

Pair-Factorised Steady State model exhibits condensate's growth on monolayers



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Models of mass transport and condensation

Motivation:

- Cluster or atomic deposition, diffusion of atomic clusters on substrate surface, thin films growth, self-assembled quantum dots, etc.

Existing models:

- Balls-in-Boxes, Zero-Range Process (**ultralocal** rules of dynamics \rightarrow stationary state factorises \rightarrow condensates occupy **single sites**), Solid-on-Solid, etc.
- The models aim at describing particle/mass condensation: critical density, size and shape of the condensate, fluctuations, etc.

Pair-Factorised Steady State (PFSS)

M particles at N sites, set of occupation numbers $\{m_1, \dots, m_N\}$.

A particle at a site i hops with a rate $u(m_i | m_{i+1}, m_{i-1})$ (can be asymmetric).

$$\text{Master eq. } \frac{dP(\vec{m}, t)}{dt} = \sum_{\vec{m}'} [W(\vec{m}' \rightarrow \vec{m})P(\vec{m}', t) - W(\vec{m} \rightarrow \vec{m}')P(\vec{m}, t)],$$

$$\text{In steady-state: } \sum_{\vec{m}'} P(\vec{m})W(\vec{m} \rightarrow \vec{m}') = \sum_{\vec{m}'} P(\vec{m}')W(\vec{m}' \rightarrow \vec{m})$$

$$\text{one obtains: } P(\vec{m}) = \frac{1}{Z} \prod_{i=1}^N g(m_i, m_{i+1}) \delta_{m_1 + \dots + m_N, M},$$

The stationary state is pair factorised.

$$\text{Assume weights: } g(m, n) = K(|m - n|) \sqrt{p(m)p(n)},$$

B Waclaw, J Sopik, W Janke and H Meyer-Ortmanns, J. Stat. Mech. (2009) P10021

B Waclaw, J Sopik, W Janke and H Meyer-Ortmanns, PRL 103 (2009) 080602

Known 1D results

Assume $P(m)$ is a **probability of a microstate**.

Canonical:
$$Z(N, M) = \sum_{\{m_i\}} \prod_i g(m_i, m_{i+1}) \delta_{\sum_i m_i, M}$$

Grand-can.:
$$Z_N(z) = \sum_M Z(N, M) z^M = \sum_{\{m_i\}} z^{\sum_i m_i} \prod_i g(m_i, m_{i+1})$$

In 1D case, the partition function takes the form:

$$Z_N(z) = \sum_{m_1, \dots, m_N} T_{m_1 m_2} T_{m_2 m_3} \cdots T_{m_N m_1} = \text{Tr} T(z)^N,$$

where $T_{mn} = z^{(m+n)/2} g(m, n)$

So solving 1D problem is equivalent to the **eigenproblem**:

$$\sum_n T_{mn} \phi_n = \lambda_{\max} \phi_m,$$

Lennard - Jones

$$g(m, n) = \exp(-J|m - n|) \exp\left(\frac{U}{4} \left[\left(\frac{\sigma}{m+1}\right)^3 + \left(\frac{\sigma}{n+1}\right)^3 - \left(\frac{\sigma}{m+1}\right)^9 - \left(\frac{\sigma}{n+1}\right)^9 \right]\right)$$

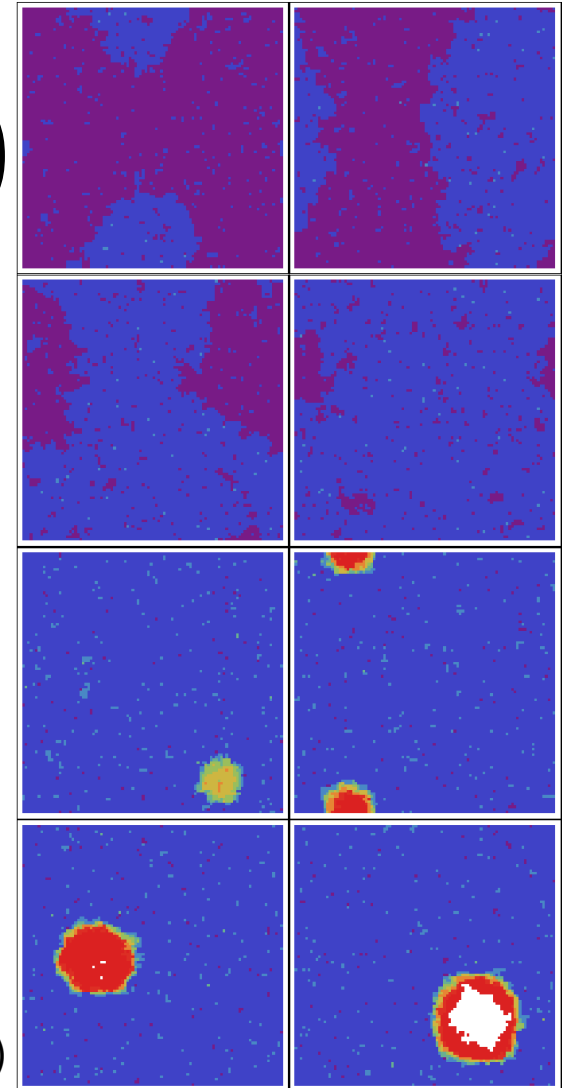
A layer of thickness $\approx \sigma$ builds up.

A condensate forms **upon** the layer.

σ tunes the thickness (almost) linearly!

Densities: 0.2, 0.5,
0.7, 0.9,
1.1, 1.2,
1.5, 1.8

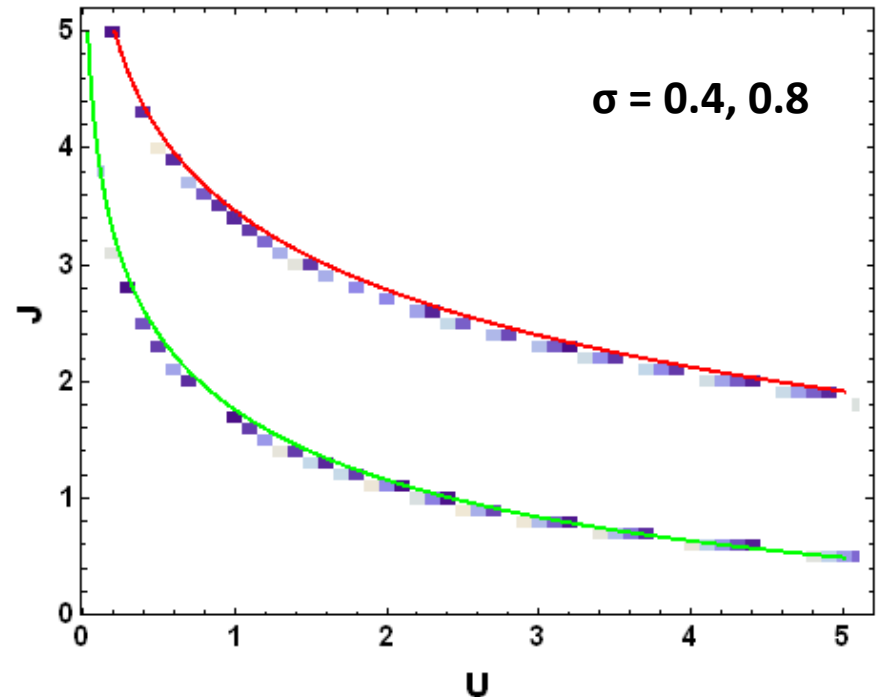
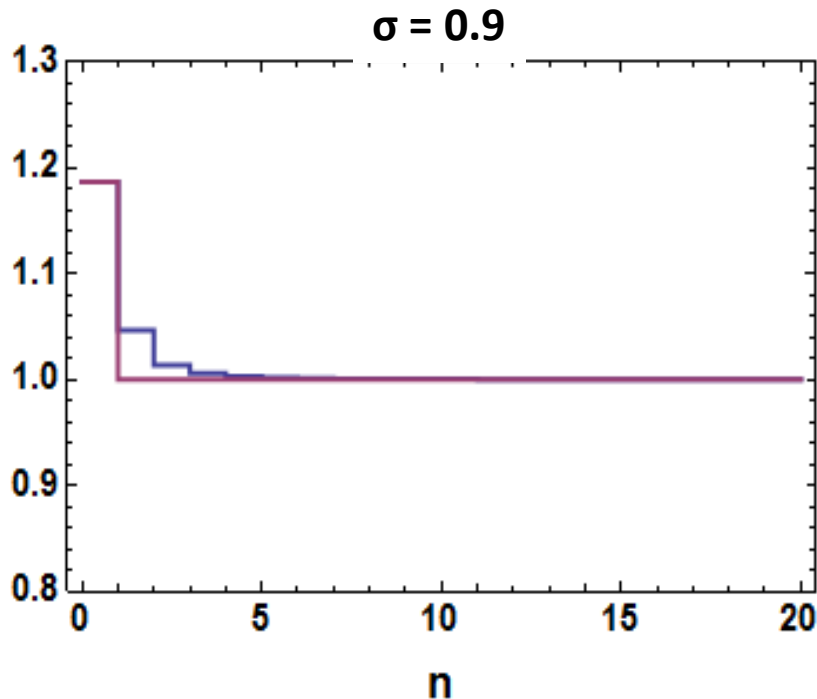
$\sigma = 1$
 $J=1, U=6$



Good 1D approximations, $\sigma < 1$

Instead of $\exp \left[\frac{U}{2} \left[\left(\frac{\sigma}{m+1} \right)^3 - \left(\frac{\sigma}{m+1} \right)^9 \right] \right]$

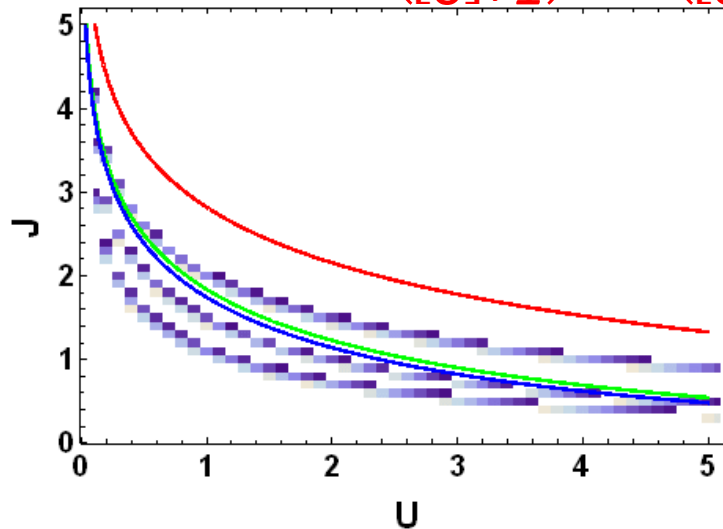
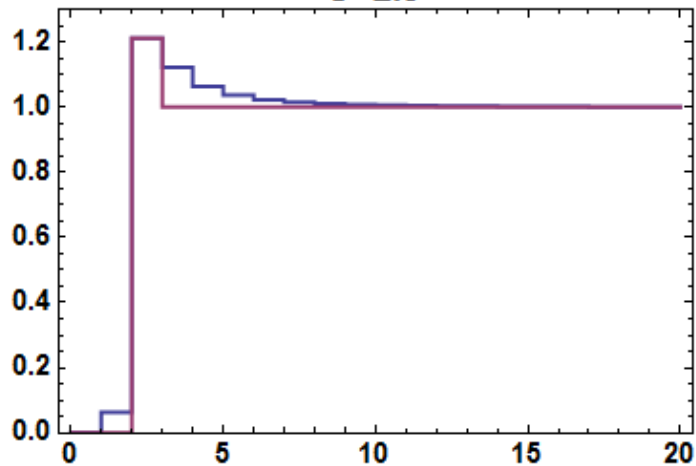
use $\exp \left[\frac{U'}{2} \delta(n, 0) \right]$ with $U' = U(\sigma^3 - \sigma^9)$



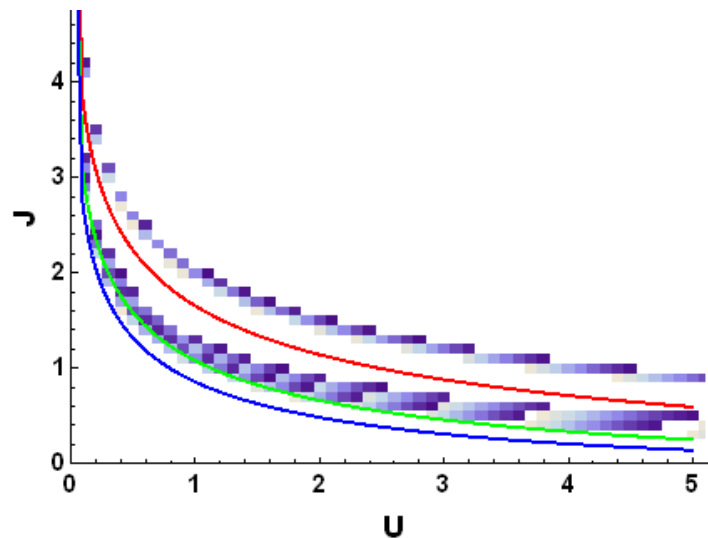
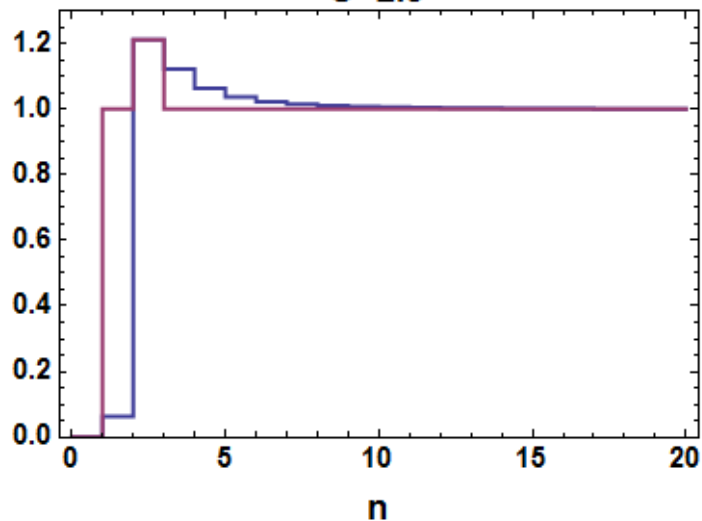
Bad 1D approximations, $\sigma > 1$

Similarly take $\exp\left[\frac{U'}{2} \delta(n, 0)\right]$ with $U' = U\left[\left(\frac{\sigma}{|\sigma|+1}\right)^3 - \left(\frac{\sigma}{|\sigma|+1}\right)^9\right]$

$S=2.5$



$S=2.5$



1D approximations

For these approximations one can solve

$$\sum_n g(m, n) \varphi_n = \lambda_{max} \varphi_m$$

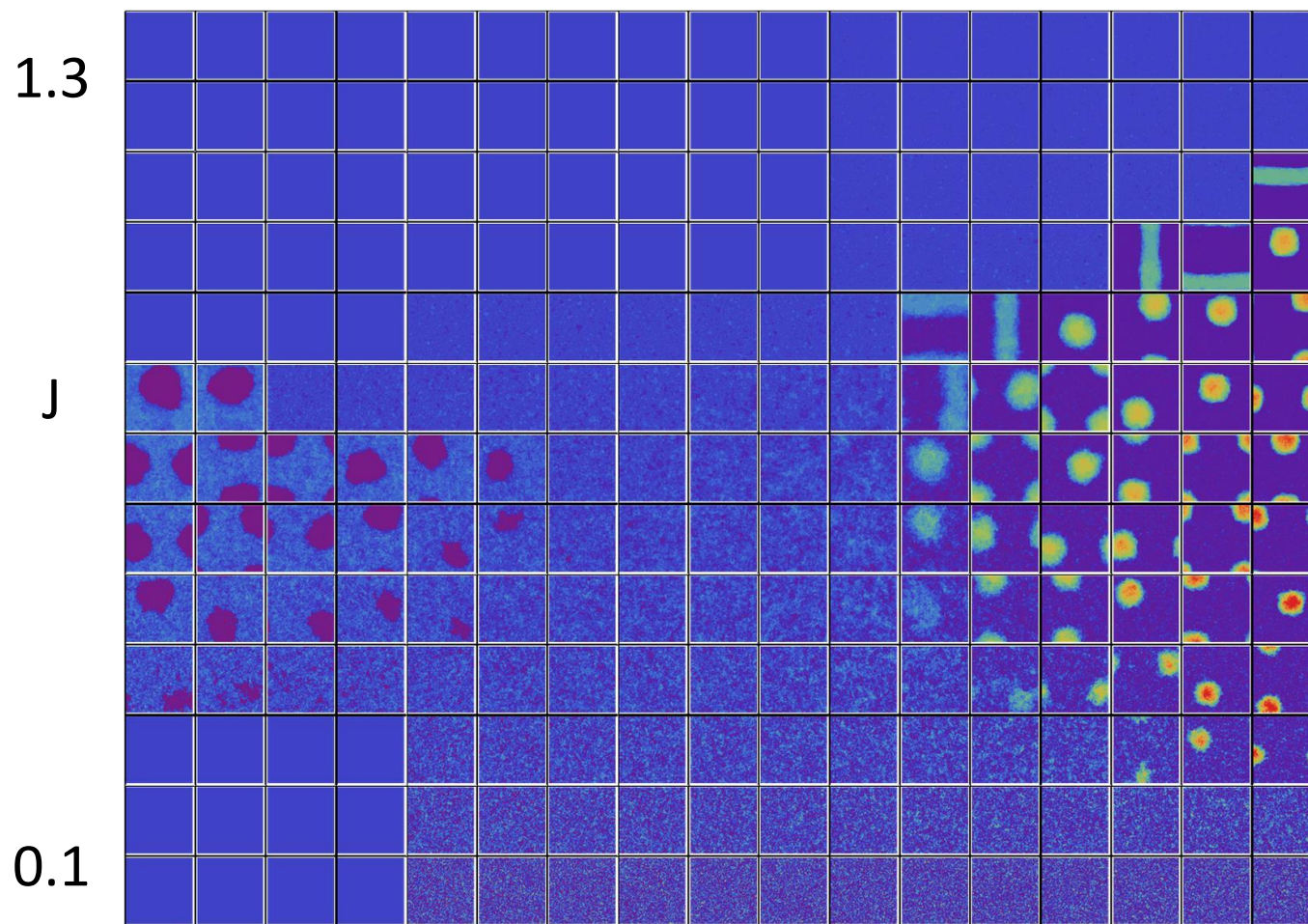
and calculate the **critical density**

$$\rho_c(\sigma) = \frac{\sum_n^\infty m \varphi_m^2}{\sum_n^\infty \varphi_m^2} = \rho_c(0) + \sigma$$

Number of monolayers grows linearly with σ

2D phase diagrams

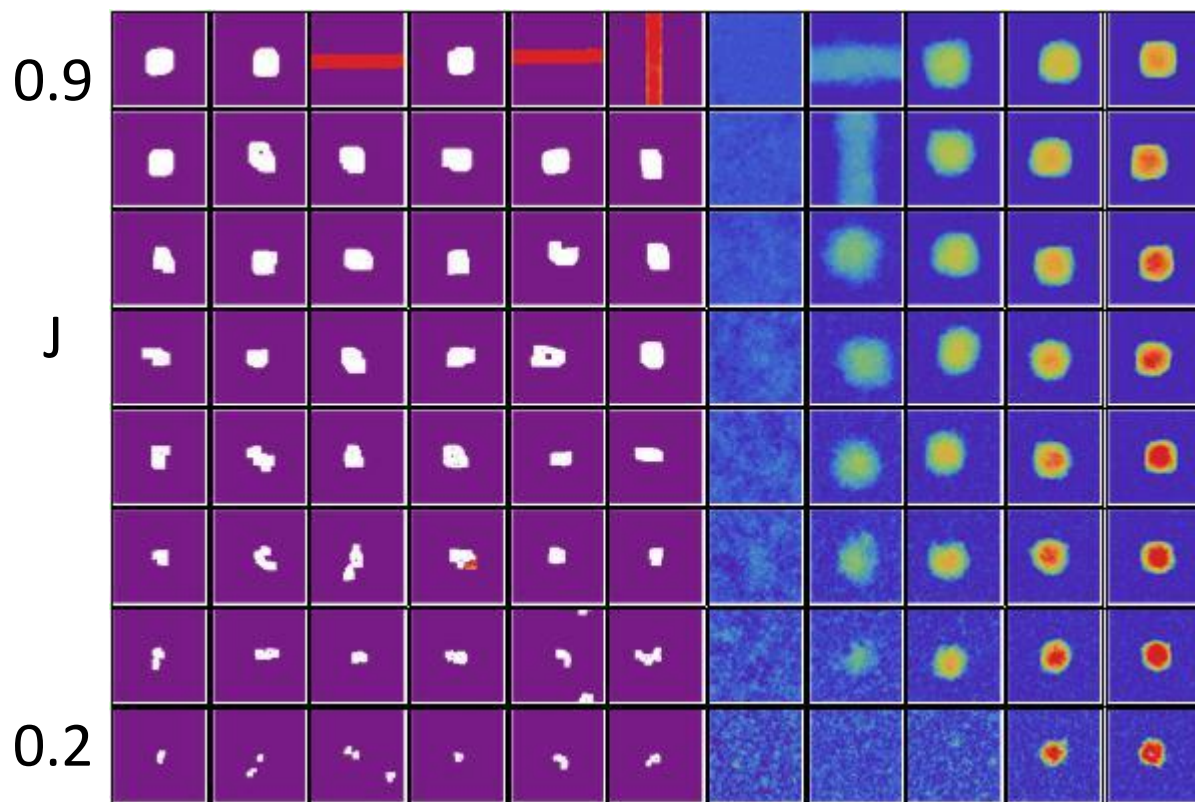
$$\sigma = 1$$



$$U = -10, \dots, 6$$

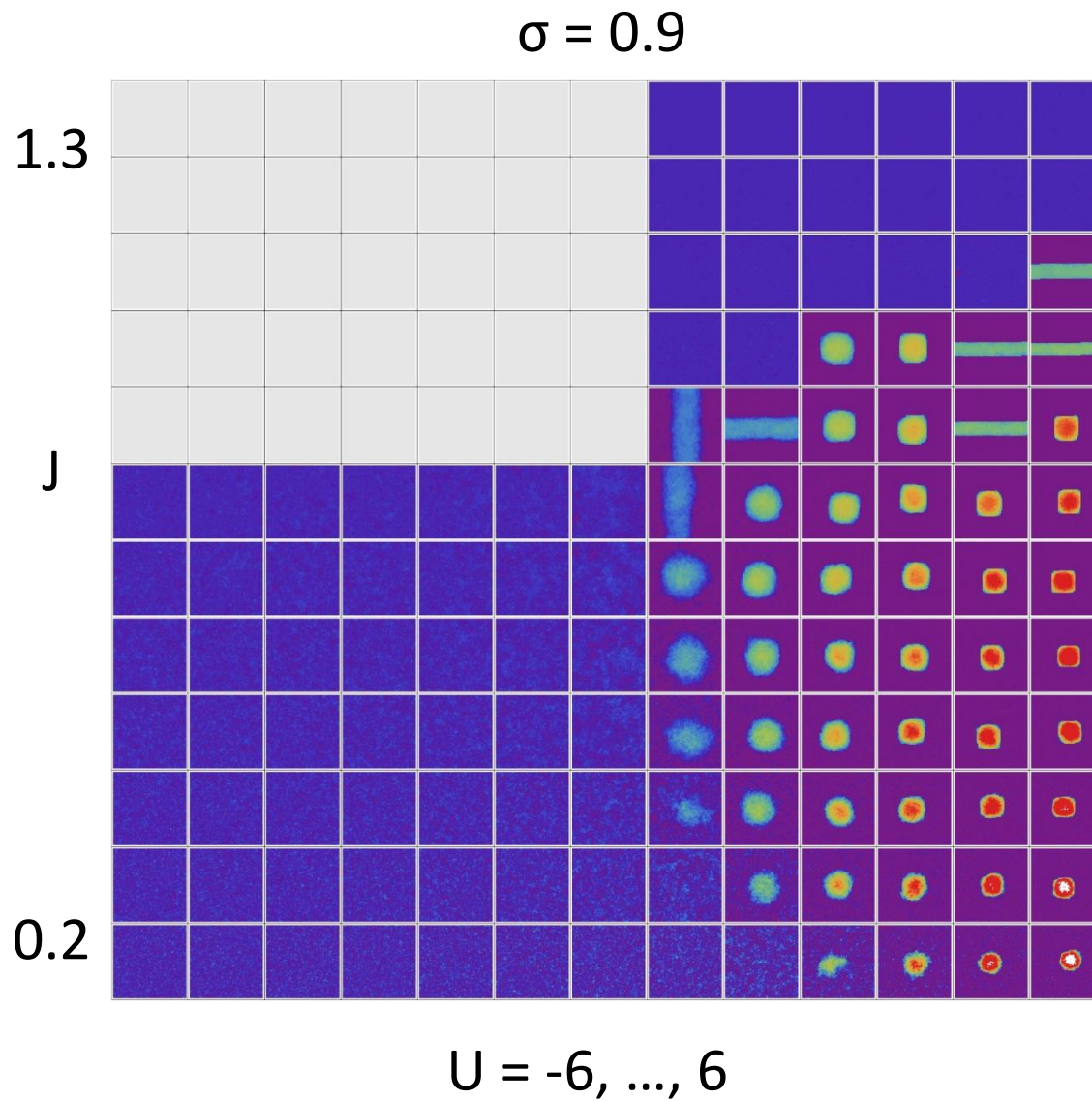
2D phase diagrams

$$\sigma = 3$$



$$U = -6, \dots, 4$$

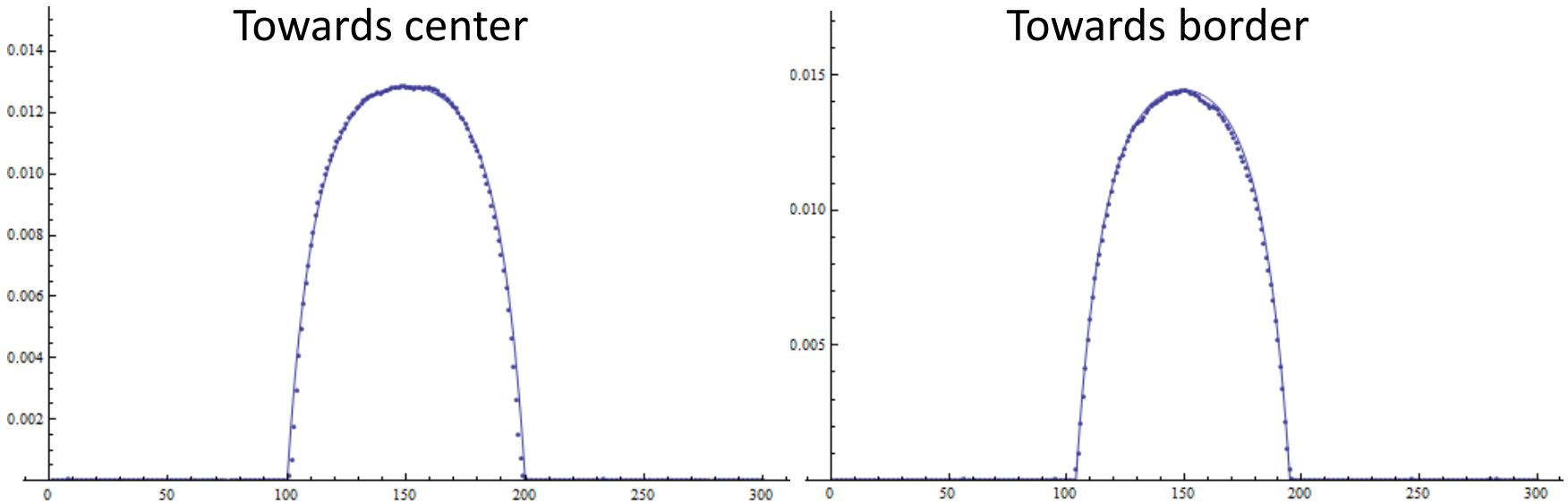
2D phase diagrams



2D condensate's shapes

In 1D width of the condensate $\sim \sqrt{N}$

And the shape
$$h(t) = \frac{w}{2v} \ln \frac{\tilde{K}(v)}{\tilde{K}(vt)}$$
 (t – normalised position,
 w and v depend on J and U)
 $K(|m - n|) = \exp(-J|m - n|)$



Sections of 2D condensate can be fitted with 1D shape.

Larger flattening of the top for the same J values than in 1D.

Conclusions

- The model can exhibit growth of the condensate on monolayers (reminding of Stranski-Krastanov growth)
- Number (thickness) of monolayers connected with the σ parameter of Lennard-Jones potential
- 2D condensates behave very much like 1D



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