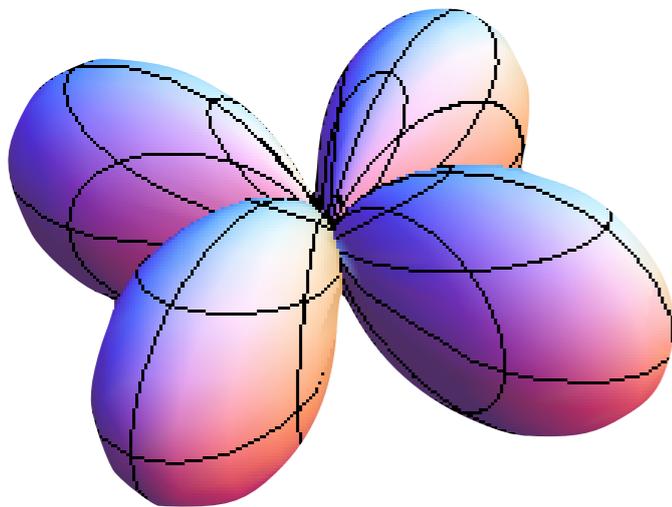


Orbital ordering in e_g orbital systems:

Critical properties of the classical 120° model

CompPhys'10, November 2010

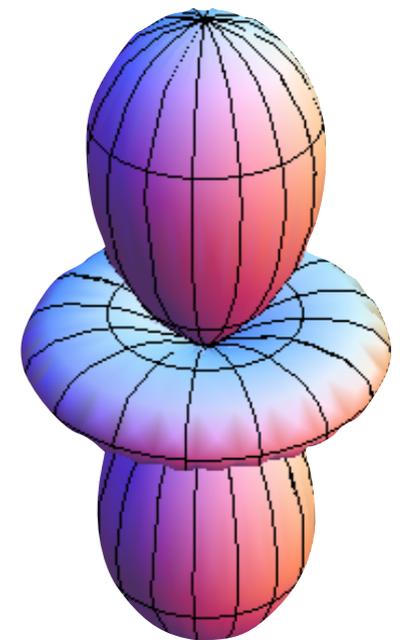


Sandro Wenzel

EPF Lausanne

Andreas Laeuchli

MPI PKS, Dresden



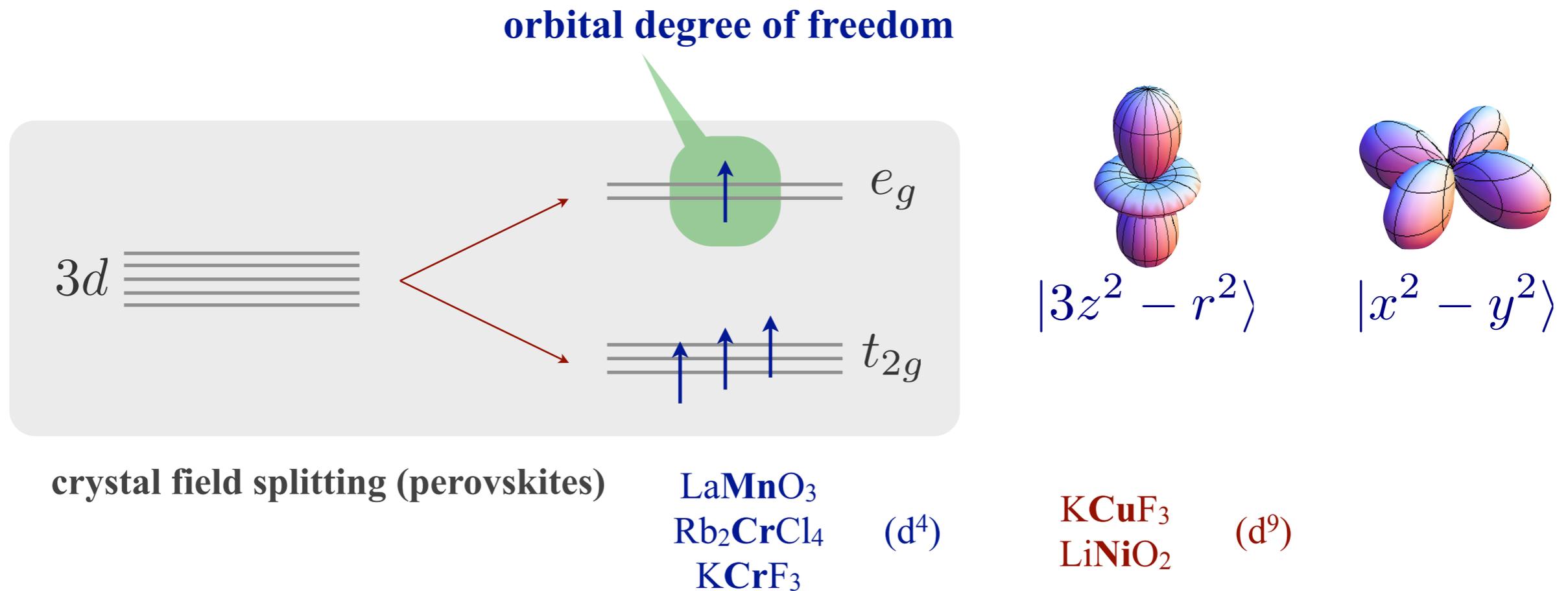
ÉCOLE POLYTECHNIQUE
FÉDÉRALE DE LAUSANNE



Mott insulators with partially filled d -shells

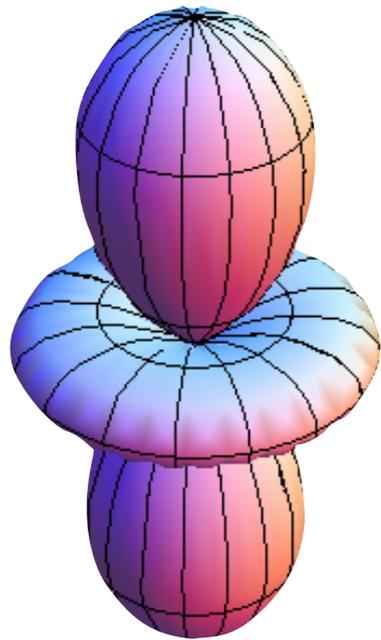
Mott insulating transition metal oxides with **partially filled $3d$ -shells** – such as the manganites – exhibit rich phase diagrams.

Non-trivial interplay of spin, charge, and orbital degrees of freedom.

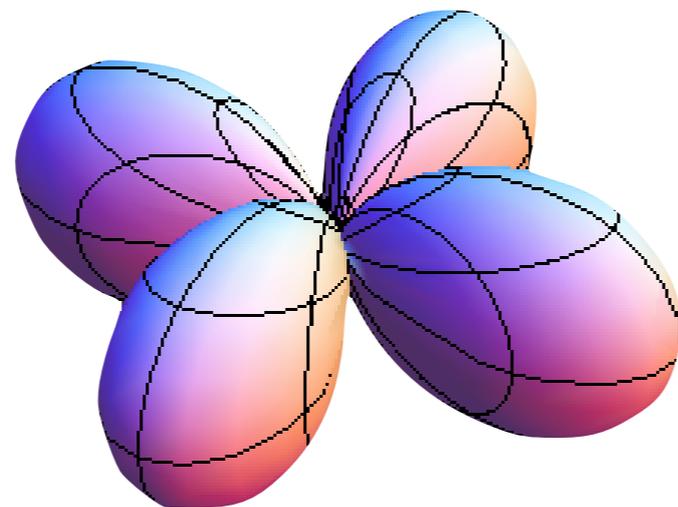


Description of orbital degree of freedom

General orbital state is superposition of e_g basis states

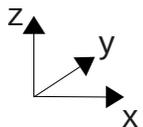


$$|3z^2 - r^2\rangle \quad |1\rangle$$



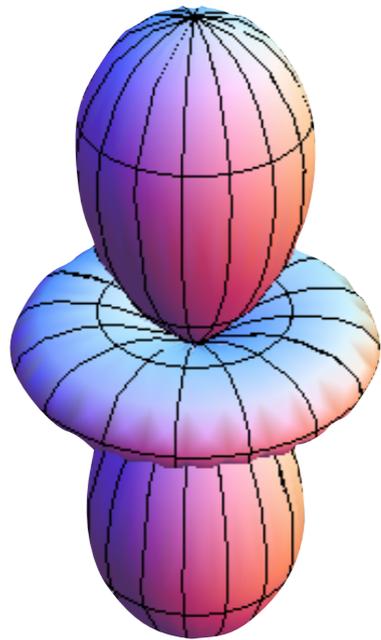
$$|x^2 - y^2\rangle \quad |2\rangle$$

$$|\psi\rangle = e^{i\phi} \cos(\theta/2)|1\rangle + e^{-i\phi} \sin(\theta/2)|2\rangle$$

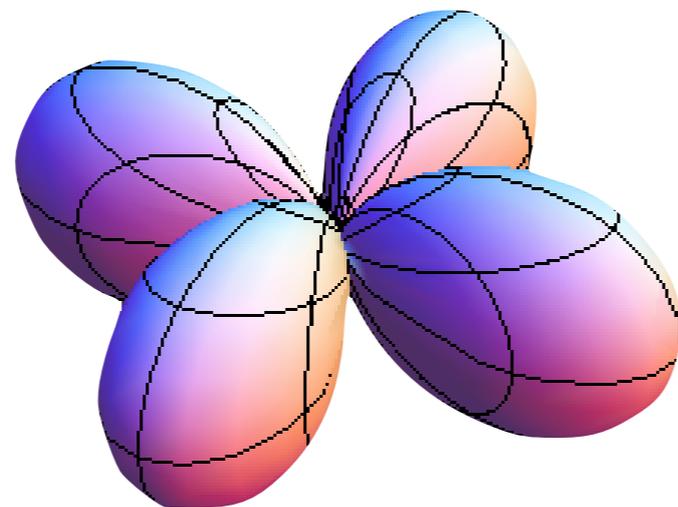


Description of orbital degree of freedom

General orbital state is superposition of e_g basis states

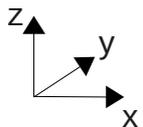


$$|3z^2 - r^2\rangle \quad |1\rangle$$



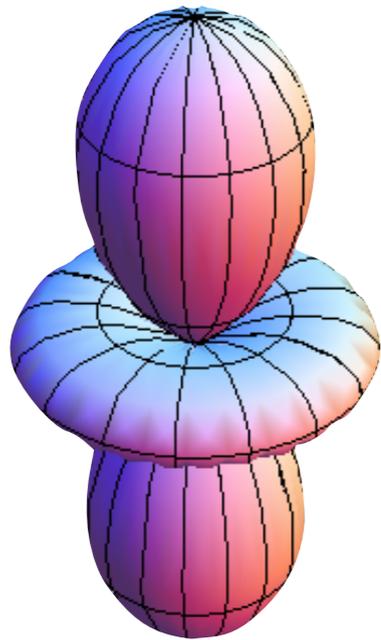
$$|x^2 - y^2\rangle \quad |2\rangle$$

$$|\psi\rangle = e^{i\phi} \cos(\theta/2) |1\rangle + e^{-i\phi} \sin(\theta/2) |2\rangle$$

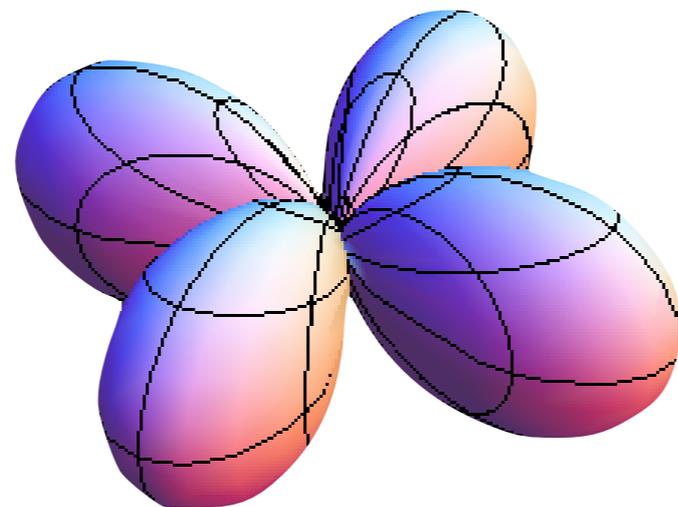


Description of orbital degree of freedom

General orbital state is superposition of e_g basis states

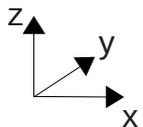


$$|3z^2 - r^2\rangle \quad |1\rangle$$



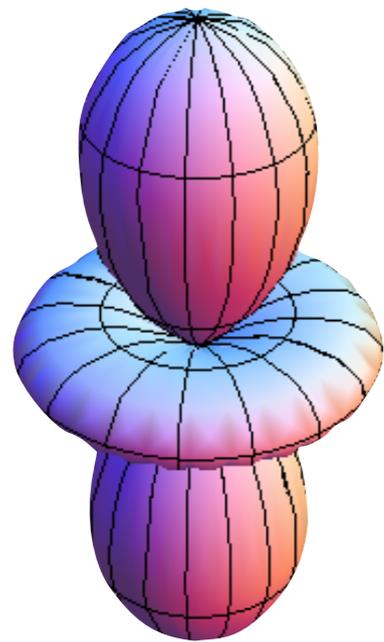
$$|x^2 - y^2\rangle \quad |2\rangle$$

$$|\psi\rangle = \cos(\theta/2)|1\rangle + \sin(\theta/2)|2\rangle$$

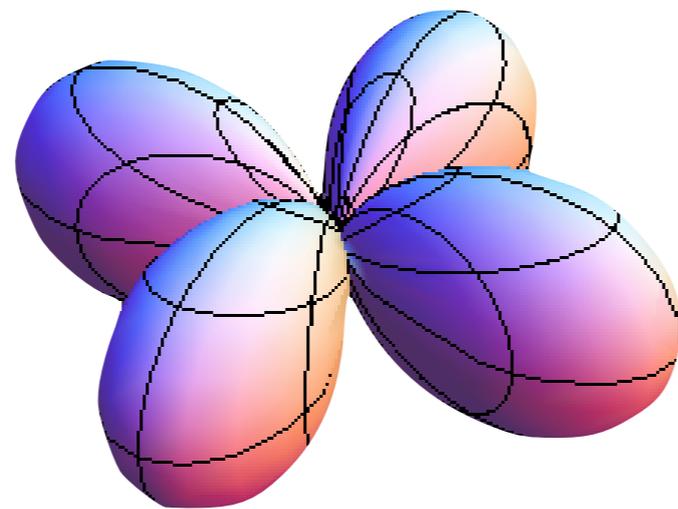


Description of orbital degree of freedom

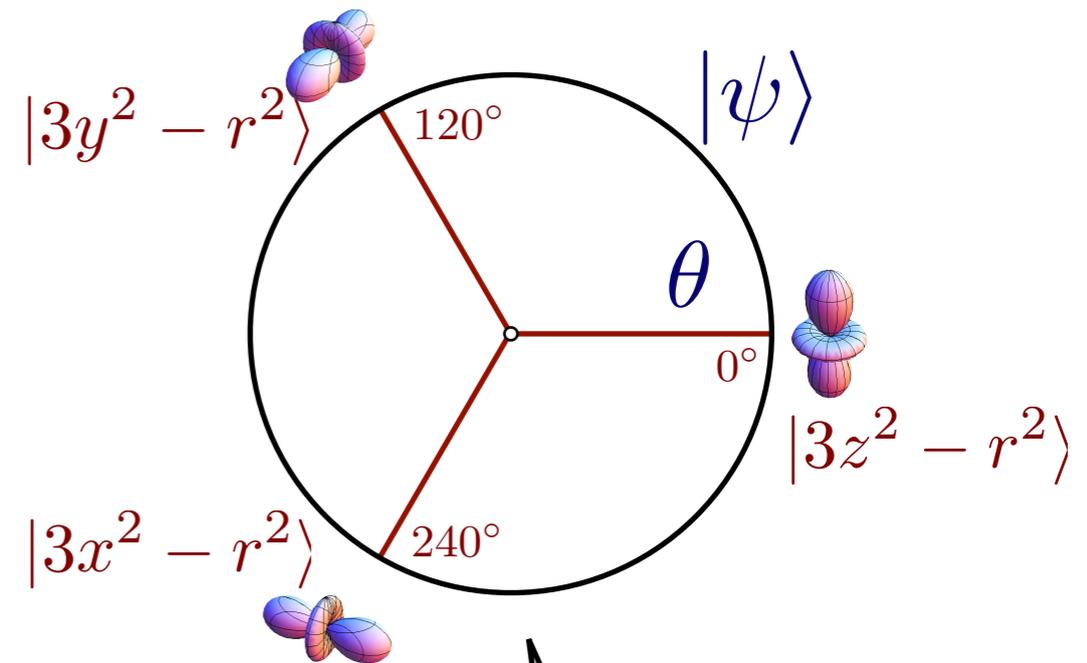
General orbital state is superposition of e_g basis states



$$|3z^2 - r^2\rangle \quad |1\rangle$$

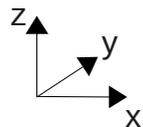


$$|x^2 - y^2\rangle \quad |2\rangle$$



$$|\psi\rangle = \cos(\theta/2)|1\rangle + \sin(\theta/2)|2\rangle$$

defines a XY spin
 $T = (T_z, T_x)$

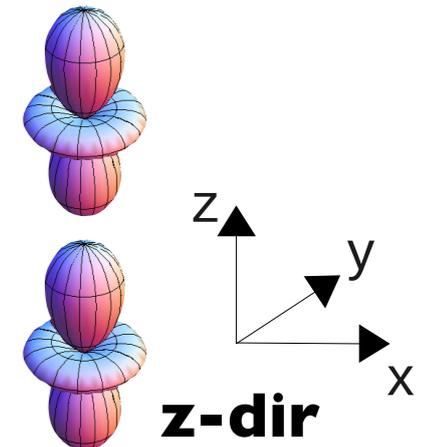
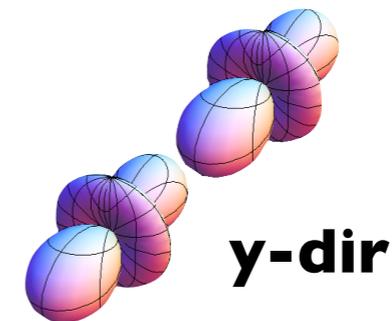
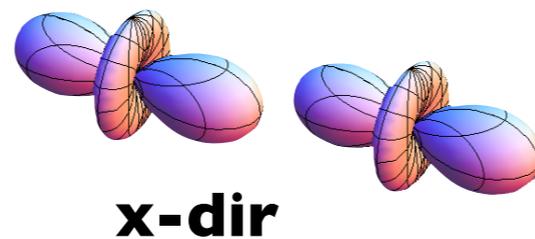


Orbital interactions: The 120° model

J. van den Brink, *New Journal of Physics* **6**, 210 (2004).

Jahn-Teller effect: orbital interactions mediated by phonons

favourable nearest neighbour configurations:

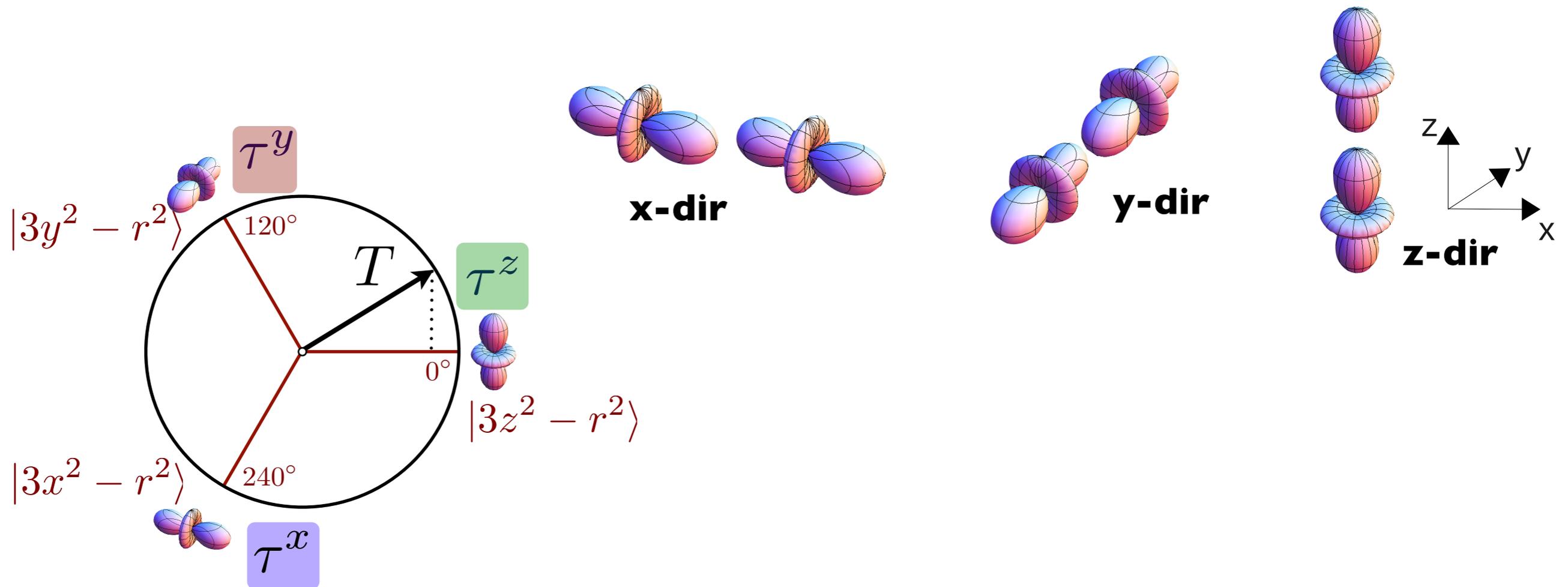


Orbital interactions: The 120° model

J. van den Brink, New Journal of Physics 6, 210 (2004).

Jahn-Teller effect: orbital interactions mediated by phonons

favourable nearest neighbour configurations:

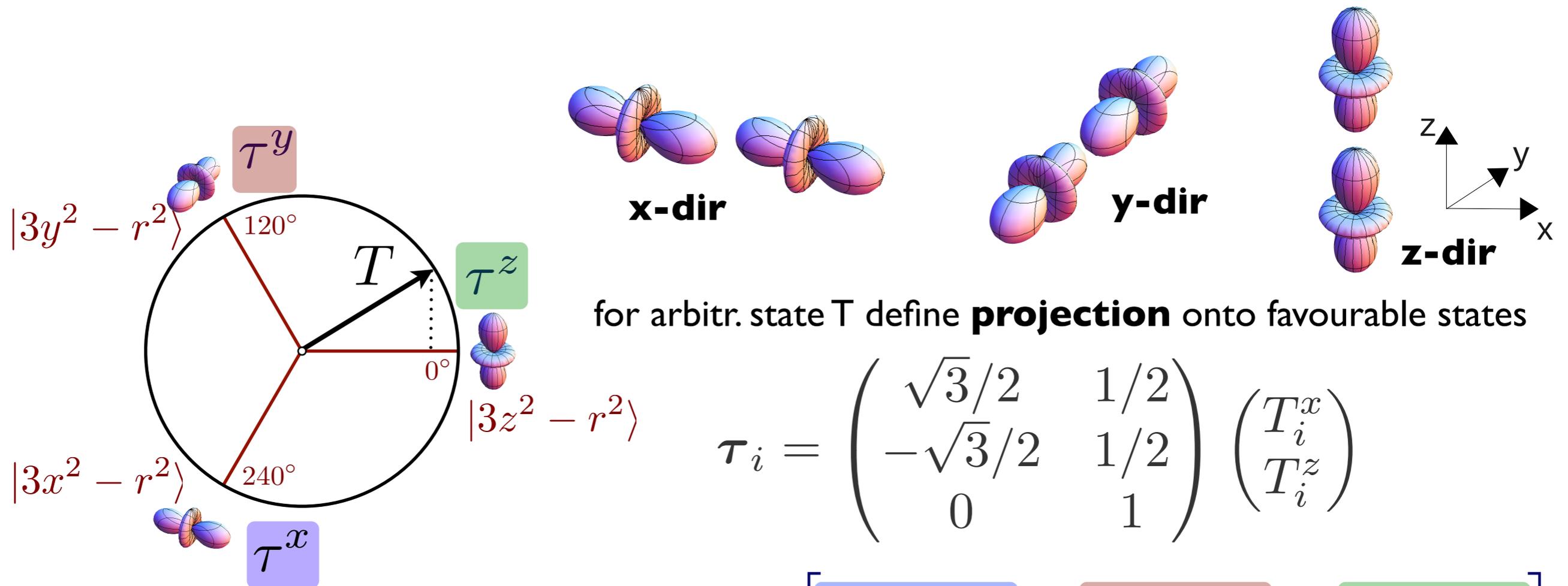


Orbital interactions: The 120° model

J. van den Brink, New Journal of Physics 6, 210 (2004).

Jahn-Teller effect: orbital interactions mediated by phonons

favourable nearest neighbour configurations:



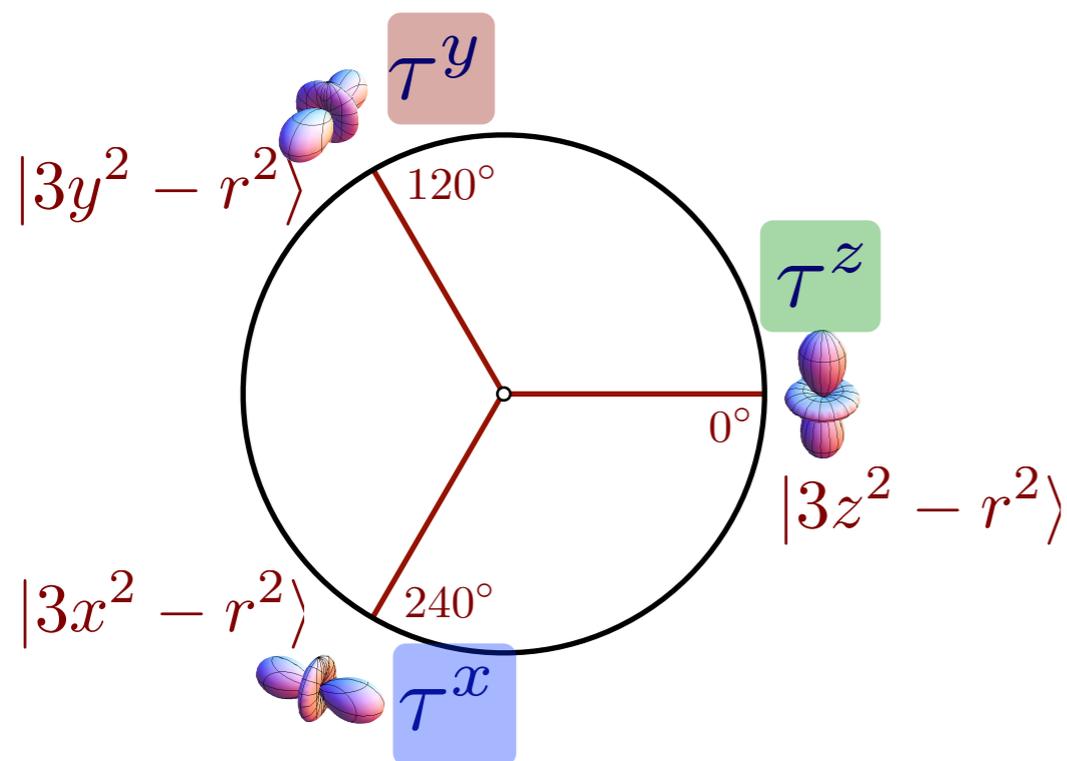
for arbitr. state T define **projection** onto favourable states

$$\tau_i = \begin{pmatrix} \sqrt{3}/2 & 1/2 \\ -\sqrt{3}/2 & 1/2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} T_i^x \\ T_i^z \end{pmatrix}$$

$$H_{120} = - \sum_i \left[\tau_i^x \tau_{i+\mathbf{e}_x}^x + \tau_i^y \tau_{i+\mathbf{e}_y}^y + \tau_i^z \tau_{i+\mathbf{e}_z}^z \right]$$

Critical properties of the classical 120° model

pseudospins are classical $O(2)$ spins



1. Is there collective ordering at finite T ?
2. What are the critical properties (universality)?

our approach: extensive Monte Carlo simulations!

$$H_{120} = \sum_i \left[\tau_i^x \tau_{i+\mathbf{e}_x}^x + \tau_i^y \tau_{i+\mathbf{e}_y}^y + \tau_i^z \tau_{i+\mathbf{e}_z}^z \right]$$

Zero temperature: degenerate ground states

M. Biskup *et al.*, *Comm. Math. Phys.* **255**, 253 (2005).
Z. Nussinov *et al.*, *Europhys. Lett.* **67**, 990 (2004).

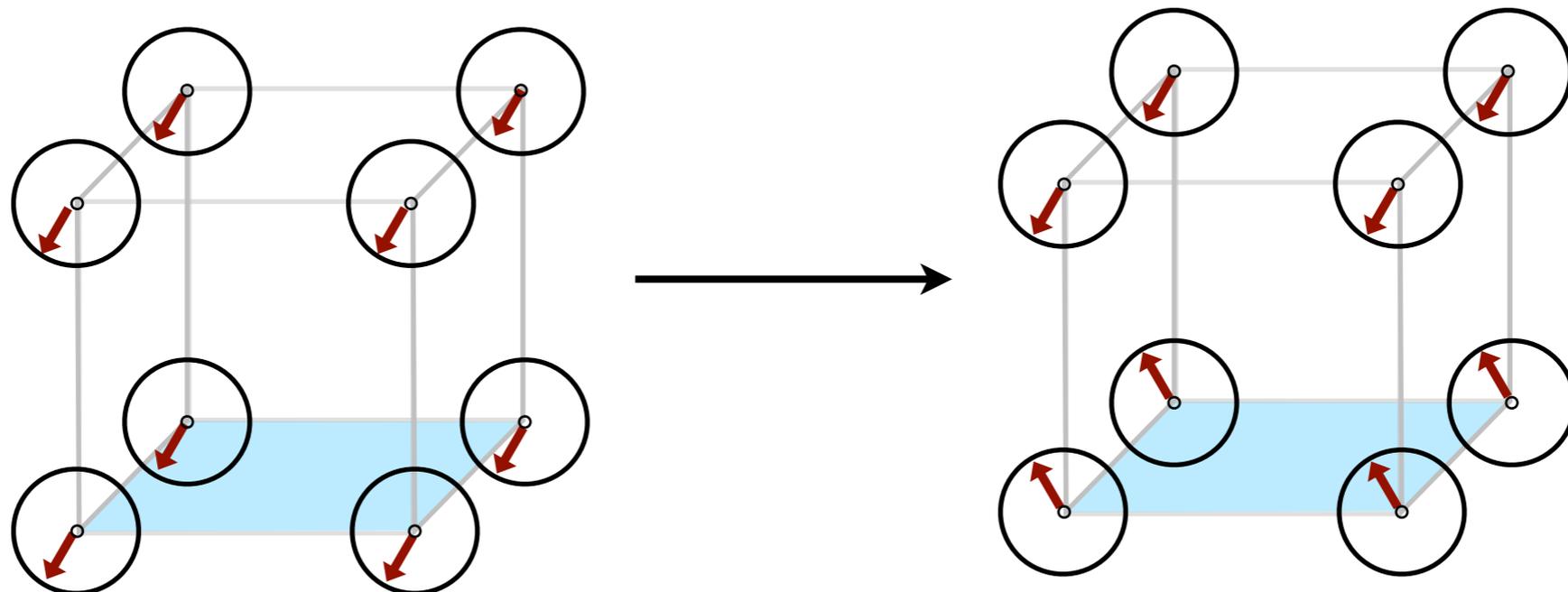
Emergent symmetries: U(1) and Z_2 symmetries

Ground-state manifold: infinite, but sub-extensive number of states

$$\mathcal{D} = 2^{3L}$$

example of Z_2 symmetry:

reflect all spins in xy plane at 0°



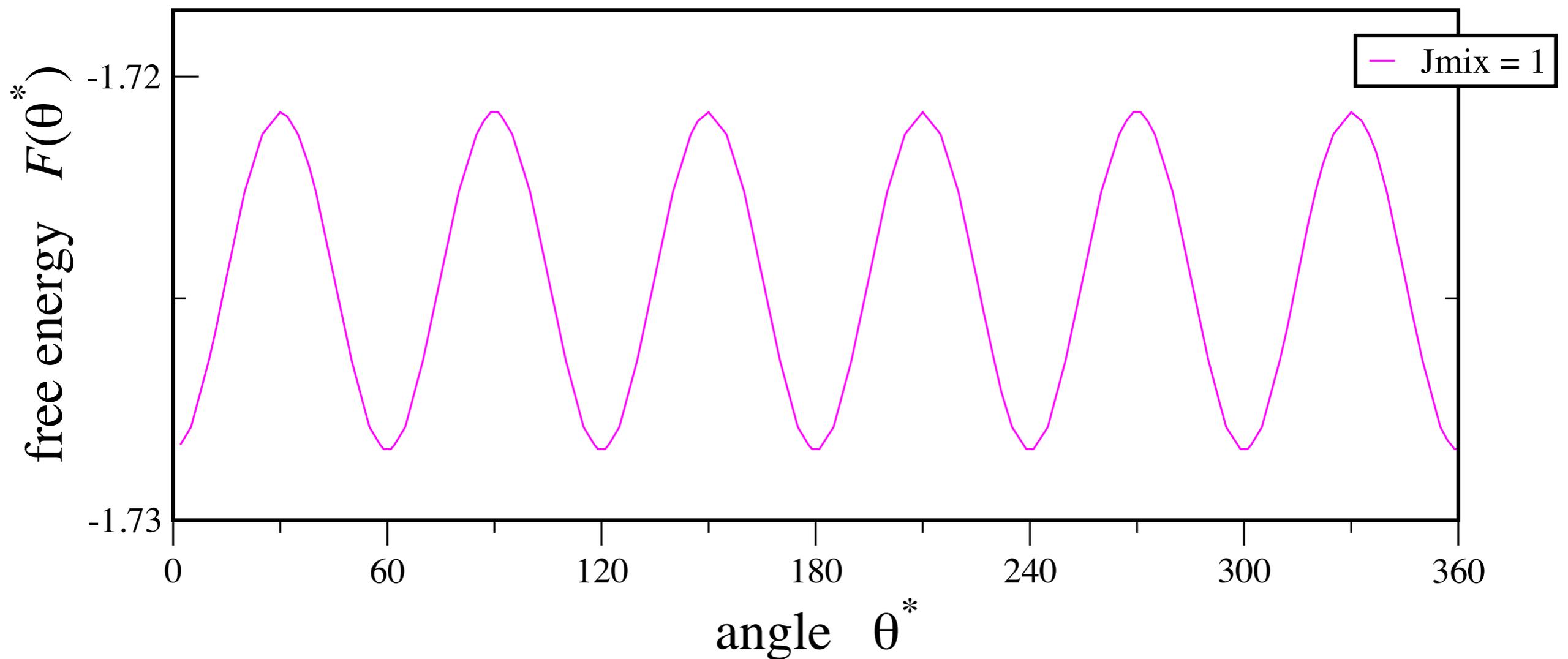
Low temperatures: Order by disorder

M. Biskup *et al.*, *Comm. Math. Phys.* **255**, 253 (2005).

Z. Nussinov *et al.*, *Europhys. Lett.* **67**, 990 (2004).

A.v. Rynbach, S. Todo, S. Trebst, *Phys. Rev. Lett.* **105**, 14640 (2010).

Spin-wave approximation: expansion in fluctuations around ordered state with $\delta\theta_i = \theta_i - \theta^*$
 $\theta_i = \theta^*$ at each site.



Low temperatures: Order by disorder

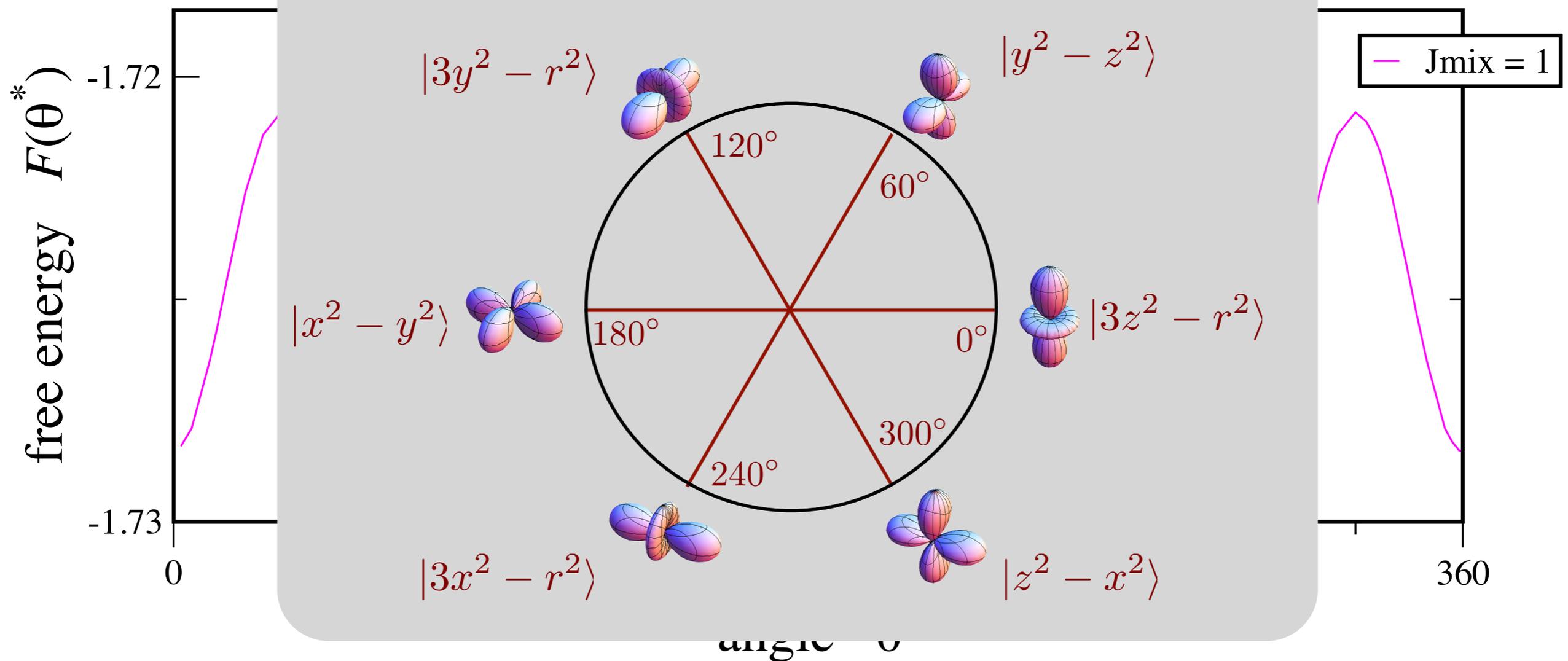
M. Biskup *et al.*, *Comm. Math. Phys.* **255**, 253 (2005).

Z. Nussinov *et al.*, *Europhys. Lett.* **67**, 990 (2004).

A.v. Rynbach, S. Todo, S. Trebst, *Phys. Rev. Lett.* **105**, 14640 (2010).

Spin-wave approximation: expansion in fluctuations around ordered state with $\delta\theta_i = \theta_i - \theta^*$
 $\theta_i = \theta^*$ at each site.

proposed collective ordering states

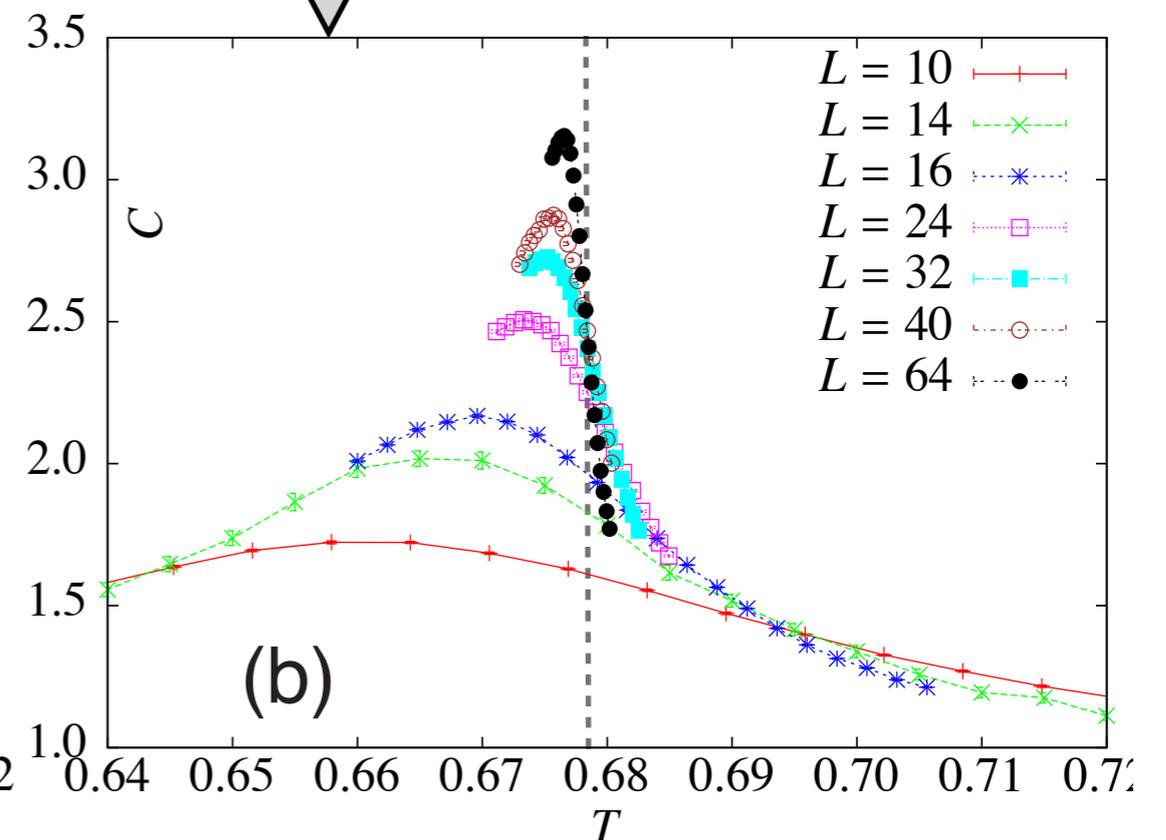
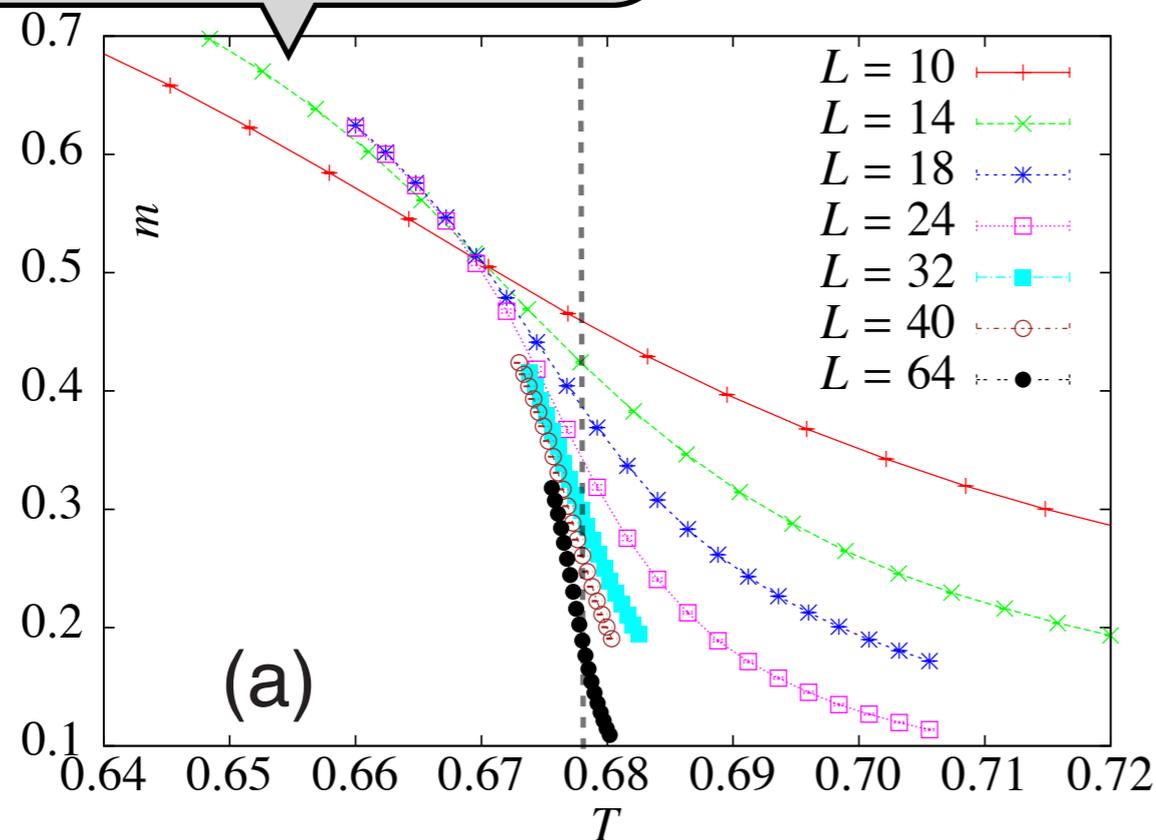


Orbital-ordering phase transition

$$m = \frac{1}{N} \sqrt{\left(\sum_i T_i^x\right)^2 + \left(\sum_i T_i^z\right)^2}$$

heat capacity

$$T_c = 0.6775(1)$$

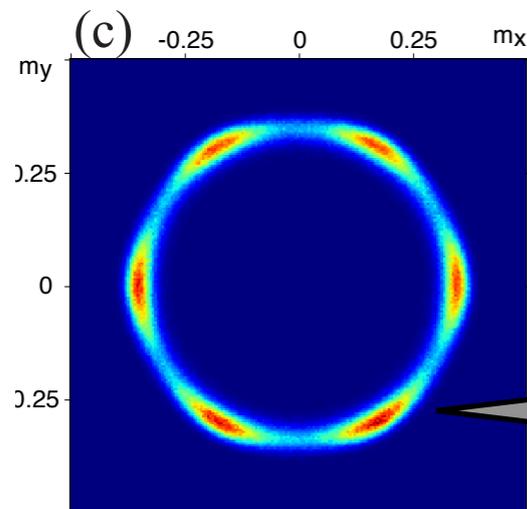
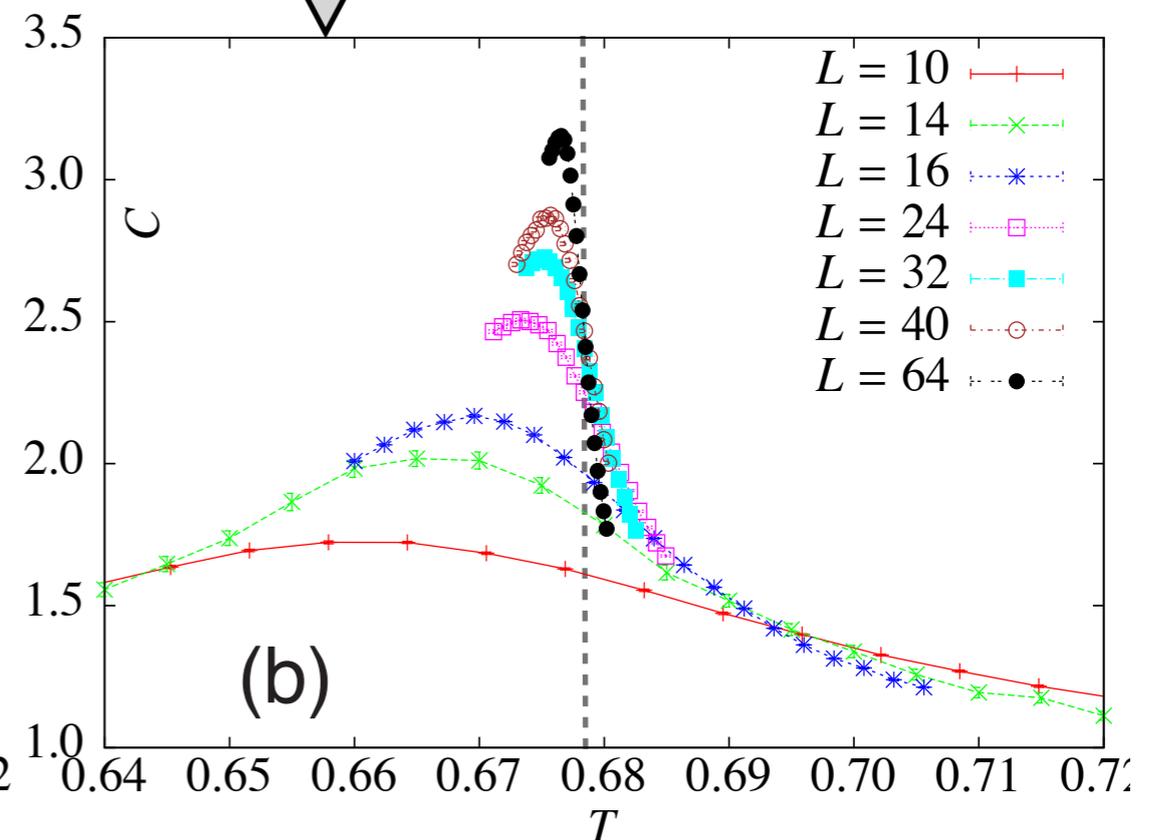
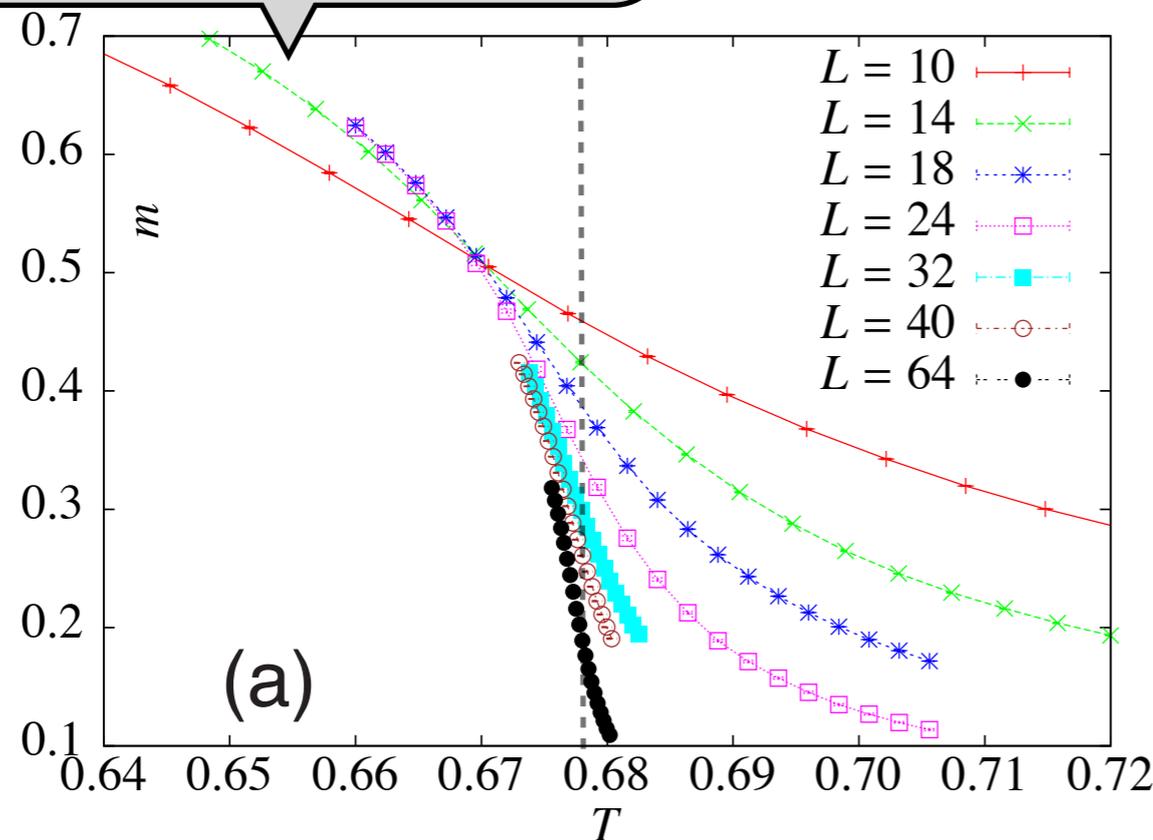


Orbital-ordering phase transition

$$m = \frac{1}{N} \sqrt{\left(\sum_i T_i^x\right)^2 + \left(\sum_i T_i^z\right)^2}$$

heat capacity

$$T_c = 0.6775(1)$$



distribution of order parameter

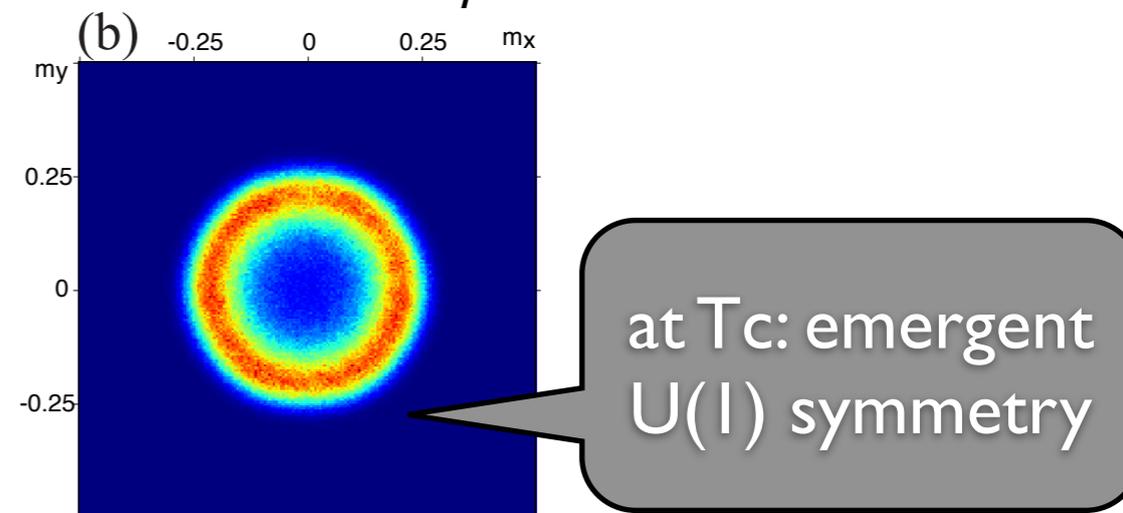
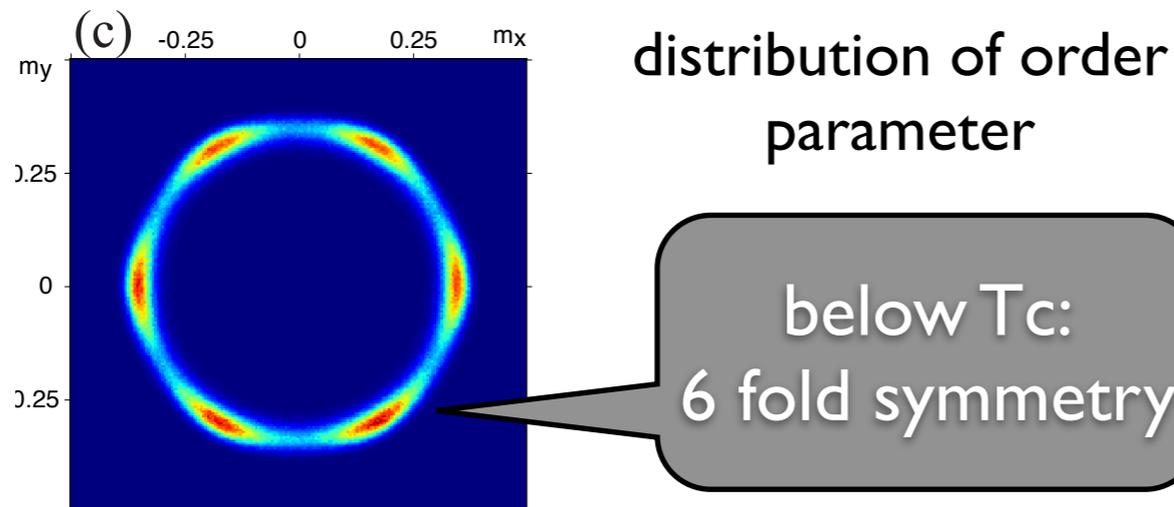
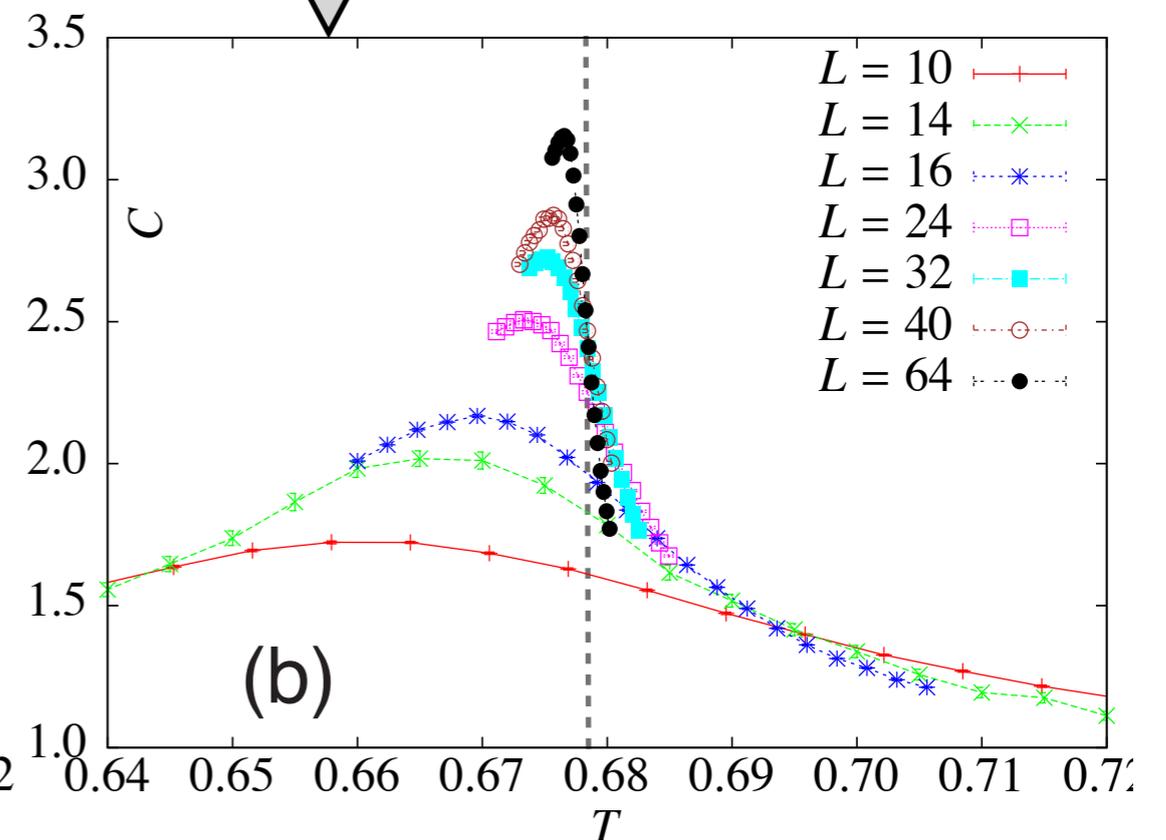
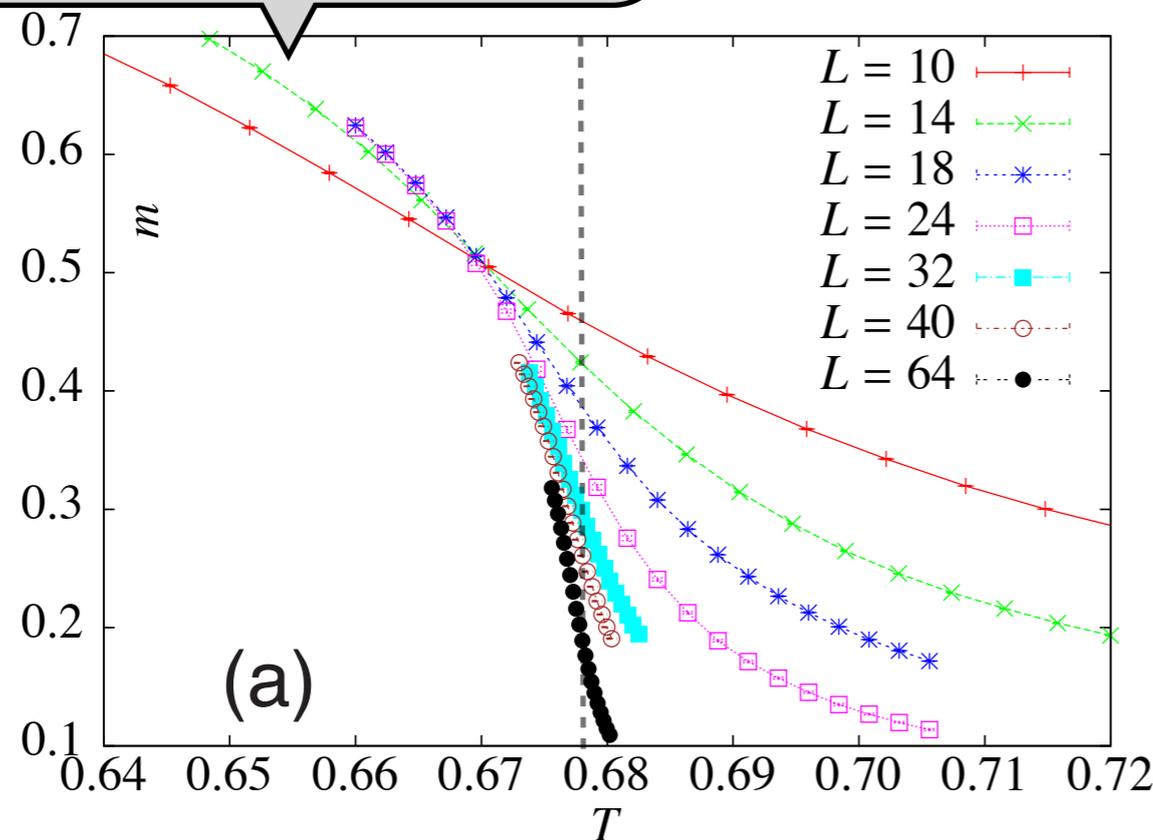
below T_c :
6 fold symmetry

Orbital-ordering phase transition

$$m = \frac{1}{N} \sqrt{\left(\sum_i T_i^x\right)^2 + \left(\sum_i T_i^z\right)^2}$$

heat capacity

$$T_c = 0.6775(1)$$



Orbital-ordering phase transition

$$m = \frac{1}{N} \sqrt{\left(\sum_i T_i^x\right)^2 + \left(\sum_i T_i^z\right)^2}$$

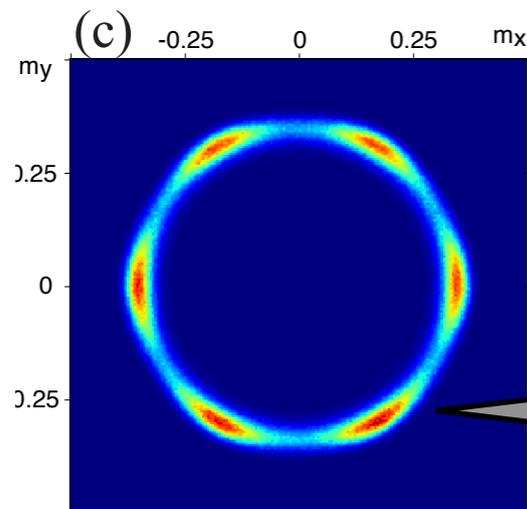
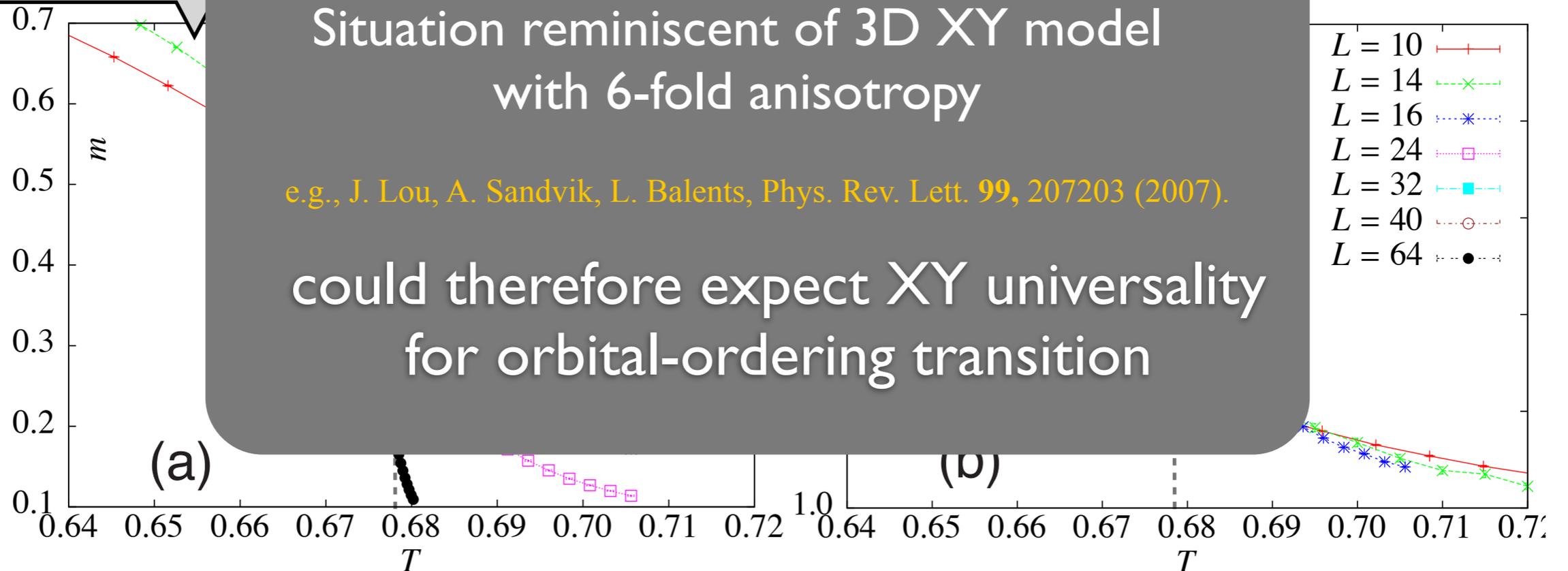
heat capacity

0.6775(1)

Situation reminiscent of 3D XY model
with 6-fold anisotropy

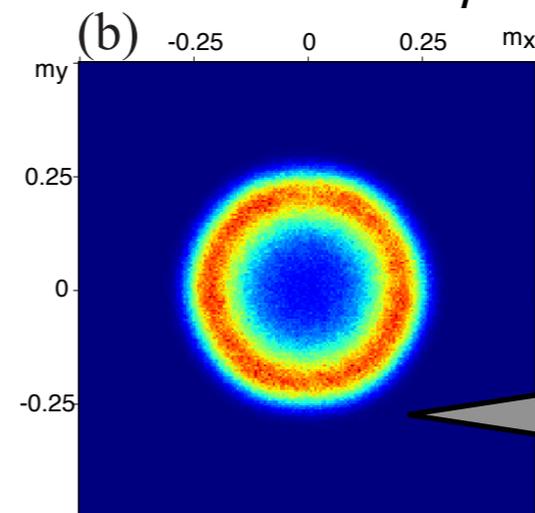
e.g., J. Lou, A. Sandvik, L. Balents, *Phys. Rev. Lett.* **99**, 207203 (2007).

could therefore expect XY universality
for orbital-ordering transition



distribution of order
parameter

below T_c :
6 fold symmetry

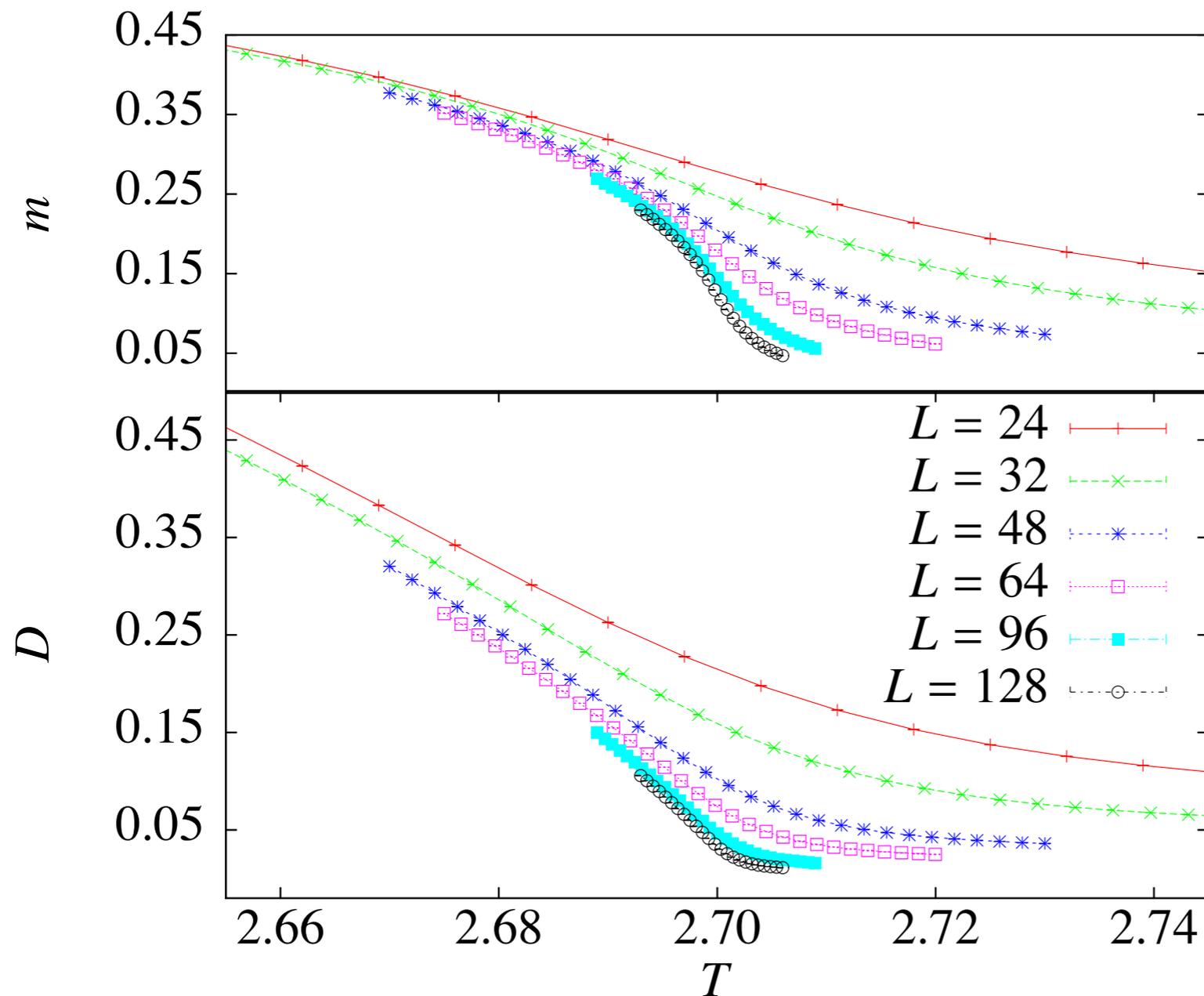


at T_c : emergent
 $U(1)$ symmetry

Additional directional ordering

a key feature not present in an ordinary **XY** spin model:

$$D = (1/N) \left([E_x - E_y]^2 + [E_y - E_z]^2 + [E_z - E_x]^2 \right)$$



simultaneous
onset of orbital and
directional order

anisotropic
fluctuations in
ordered state

Critical exponents of the 120° model

$$\xi \sim |t|^{-\nu}$$

orbital-orbital
correlation length

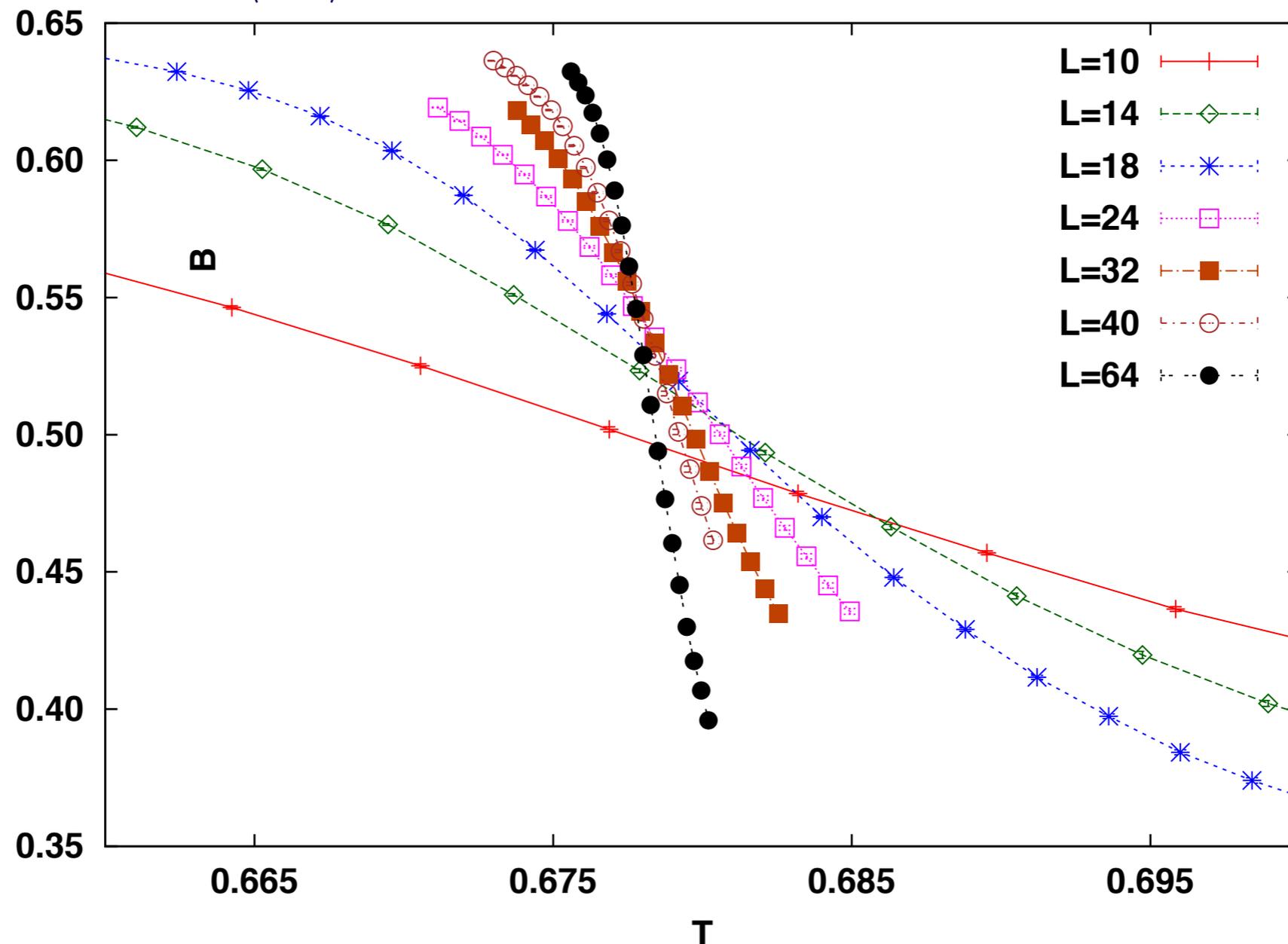
$$G(r) = r^{-2+d-\eta}$$

decay of orbital-orbital correlations
at critical temperature

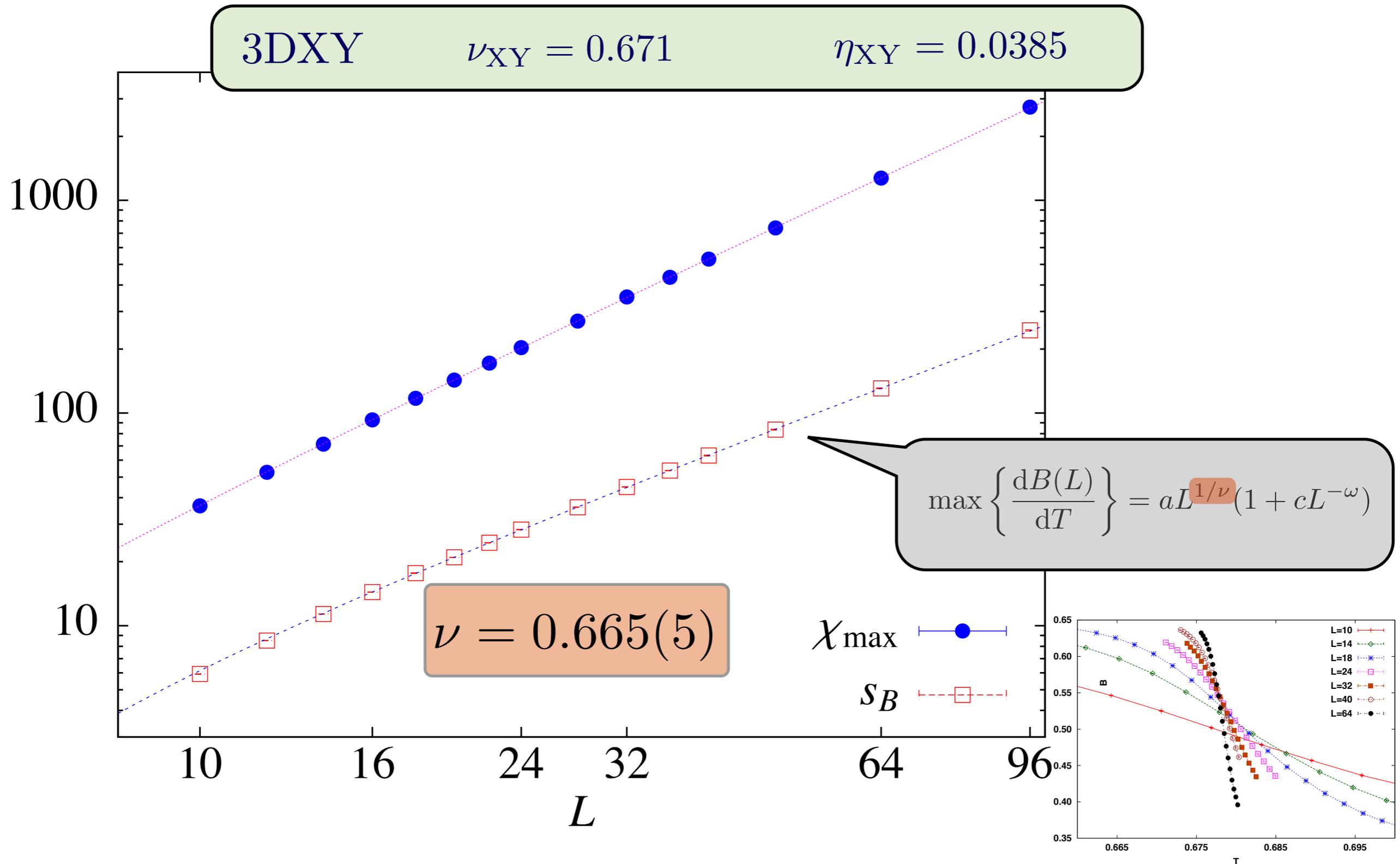
Critical exponents of the 120° model

$$\xi \sim |t|^{-\nu}$$

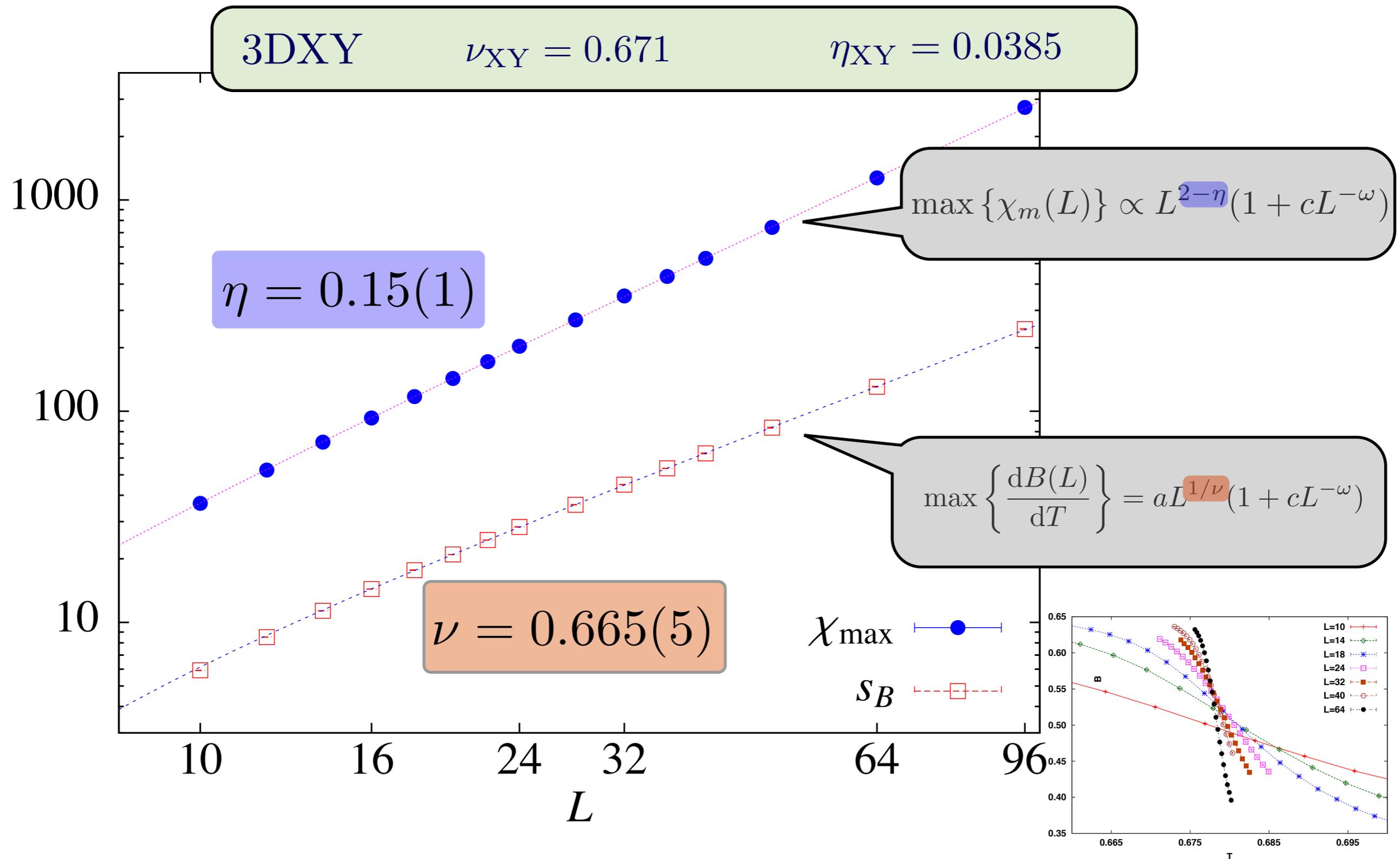
$$B = \frac{\langle m^4 \rangle}{\langle m^2 \rangle^2} \quad \max \left\{ \frac{dB(L)}{dT} \right\} = aL^{1/\nu} (1 + cL^{-\omega})$$



Critical exponents of the 120° model



Critical exponents of the 120° model



The discrete 120° (clock) model

Want a
simplified
model ...

**standard XY
model**

$$H_{XY} = \sum_{\langle ij \rangle} \mathbf{S}_i \mathbf{S}_j$$

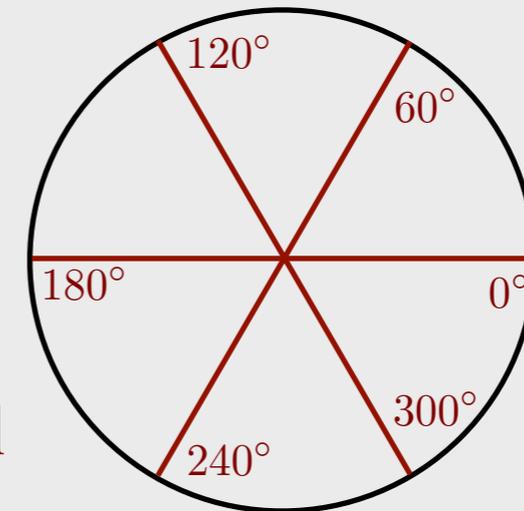
**continuous
spins**

**120°
model**

T

XY Universality
class

**6-state
clock model**



discrete spins

**120°-clock
model**

The discrete 120° (clock) model

Want a
simplified
model ...

discrete models have the same
ordered state on spin level!

**standard XY
model**

$$H_{XY} = \sum_{\langle ij \rangle} S_i S_j$$

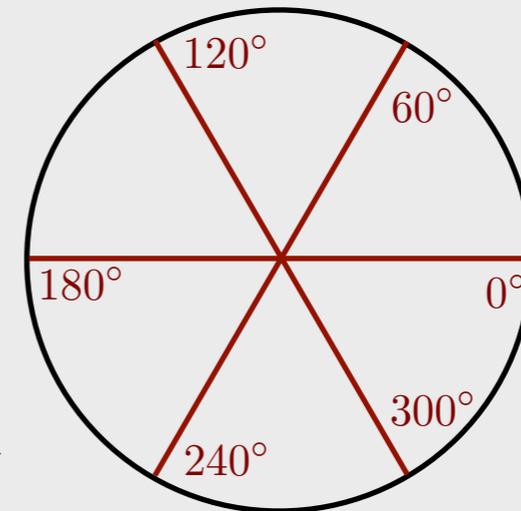
**continuous
spins**

**120°
model**

T

XY Universality
class

**6-state
clock model**



discrete spins

**120°-clock
model**

The discrete 120° (clock) model

Want a
simplified
model ...

**standard XY
model**

$$H_{XY} = \sum_{\langle ij \rangle} \mathbf{S}_i \mathbf{S}_j$$

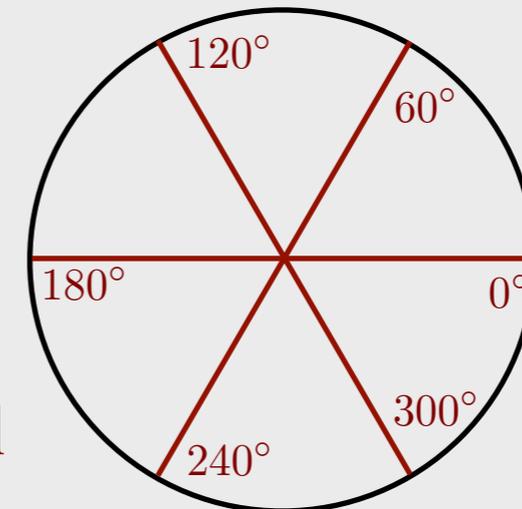
**continuous
spins**

**120°
model**

T

XY Universality
class

**6-state
clock model**



**120°-clock
model**

discrete spins

discrete models have the same
ordered state on spin level!

discrete structure allows for
VERY efficient MC simulations:
can reach $N=128^3$ sites
(2.09 million spins !)

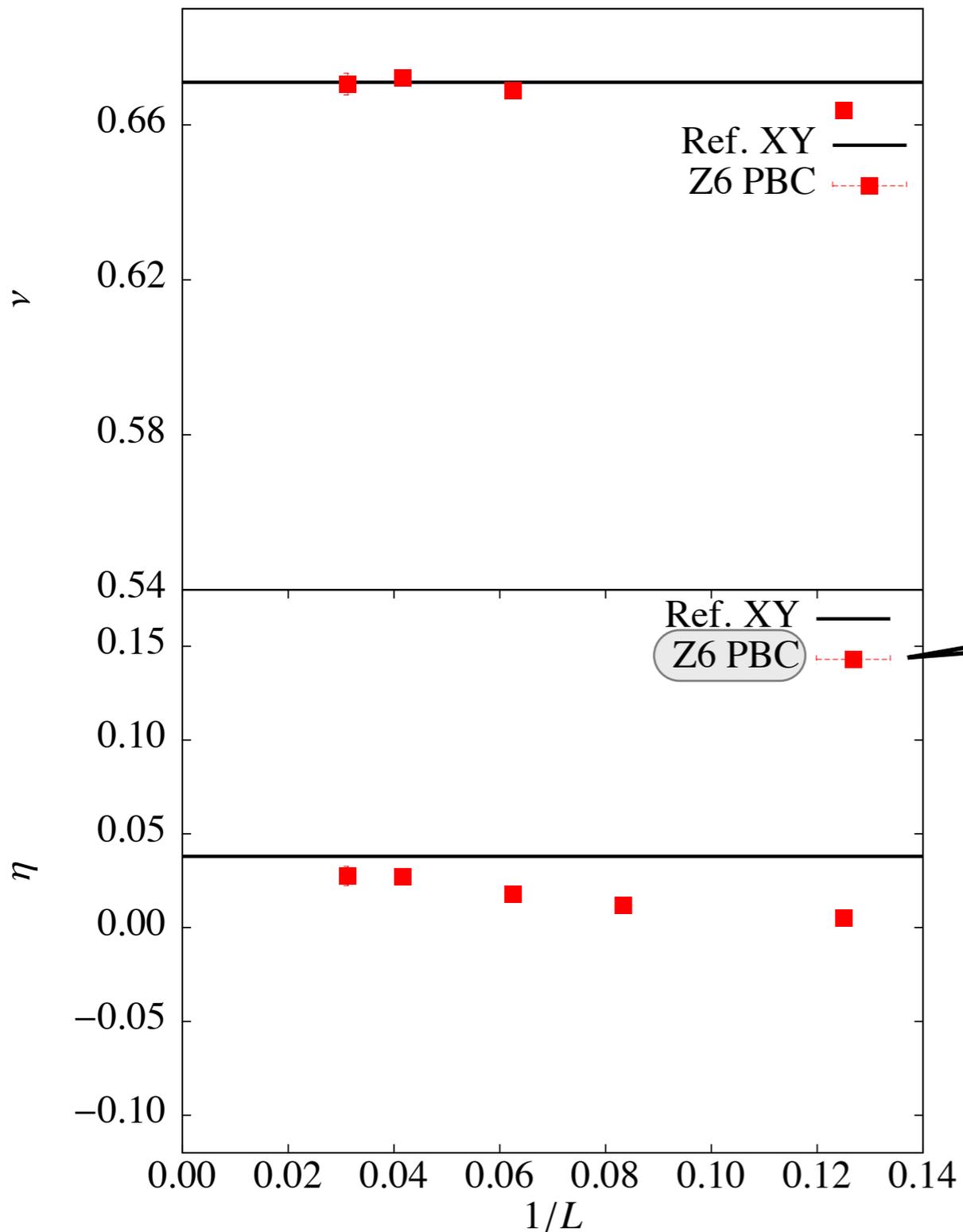
Detailed comparison of FSS

systematic comparison of critical scaling:
scaling of **effective** critical exponents

$$\chi(L) \propto L^{2-\eta} (1 + \dots)$$



$$\eta = 2 - \ln(\chi(2L)/\chi(L)) / \ln(2)$$



6 state clock model

XY?



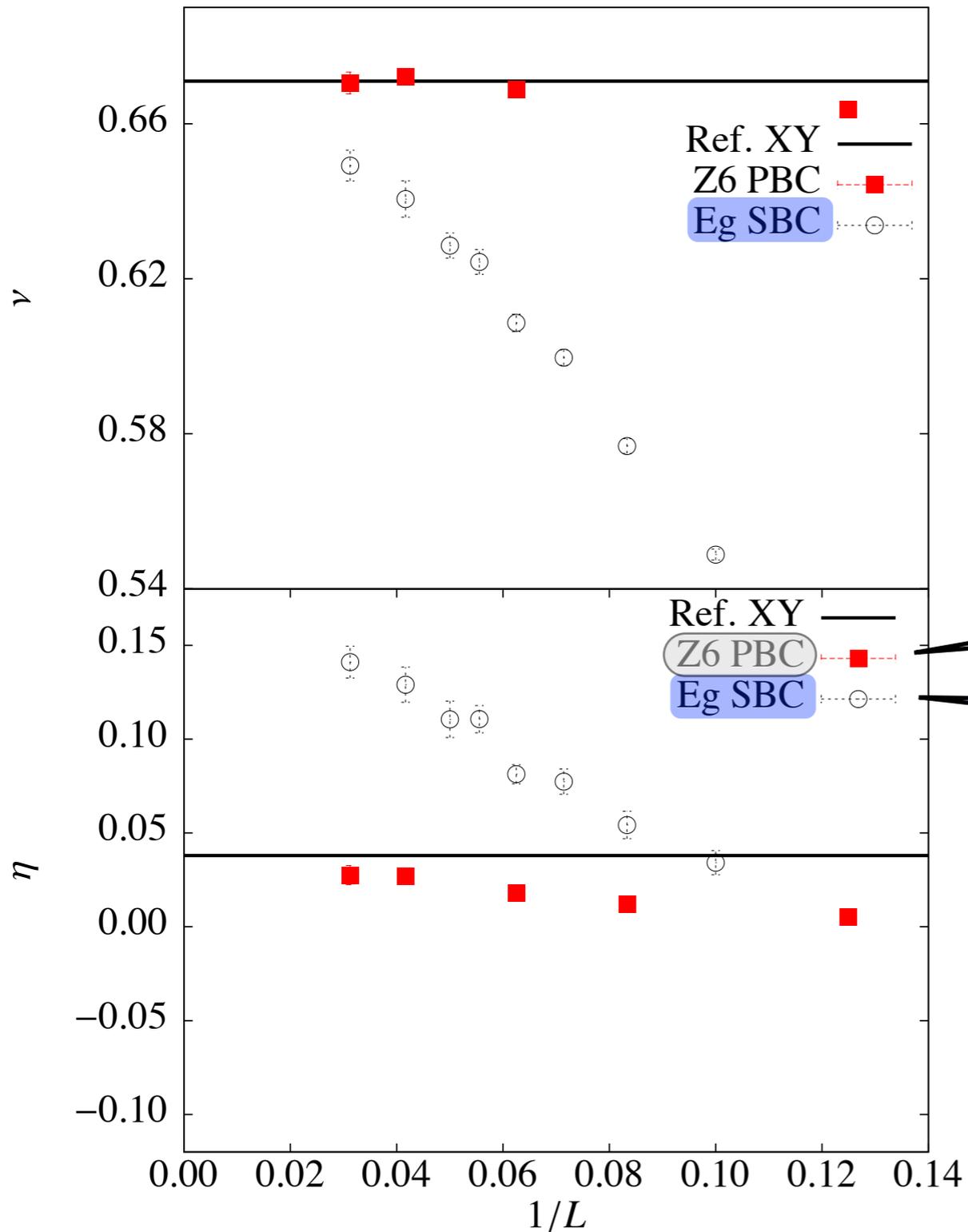
Detailed comparison of FSS

systematic comparison of critical scaling:
scaling of **effective** critical exponents

$$\chi(L) \propto L^{2-\eta} (1 + \dots)$$



$$\eta = 2 - \ln(\chi(2L)/\chi(L)) / \ln(2)$$



6 state clock model

XY?



120° model



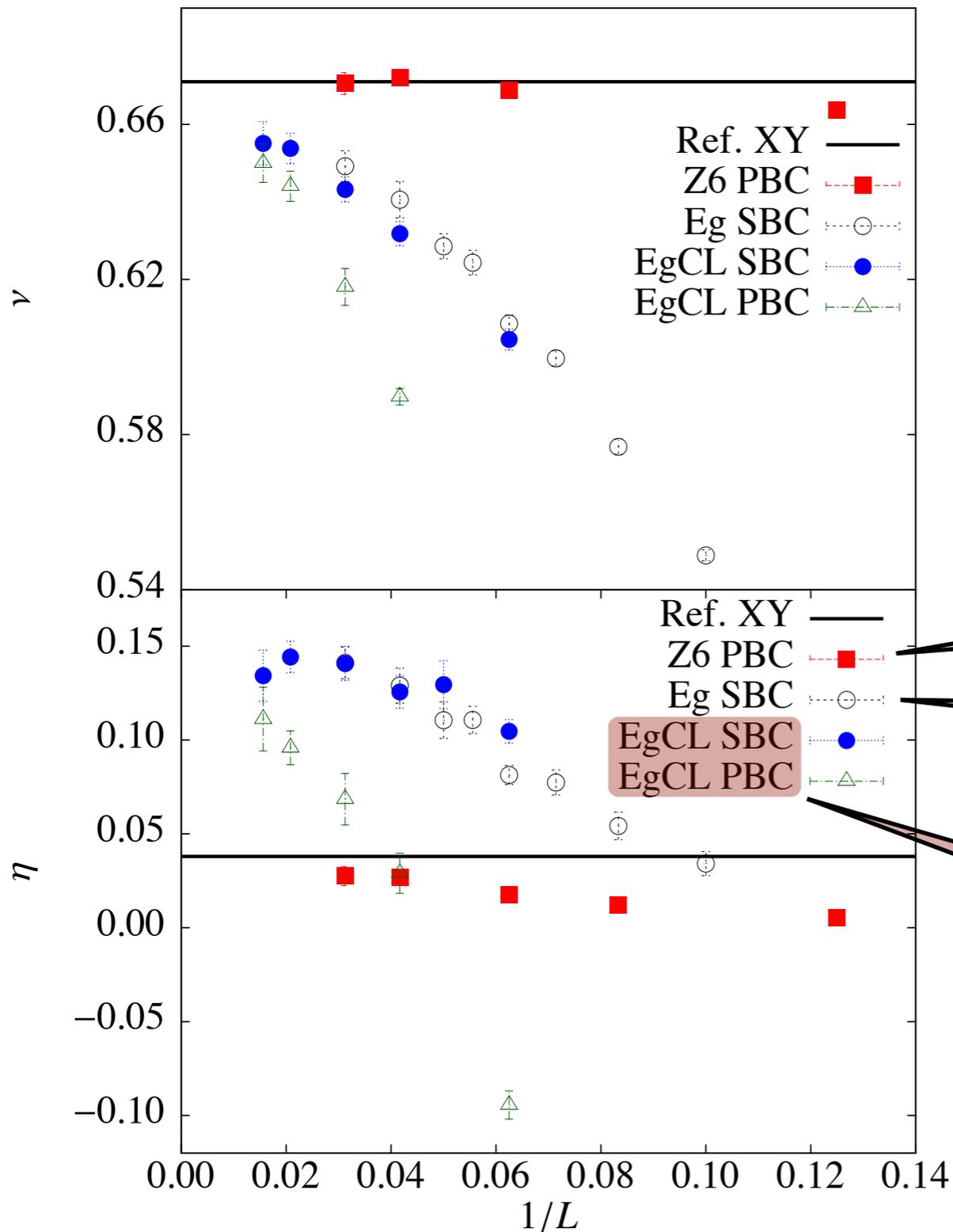
Detailed comparison of FSS

systematic comparison of critical scaling:
scaling of **effective** critical exponents

$$\chi(L) \propto L^{2-\eta} (1 + \dots)$$



$$\eta = 2 - \ln(\chi(2L)/\chi(L)) / \ln(2)$$



6 state clock model

XY?



120° model



120°-clock model
(for multiple bc)



Summary

first large scale Monte Carlo study of the prototypical classical 120° model for orbital ordering

- **ordering nature:** verified orbital ordering with 6-fold degeneracy
- **criticality:** obtained critical exponents which differ from standard magnetic universality classes (XY)
- **consequence:** expect to see **novel physics** different from ordinary spin models of magnetisms (e.g. disorder, ...)
- **outlook:** write down a Ginzburg-Landau theory..., simulations on GPUs, ...

T. Tanaka, M. Matsumoto, and S. Ishihara, Phys. Rev. Lett. **95**, 267204 (2005)

Summary

first large scale Monte Carlo study of the prototypical classical 120° model for orbital ordering

- **ordering nature:** verified orbital ordering with 6-fold degeneracy
- **criticality:** obtained critical exponents which differ from standard magnetic universality classes (XY)
- **consequence:** expect to see **novel physics** different from ordinary spin models of magnetisms (e.g. disorder, ...)
- **outlook:** write down a Ginzburg-Landau theory..., simulations on GPUs, ...

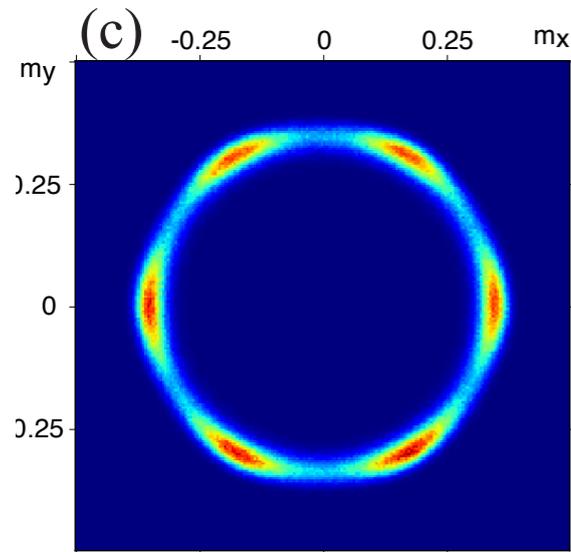
T. Tanaka, M. Matsumoto, and S. Ishihara, Phys. Rev. Lett. **95**, 267204 (2005)

!! THANKS !!

- ... for your attention
- ... to Simon Trebst for providing some slides/figures

Influence of U(1) length-scale for $T < T_c$

J. Lou, A. Sandvik, L. Balents, Phys. Rev. Lett. **99**, 207203 (2007).



consider order distribution function for $T < T_c$:

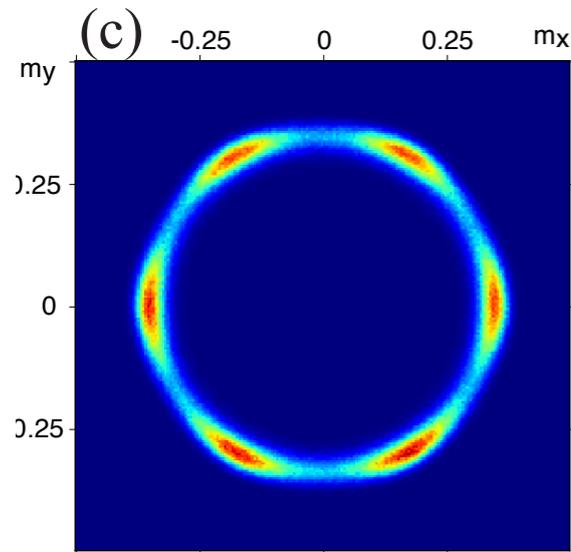
- $T = T_c$: U(1) symmetry
- $T < T_c$:
 - true 6-fold structure
 - continuous U(1)

$$L > \Lambda$$

$$L < \Lambda$$

Influence of U(1) length-scale for $T < T_c$

J. Lou, A. Sandvik, L. Balents, Phys. Rev. Lett. **99**, 207203 (2007).



consider order distribution function for $T < T_c$:

- $T = T_c$: U(1) symmetry
- $T < T_c$:
 - true 6-fold structure
 - continuous U(1)

$$L > \Lambda$$

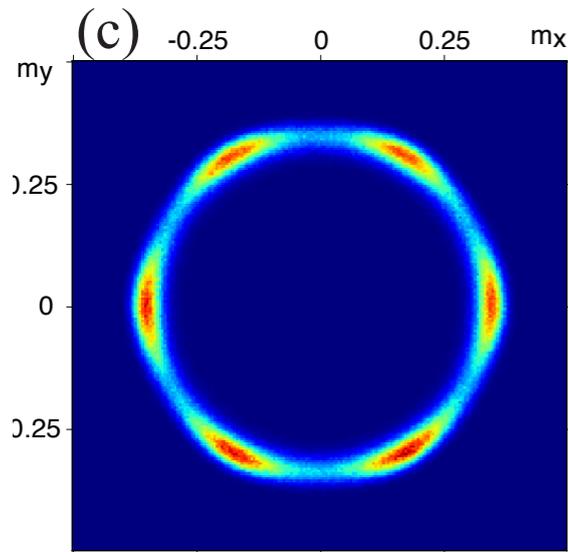
$$L < \Lambda$$

$$\Lambda \sim \xi^a, \quad a > 1$$

... a is universal

Influence of U(1) length-scale for $T < T_c$

J. Lou, A. Sandvik, L. Balents, Phys. Rev. Lett. **99**, 207203 (2007).



consider order distribution function for $T < T_c$:

- $T = T_c$: U(1) symmetry
- $T < T_c$:
 - true 6-fold structure
 - continuous U(1)

$$L > \Lambda$$

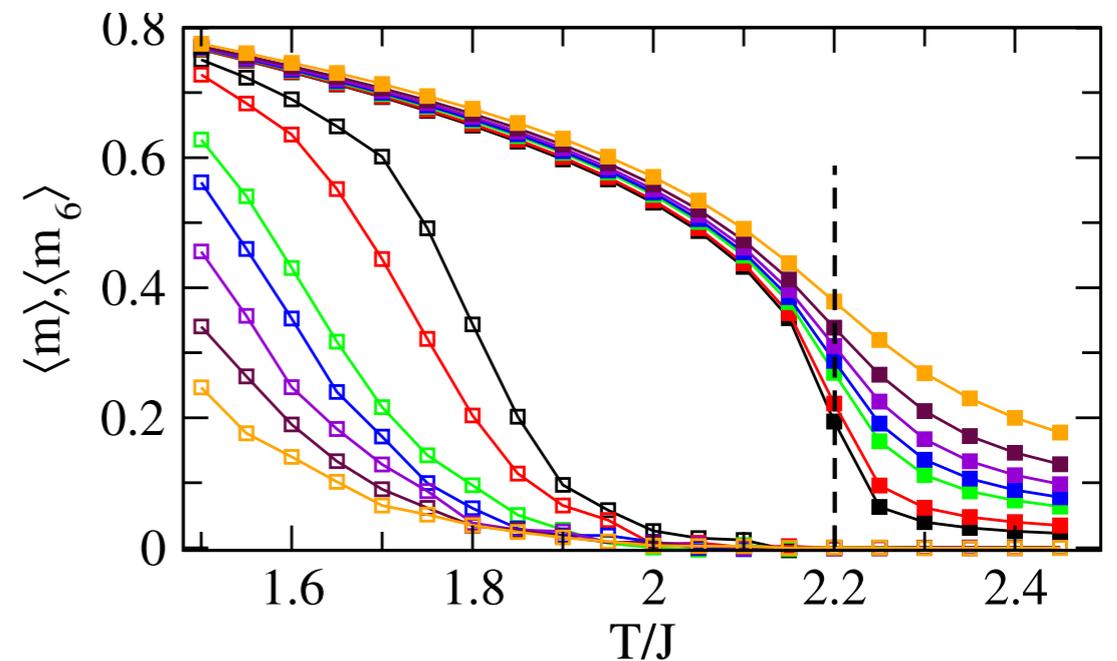
$$L < \Lambda$$

$$\Lambda \sim \xi^a, \quad a > 1$$

... a is universal

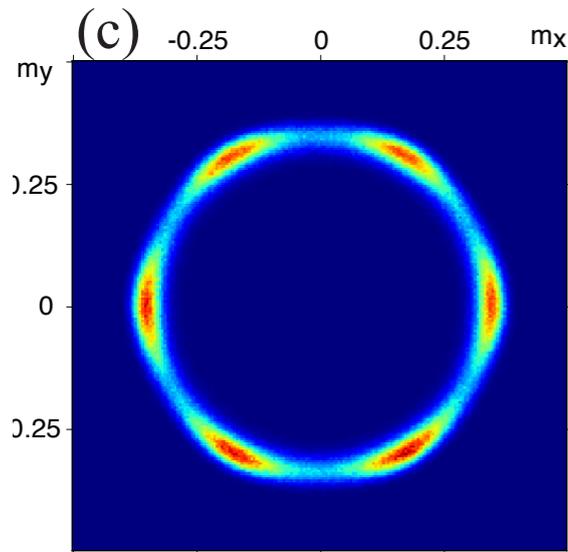
$$\langle m \rangle = \int_0^1 dr \int_0^{2\pi} d\theta r^2 P(r, \theta) \quad \xi$$

$$\langle m_6 \rangle = \int_0^1 dr \int_0^{2\pi} d\theta r^2 P(r, \theta) \cos(6\theta) \quad \Lambda$$



Influence of U(1) length-scale for $T < T_c$

J. Lou, A. Sandvik, L. Balents, Phys. Rev. Lett. **99**, 207203 (2007).



consider order distribution function for $T < T_c$:

- $T = T_c$: U(1) symmetry
- $T < T_c$:
 - true 6-fold structure
 - continuous U(1)

$$L > \Lambda$$

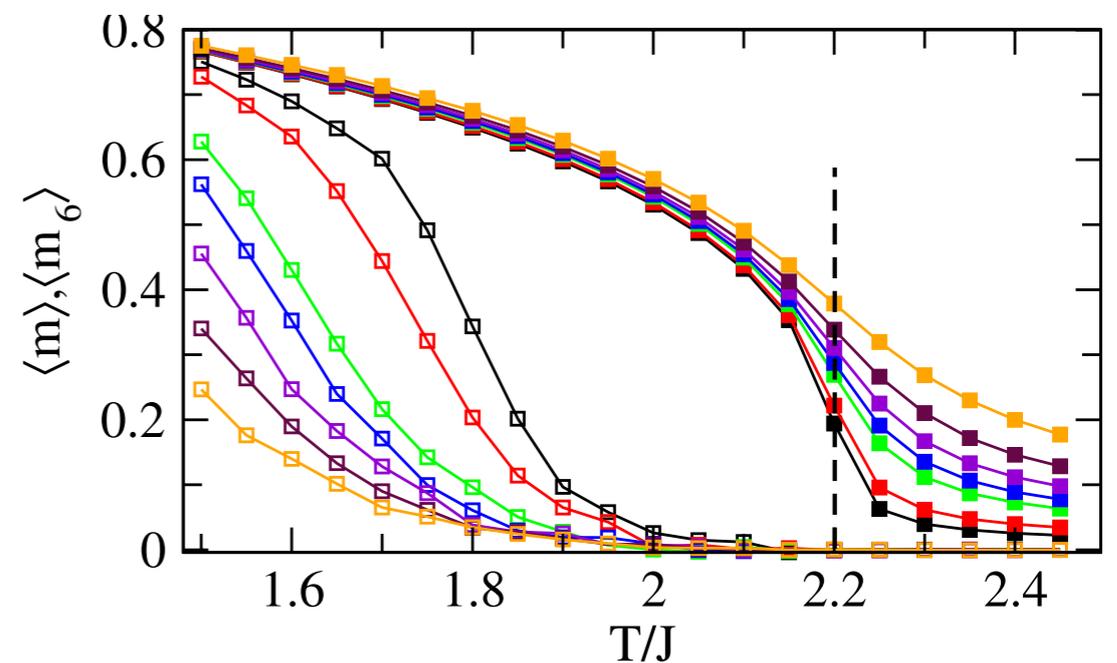
$$L < \Lambda$$

$$\Lambda \sim \xi^a, \quad a > 1$$

... a is universal

$$\langle m \rangle = \int_0^1 dr \int_0^{2\pi} d\theta r^2 P(r, \theta) \quad \xi$$

$$\langle m_6 \rangle = \int_0^1 dr \int_0^{2\pi} d\theta r^2 P(r, \theta) \cos(6\theta) \quad \Lambda$$



$$a_{q=6;XY} = 2.2(1)$$

$$a_{120^\circ} = 1.2(1)$$