

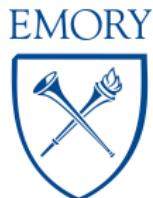
Renormalization Group for Quantum Walks

CompPhys13

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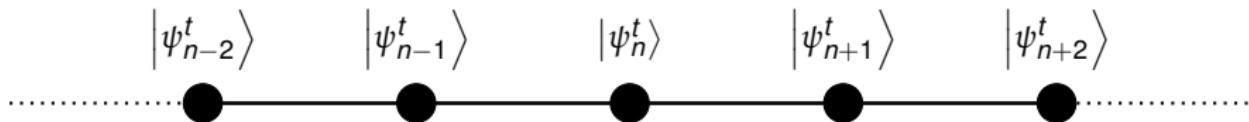


arXiv:1311.3369

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Discrete time walks in one dimension

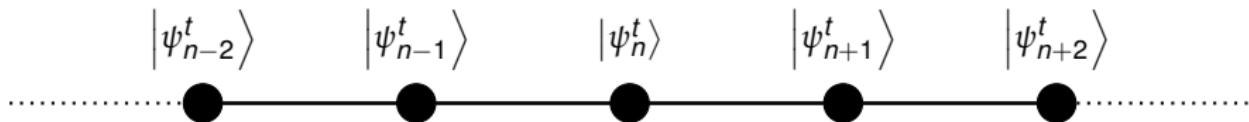


Discrete Time Evolution

$$|\psi^{t+1}\rangle = \mathcal{U} \cdot |\psi^t\rangle$$

$$\mathcal{U} = \sum_n \mathcal{P} \cdot |n\rangle \langle n-1| + \mathcal{Q} \cdot |n\rangle \langle n+1|$$

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Classical

$$p_n^t = \||\psi_n^t\rangle\|_1$$

$\mathcal{P} + \mathcal{Q}$ stochastic

Discrete time walks in one dimension



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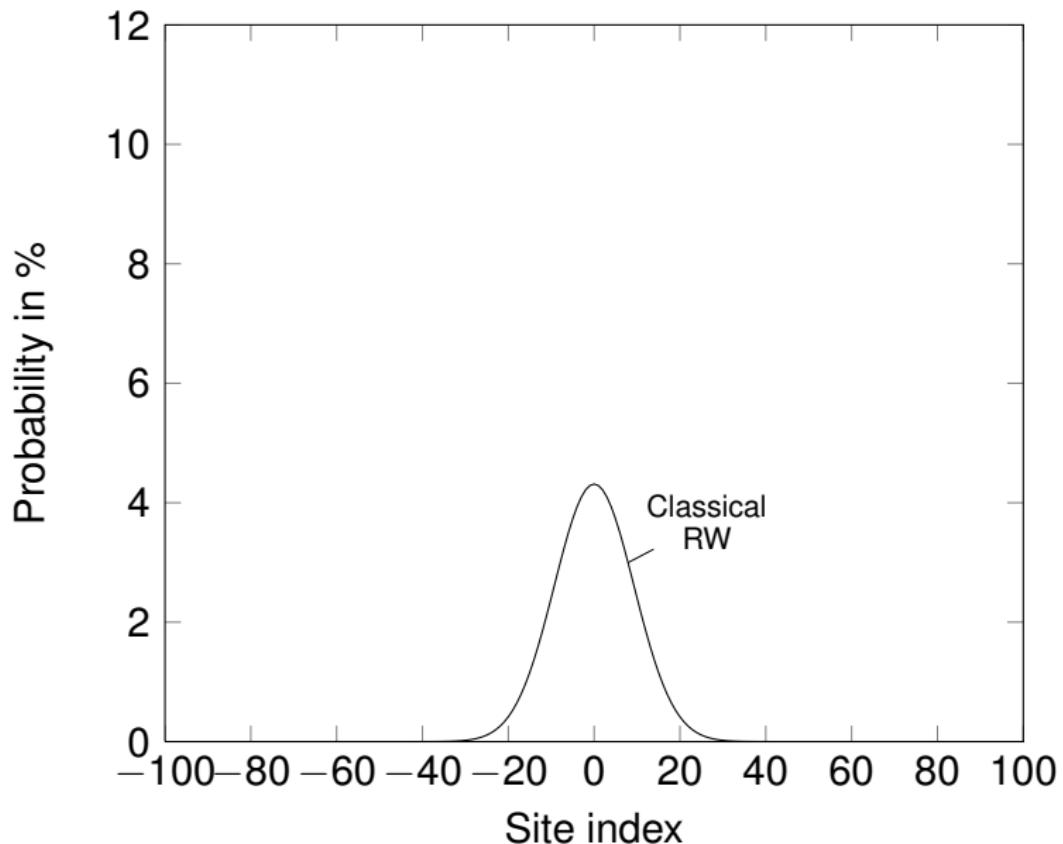
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Quantum Mechanical

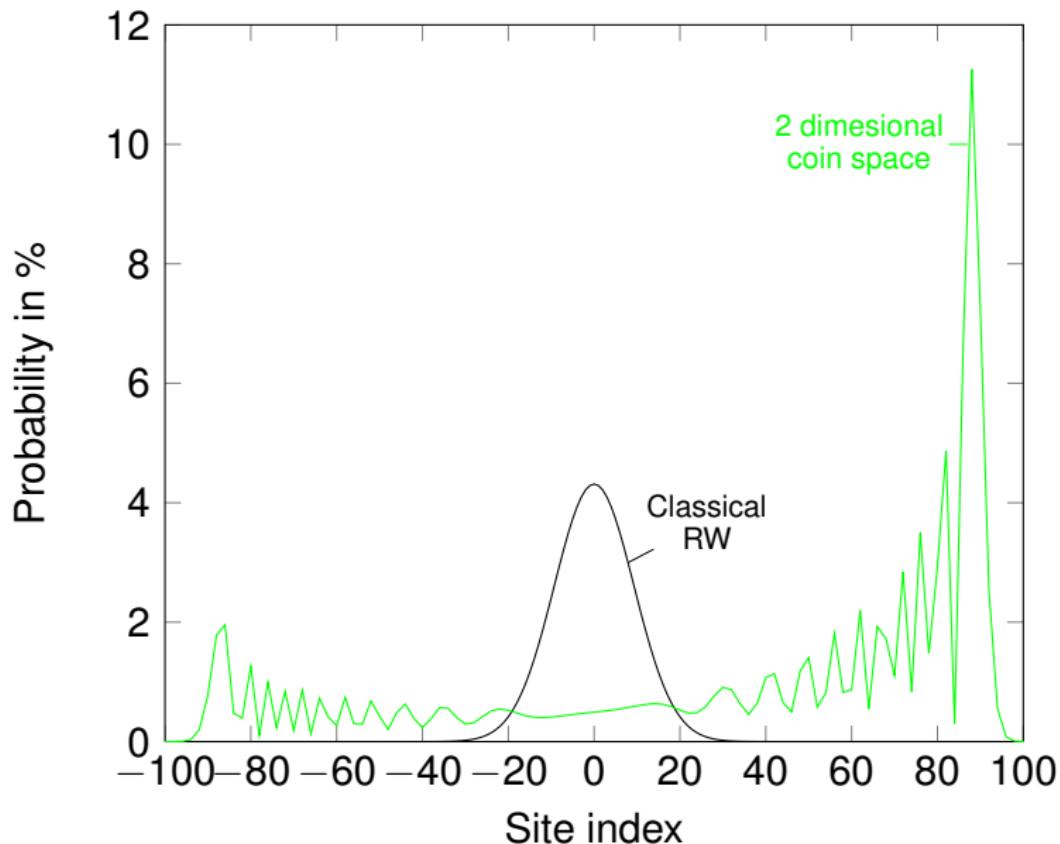
$$p_n^t = \langle \psi_n^t | \psi_n^t \rangle \quad \mathcal{U} = \mathcal{S} \cdot (\mathbb{1} \otimes \mathcal{C})$$

$$\mathcal{P}\mathcal{P}^\dagger + \mathcal{Q}\mathcal{Q}^\dagger = \mathbb{1} \quad \mathcal{P}\mathcal{Q}^\dagger = \mathcal{Q}\mathcal{P}^\dagger = 0$$

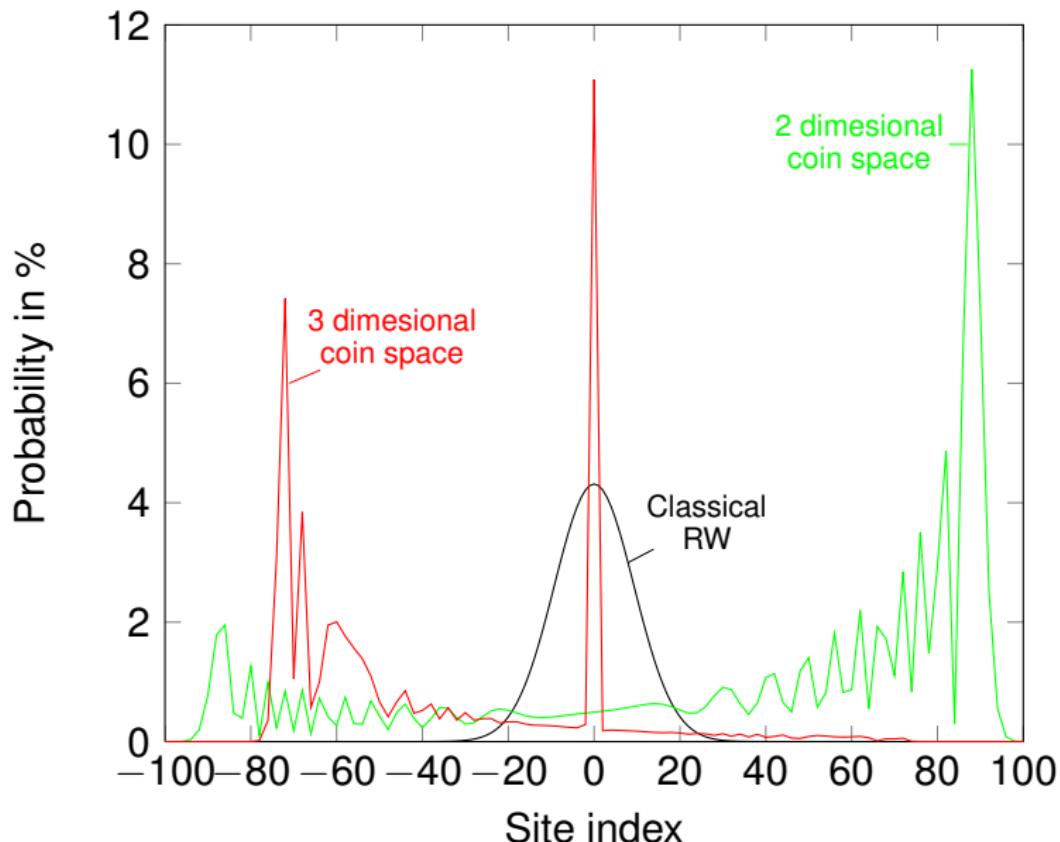
One dimensional walks



One dimensional walks



One dimensional walks



Motivation

- model for quantum transport (Aharonov PRA 1993)
- Grover's search algorithm (Grover PRL 1997)
find a marked vertex in a graph in $T \sim \mathcal{O}(\sqrt{N})$
- more general quantum search algorithm
- general quantum computation (Lovett PRA 2010)
- generally faster spreading compared to the random walk

$$\langle r^2 \rangle = t^{\frac{2}{d_w}} \quad \text{with } d_w^{QW} \stackrel{?}{=} \frac{1}{2} d_w^{RW}$$

- analytic results only on translational invariant lattice

The generating function

The discrete Laplace transform/ The z -Transform

The generating function and its backtransform

$$|\tilde{\psi}(z)\rangle = \sum_{t=0}^{\infty} z^t \cdot |\psi^t\rangle \quad |\psi^t\rangle = \frac{1}{2\pi i} \oint_{|z|=1} z^{-t-1} \cdot |\tilde{\psi}(z)\rangle dz$$

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$$|\psi^{t+1}\rangle = \mathcal{U} \cdot |\psi^t\rangle \quad \implies \quad |\tilde{\psi}(z)\rangle = \tilde{\mathcal{U}}(z) \cdot |\tilde{\psi}(z)\rangle + |\psi^0\rangle$$

Recurrence equations become algebraic equations.

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Generating function for p_n^t in quantum walks

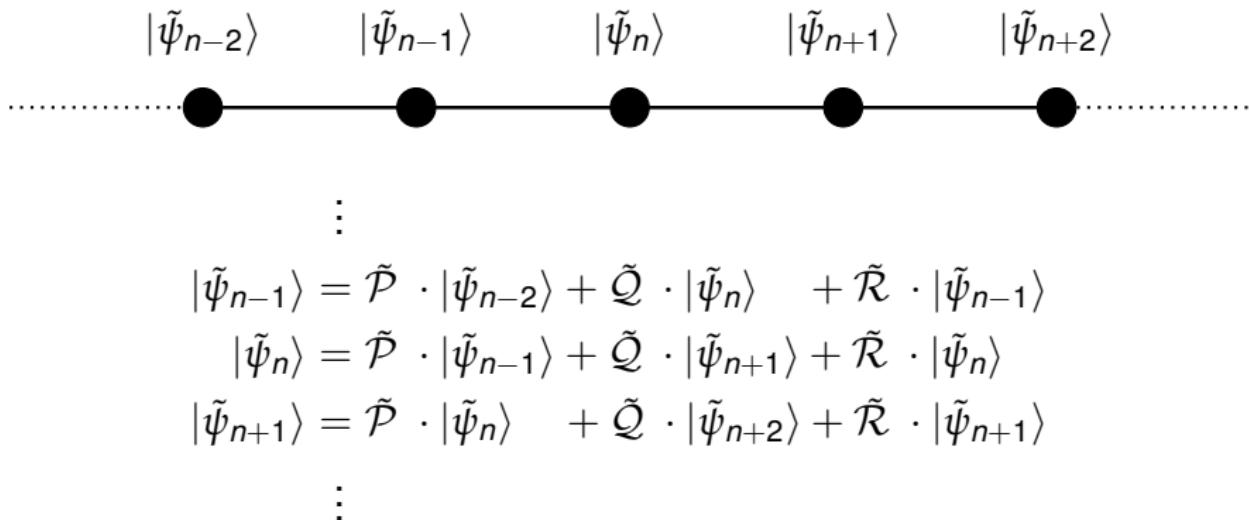
$$\tilde{p}_n(z) = \sum_{t_0}^{\infty} z^t \cdot p_n^t = \frac{1}{2\pi i} \oint \langle \tilde{\psi}(z/y) | \tilde{\psi}(y) \rangle \frac{dy}{y}$$

Renormalization group in one dimension

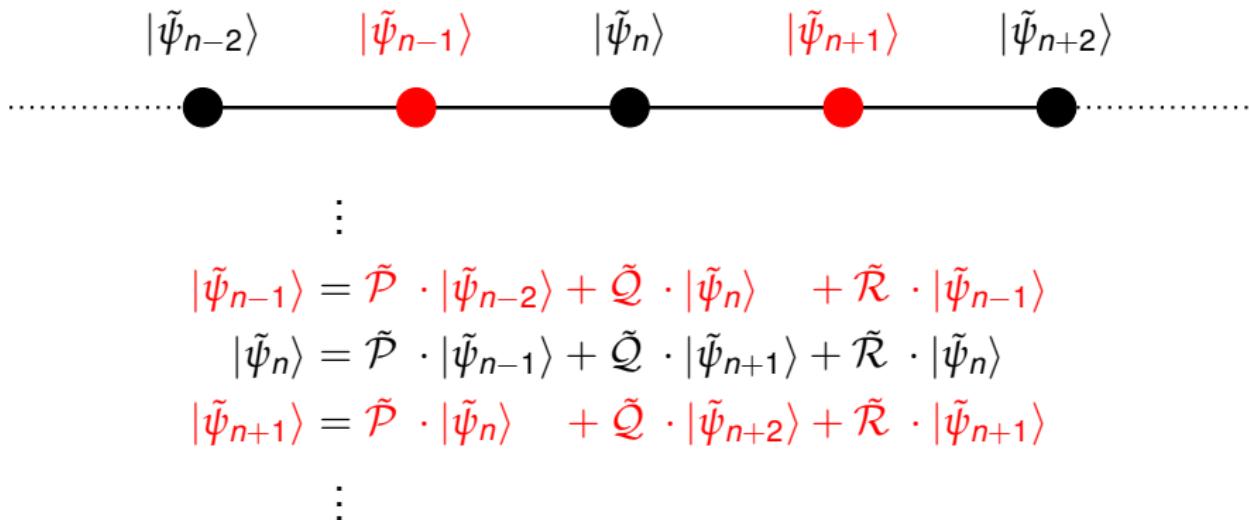
$|\tilde{\psi}_{n-2}\rangle$ $|\tilde{\psi}_{n-1}\rangle$ $|\tilde{\psi}_n\rangle$ $|\tilde{\psi}_{n+1}\rangle$ $|\tilde{\psi}_{n+2}\rangle$



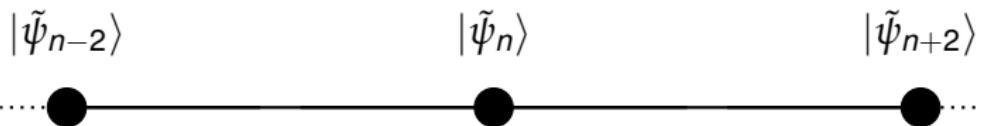
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$$|\tilde{\psi}_n\rangle = \tilde{\mathcal{P}}' \cdot |\tilde{\psi}_{n-2}\rangle + \tilde{\mathcal{Q}}' \cdot |\tilde{\psi}_{n+2}\rangle + \tilde{\mathcal{R}}' \cdot |\tilde{\psi}_n\rangle$$

The one dimensional classical random walk

- the recursion equations:

$$p' = \frac{p^2}{1-r} \quad q' = \frac{q^2}{1-r} \quad r' = r + \frac{2pq}{1-r}$$

- three fixed points:

$$(p^*, q^*, r^*) = (1 - r^\infty, 0, r^\infty) \quad t_k \sim L_k^{d_w} \quad \text{with} \quad d_w = 1$$

$$(p^*, q^*, r^*) = (0, 1 - r^\infty, r^\infty) \quad d_w = 1$$

$$(p^*, q^*, r^*) = (0, 0, 1) \quad d_w = 2$$

The one dimensional quantum walk

- the recursion (matrix) equations:

$$\mathcal{P}' = \mathcal{P} \cdot (\mathbb{1} - \mathcal{R})^{-1} \cdot \mathcal{P} \quad \mathcal{Q}' = \mathcal{Q} \cdot (\mathbb{1} - \mathcal{R})^{-1} \cdot \mathcal{Q}$$

$$\mathcal{R}' = \mathcal{R} + \mathcal{P} \cdot (\mathbb{1} - \mathcal{R})^{-1} \cdot \mathcal{Q} + \mathcal{Q} \cdot (\mathbb{1} - \mathcal{R})^{-1} \cdot \mathcal{P}$$

- three different regimes:
 - a non-trivial fixed point (despite unitarity)!
 - quasi periodic/chaotic behavior of the parameters
 - subspace of trivial fixed points: $\mathcal{P} = \mathcal{Q} = 0$
- we need a generalization of the Jacobian

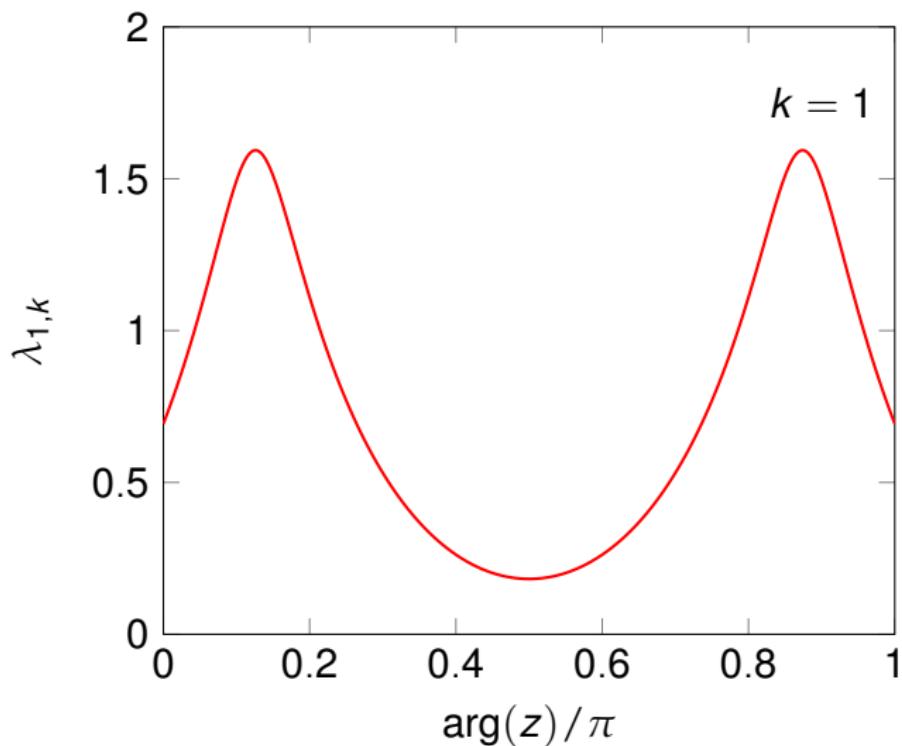
The Lyapunov exponents

- given a dynamical system evolving according to

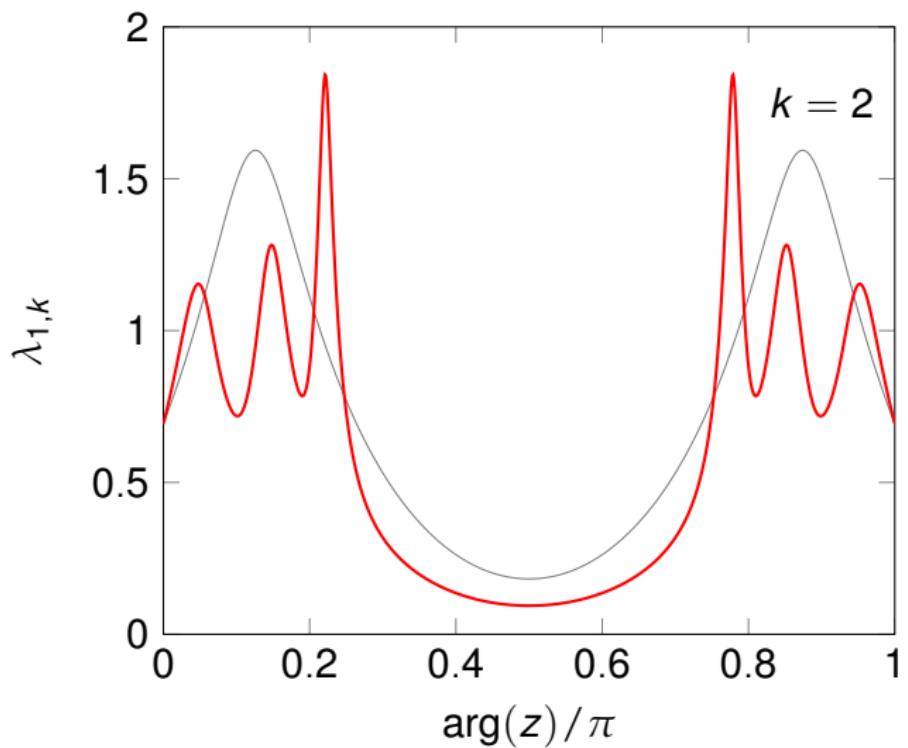
$$\underline{x}(t+1) = F(\underline{x}(t))$$

- J_t - Jacobian at $x(t)$, i.e. $F(x(t) + \delta x) \approx x(t+1) + J_t \cdot \delta x$
- $\mathcal{J}_t = \prod_{t'=0}^t J_{t'}$ - relates the change in $x(t+1)$ when changing $x(0)$
- $\{e^{\lambda_{i,t}}\}$ - eigenvalue spectrum of $\Lambda = (\mathcal{J}_t^\dagger \cdot \mathcal{J}_t)^{\frac{1}{2(t+1)}}$
- $\lambda_{1,\infty}$ - (largest) Lyapunov exponent dominates long term behavior

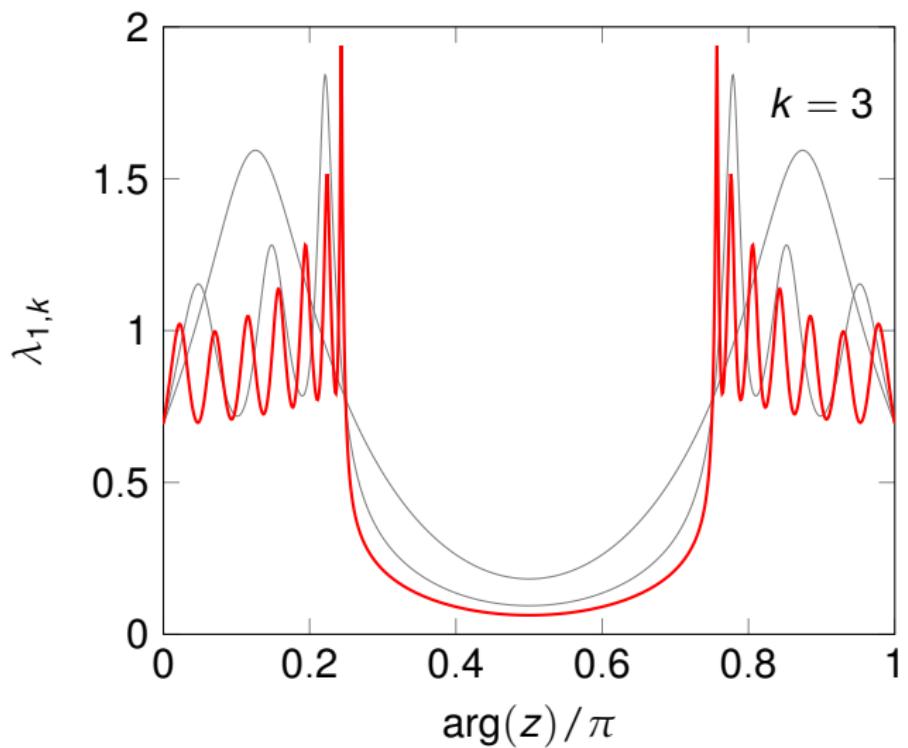
Lyapunov exponent for the line



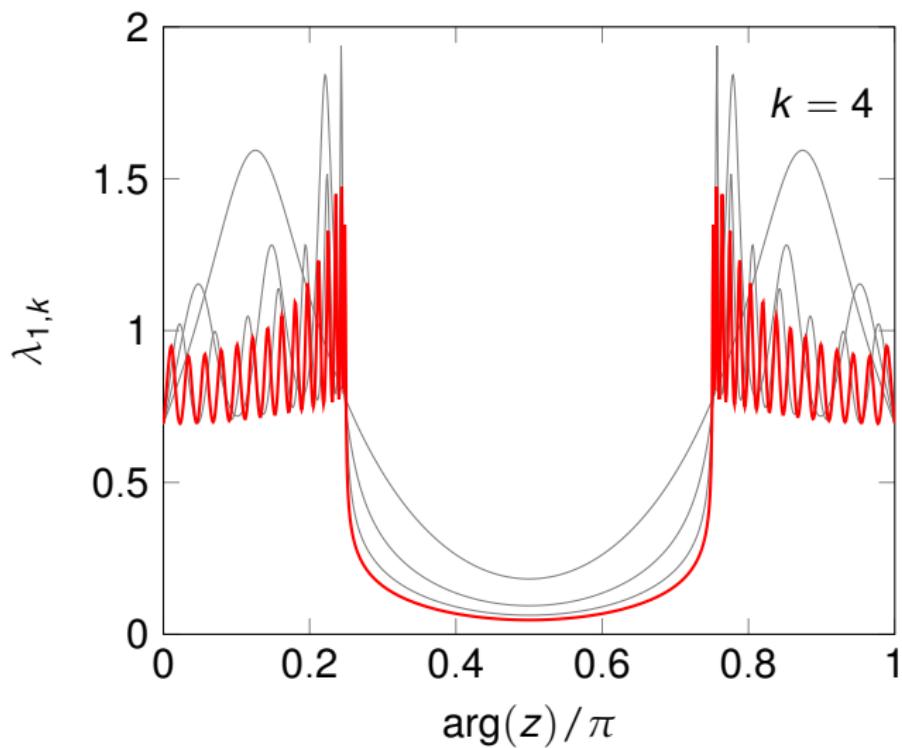
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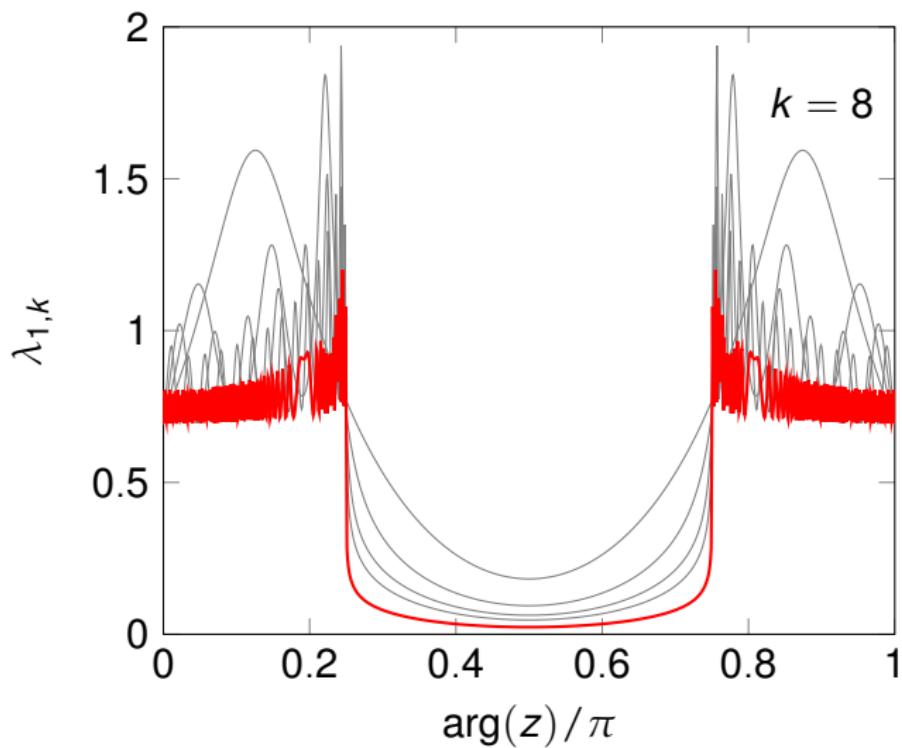
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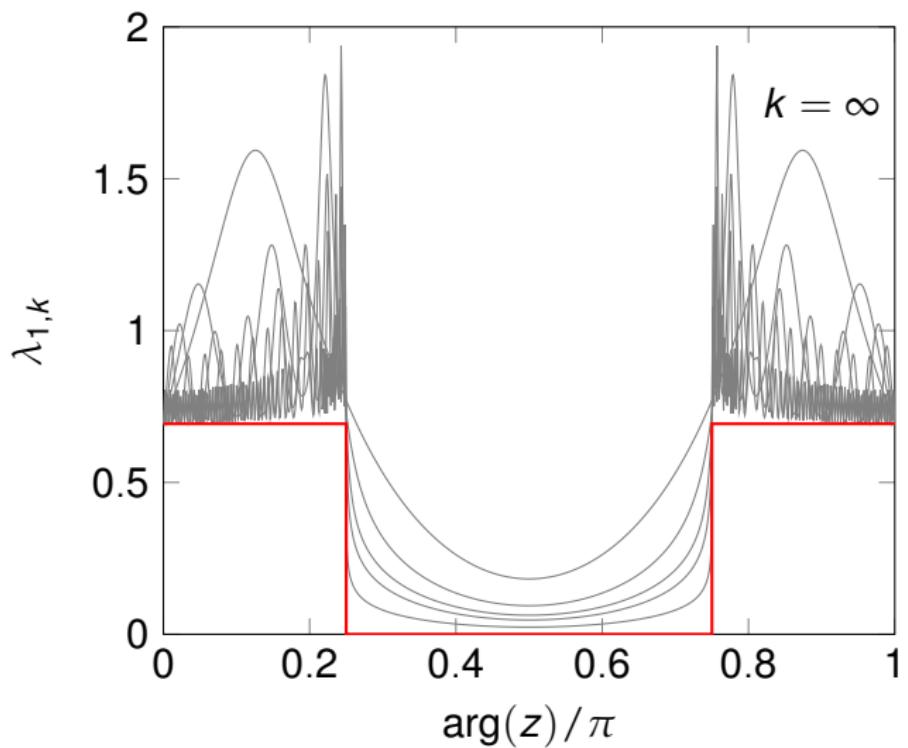
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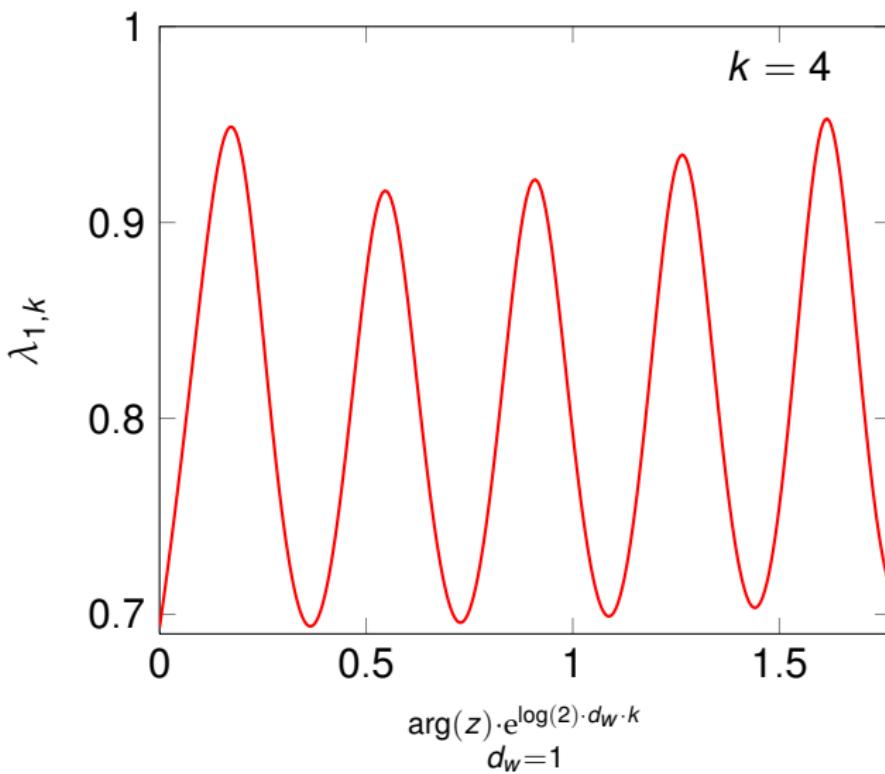
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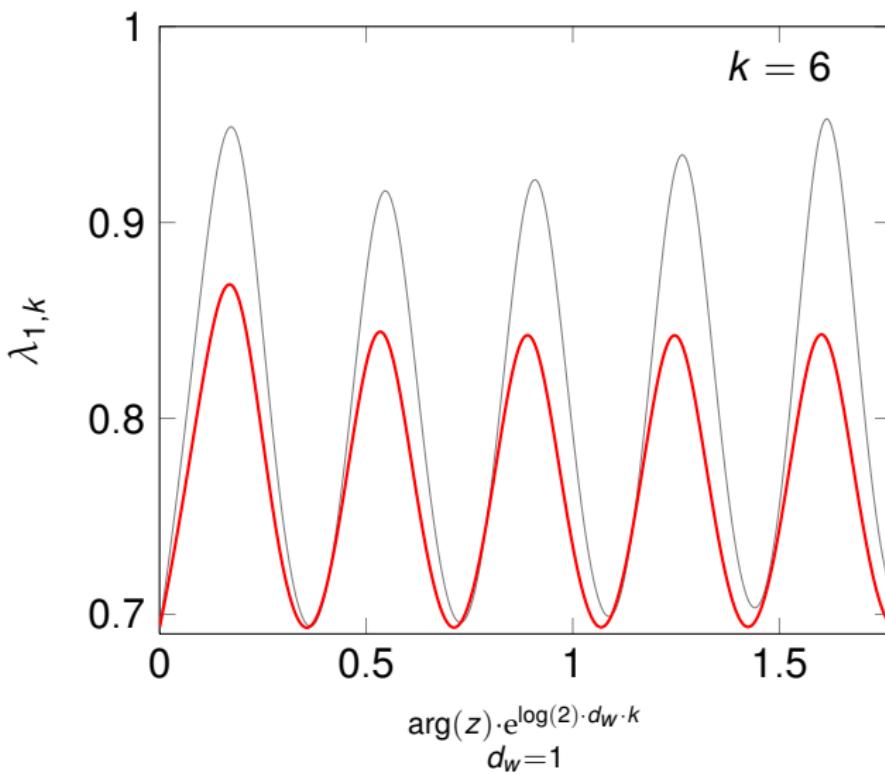
Lyapunov exponent for the line



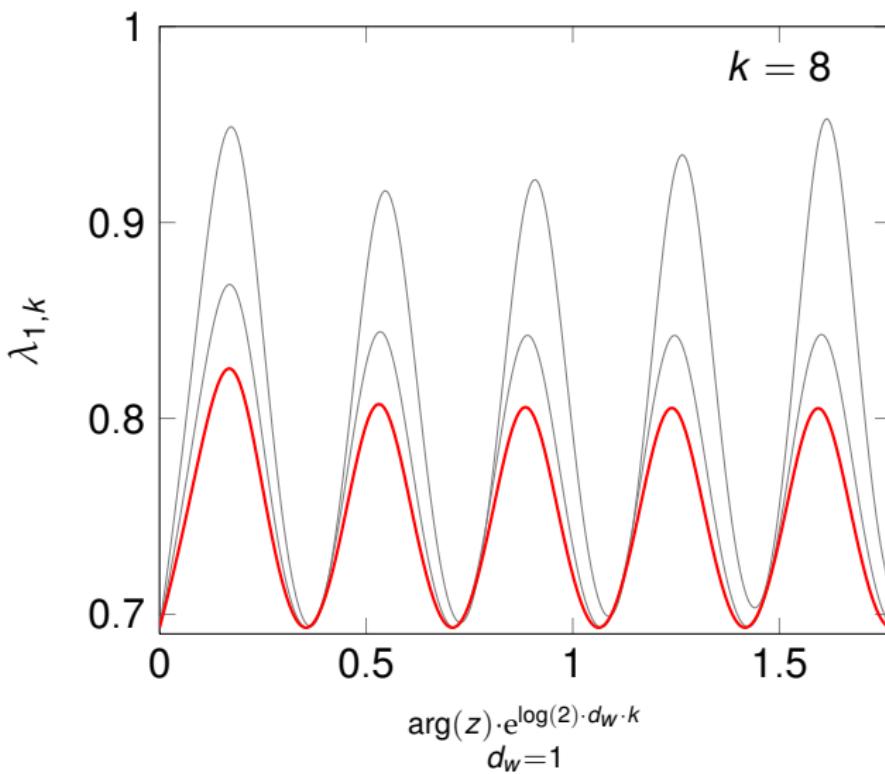
Scaling of the Lyapunov exponent for the line



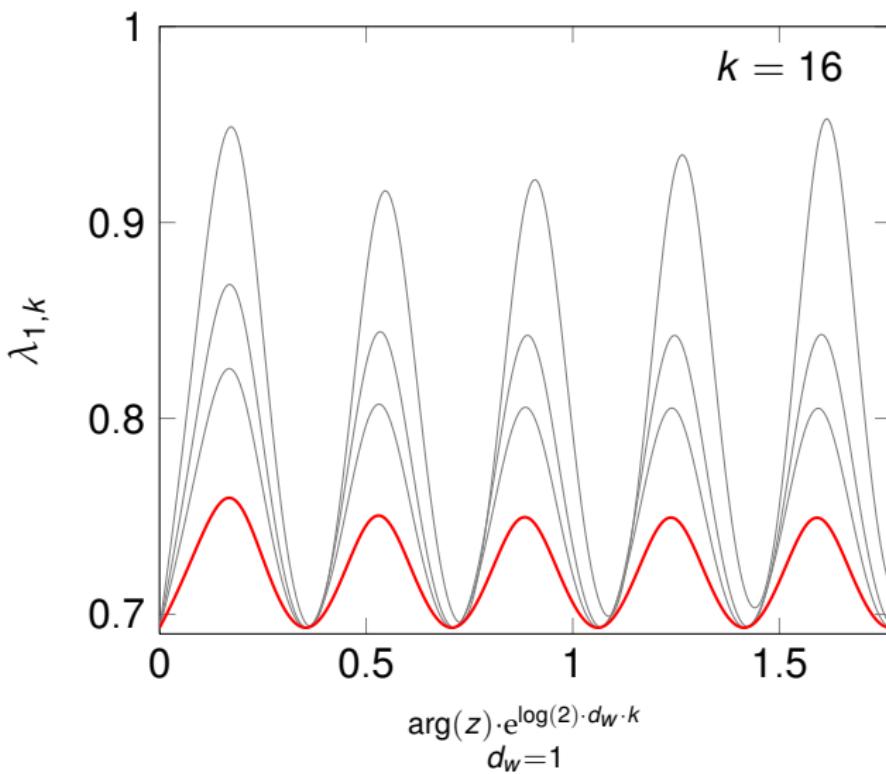
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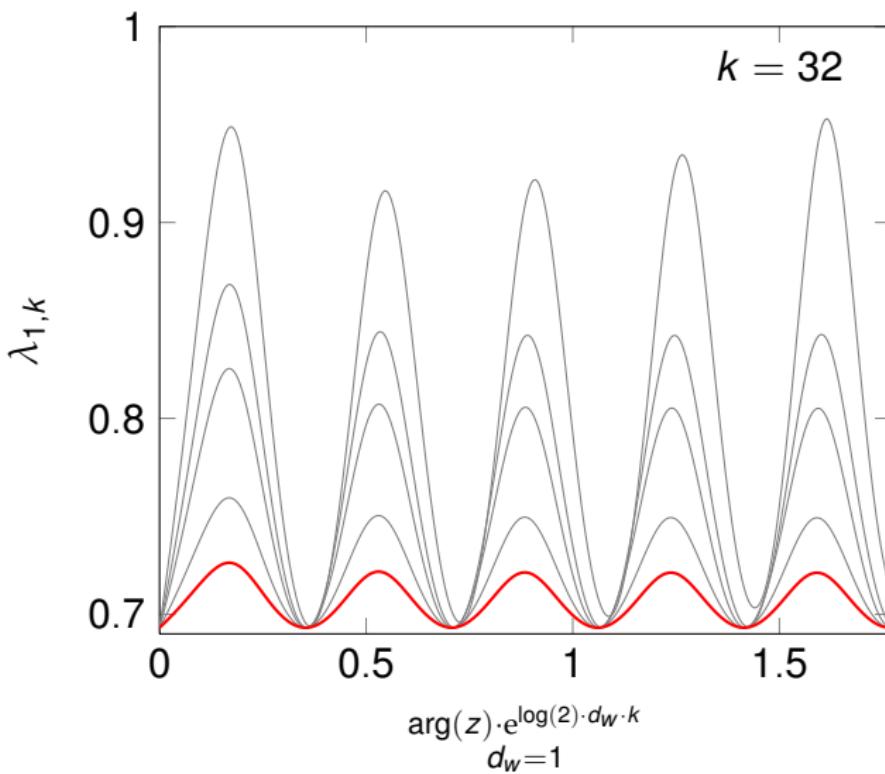
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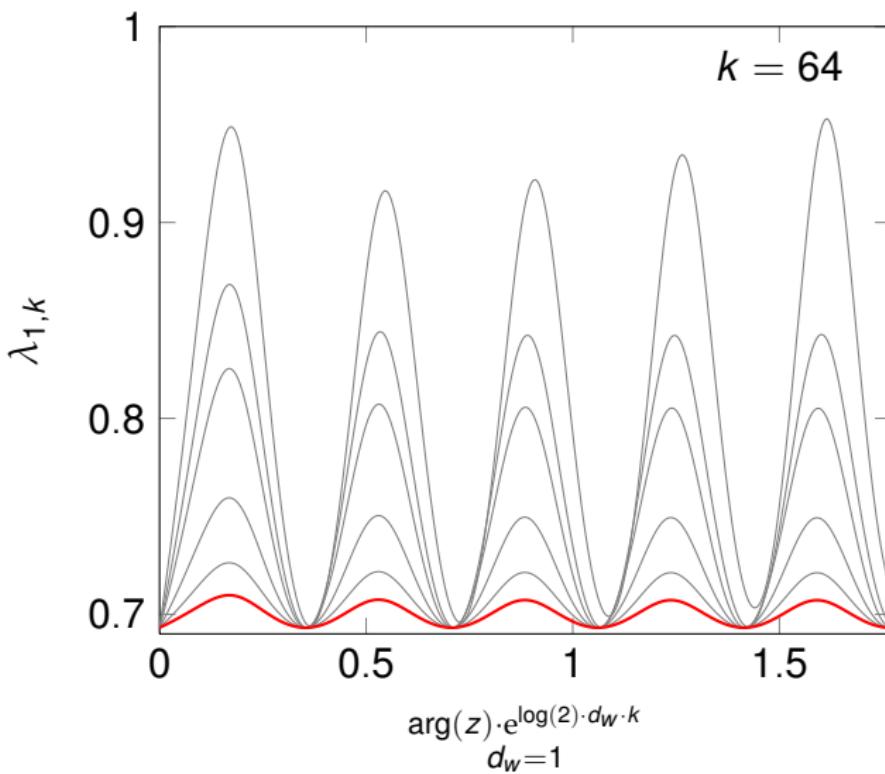
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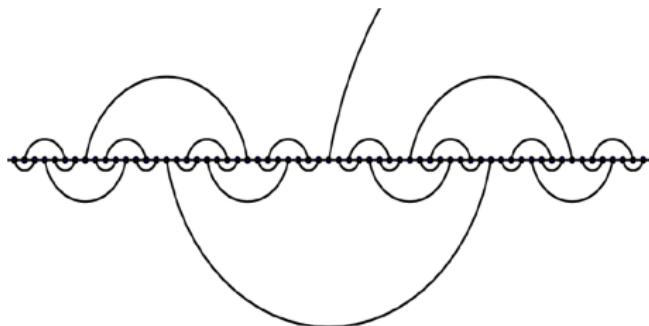


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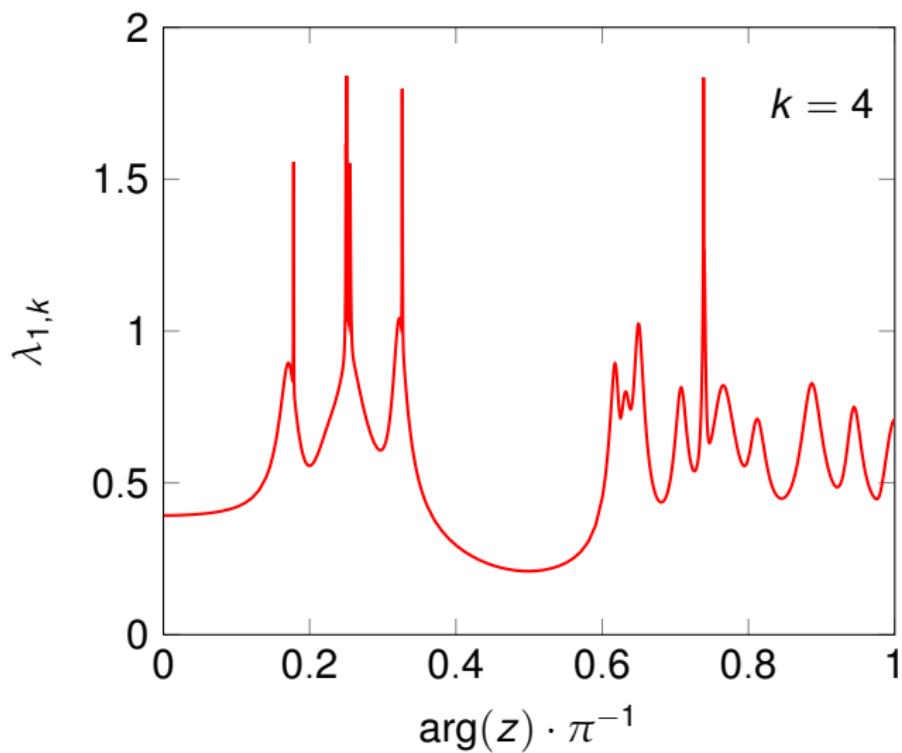


3 regular Hanoi network HN3

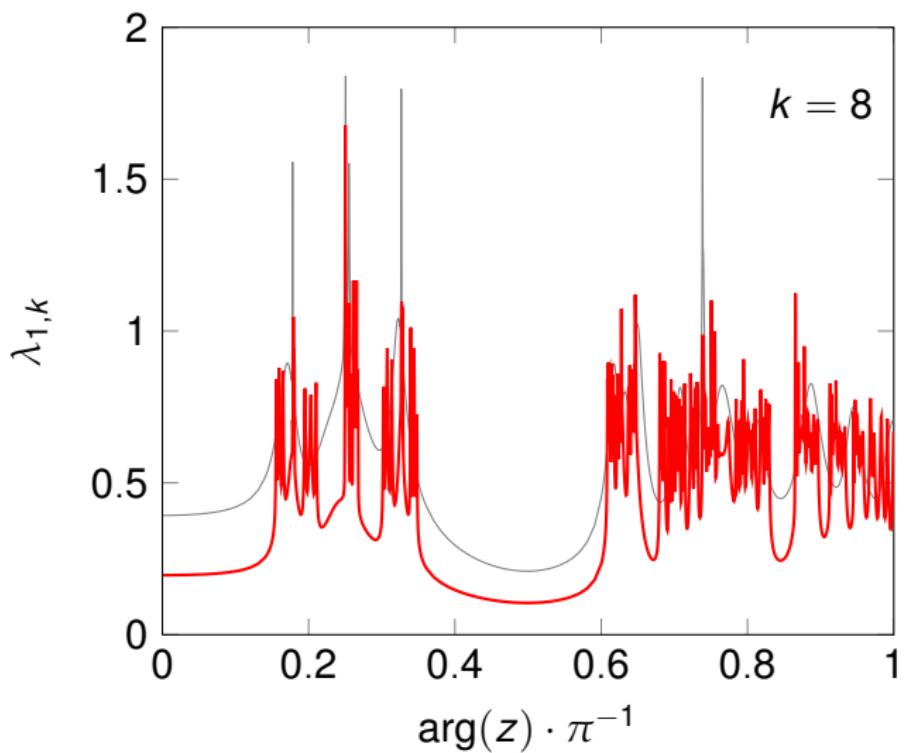
- hierarchical network
(Boettcher EPL 2008)
- exact renormalization
- interpolates between lattices and small-world networks
- $d_w^{RW} = 2 - \log_2(\phi) = 1.30576\dots$



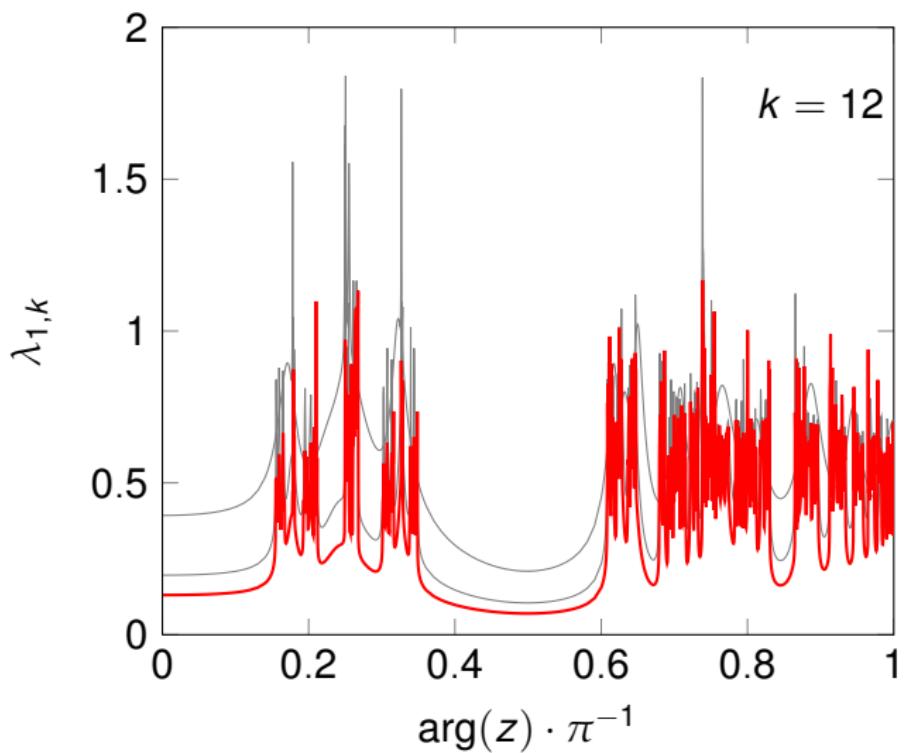
Lyapunov exponent for HN3



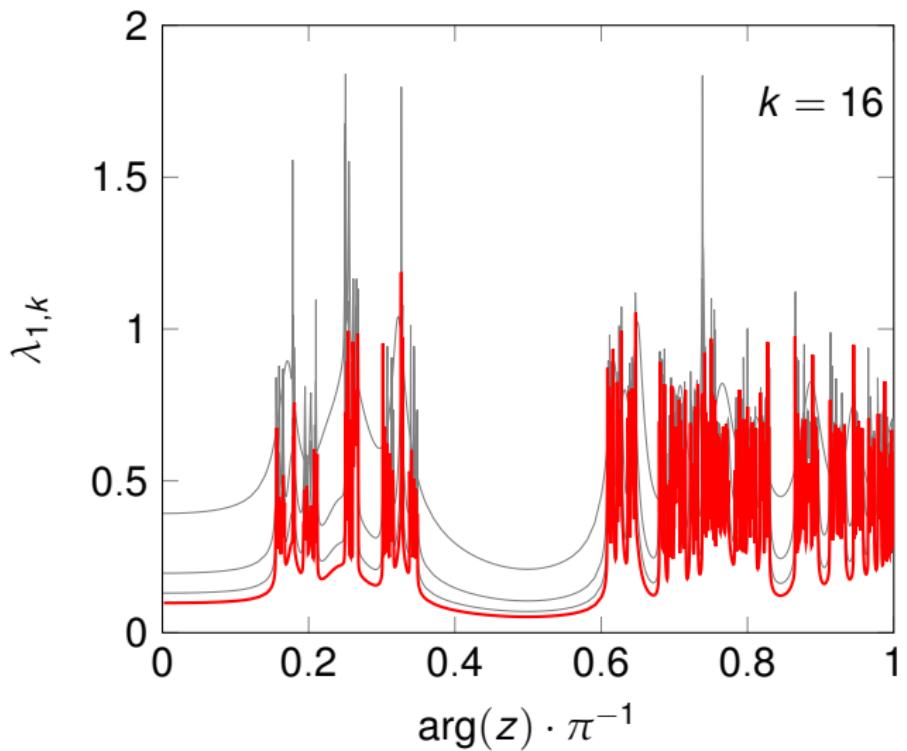
Lyapunov exponent for HN3



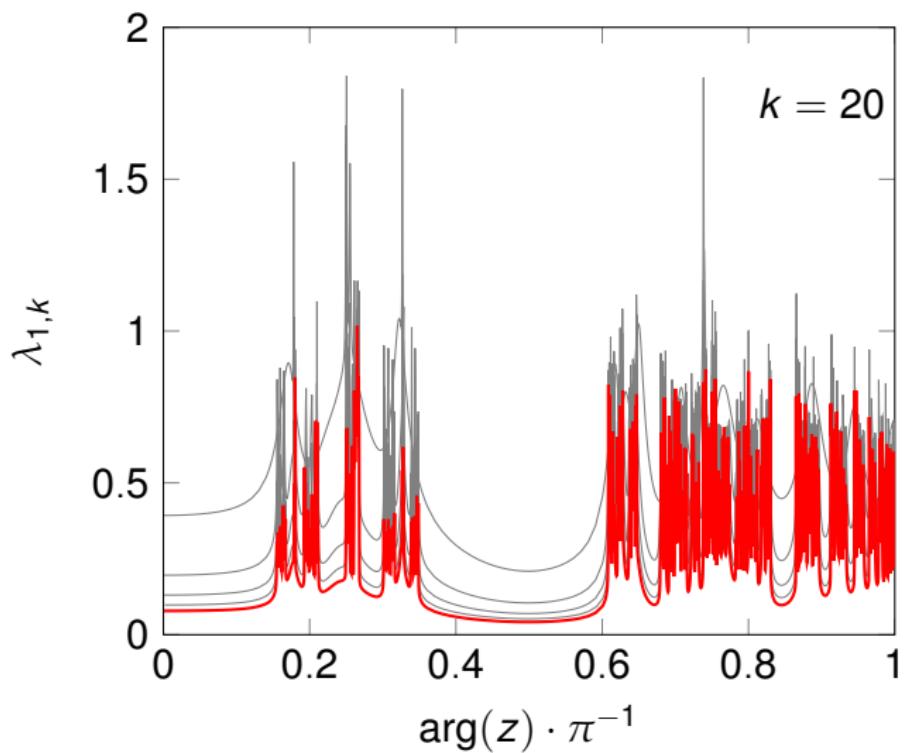
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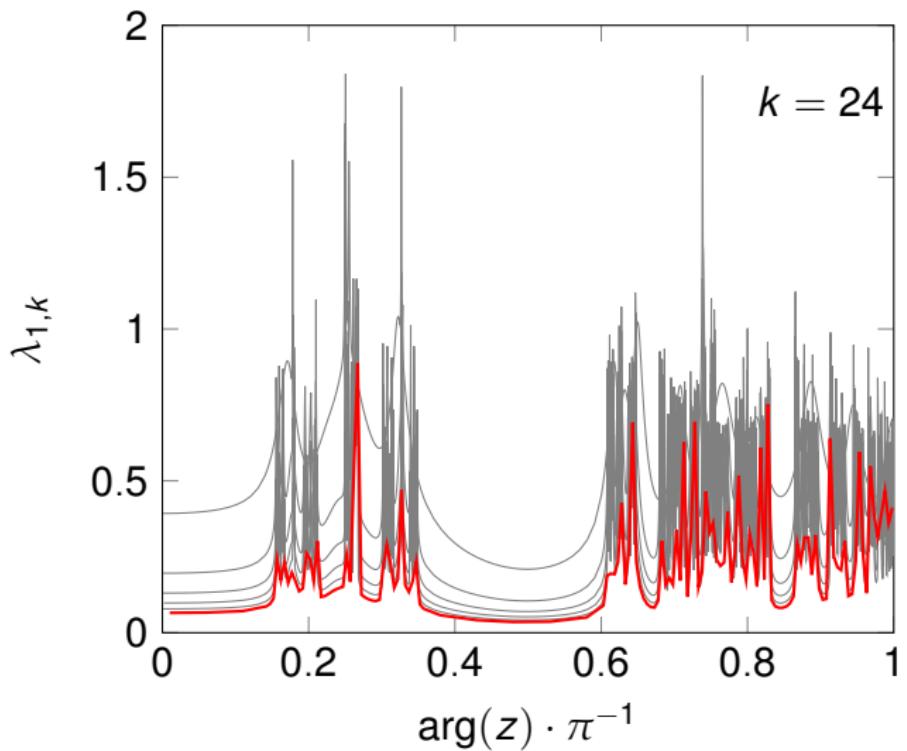
Lyapunov exponent for HN3



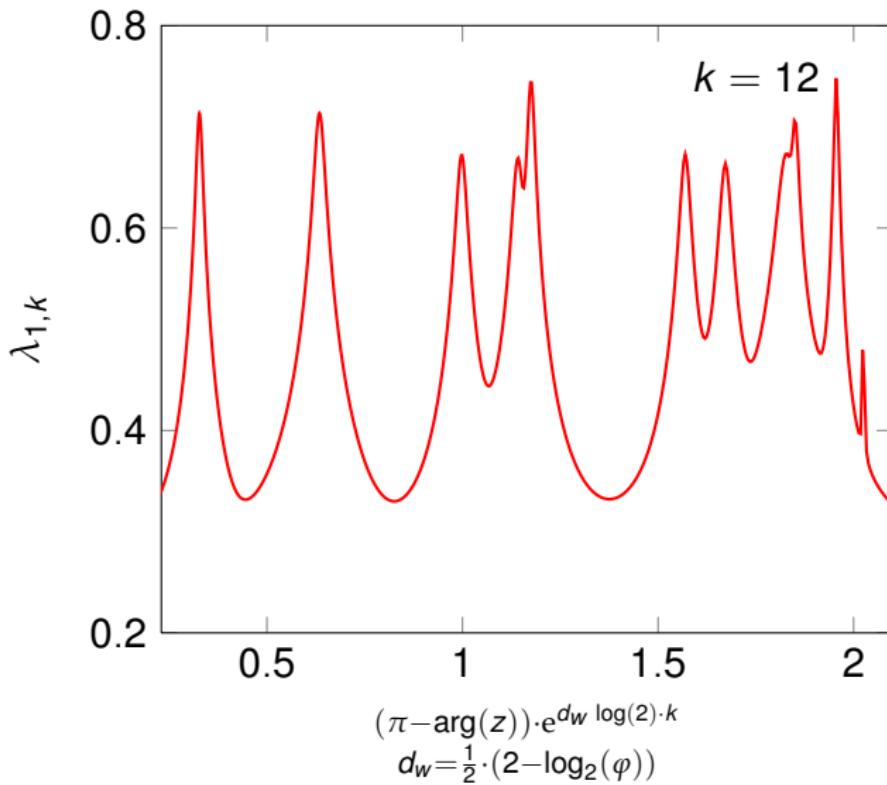
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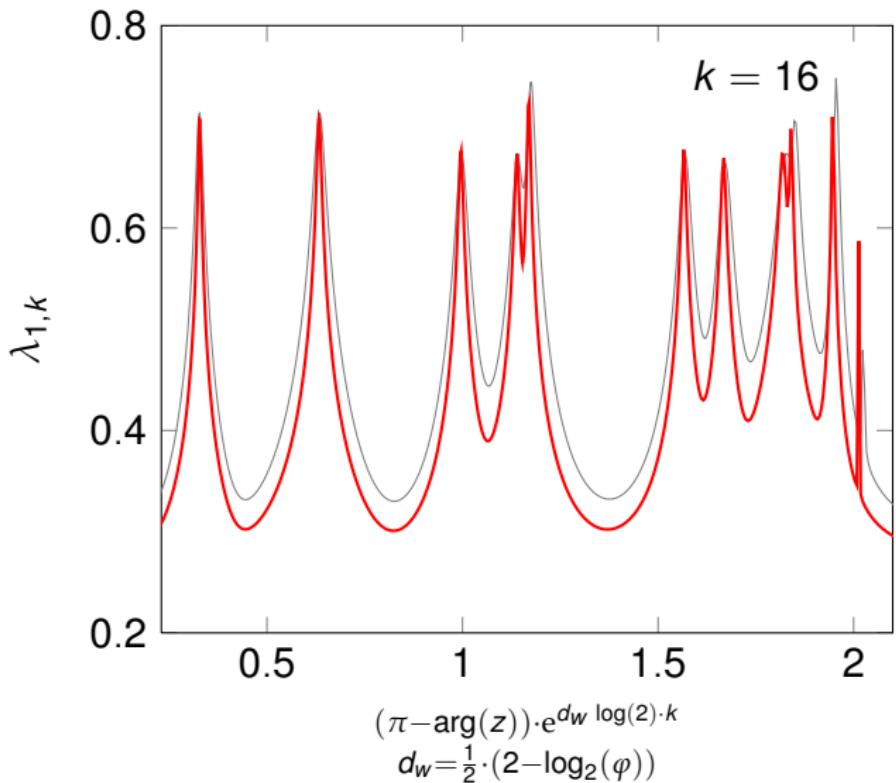
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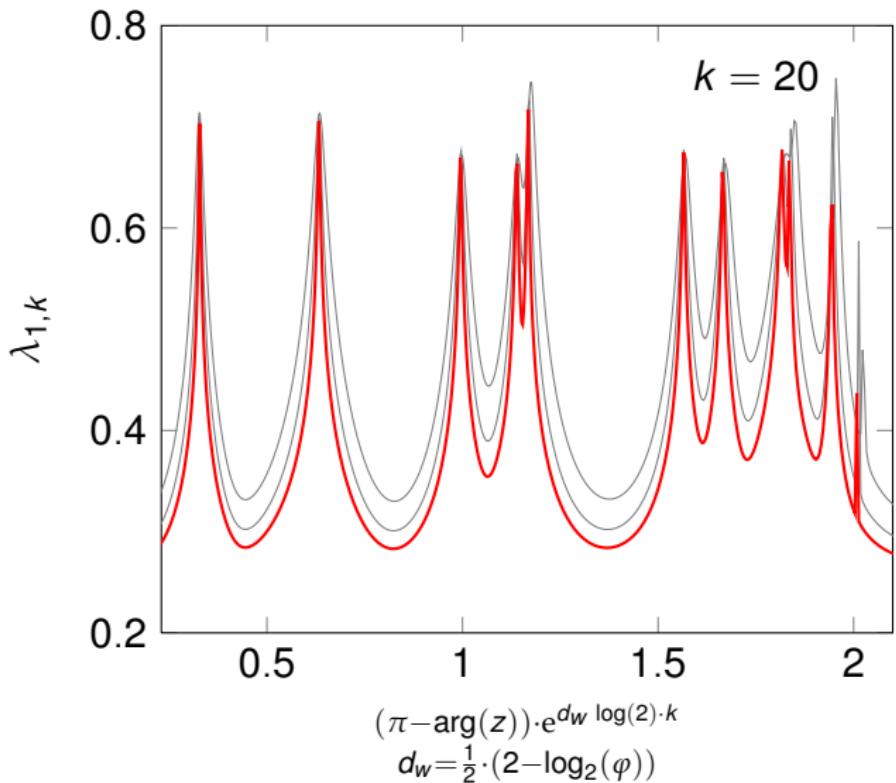
Scaling of the Lyapunov exponent for HN3



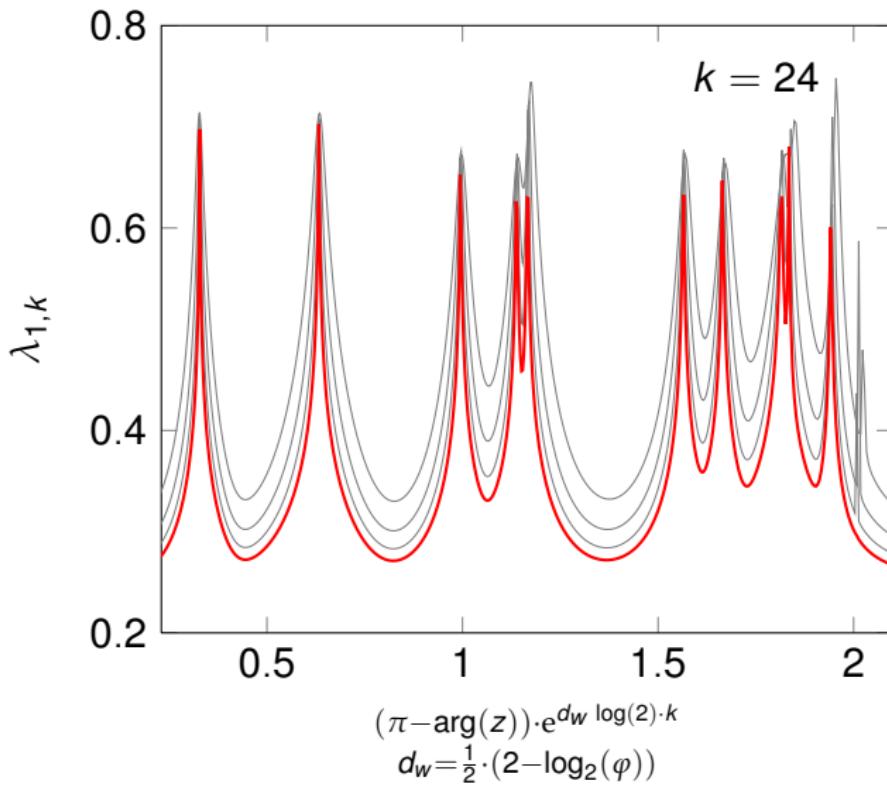
Scaling of the Lyapunov exponent for HN3



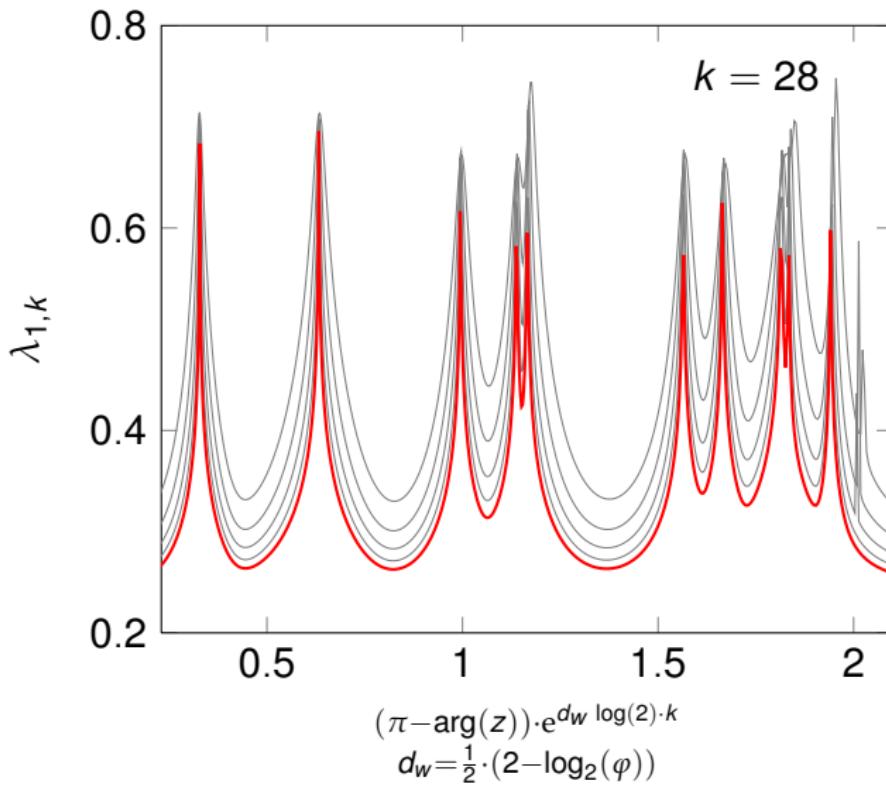
Scaling of the Lyapunov exponent for HN3



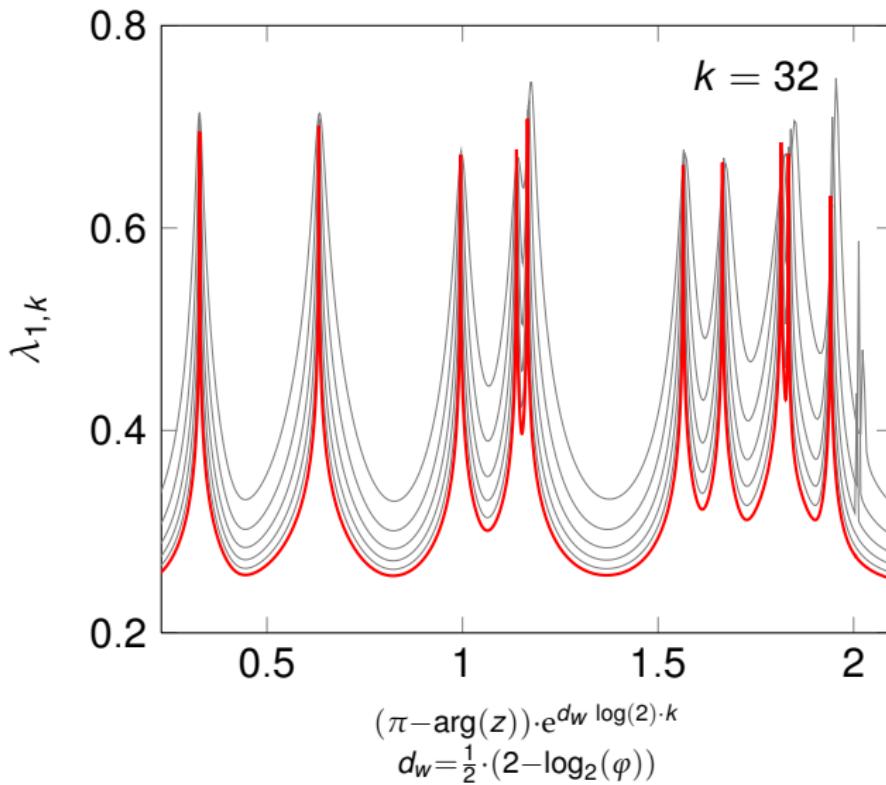
Scaling of the Lyapunov exponent for HN3



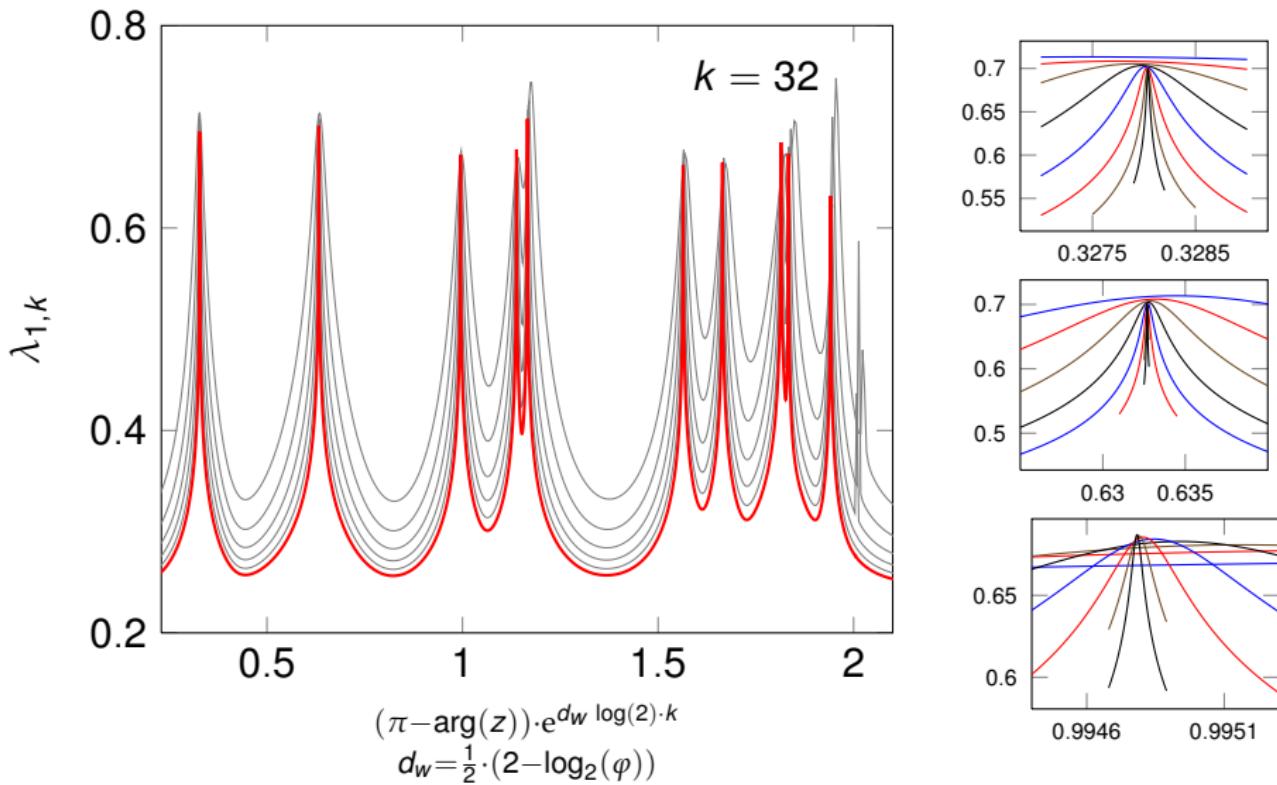
Scaling of the Lyapunov exponent for HN3



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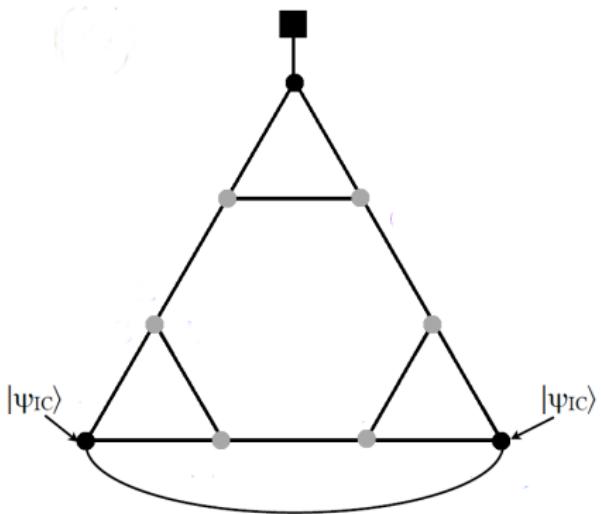


Scaling of the Lyapunov exponent for HN3

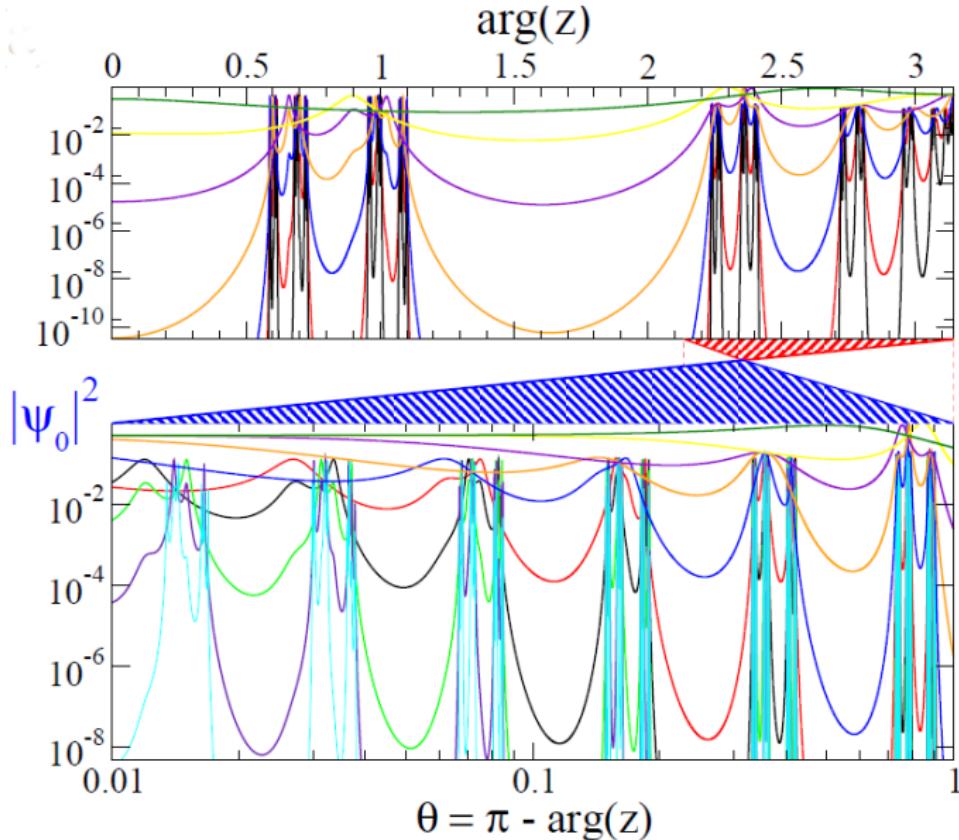


The dual Sierpinsky gasket

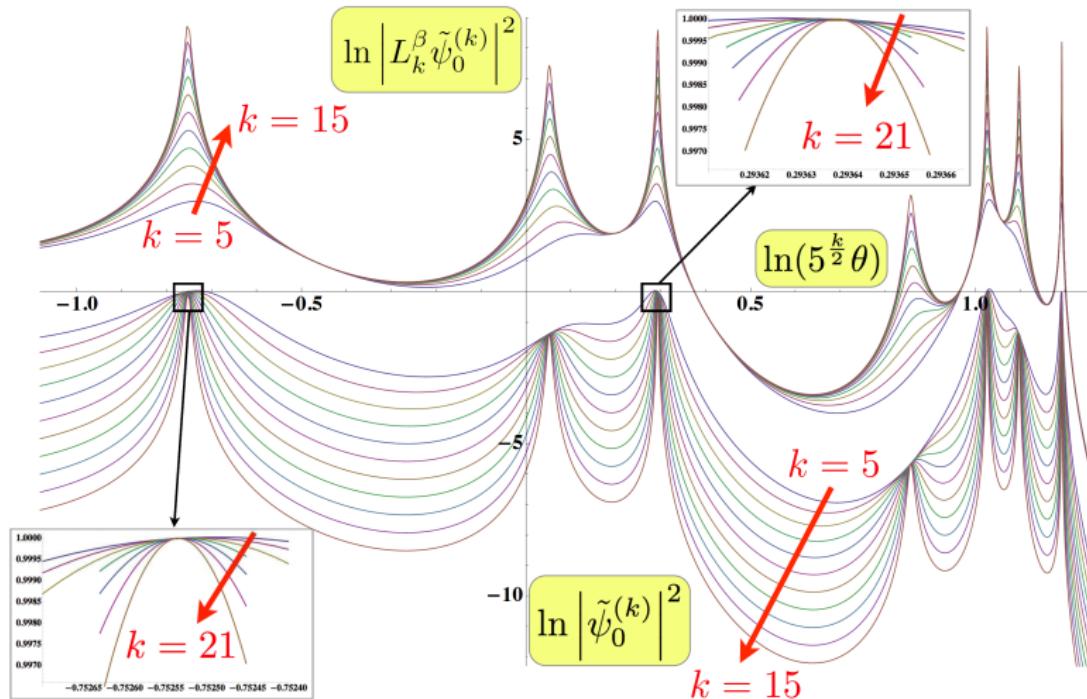
- degree 3 fractal
- exact renormalization
- $d_w^{RW} = \log_2 5$
- one absorbing boundary (square)



Lyapunov exponent for the dual Sierpinsky gasket



Scaling of the Lyapunov exponent for the DSG



Conclusions

- first attempt to study quantum walks on graphs without translational invariance
- renormalization group leads to potentially chaotic recursion equations
- scaling determined by non-local properties rather than simple fixed points
- Lyapunov exponents and observables scale similarly
- hints for $d_w^{QW} = d_w^{RW}$

- How to analyze situation where no fixed point on the unit circle exists?
- universality among different coins?
- problem of localization?