

Corner contribution to percolation cluster numbers

Ferenc Iglói (Budapest)

in collaboration with

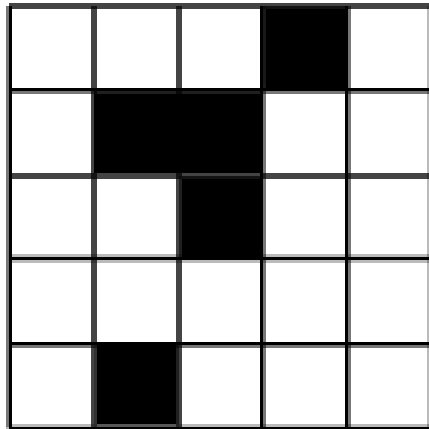
John Cardy (Oxford) and *István Kovács* (Budapest)

arXiv:1210.4671

Leipzig, 30th November 2012

Percolation

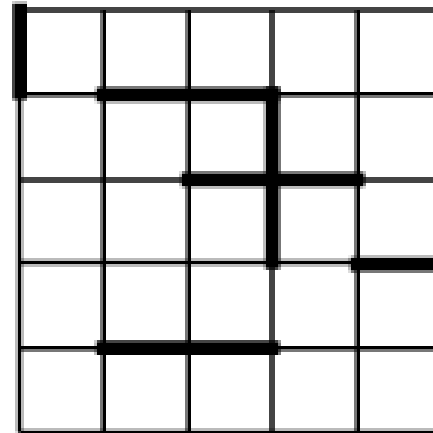
site percolation



2D: $p_c \approx 0.592746$

3D: $p_c \approx 0.3116$

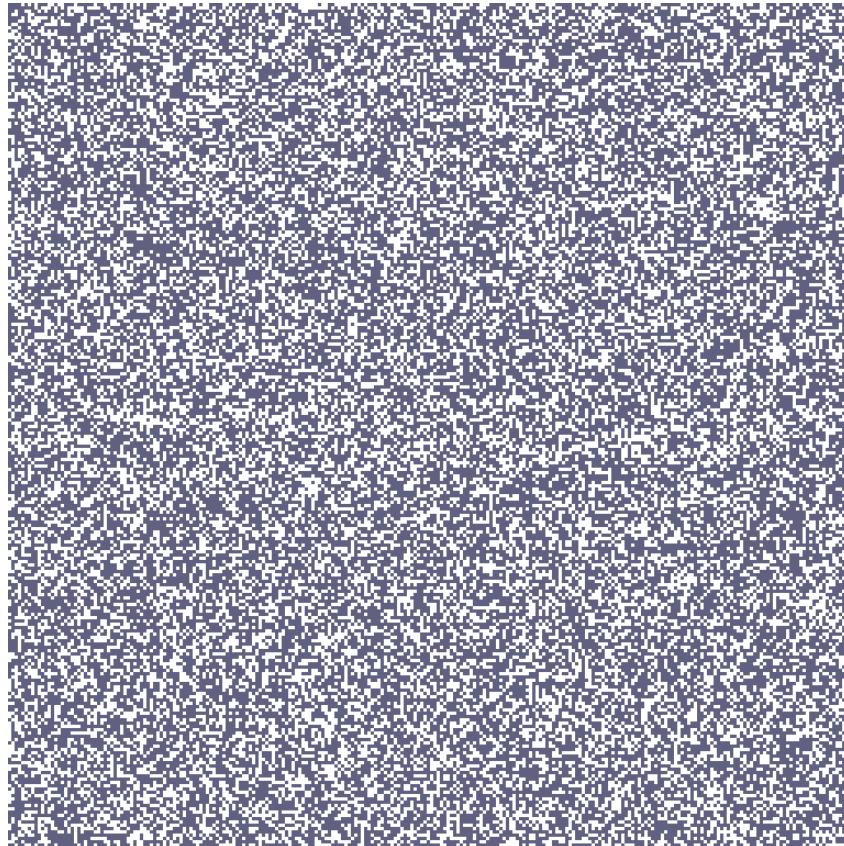
bond percolation



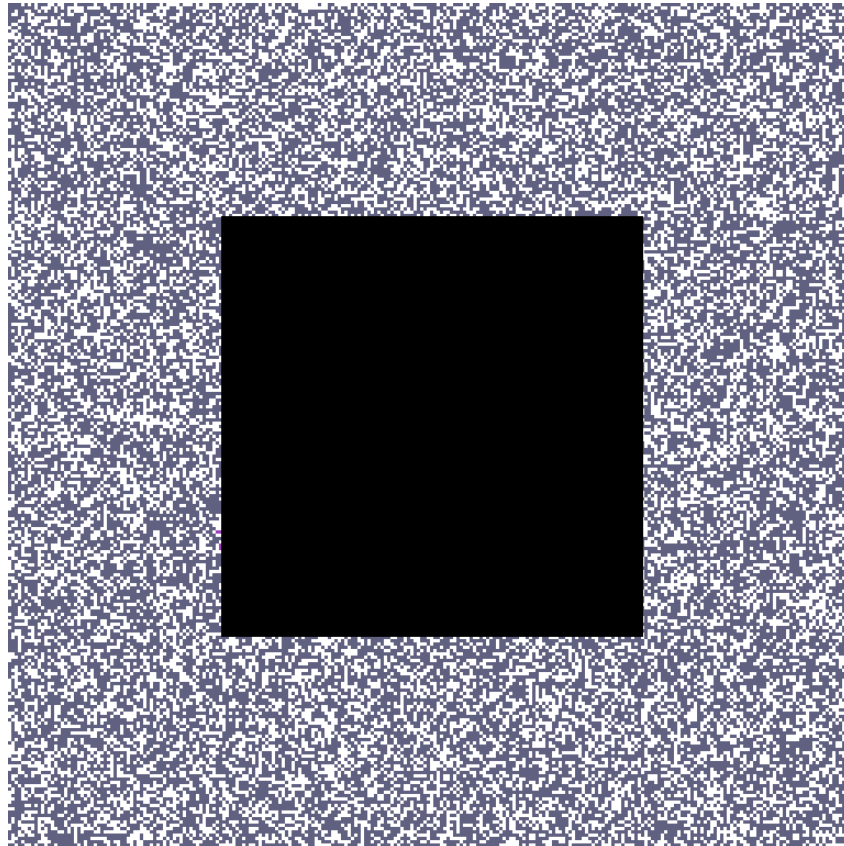
$p_c = 0.5$

$p_c \approx 0.24881$

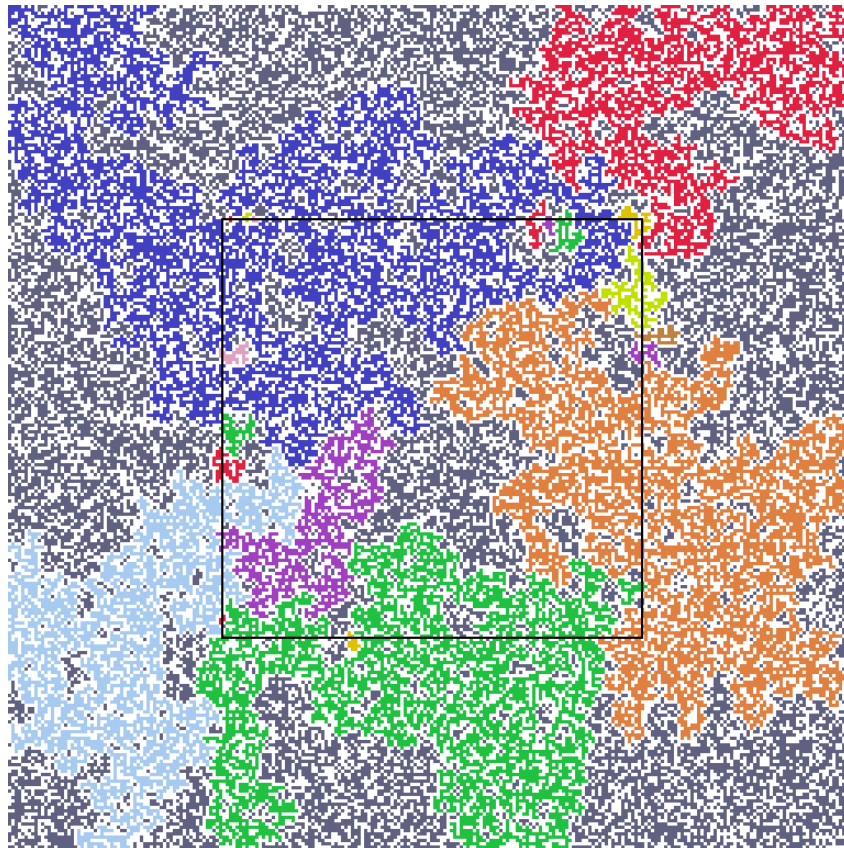
A critical system



Square subsystem



What is the number of crossed clusters?



Entanglement entropy of the diluted quantum Ising model

- Hamiltonian:

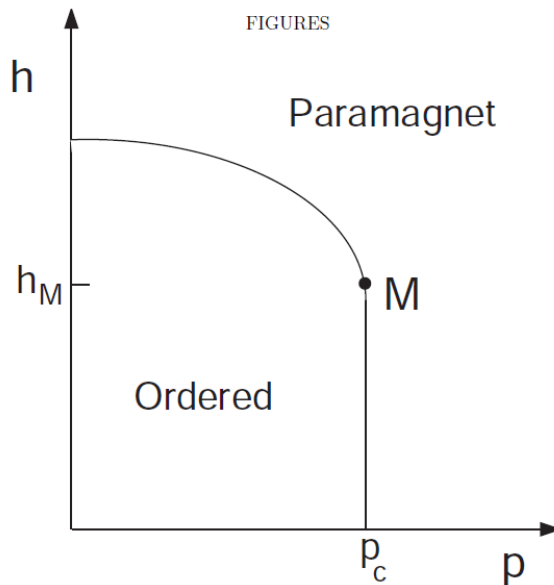
$$\mathcal{H} = -\sum_{\langle ij \rangle} J_{ij} \sigma_i^x \sigma_j^x - \sum_i h \sigma_i^z$$

- with diluted couplings

$J_{ij} = 0$ with proba. p

$J_{ij} = J > 0$ with proba. $1 - p$

- phase diagram:



- for small h percolation transition at p_c

- ground-state of the system:
each connected cluster with c number of spins is in a GHZ-state:

$$\frac{1}{\sqrt{2}} (|\underbrace{\uparrow\uparrow\cdots\uparrow}_{c \text{ times}}\rangle + |\underbrace{\downarrow\downarrow\cdots\downarrow}_{c \text{ times}}\rangle)$$

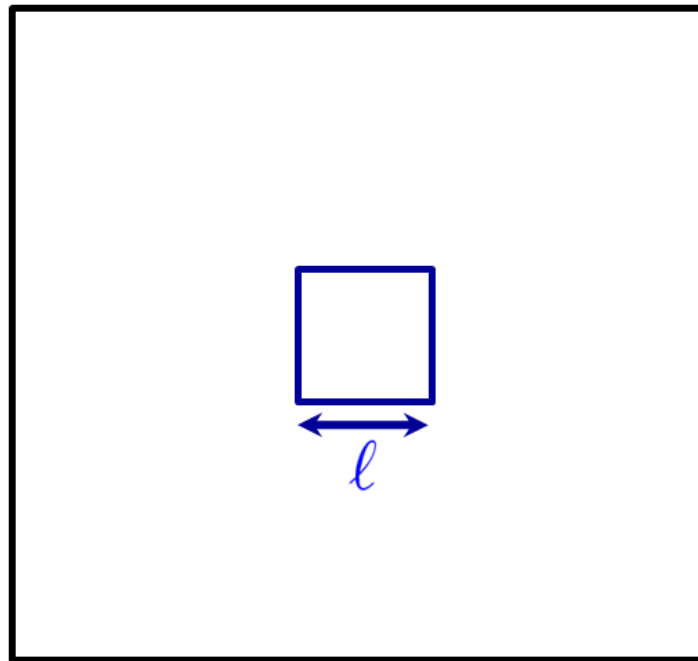
- entanglement entropy between a subsystem: A and the environment: B , having a boundary: Γ :

$$\mathcal{S}_\Gamma = -\text{Tr}(\rho_A \log_2 \rho_A)$$

- $\rho_A = \text{Tr}_B |0\rangle\langle 0|$: reduced density matrix
- $|0\rangle$ ground state of $A \cup B$
- for the GHZ-state the entanglement entropy is $\mathcal{S}_{\text{GHZ}} = 1$
- each crossed cluster gives 1 contribution to $\mathcal{S}_\Gamma$
- asymptotically $\mathcal{S}_\Gamma$ and N_Γ are equivalent

Expectation: area law

$$\mathcal{S}(\ell) \sim \ell^{d-1}$$

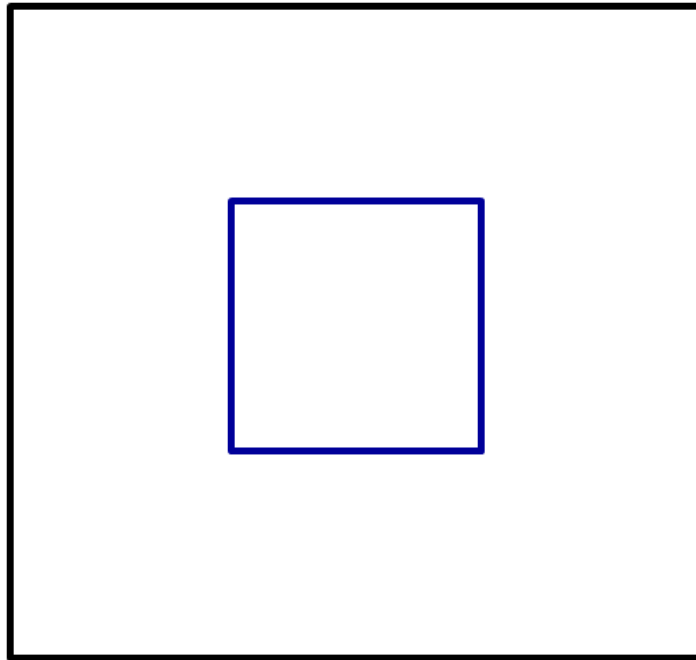


How does it grow with the subsystems size?

Expectation: area law

$$\mathcal{S}(\ell) \sim \ell^{d-1}$$

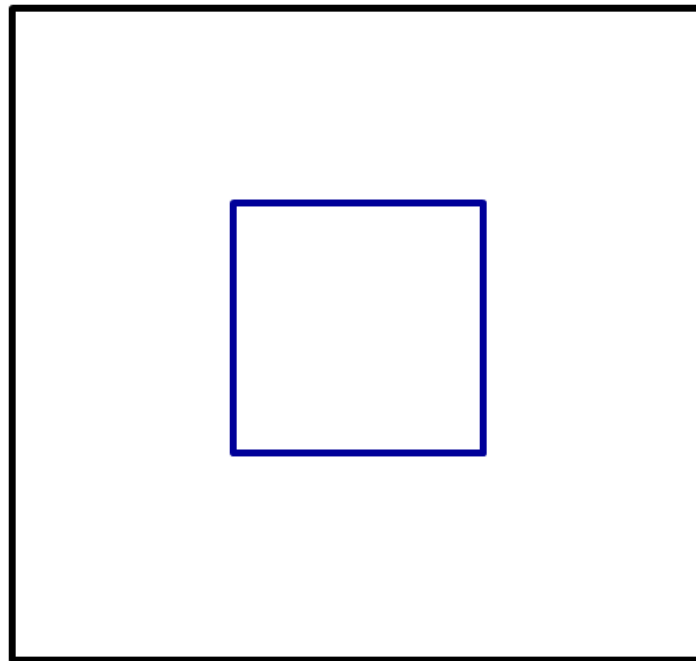
•



How does it grow with the subsystems size?

Expectation: area law

$$\mathcal{S}(\ell) \sim \ell^{d-1}$$



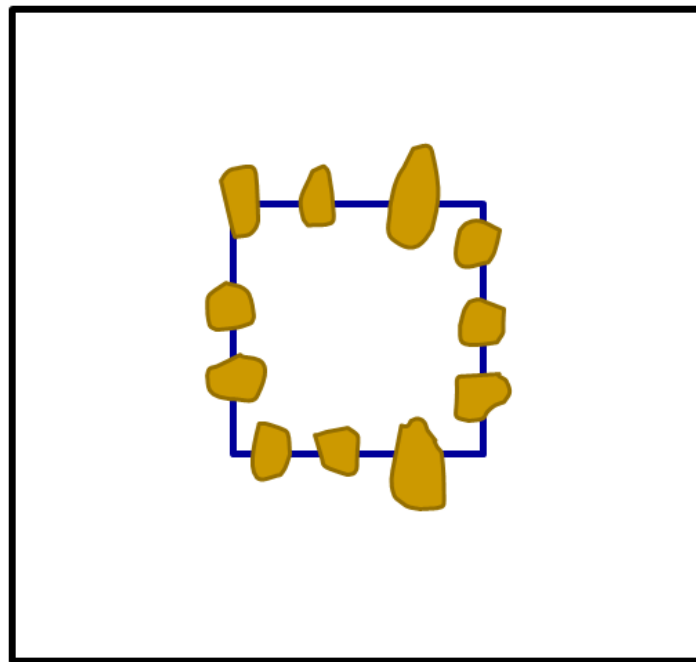
spin-cluster



How does it grow with the subsystems size?

Expectation: area law

$$\mathcal{S}(\ell) \sim \ell^{d-1}$$

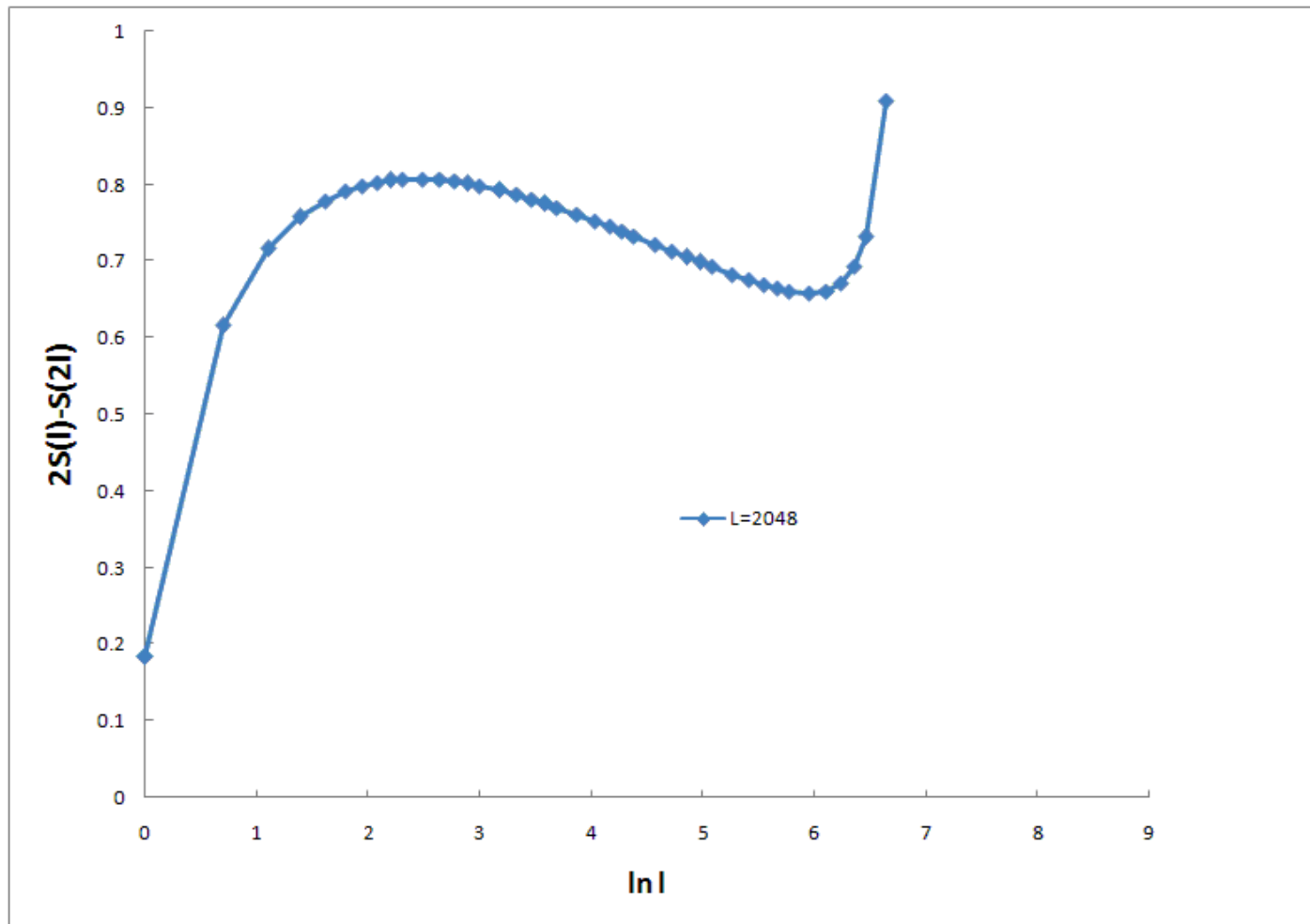


spin-cluster



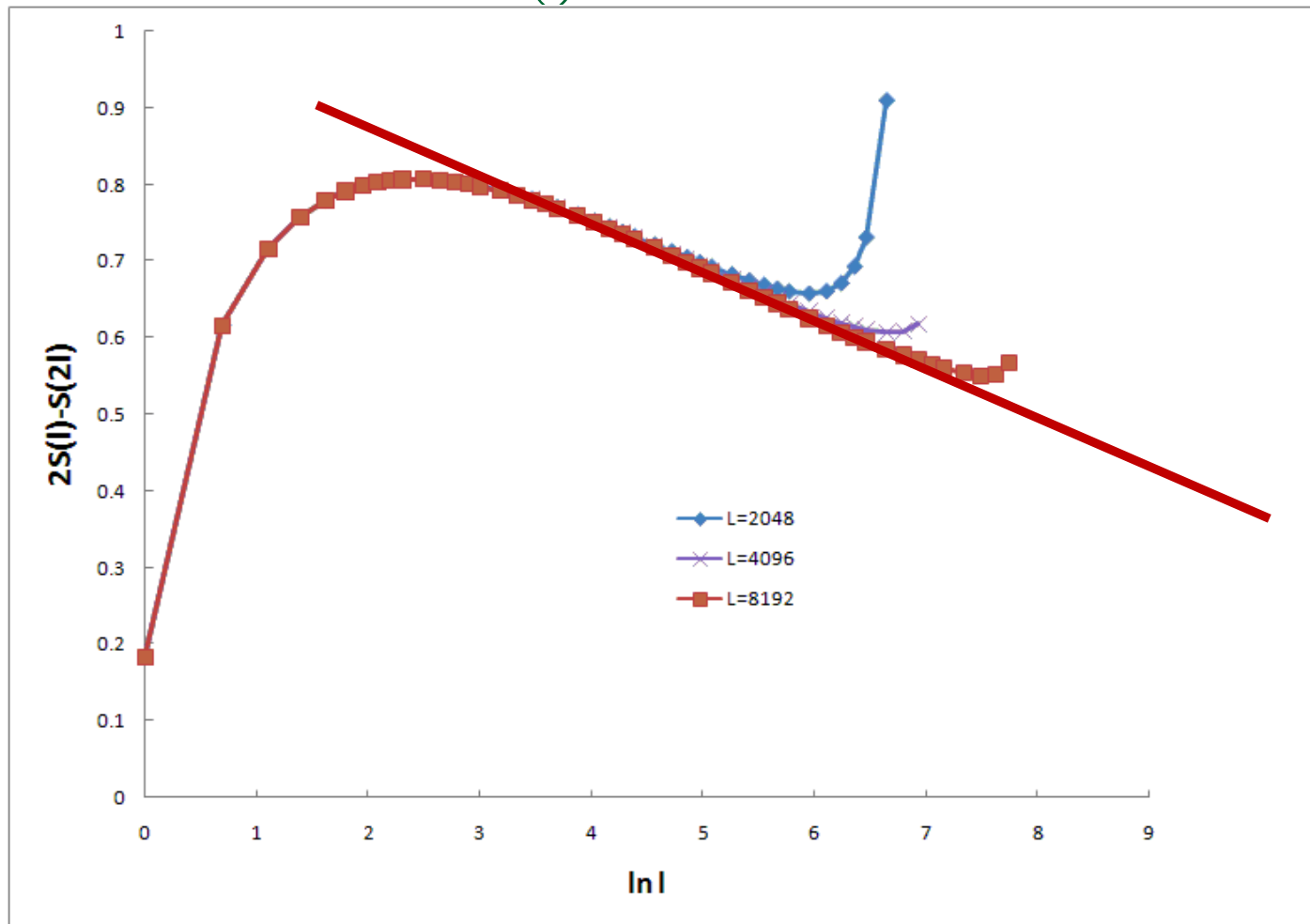
How does it grow with the subsystems size? $\mathcal{S}_{2D}(\ell) \sim \ell$

Corrections?



First study: Yu et al. Phys. Rev. B 77, 140402 (2008)

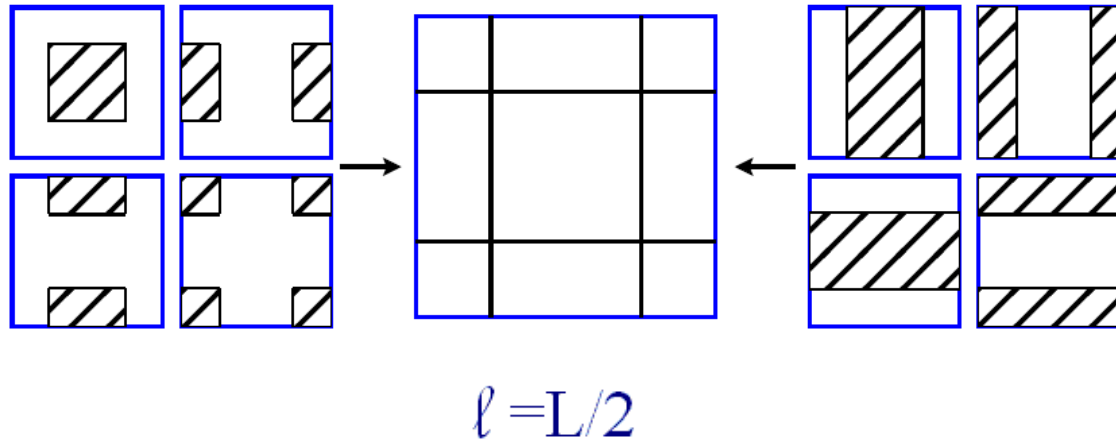
$$S(l) = a + b \ln l$$



First study: Yu et al. Phys. Rev. B 77, 140402 (2008)

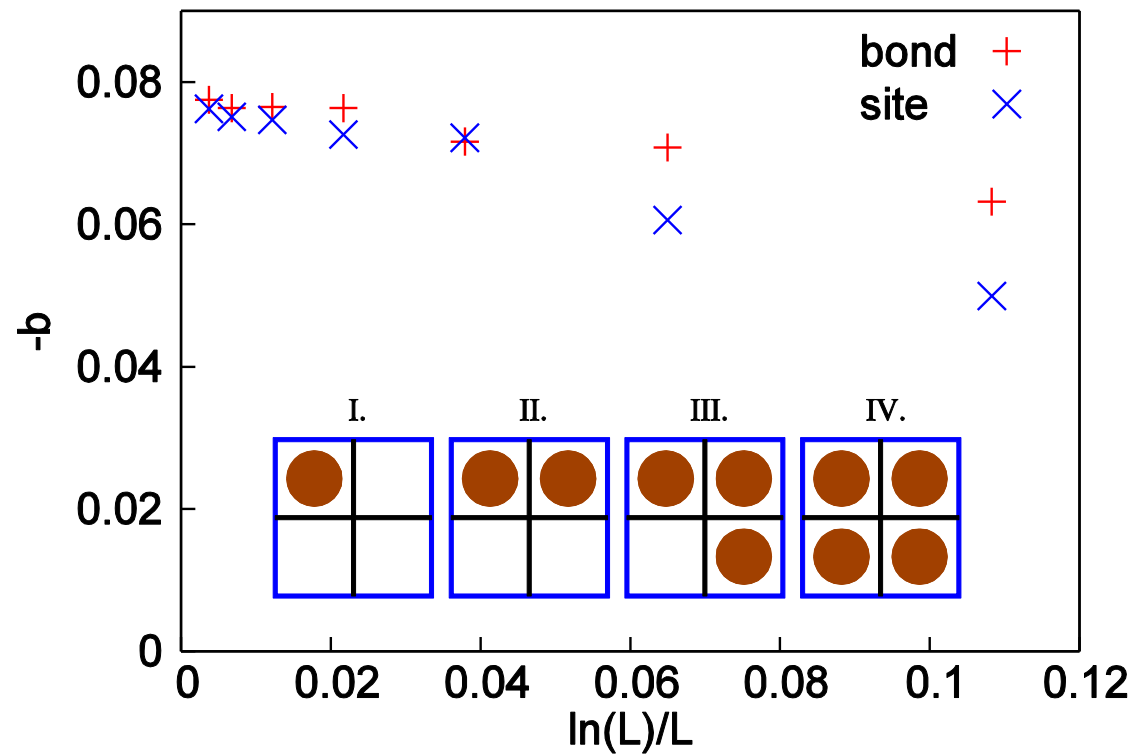
Idea: it may come from the corners

How to check this?



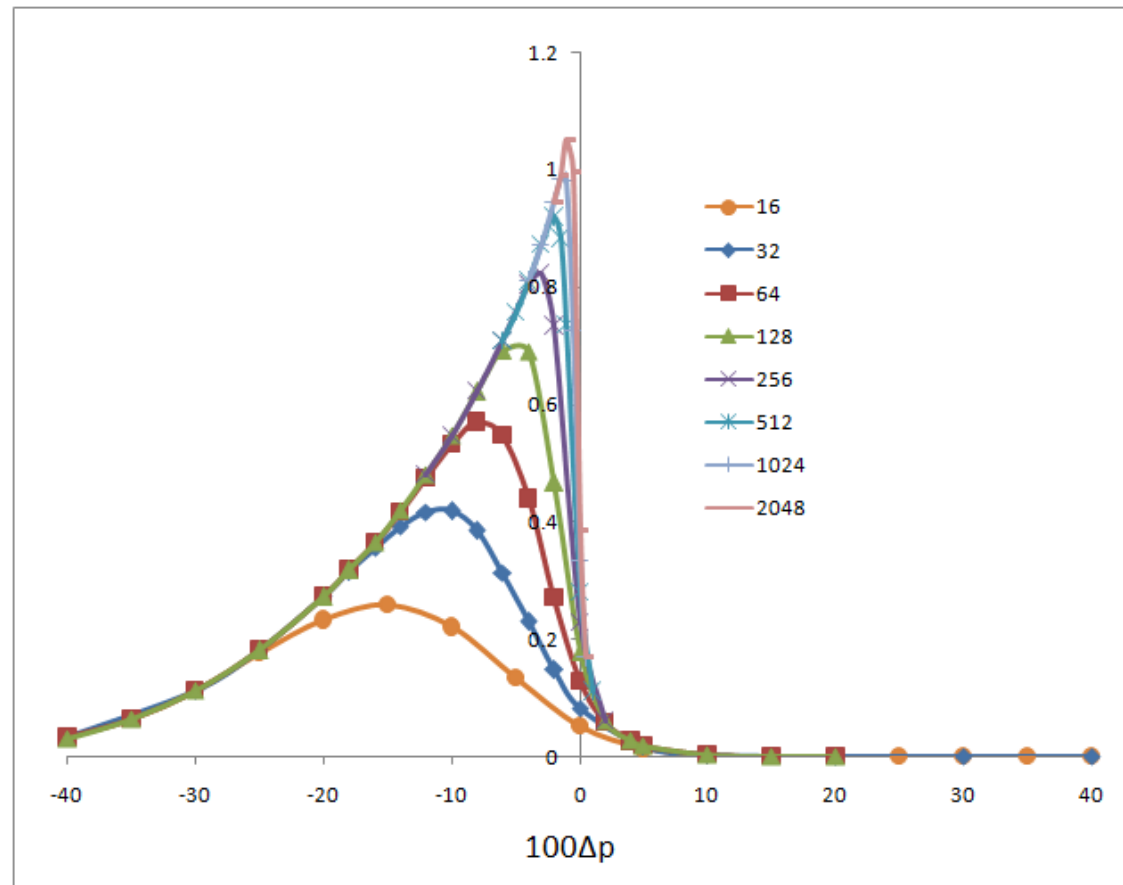
The difference is caused by the corners

The logarithmic correction is universal

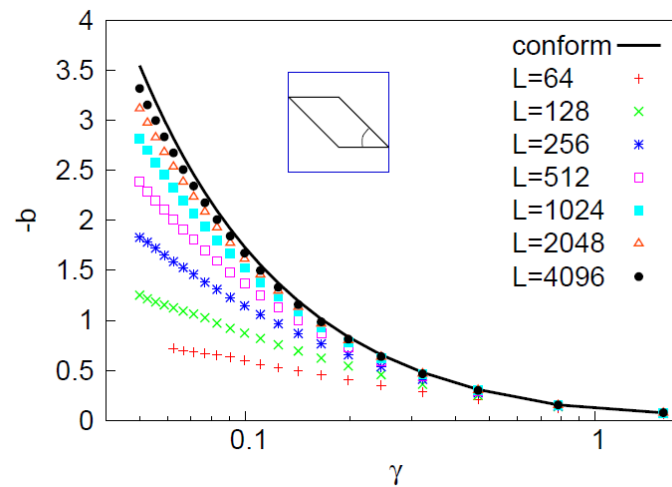
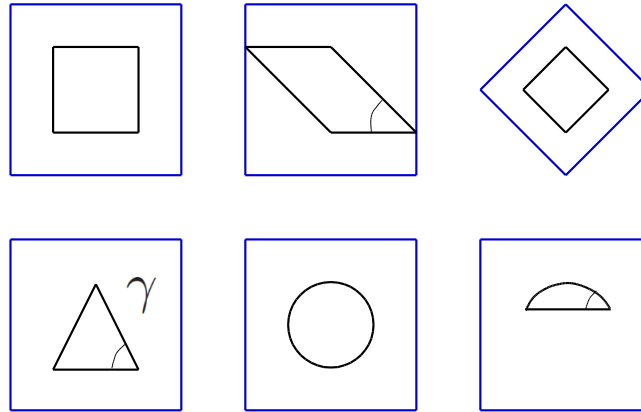


$$b = -0.077(1)$$

It is a consequence of diverging correlation length



Other shapes: angle-dependency



Potts model - random cluster representation

- Q -state Potts model
- $s_i = 1, 2, \dots, Q$
- lattice with $i = 1, 2, \dots, n$ sites and m edges
- $\langle ij \rangle$: nearest neighbours

$$Z(Q) = \sum_s \prod_{\langle ij \rangle} \exp(K \delta_{s_i, s_j})$$

- $\exp(K \delta_{s_i, s_j}) = 1 + \frac{p}{1-p} \delta_{s_i, s_j}$
- $p = 1 - e^{-K}$

$$Z(Q) = \sum_s \prod_{\langle ij \rangle} \left(1 + \frac{p}{1-p} \delta_{s_i, s_j} \right)$$

- expand the products: the $\frac{p}{1-p} \delta_{s_i, s_j}$ terms define the “F” Fortuin-Kasteleyn clusters

- with $N_{tot}(F)$ number of connected components and $M(F)$ edges

$$\begin{aligned} Z(Q) &= \sum_F Q^{N_{tot}(F)} \left(\frac{p}{1-p} \right)^{M(F)} \\ &\sim \sum_F Q^{N_{tot}(F)} p^{M(F)} (1-p)^{m-M(F)} \\ &= \langle Q^{N_{tot}} \rangle \end{aligned} \quad (1)$$

- Q real parameter, $Q \rightarrow 1$ percolation
- critical point: $p_c = \frac{1}{2}$
- mean number of clusters:

$$\langle N_{tot} \rangle = \left. \frac{\partial \ln Z(Q)}{\partial Q} \right|_{Q=1}.$$

Number of crossing clusters

- impose a boundary condition
- all the Potts spins are fixed (say, in state 1) on Γ

$$Z_{\Gamma}(Q) \sim \langle Q^{N_{tot}-N_{\Gamma}} \rangle ,$$

where N_{Γ} is the number of clusters which intersect Γ .

$$\langle N_{tot} - N_{\Gamma} \rangle = \left. \frac{\partial \ln Z_{\Gamma}(Q)}{\partial Q} \right|_{Q=1}$$

- at the critical point $p = p_c$

$$\begin{aligned} \ln Z(Q) &\sim A f_b(Q) \\ \ln Z_{\Gamma}(Q) &\sim A f_b(Q) + L_{\Gamma} f_s(Q) + C_{\Gamma}(Q) \ln L_{\Gamma} \end{aligned} \quad (2)$$

- $A \propto n$ total area
- L_{Γ} length of Γ
- f_b : bulk free-energy density

- f_s surface free-energy density

- last term: corner contribution

$$\langle N_{\Gamma} \rangle = -f'_s(1)L_{\Gamma} - C'_{\Gamma}(1) \ln L_{\Gamma}$$

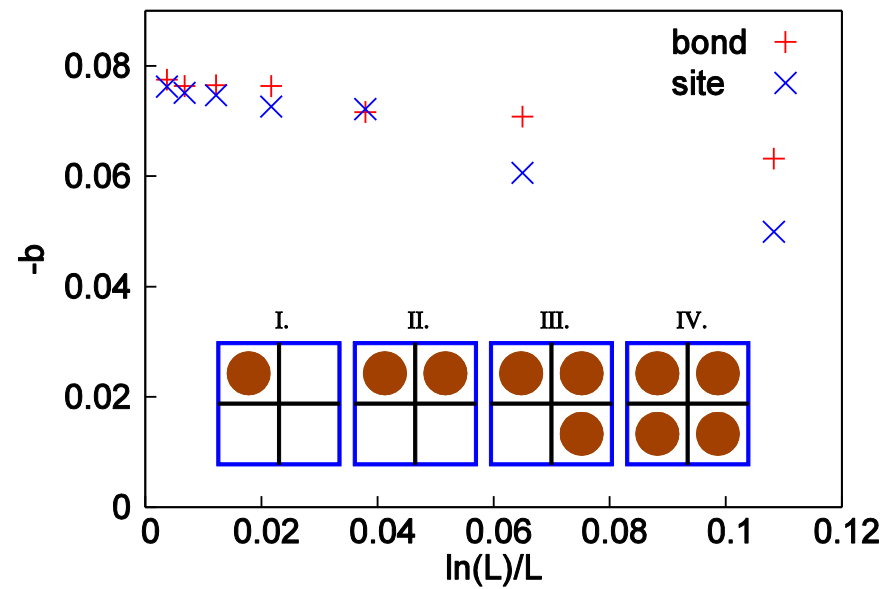
- Cardy-Peschel formula for the corner contribution of conformally invariant systems:

$$\begin{aligned} C_{\Gamma}(Q) &= \frac{c(Q)}{24} \sum_k \left[\left(\frac{\pi}{\gamma_k} \right) - \left(\frac{\gamma_k}{\pi} \right) \right. \\ &\quad \left. + \left(\frac{\pi}{2\pi - \gamma_k} \right) - \left(\frac{2\pi - \gamma_k}{\pi} \right) \right] \end{aligned} \quad (3)$$

- γ_k ($2\pi - \gamma_k$) interior (exterior) angle of Γ
- $c(Q)$: central charge of the Q -state Potts model

$$c'(1) = \frac{5\sqrt{3}}{4\pi}$$

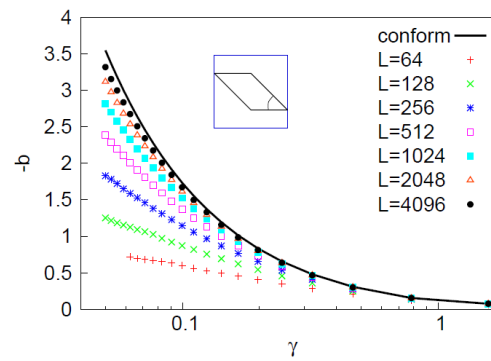
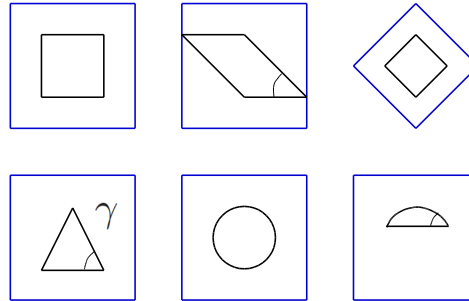
The logarithmic correction is universal



$$b = -0.077(1)$$

$$-C'(1) = -5\sqrt{3}/(36\pi) = -0.07657$$

Other shapes: angle-dependency



$$C'(1) = -\frac{c'(1)}{12} \left[4 - \pi \left(\frac{1}{\gamma} + \frac{1}{\pi - \gamma} + \frac{1}{\pi + \gamma} + \frac{1}{2\pi - \gamma} \right) \right]$$

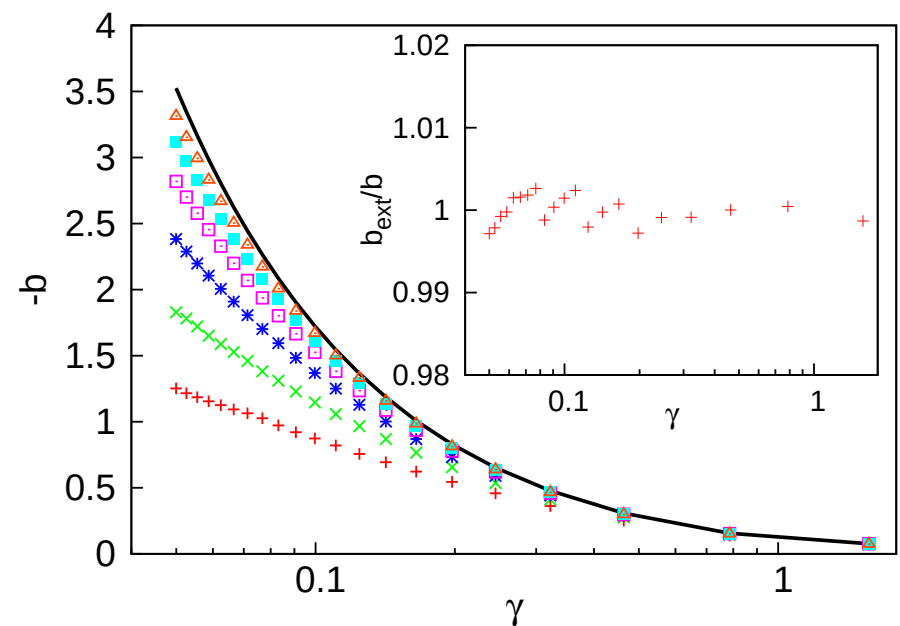
Anisotropic percolation

- different probabilities:
 p_x horizontally, p_y vertically
- at the critical point: $p_x + p_y = 1$
- diagonally oriented square
- with symmetric angles:
 $\gamma = \pi - \gamma = \pi/2$)

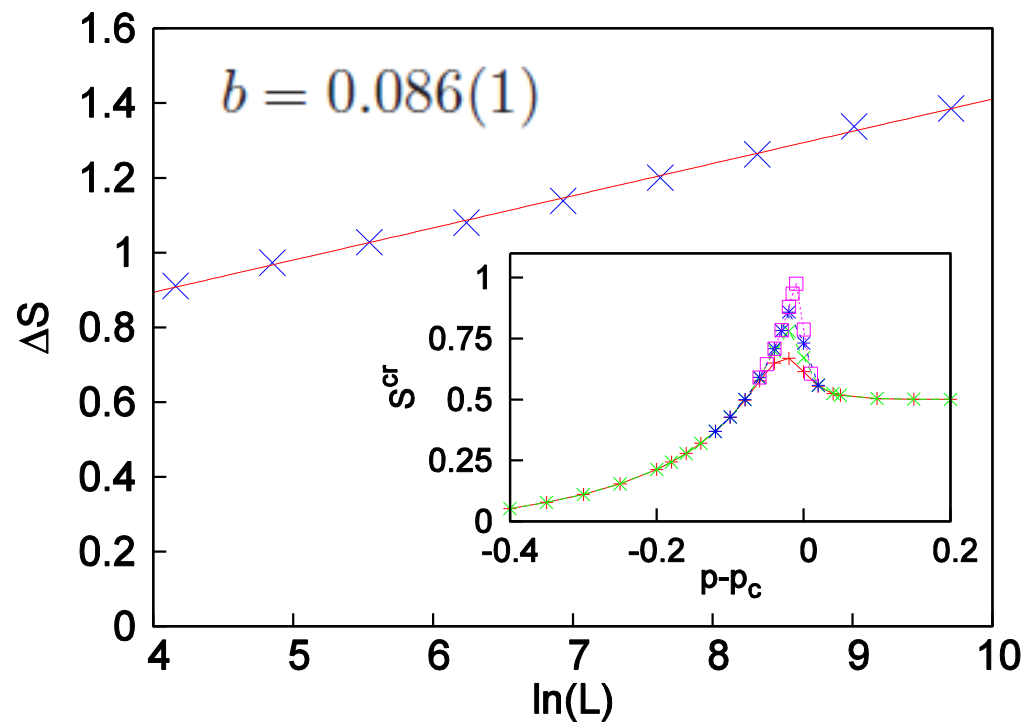
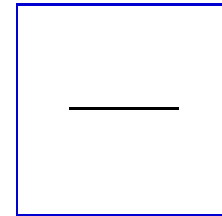
according to discrete holomorphicity
in the scaling limit equivalent to an

- isotropic system: $p_x = p_y = 1/2$
- with asymmetric angles

$$\frac{p_y}{p_x} = \frac{\sin((\pi - \gamma)/3)}{\sin(\gamma/3)}$$



In 2D we have a complete understanding...
 ...it can be any open curve as well!



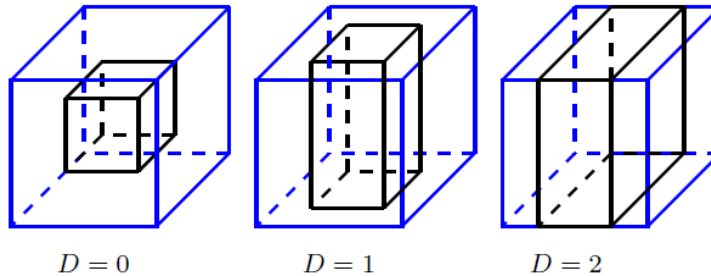
$$-C'_{\Gamma}(1) = \frac{5\sqrt{3}}{32\pi} \approx 0.08615$$

joint work with John Cardy (Oxford)

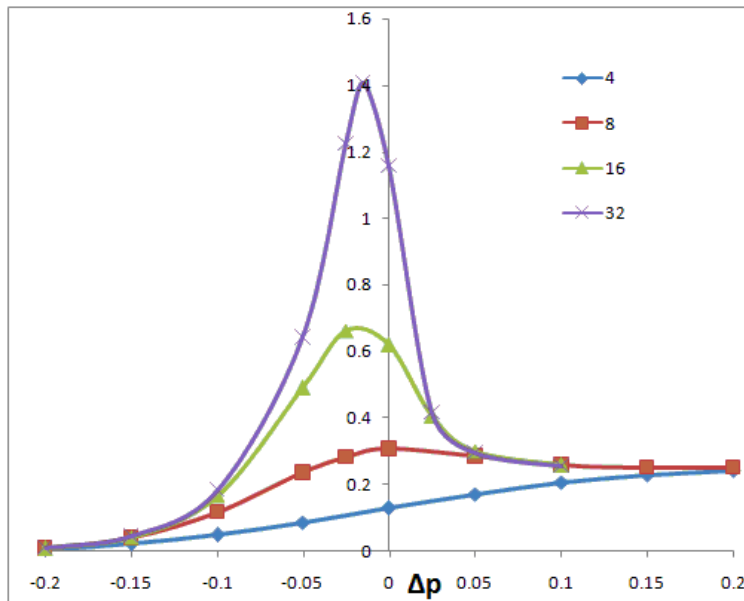
Conclusion - Outlook

- percolation cluster numbers: N_{Γ}
- entanglement entropy of diluted quantum Ising model: $\mathcal{S}_{\Gamma}$
- asymptotically: $N_{\Gamma} \approx \mathcal{S}_{\Gamma}$
- in $2d$ corner contributions at the critical point are
 - universal
 - exactly known from the Cardy-Peschel formula
- Related problems:
 - percolation cluster numbers in $3d$
 - entanglement entropy of the random quantum Ising model for $d \geq 2$

3D: There is a similar corner contribution

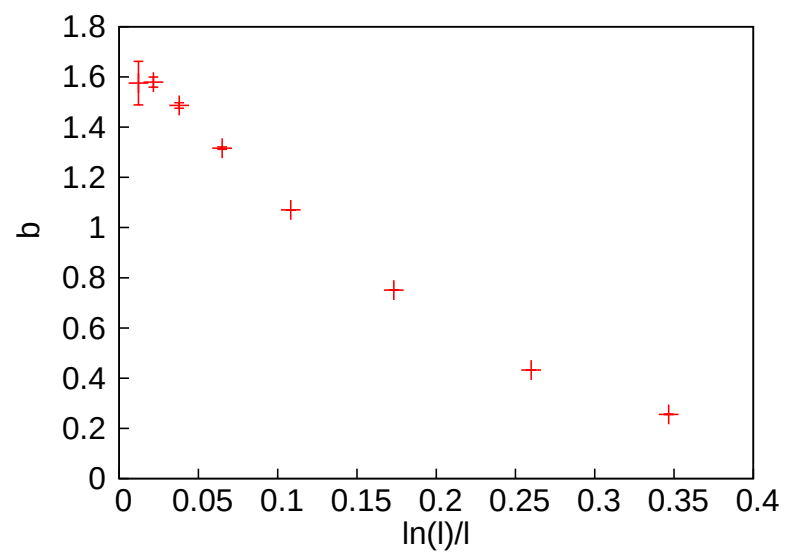


$$\mathcal{S}_0^{(d)}(0) = \sum_{D=0}^{d-1} \left(-\frac{1}{2}\right)^D \binom{d}{D} \mathcal{S}_D^{(d)}$$

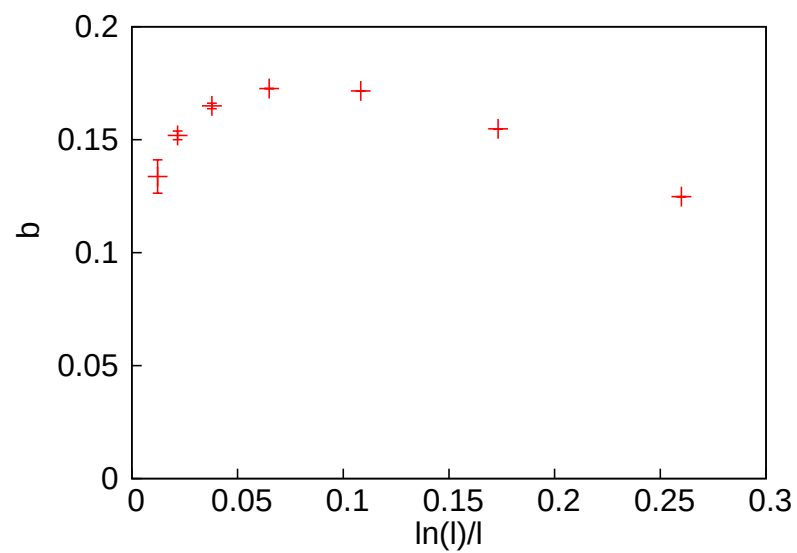


The sign is different!

Is it a log or $\log^2$?



prefactor of $\log l$



prefactor of $\log^2 l$