

Magnetic properties of two-color QCD at finite temperature

work with Michael Ilgenfritz, Martin Kalinowski, Michael Muller-Preussker
and Alexander Schreiber

arXive 1203.3360 1310.7876

I. Motivations

II. Formalism

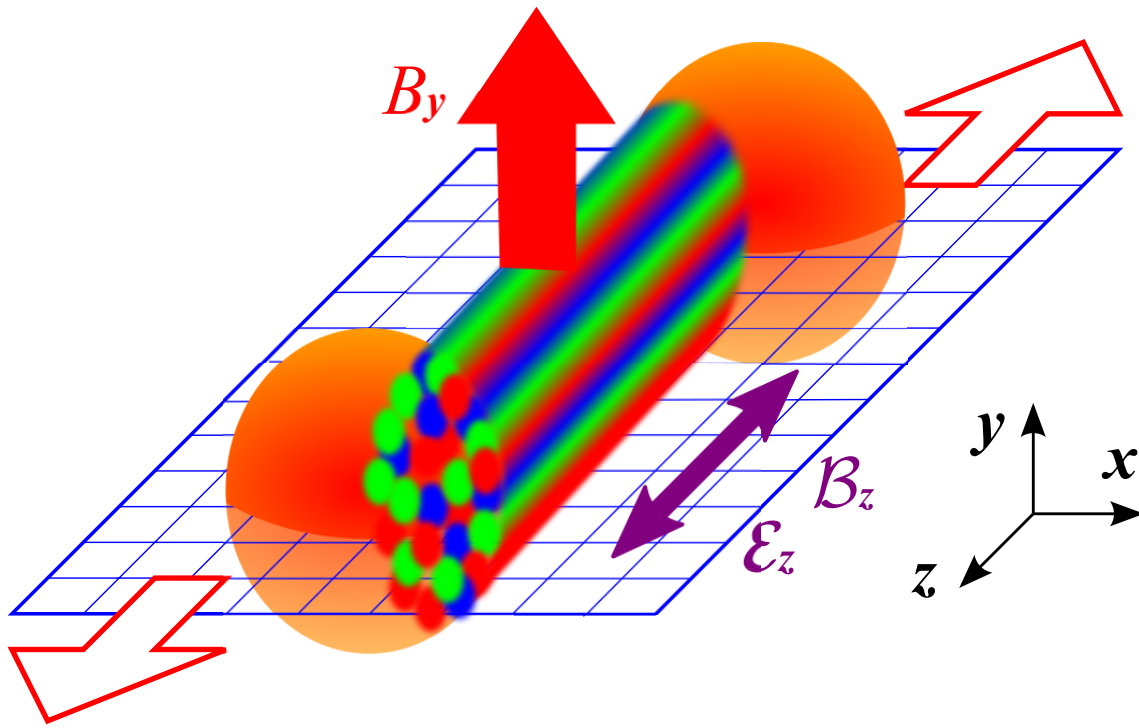
III. Magnetic catalysis

IV. The phase diagram

V. Conclusions

(VI. Computational method)

Motivations: Non-central heavy ion collisions at RHIC and LHC



$$B_y \sim m_\pi^2 \sim 10^{18} \text{ Gauss}$$

$$T \sim m_\pi$$

Furthermore

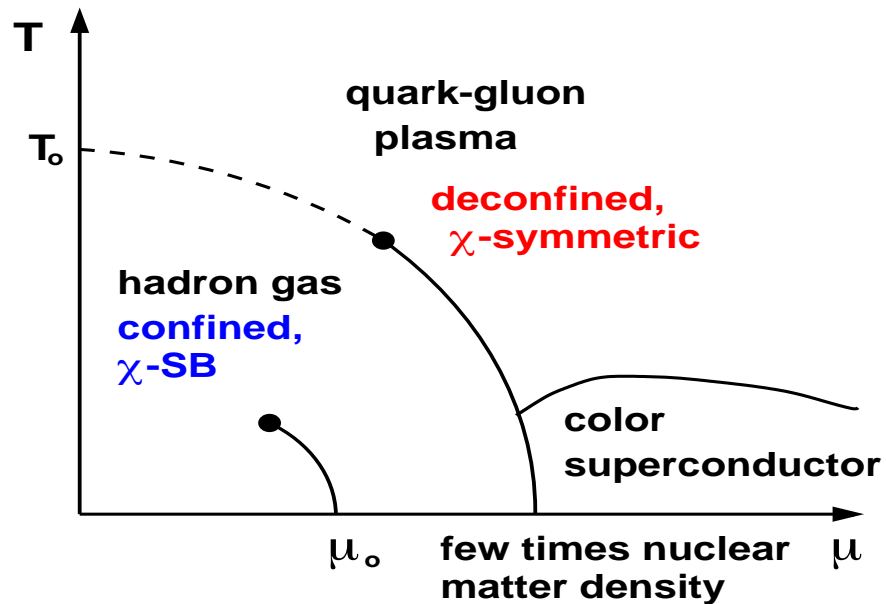
- Very strong magnetic fields in the early universe $\sim 10^{23}$ Gauss
- Strong magnetic fields in compact stars (magnetars)
 $\sim 10^{13}$ Gauss

(Largest man-made magnet = 10^6 Gauss)

Non perturbative methods necessary

- Models
- Lattice simulations

Usually: (T, μ) Phase diagram $(\mu = \text{Baryon chem. pot.})$



This talk: (T, B) Phase Diagram

$\vec{B} = B e_z$ B constant and homogeneous

Lattice formalism

Lattice size: $N_\sigma^3 \times N_\tau \equiv \mathcal{V}$

$$V = (N_\sigma a)^3$$

$$T = \frac{1}{N_\tau a} \quad a = \text{lattice spacing}$$

$$Z(T, B, \mu) = \int dP(U) \quad \langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{O}(U) dP(U)$$

U stands for $U_\mu(x) = e^{iA_\mu(x)} \in SU(N)$

$$dP(U) = \prod dU e^{-\beta S_g(U)} \text{Det}(D(e^{iqBa^2}, e^{\mu a}) + ma) > 0 \text{ only if } \mu = 0 \text{ for } N > 2$$

$\beta =$ bare coupling constant, $ma =$ bare quark mass, $a = a(\beta)$

Our calculations

$$SU(2) \quad \mu = 0 \quad N_f = 4$$

$$32^3 \times N_\tau \quad N_\tau = 10, 8, 6, 4 \quad \beta = 1.8 \quad \text{fixed scale approach}$$

Implementing a constant magnetic field on a periodic lattice \Rightarrow

$$\phi = a^2 q B = 2\pi N_b / N_\sigma^2, \quad N_b \in Z \quad \phi \text{ periodic} \Rightarrow N_b < N_\tau^2 / 4$$

Reasons for our choice

Simpler theory but similar chiral properties as in QCD

Action local on the lattice, some chiral symmetry left

Can be extended to $\mu \neq 0$

Simulations with CUDA Fortran on GPU:s

Chiral order parameter (Exact for $m=0$)

$$a^3 \langle \bar{\chi}\chi \rangle = \frac{1}{4\mathcal{V}} \langle \text{Tr}(D + m)^{-1} \rangle = -\frac{1}{4\mathcal{V}} \frac{\partial}{\partial m} \log(Z)$$

$\langle \bar{\chi}\chi \rangle$ is the chiral condensate

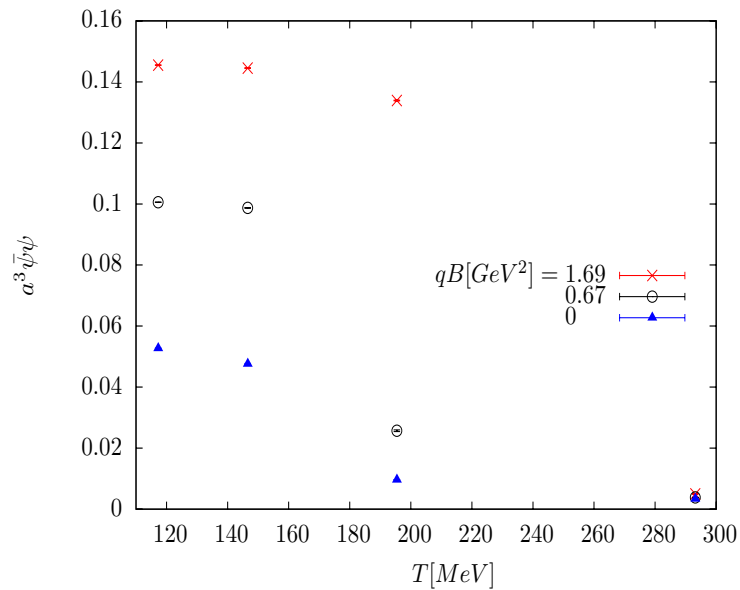
Confinement order parameter (Exact for $m \rightarrow \infty$)

$$\langle L \rangle = \frac{1}{N_\sigma^3} \sum_{\bar{n}} \langle \frac{1}{2} \text{Tr} \prod_{n_4=1}^{N_\tau} U(\bar{n}, n_4) \rangle \sim e^{-\frac{F_q}{T}}$$

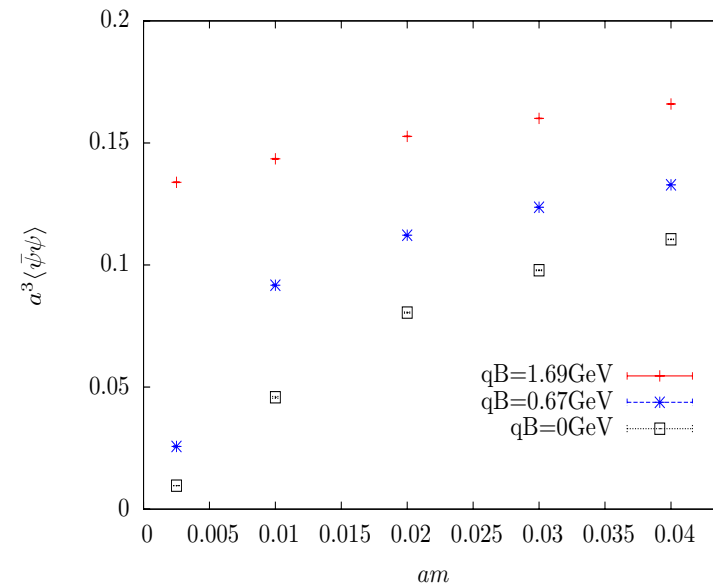
F_q is the free energy of a static quark

Magnetic catalysis: The chiral condensate increases with B

Scale from Static potential $ma = 0.0025 \Rightarrow m_\pi = 175 \text{ MeV}$ at $B = 0$



$$ma = 0.0025$$



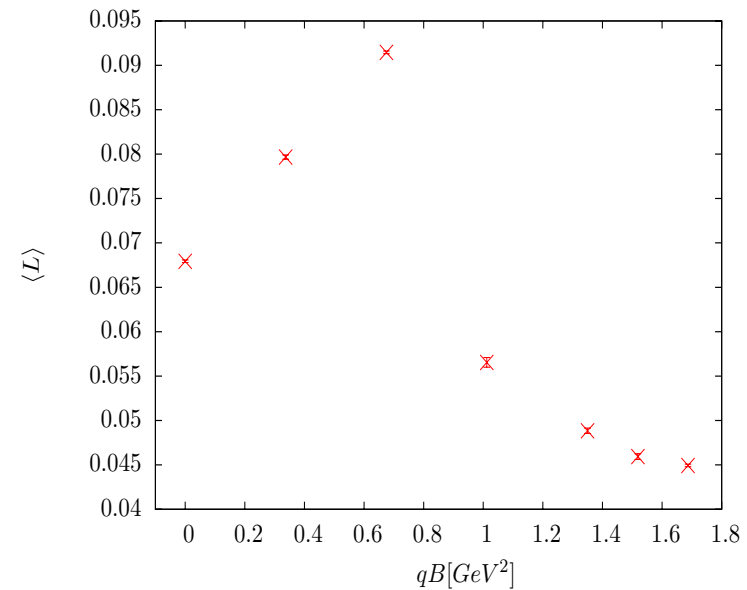
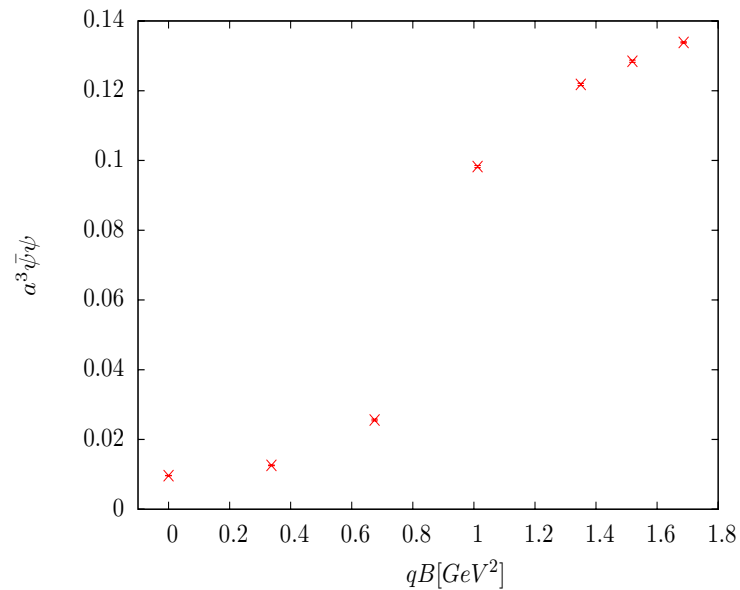
$$T = 193 \text{ MeV}$$

$$T_c(B = 0) \approx m_\pi$$

Magnetic catalysis in the transition region?

(Regensburg-Wuppertal collaboration claims the opposite!)

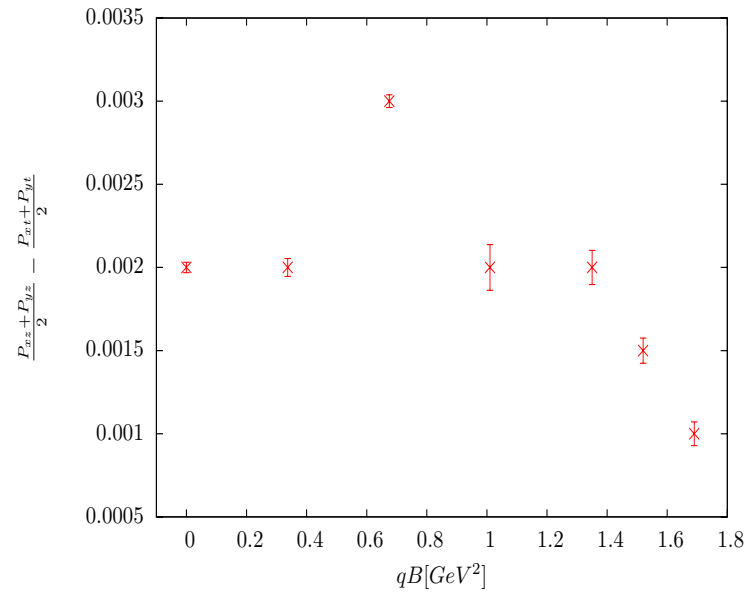
Our data



$$ma = 0.0025$$

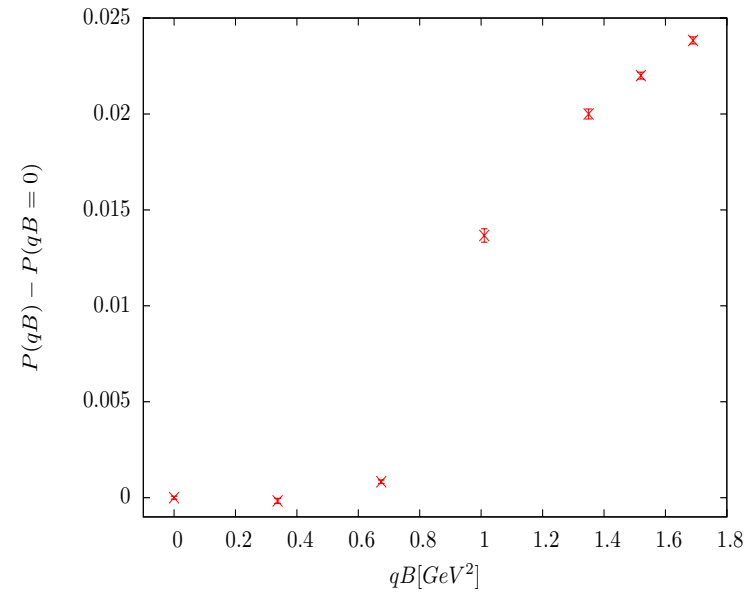
$$T = 193 \text{ MeV}$$

The indirect influence of B on the gluons



Transverse part of the
gluonic energy density

$$ma = 0.0025$$

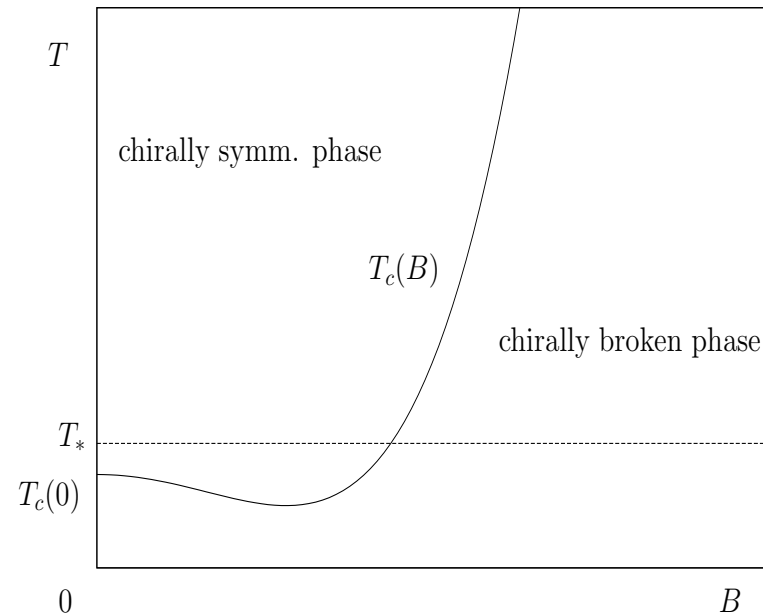


Gluonic part of the
conformal anomaly

$$T = 193 \text{ MeV}$$

Conclusions

Phase Diagram



Next: Simulations nearer T_c , smaller masses, $\mu \neq 0$
cme-effect

$$D + ma \equiv M \quad \text{Det}M = \text{Det}M^\dagger \quad M^\dagger M > 0$$

$$dP = e^{-\beta S_g} (\text{Det}M)^2 = e^{-\beta S_g} (\text{Det}M^\dagger M) = d\phi^\dagger d\phi dP e^{-\beta S_g + \phi^\dagger (M^\dagger M)^{-1} \phi}$$

$$(M^\dagger M)_{eo} = 0 \quad \text{choose} \quad (M^\dagger M)_{ee}$$

Molecular dynamics:

$$dP dp d\phi^\dagger d\phi e^{-H} \quad H = \sum_n p_\mu^2(n) + S_{eff}(U, \phi)$$

Discretized Eq. of motion, Omelyan integration

$$\text{Solve } (M^\dagger M)\eta = \phi \quad M^\dagger M \text{ sparse} \quad ma < \lambda < 1 \Rightarrow$$

Precond. Conj. Gradient on GPU

At end of trajectory Metropolis update \Rightarrow **Algorithm exact**