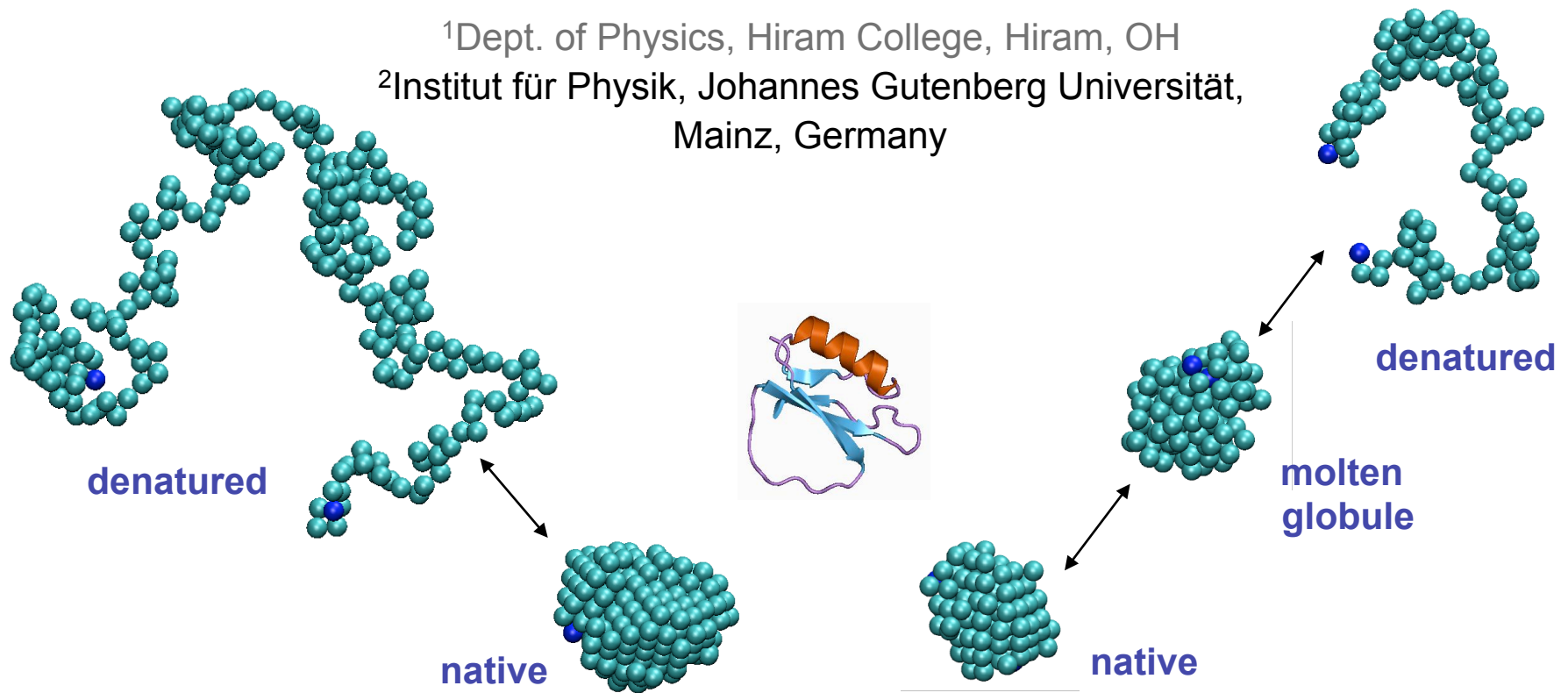


All-or-none protein-like folding of a homopolymer chain

Mark P. Taylor¹, Wolfgang Paul², and Kurt Binder²

¹Dept. of Physics, Hiram College, Hiram, OH

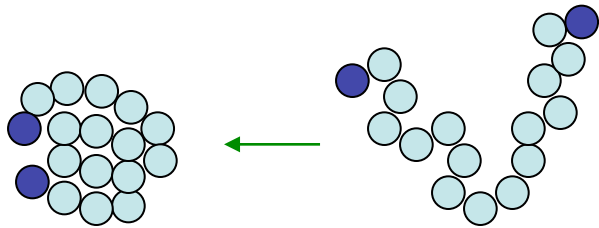
²Institut für Physik, Johannes Gutenberg Universität,
Mainz, Germany



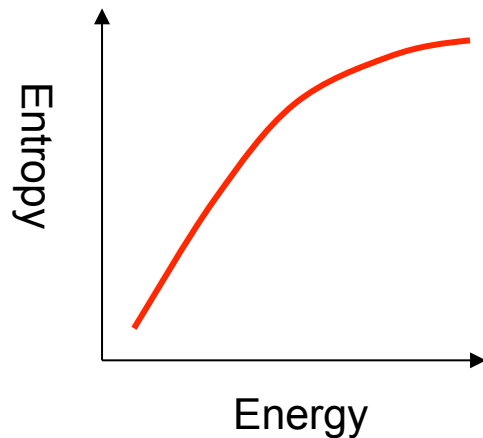
Can a simple homopolymer model capture some essentials of protein folding?

Single Chain Conformational Transitions

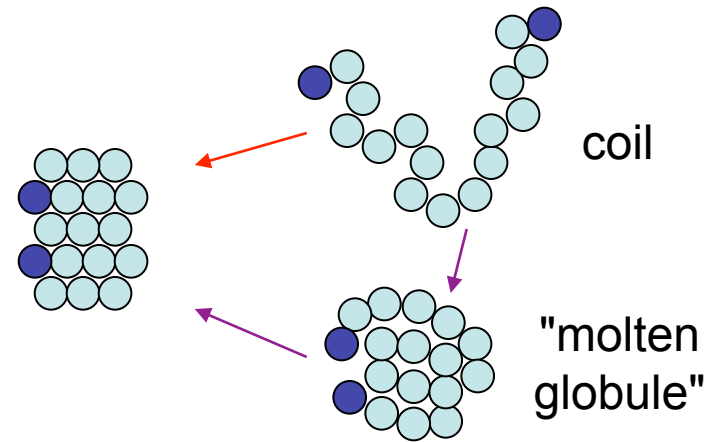
Homopolymer
Globule \leftrightarrow Coil



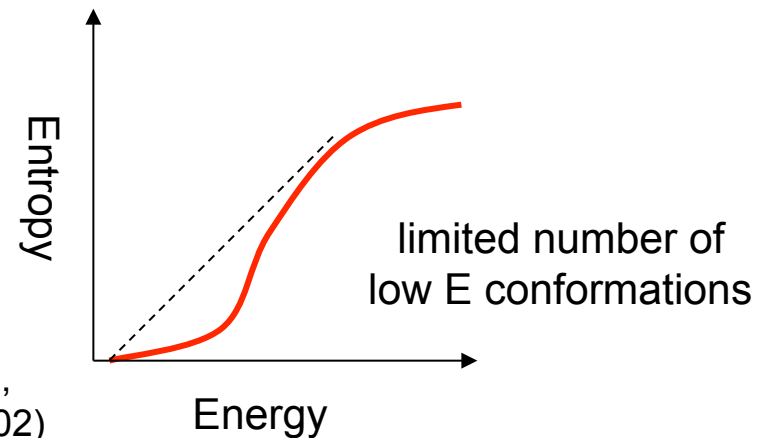
compact \leftarrow expanded
(both disordered)
continuous transition



Protein
Native \leftrightarrow Denatured



order \leftarrow disorder
discontinuous transition



Finkelstein & Ptitsyn,
"Protein Physics" (2002)

Multiple Transitions in a Simple Chain Model

Zhou, Hall, and Karplus found a **"first-order"** transition in a model homopolymer

Model: Flexible Square-Well Chain
 Length: $N = 64$
 Well Diameter: $\lambda = 1.5$
 MC and DMD simulations

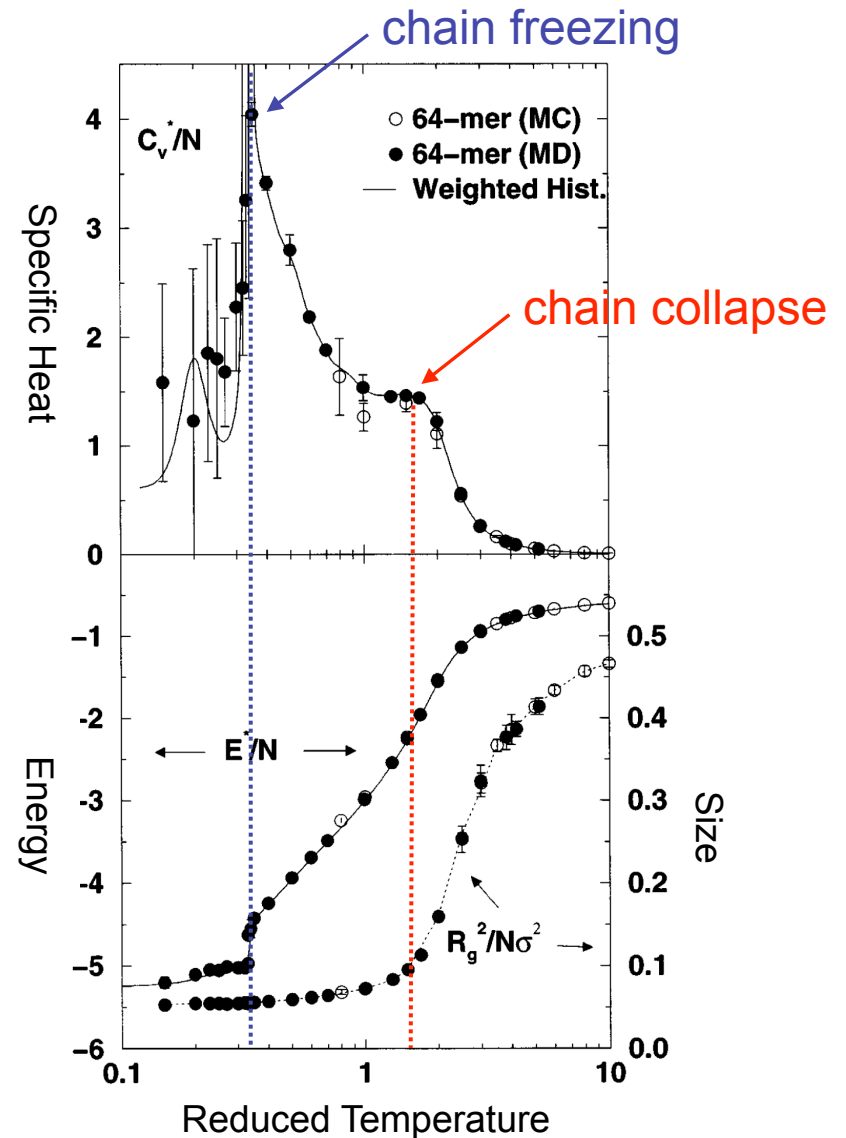
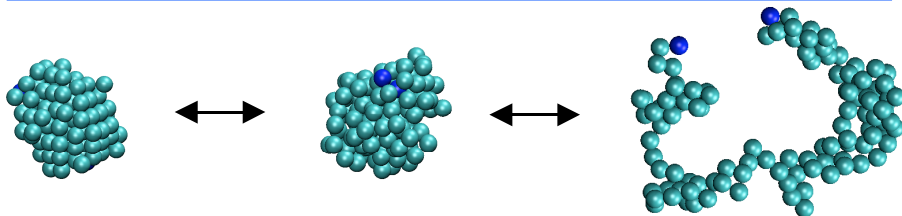
PRL 77, 2822 (1996)

Analogy with Bulk Fluid:

Solid \Leftrightarrow **Liquid** \Leftrightarrow **Gas**

Similar to Protein:

Native \Leftrightarrow **Globule** \Leftrightarrow **Denatured**



Protein-like "all-or-none" transition in a simple model?

Ten years later ... a new result:

Rampf, Paul, and Binder find evidence for a **"direct freezing"** transition in a simple model system.

Lattice Bond-Fluctuation Model

Chain Length: $N = 32-512$

Interaction Range: $\lambda = 1.225$

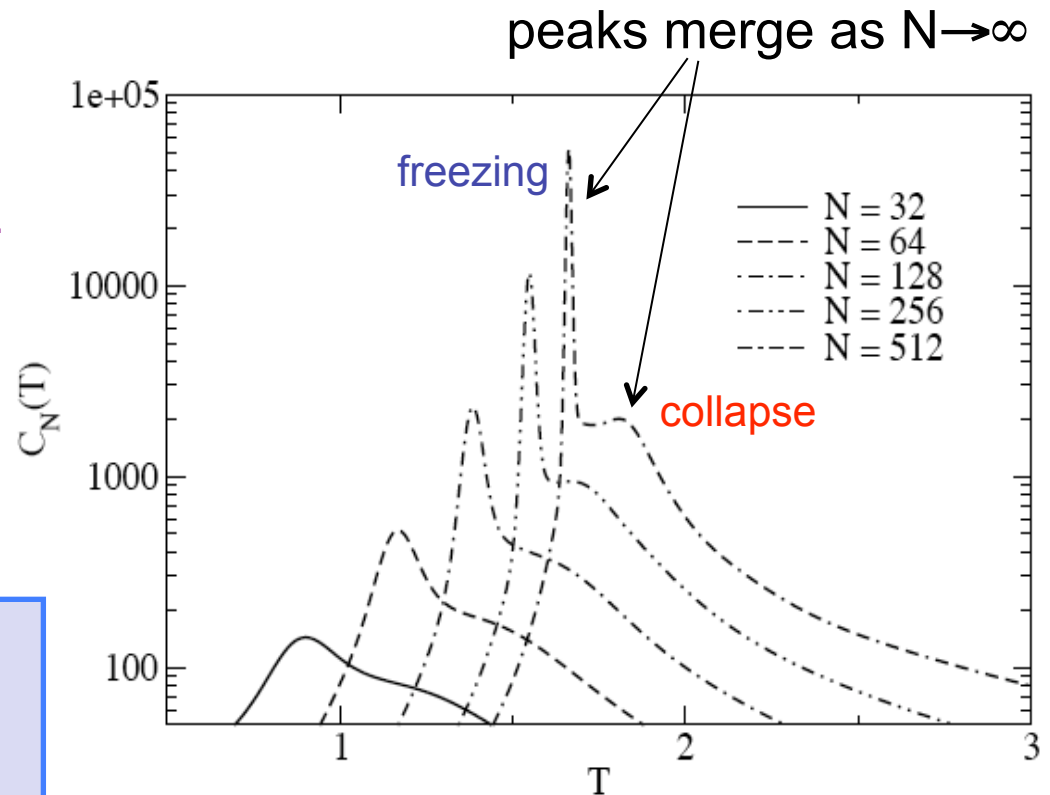
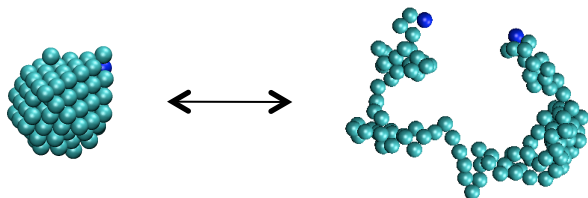
Wang-Landau Simulation Method

Analogous to simple liquid:

Solid \Leftrightarrow **Gas**

and single-step protein folding?:

Native \Leftrightarrow **Denatured**

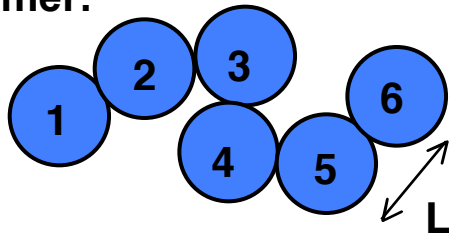


Europhys Lett. 70, 628 (2005); J. Polym. Sci. B 44, 2542 (2006);
PRE 75, 060801(R) (2007)

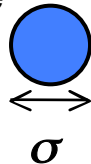
Question: can a finite-length homopolymer exhibit such an all-or-none transition?

SW Chain Model

Polymer:



built from simple monomers:



Model Parameters:

ϵ = well depth (sets energy scale)

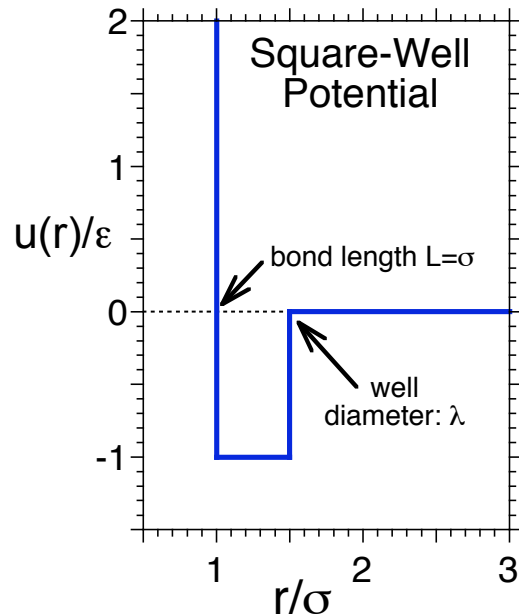
σ = hard-sphere diameter

L = fixed bond length ($L = \sigma$)

λ = interaction range/ σ

$T^* = k_B T / \epsilon$ = reduced temperature

monomer-monomer interaction



Model has a discrete energy spectrum: $E_n = n\epsilon$
(n = number of monomer-monomer interactions)

Can study this model for a continuous range of λ

For SW monomers: liquid phase is unstable if $\lambda \leq 1.25$

Question: Does the SW chain exhibit similar behavior?

Note: chain connectivity places an upper limit on entropy of "gas" phase

Density of States and Wang-Landau Sampling I

Density of States:

$g(E_n)$ = volume of configurational phase space for energy state E_n



Thermodynamics:

microcanonical entropy:

$$S(E) = k_B \ln g(E)$$

canonical partition function:

$$Z(T) = \sum g(E) \exp(-E/k_B T)$$

Wang-Landau algorithm* ... an iterative simulation method to compute $g(E_n)$:

Starting w/ $g(E_n)=1$, $H(E_n)=0 \forall n$, $f_0 = e$

Generate sequence of chain conformations using acceptance criteria:

$$P_{acc}(a \rightarrow b) = \min\left(1, \frac{g(E_a)}{g(E_b)}\right)$$

Update DOS: $g(E_n) \rightarrow f_m g(E_n)$

Update visitation

histogram: $H(E_n) \rightarrow H(E_n)+1$

When histogram \sim flat ...

reduce modification factor: $f_{m+1} = (f_m)^{1/2}$

reset histogram to zero: $H(E_n) = 0 \forall n$

iterate
m levels

m=20 is
standard

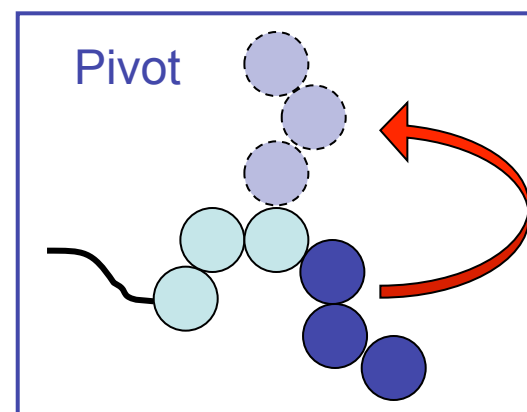
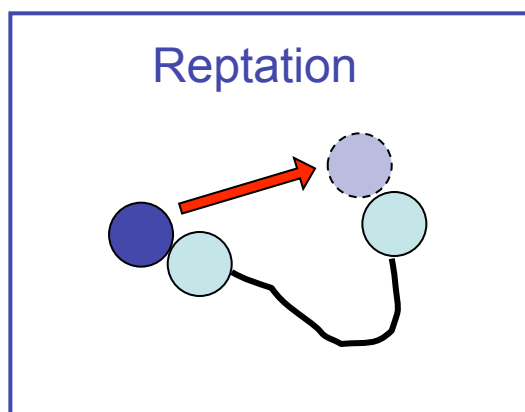
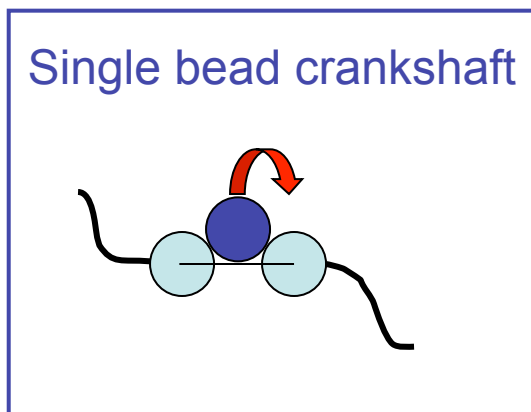
we need
m>25

*Wang & Landau, PRL 86, 2050 (2001); PRE 64, 056101 (2001).

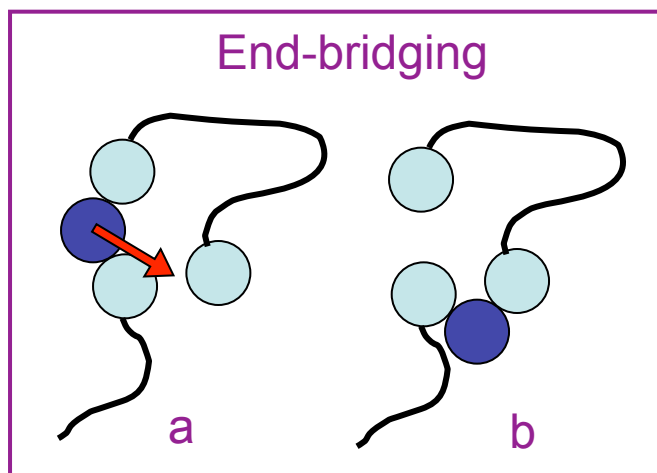
Wang-Landau Sampling II

Success of the WL methods depends critically on underlying MC move set

These "standard" moves easily sample most of configuration space:



... However, we need **this move** to access the lowest energy regions of phase space:



This move requires weight factors in the acceptance criteria:

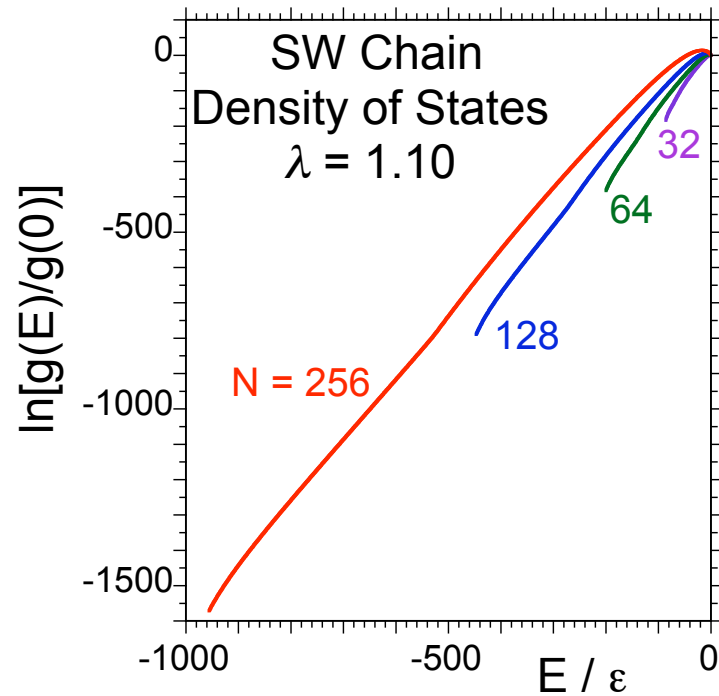
$$w_a = n_a J_a$$

of bridgable sites in state a

Jacobian factor for state a

Escobedo & de Pablo, JCP **102**, 2636 (1995)

Single Chain DOS and Canonical Analysis



For $N = 256$: $g(E)$ spans ~700 orders of magnitude!

Taylor, Paul, & Binder,
PRE **79**, 050801(R) (2009)

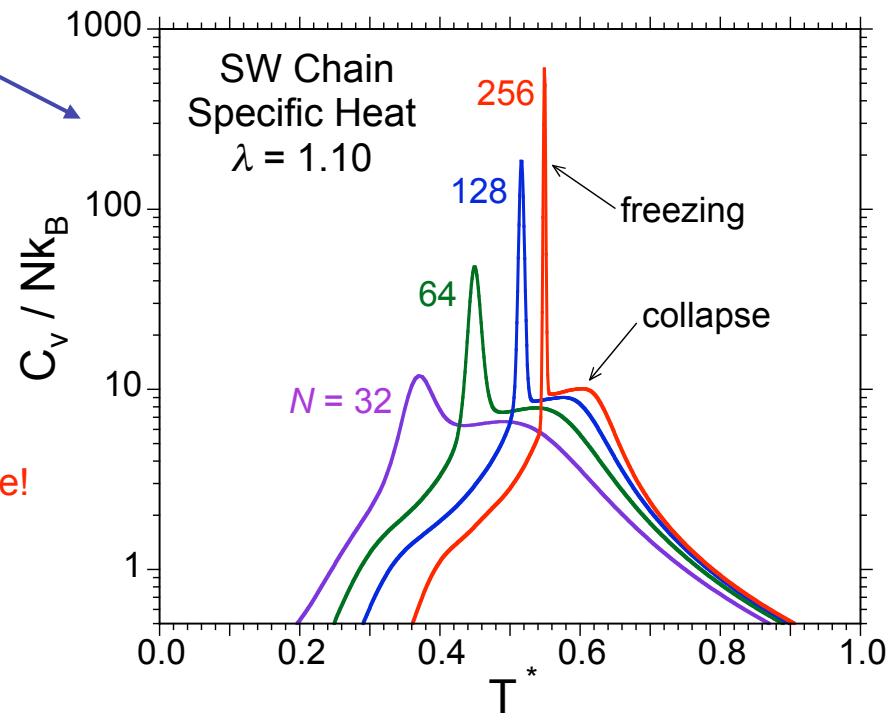
Canonical Analysis

Partition Function: $Z = \sum g(E) e^{-E/kT}$

Probability: $P(E,T) = g(E) e^{-E/kT} / Z$

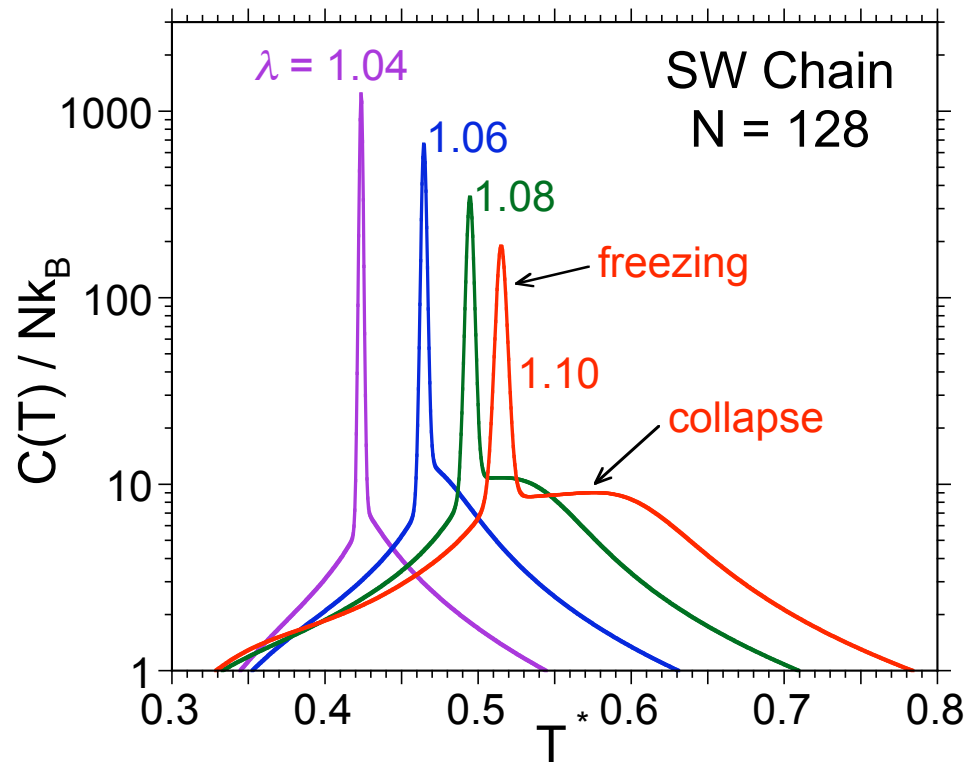
Average Energy: $\langle E(T) \rangle = \sum EP(E,T)$

Specific Heat: $C(T) = d\langle E(T) \rangle / dT$



Phase Behavior for Finite Length Chain?

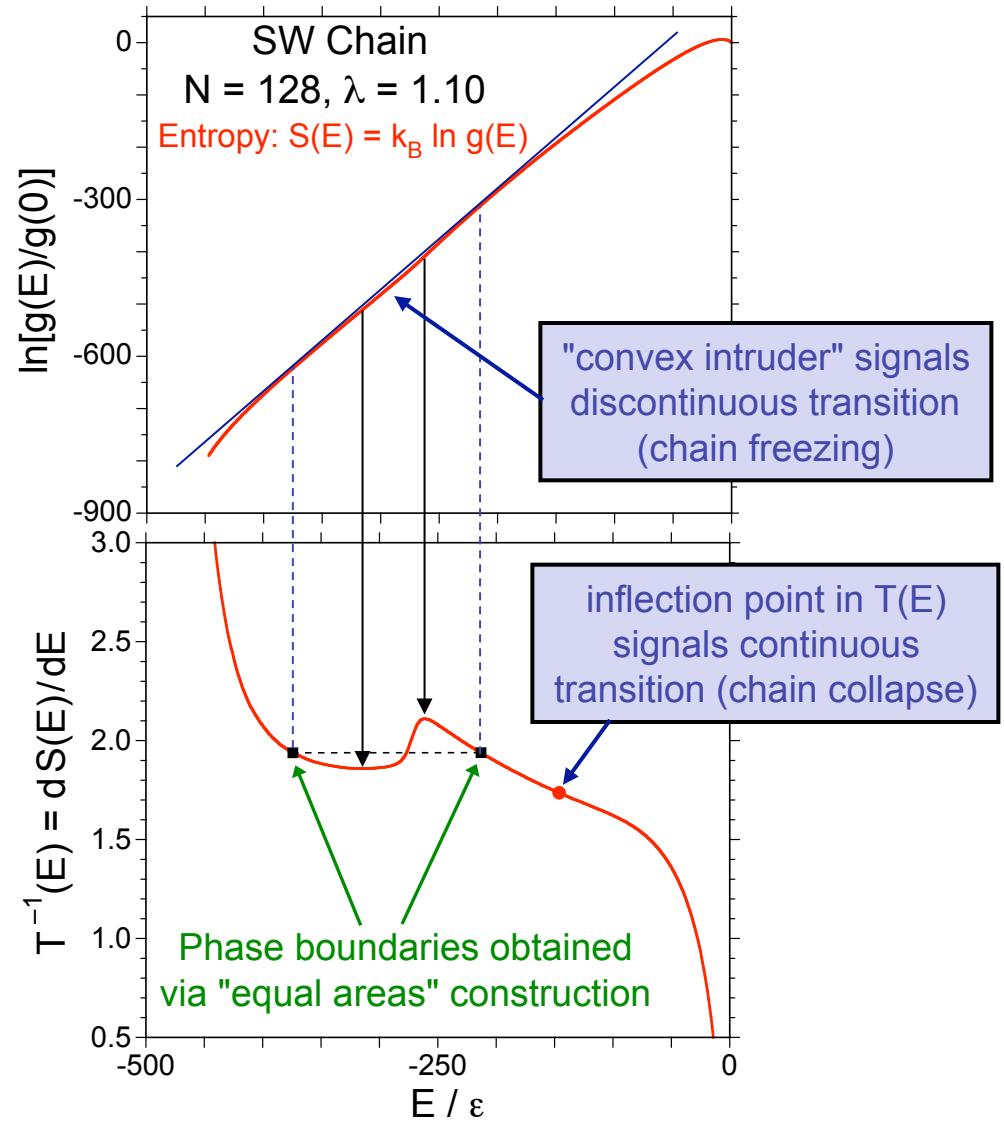
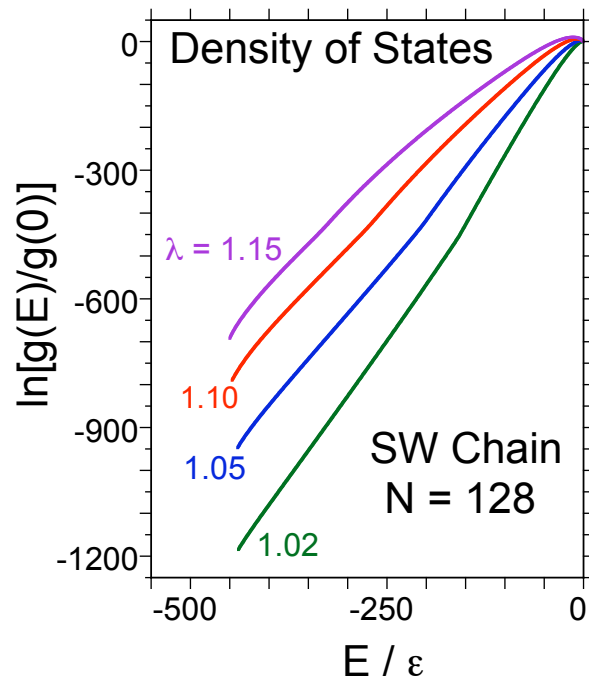
In the "canonical analysis", collapse and freezing specific heat peaks merge for small λ ...



... a "microcanonical analysis" can be used to distinguish these transitions

Microcanonical Analysis I

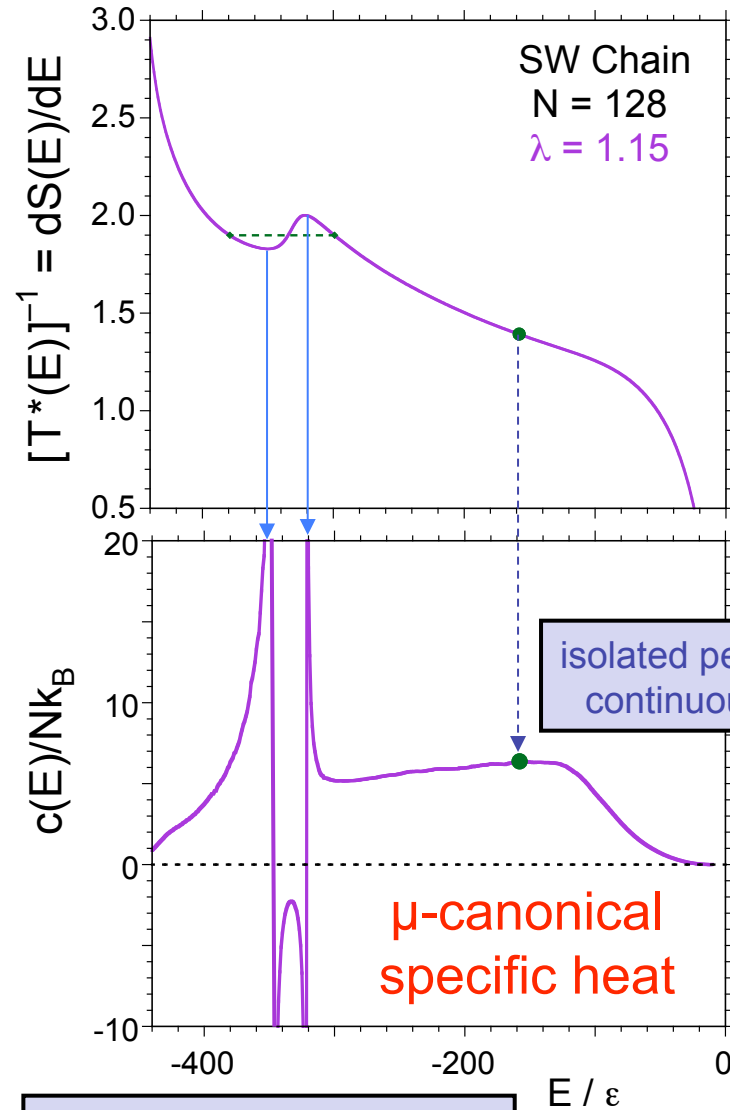
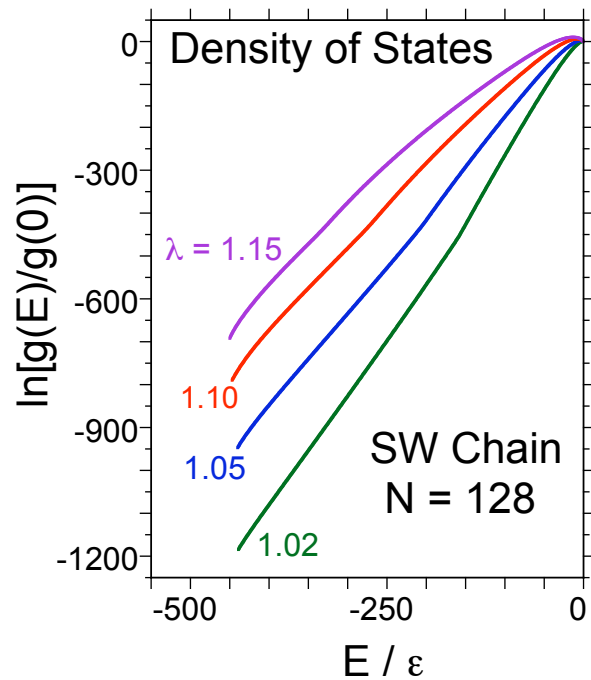
Phase transitions in a finite system determined from **curvature** of the microcanonical entropy:
 $S(E) = k_B \ln g(E)$



Gross, "Microcanonical Thermodynamics" (2001)
 Junghans, Bachmann, & Janke, PRL **97**, 218103 (2006)
 Taylor, Paul, & Binder, PRE **79**, 050801(R) (2009)

Microcanonical Analysis II

Phase transitions in a finite system determined from **curvature** of the microcanonical entropy:
 $S(E) = k_B \ln g(E)$



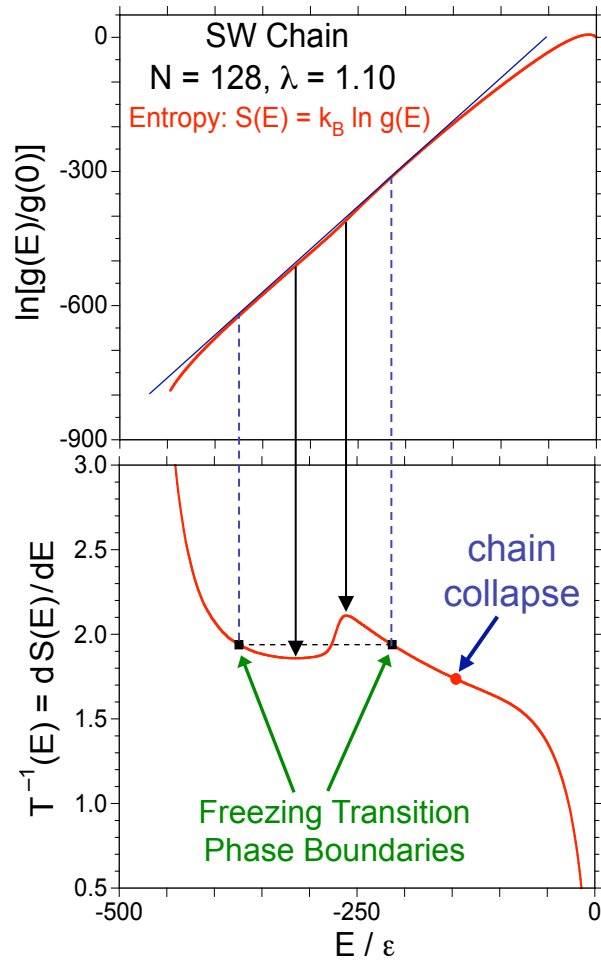
Gross, "Microcanonical Thermodynamics" (2001)

Junghans, Bachmann, & Janke, PRL **97**, 218103 (2006)

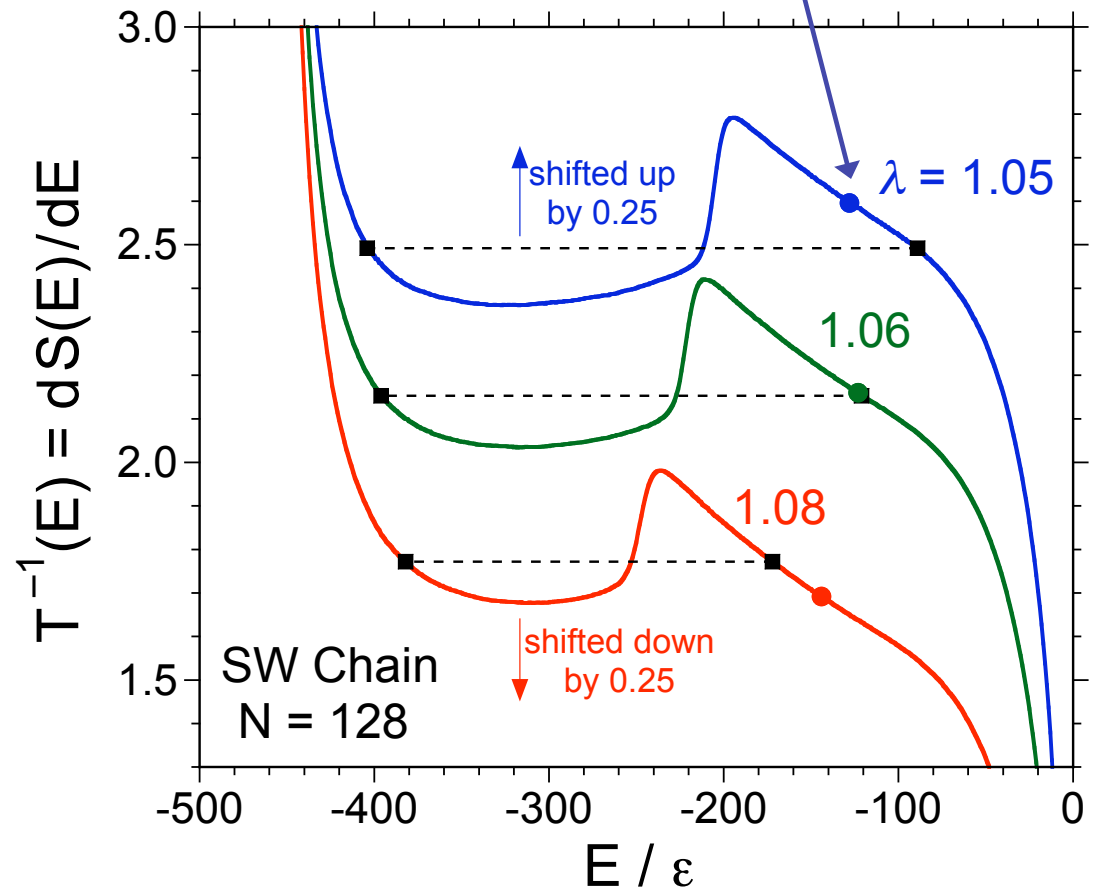
Taylor, Paul, & Binder, JCP **131**, 114907 (2009)

Microcanonical Analysis III

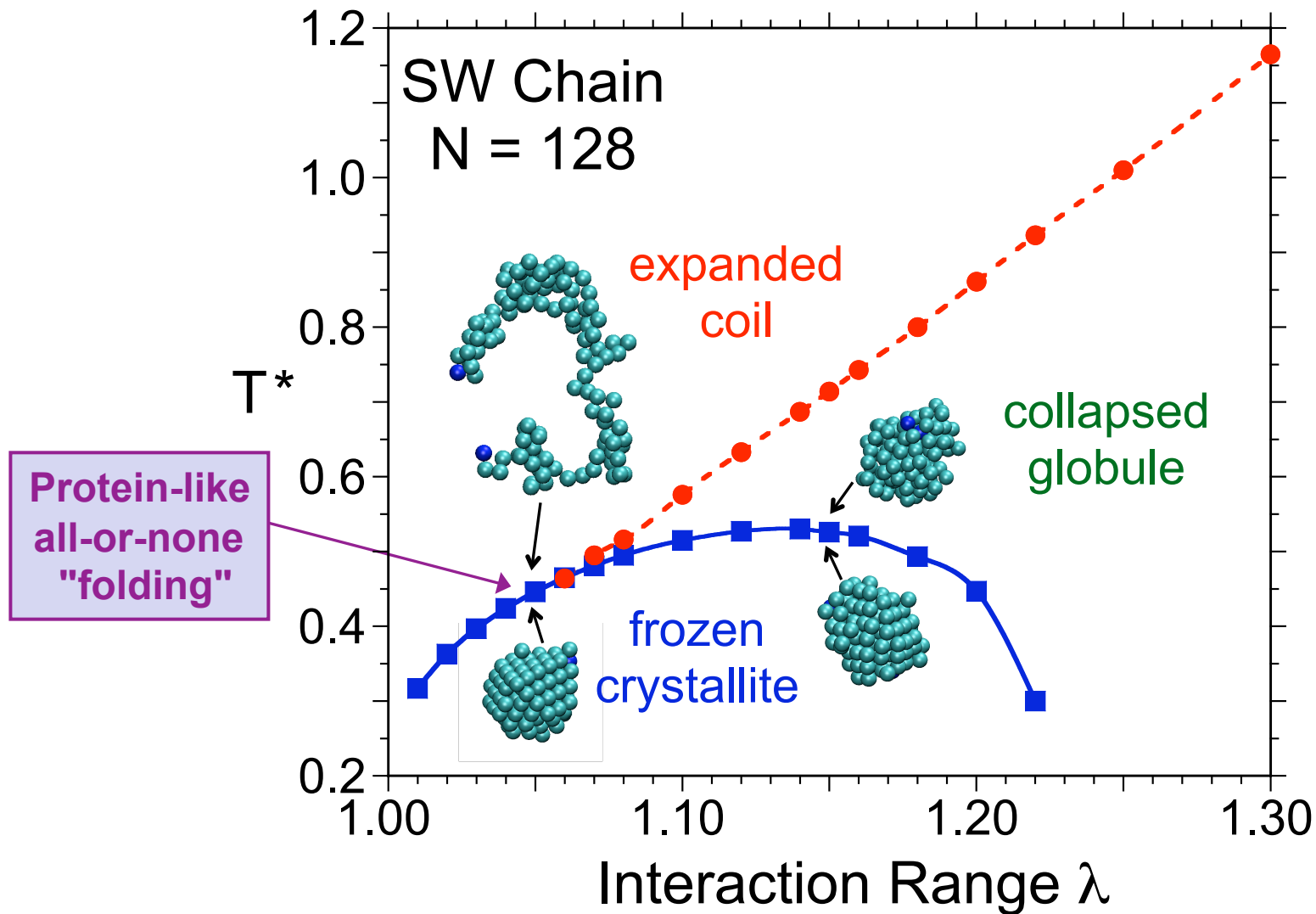
Phase transitions in a finite system determined from **curvature** of $S(E)$



Collapse transition is preempted by freezing for short-range interaction!

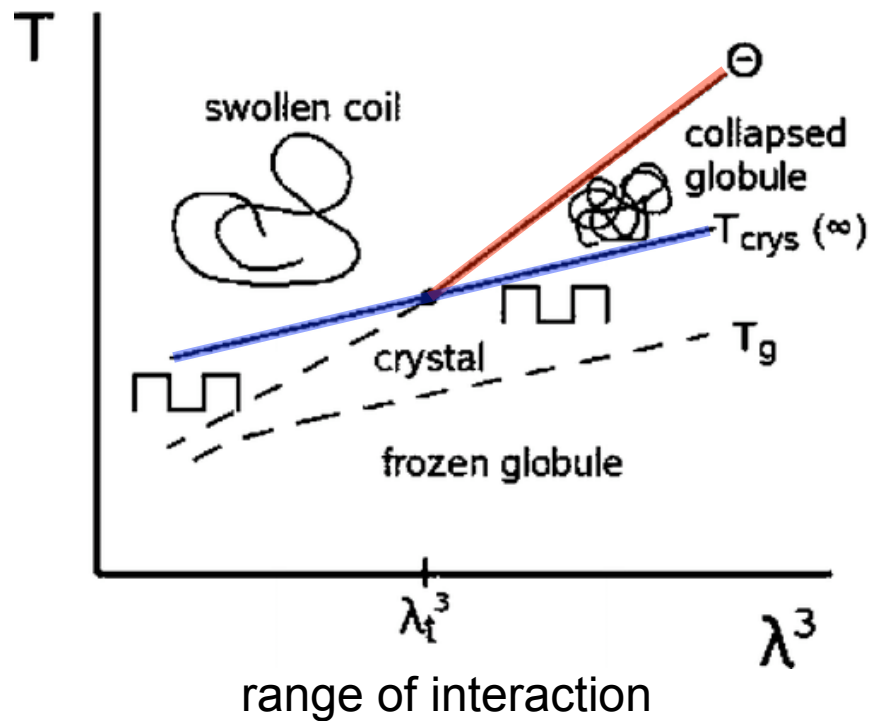


Single Chain Phase Diagram: T- λ Version



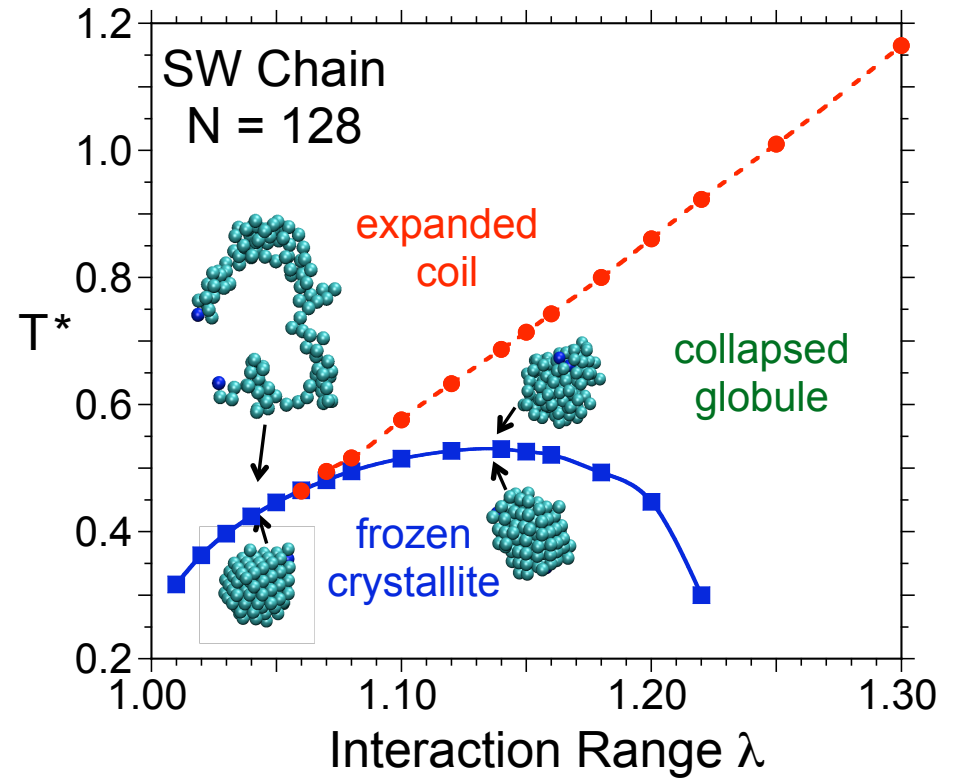
Taylor, Paul, & Binder, J. Chem. Phys. **131**, 114907 (2009)

Prediction:



Binder and Paul
 Macromolecules **41**, 4537 (2008)

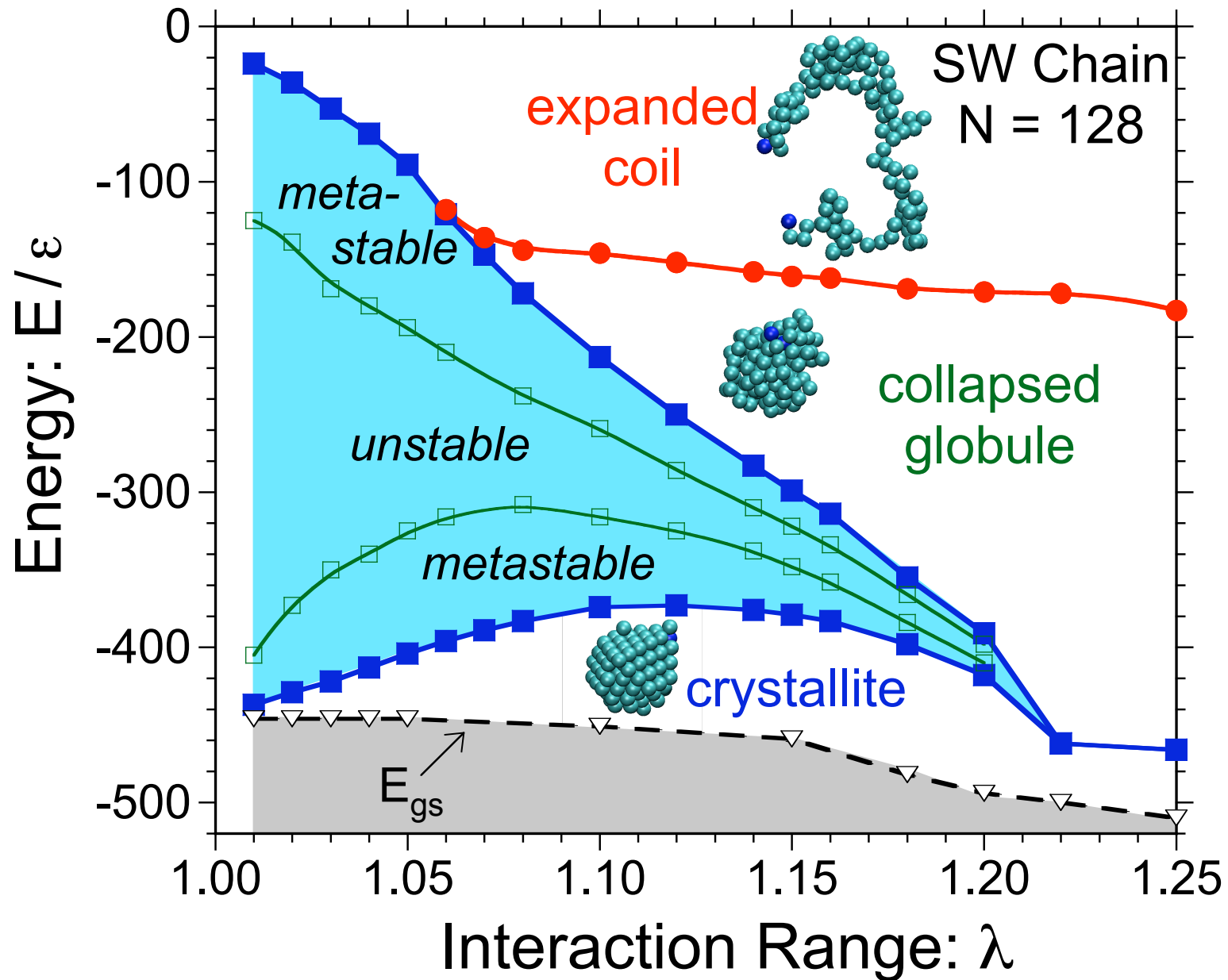
Calculation:



triple point moves to
 $\lambda = 1.15$ in $N \rightarrow \infty$ limit

Taylor, Paul, & Binder
 J. Chem. Phys. **131**, 114907 (2009)

Single Chain Phase Diagram: E- λ Version



Protein Thermodynamics: Free Energy Landscape

All-or-none folding
of a $N = 128$ SW Chain
with $\lambda = 1.05$

Defining a physical **folding temperature** sets the **model energy scale**

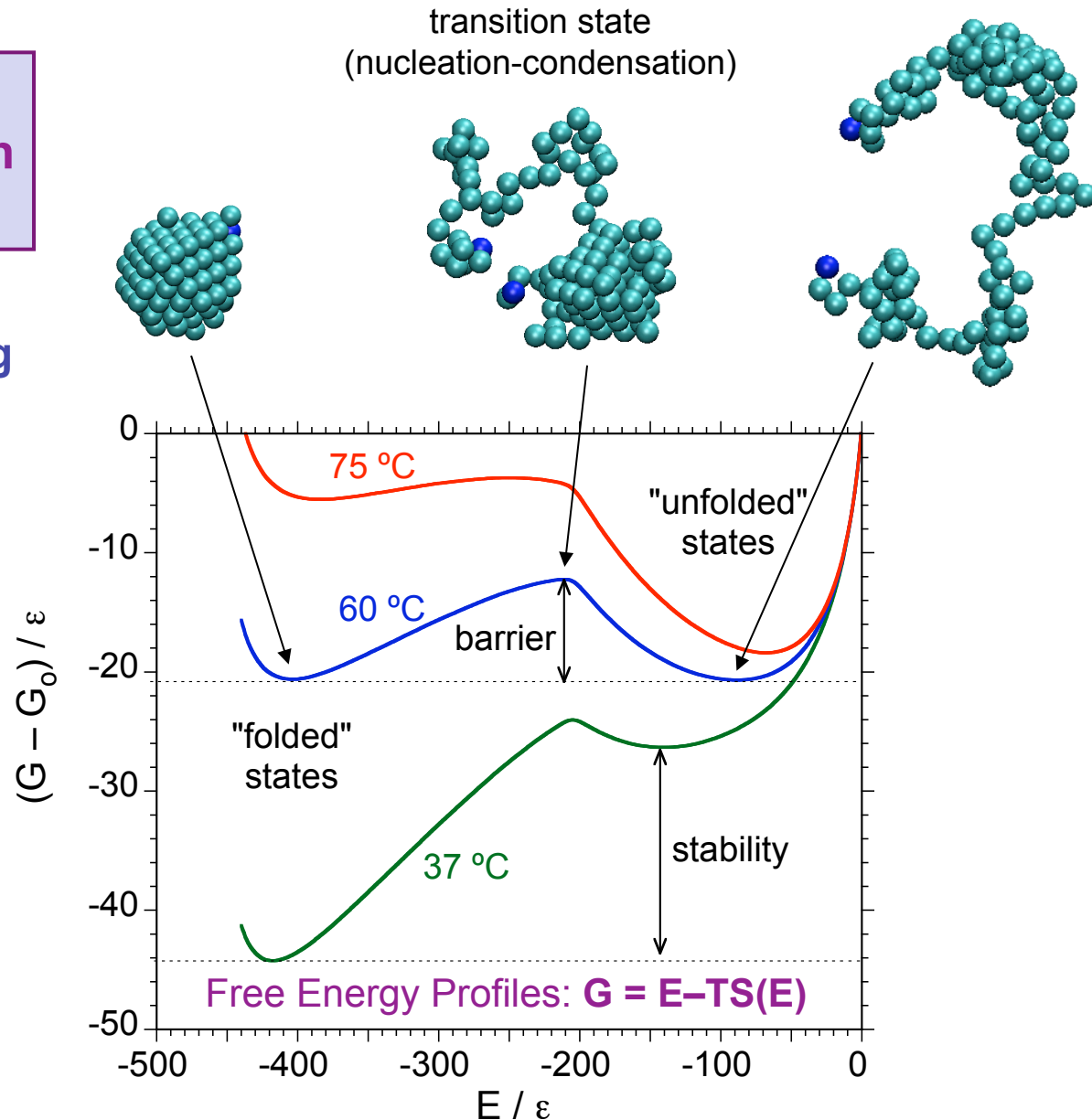
Here we take

$$T_{\text{fold}} = 333 \text{ K (60 } ^\circ\text{C)}$$

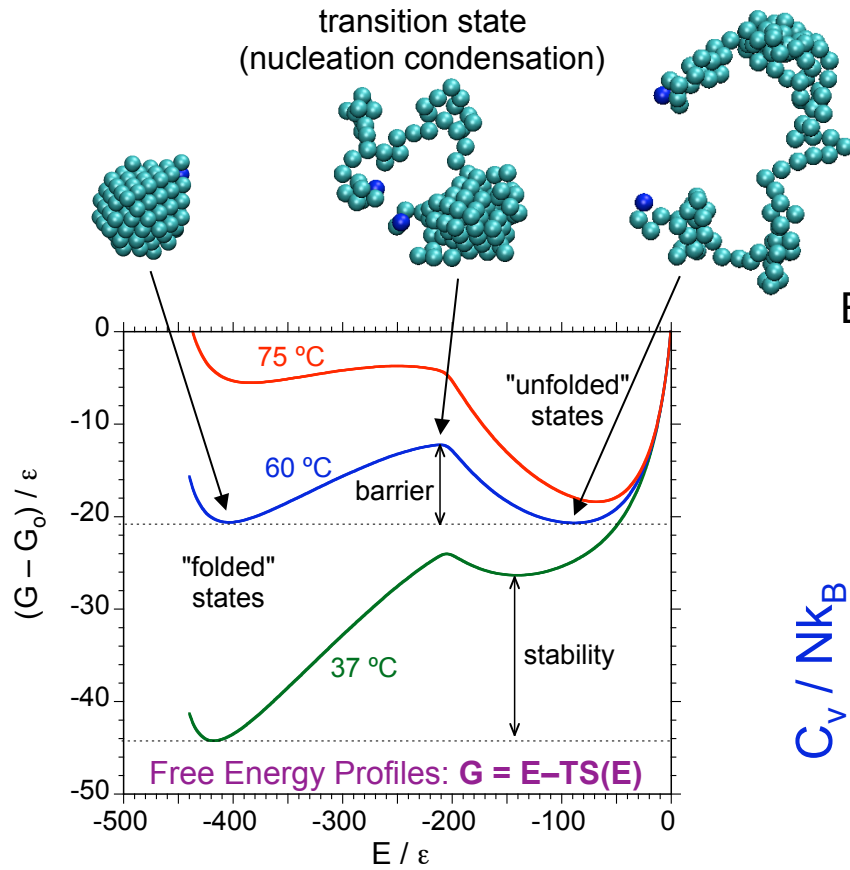
$$\Rightarrow \epsilon = 1.5 \text{ kcal/mol}$$

Taylor, Paul, & Binder,
PRE **79**, 050801(R) (2009)

See Chan and Kaya:
Proteins **40**, 543 & 637
(2000); **52**, 510 (2003).

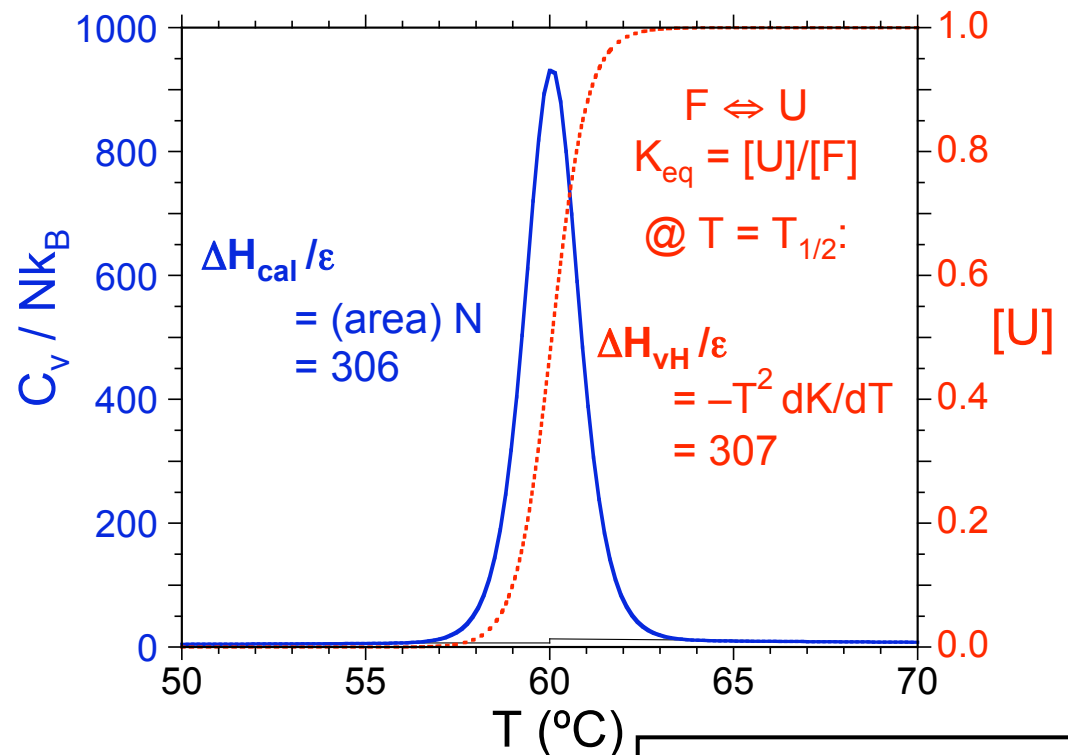


Protein Thermodynamics: Calorimetry



Transition enthalpy gives "distance" between folded/unfolded states.

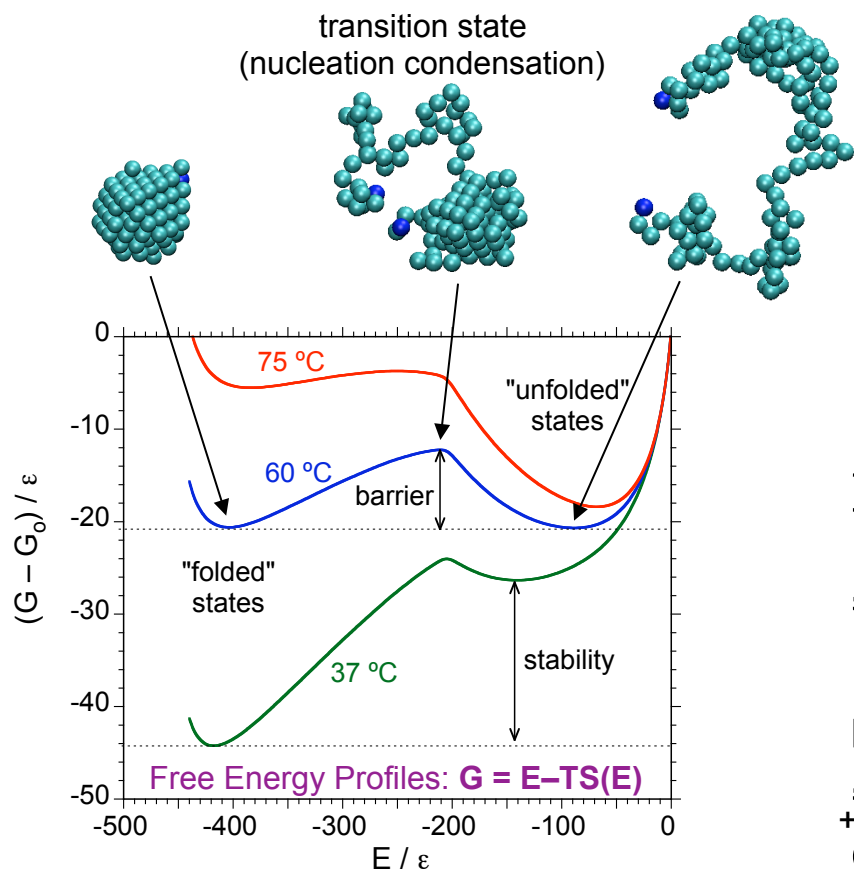
Experimental test for "two-state" folding:
Equality of **calorimetric** and **van't Hoff*** enthalpies



*experimentally this is determined from the height of C_v

See Chan and Kaya:
Proteins **40**, 543 & 637 (2000); **52**, 510 (2003).

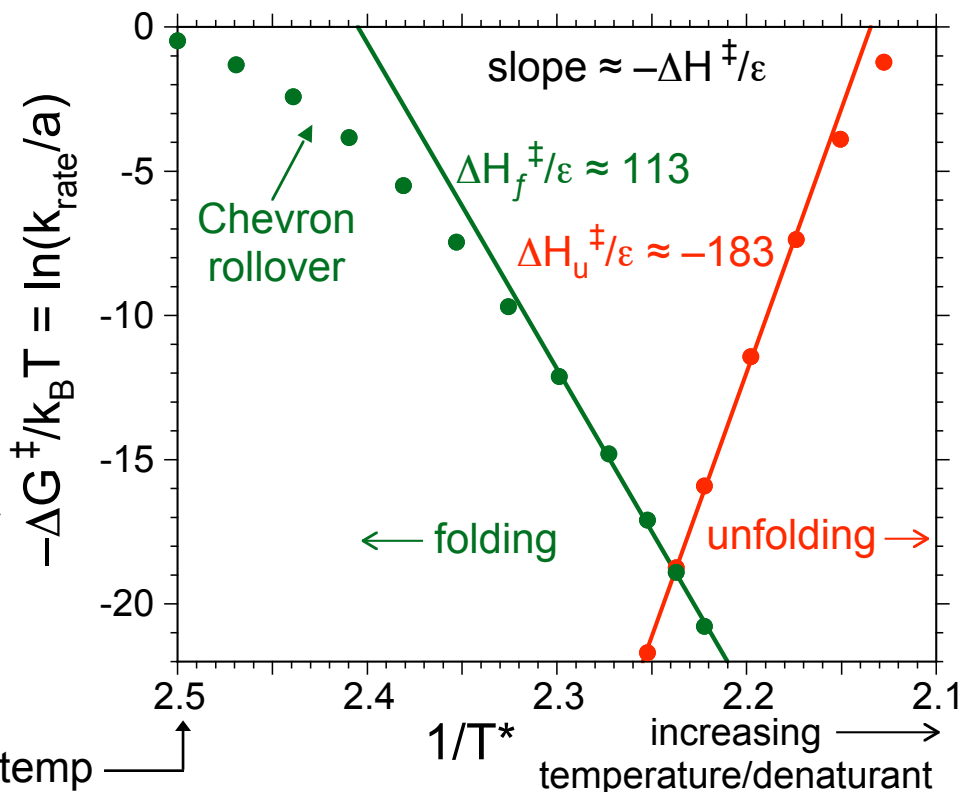
Protein Thermodynamics: Chevron Plot



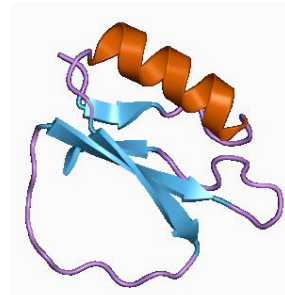
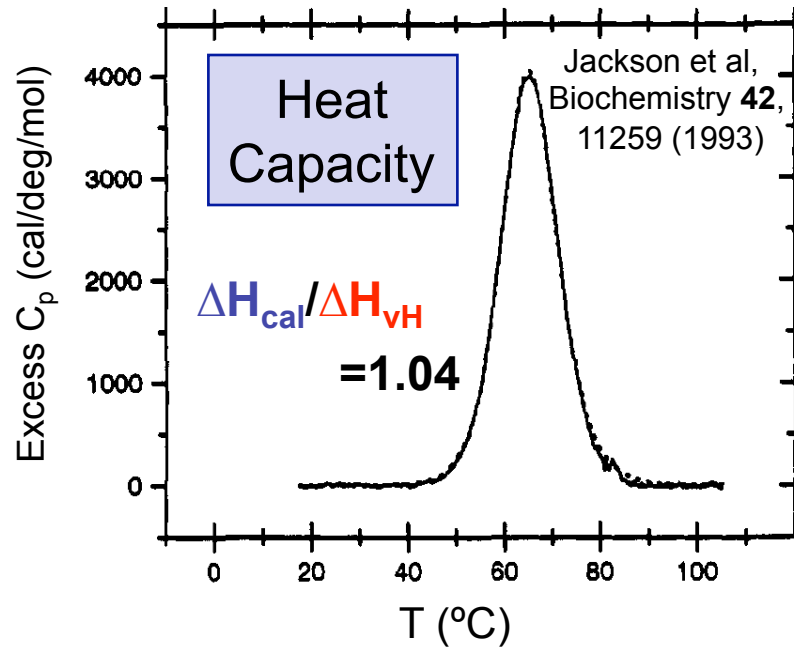
Chevron slopes give "distance" from folded/unfolded states to transition state.

Experimental test for "two-state" folding:

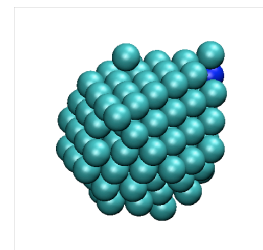
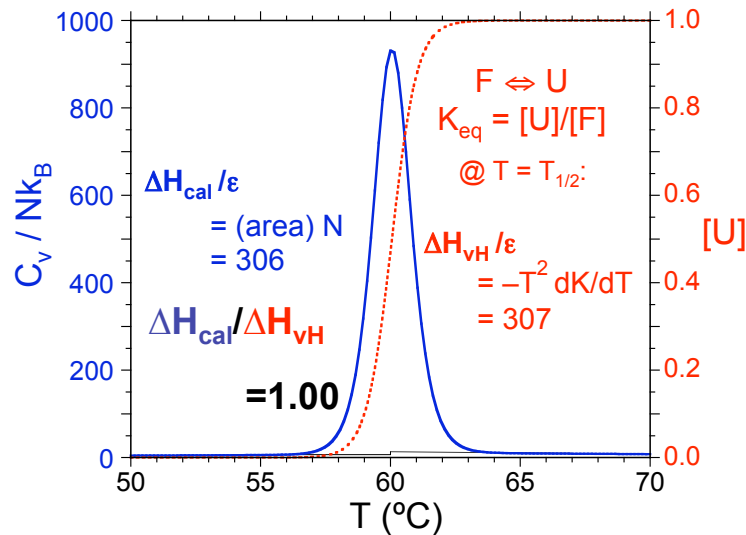
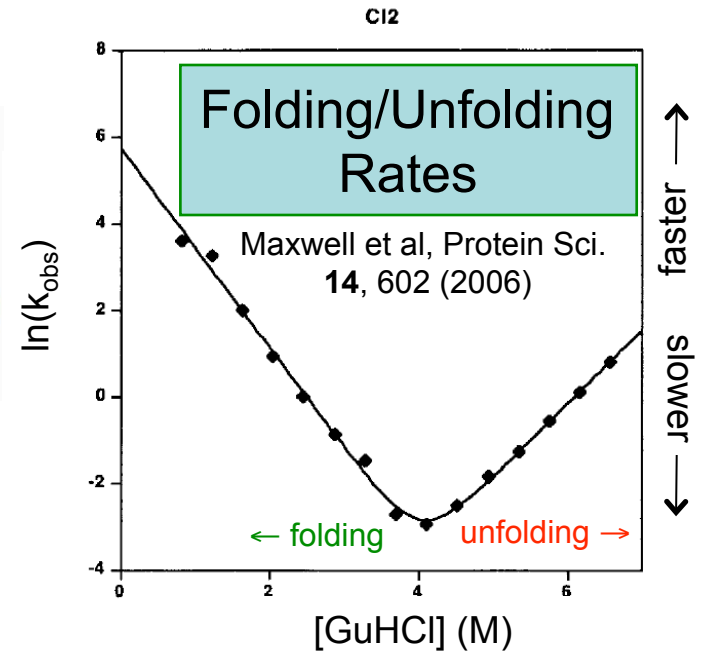
Barrier height via kinetics ...
 Anti-Arrhenius/Arrhenius
 behavior for folding/unfolding.



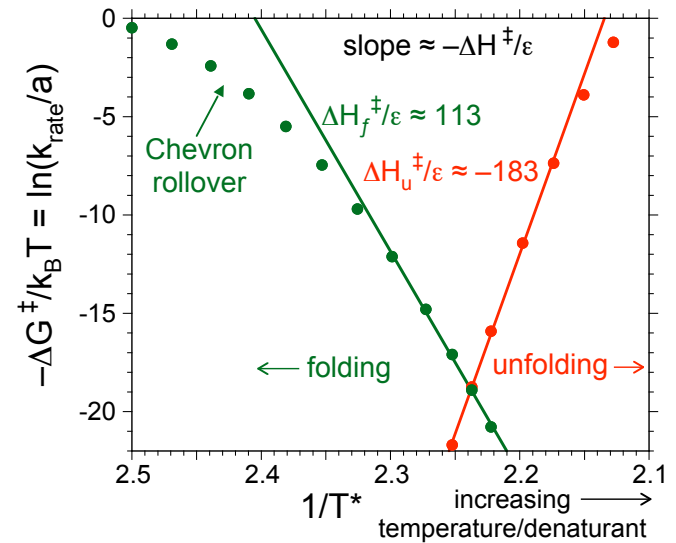
Experimental Data for Chymotrypsin Inhibitor 2 (CI2)



83 residue protein



128 bead SW chain



Summary and Outlook

Flexible SW Chain Model

Findings: Short range interaction results in "all-or-none" folding.
see: PRE **79**, 050801(R) (2009); JCP **131**, 114907 (2009)
Reproduces many key aspects of protein thermodynamics.

To do: Heteropolymer model (such as HP type).
[difficulty: bridging moves change sequence]
Explore free energy landscapes.

Funding: DFG (SFB 625-A3)

NSF (DMR-0804370)

Hiram College

Special thanks to the Binder group for their hospitality!

Happy "American" Thanksgiving