chains the leading zeros are expected to show a linear approach the $\text{Re}(w)$ -axis (or the β -axis)

$$y = \text{Re}(w) = e^{1/T}$$

$$\beta = 1/T$$



Result for transition temperature:

$$T_c = 1.016(2)$$

Finite size scaling for thermodynamics

PFZ scaling form in
complex temperature plane:

$$\ln(w_1) - \beta_c \sim (D/y_c)N^{-\phi}$$

critical point (real) → β_c

crossover exponent → $N^{-\phi}$

complex constant ($D=d_a+id_b$) → (D/y_c)

leading root (complex: $w=\beta+i\tau$) → w_1

Real part gives: $\beta_1 - \beta_c \sim N^{-\phi} \longrightarrow \boxed{|1 - T_c/T_1| \sim N^{-\phi}}$

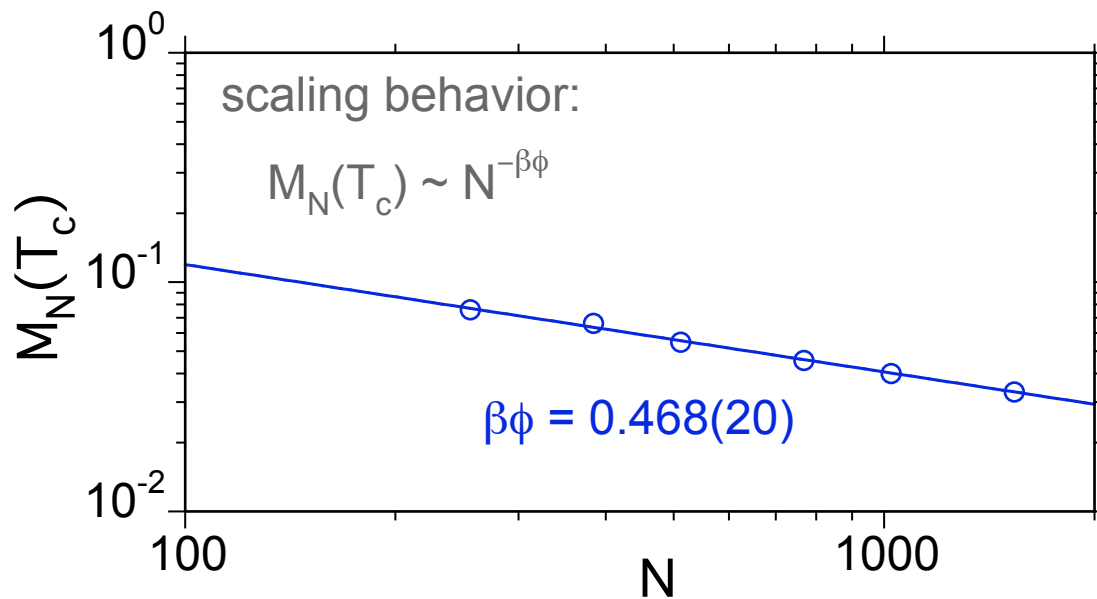
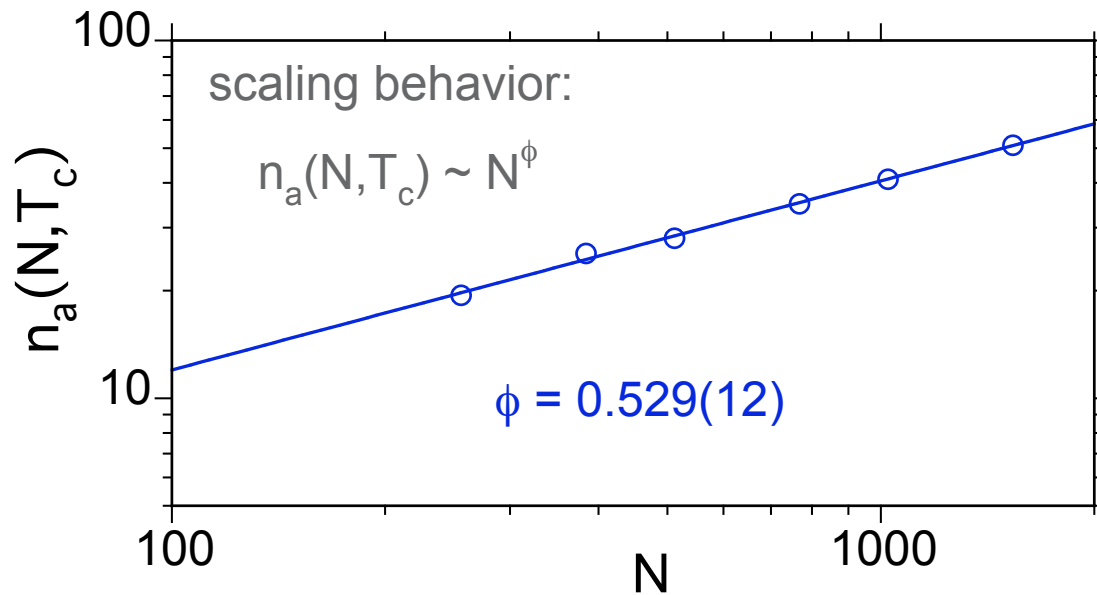
Leads to scaling relations:

$$\begin{cases} M_N(T_c) \sim N^{-\tilde{\beta}\phi} & \text{order parameter} \\ n_s(T_c) = NM_N(T_c) \sim N^\phi & \text{number adsorbed} \\ C_N(T_c)/Nk_B \sim N^{\alpha\phi} & \text{specific heat} \end{cases}$$

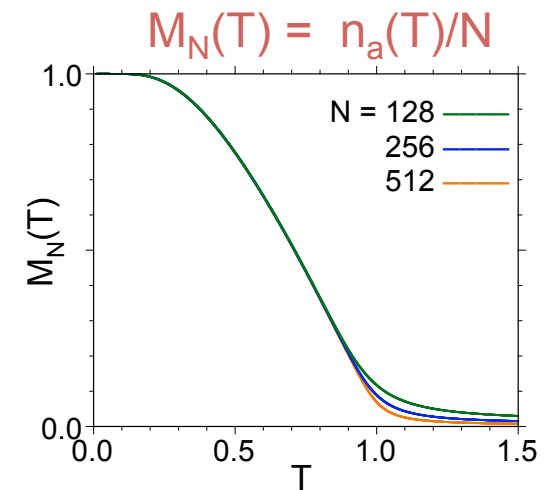
And exponent identities:

$$\begin{cases} \tilde{\beta} = -1 + 1/\phi \\ \alpha = 2 - 1/\phi \end{cases}$$

Finite size scaling at T_c : Number absorbed and order parameter



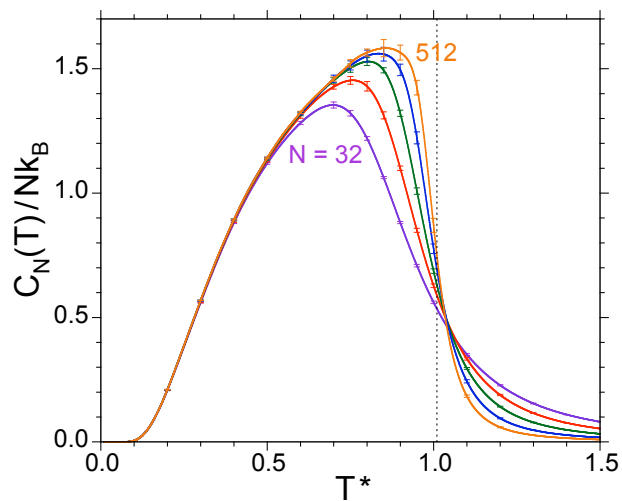
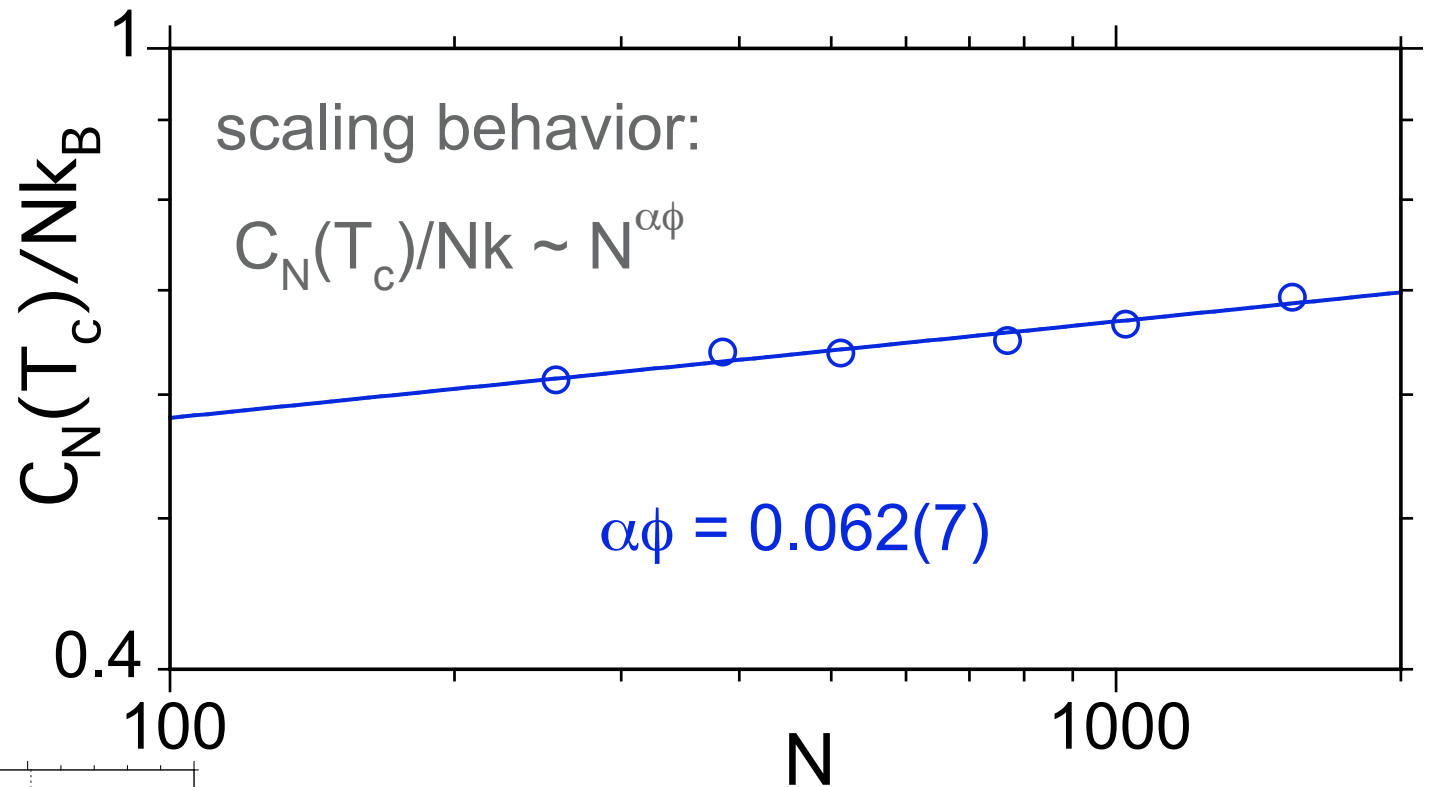
Agrees with Result
 from $\text{Im}(w_1)$:
 $\phi = 0.535(11)$



Result for order parameter
 exponent: $\beta = 0.875(41)$

Agrees with: $\beta = -1 + 1/\phi$

Finite size scaling at T_c : Specific Heat



Result for specific heat exponent: $\alpha = 0.116(13)$

Agrees with exponent identity: $\alpha = 2 - 1/\phi$

Polymer Adsorption: Summary of Results

Transition temperature: $T_c = 1.016(2)$

Crossover exponent: $\phi = 0.535(11)$

Specific heat exponent: $\alpha = 0.116(13)$

Order Parameter exponent: $\beta = 0.875(41)$

Satisfy exponent identities: $\alpha = 2 - 1/\phi$
 $\beta = -1 + 1/\phi$

Caveat: Grassberger and Hsu et al. find $\phi = 0.48$ studying very long lattice chains via the PERM algorithm. They suggest larger values of ϕ are caused by strong corrections to scaling. Note that $\phi = 0.48$ gives a negative α , indicating a non-diverging specific heat.

Grassberger, J Phys A: Math Gen **38**, 323 (2005).

Klushin, Polotsky, Hsu, Markelov, Binder, Skvortsov, PRE **87** 022604 (2013).

Summary and Outlook

Adsorption transition for a tethered polymer chain

Findings: Partition function zeros display Lee-Yang behavior.

Finite size scaling (FSS) of zeros locates T_c and determines ϕ .

FSS determines specific heat and order parameter exponents.

To do: Include correction to scaling terms in this analysis.

Carry out same analysis for a continuum chain model.

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Special thanks to Wolfgang Paul's group for their hospitality!

Happy "American" Thanksgiving

