

# Anisotropic three-dimensional Heisenberg antiferromagnets in a field

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## XXZ model on a cubic lattice

$$\mathcal{H} = J \sum_{\langle i,j \rangle} [\Delta(S_i^x S_j^x + S_i^y S_j^y) + S_i^z S_j^z] - H \sum_i S_i^z$$

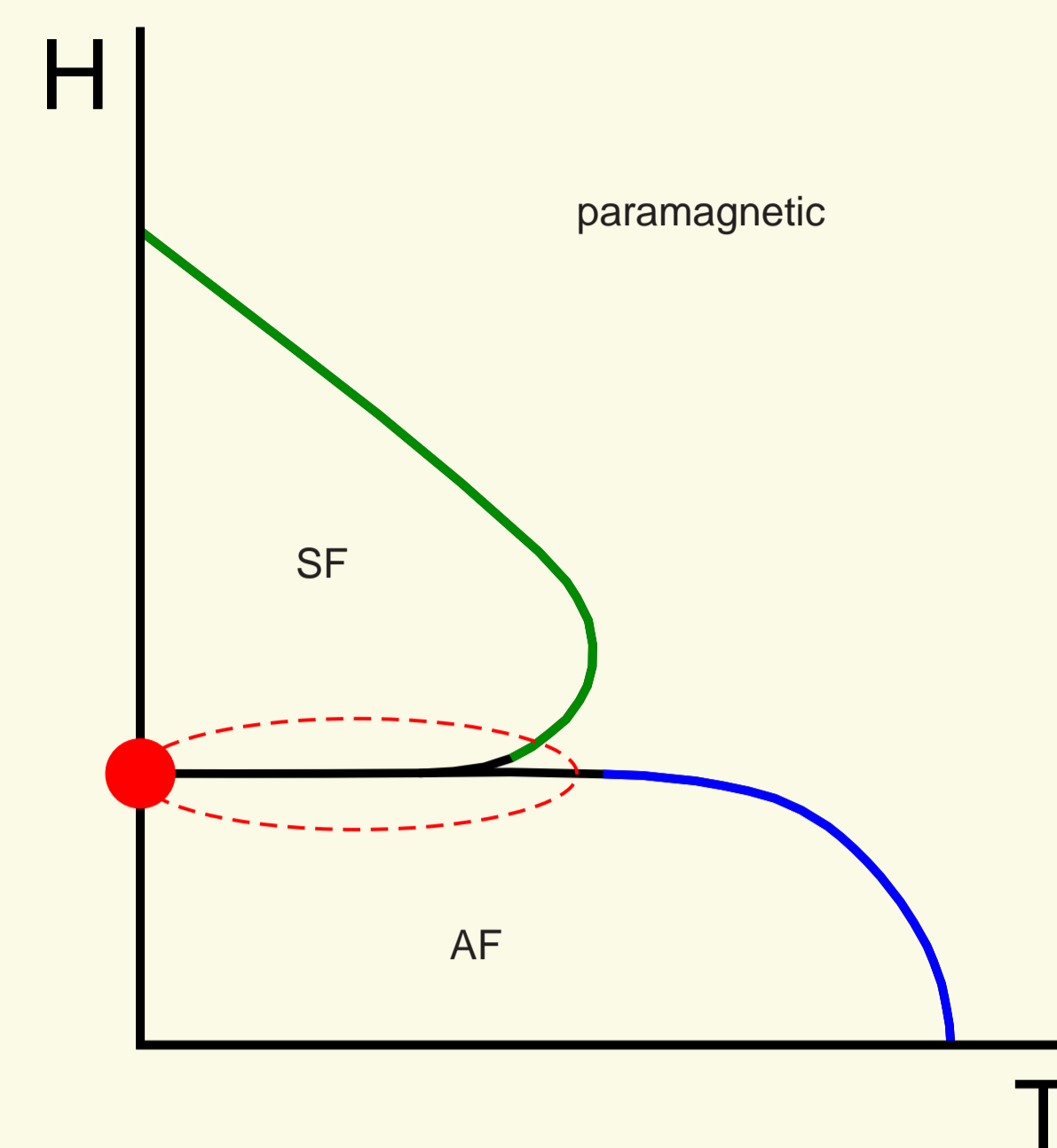
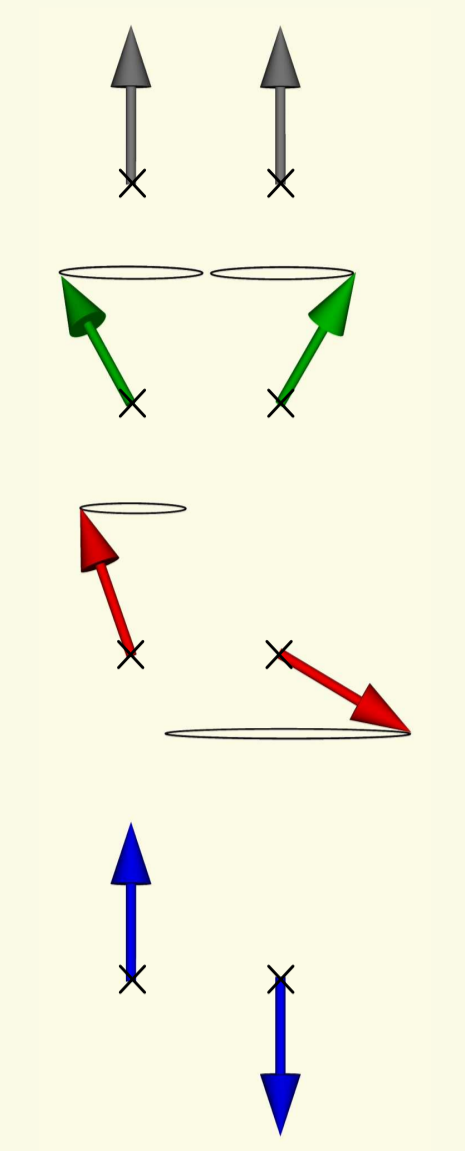
nearest-neighbour pairs	$\langle i,j \rangle$
coupling strength	$J > 0$
easy-axis anisotropy	$\Delta < 1$
field in z-direction	$H$
classical spin-vectors	$\vec{S}_i = (S_i^x, S_i^y, S_i^z);  \vec{S}_i  = 1$

## Methods

- Monte-Carlo simulations
- Ground state considerations

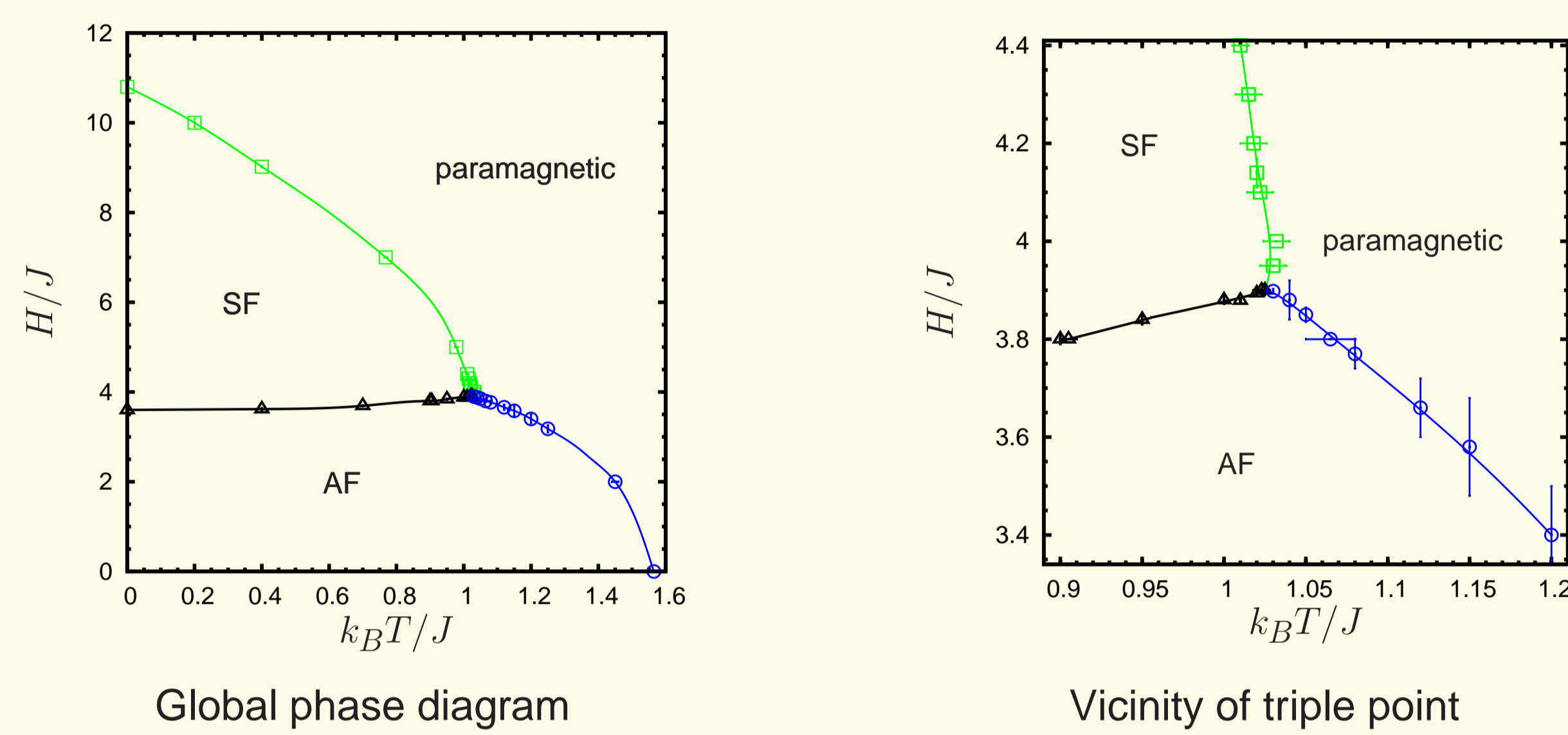
## Structures

ground states



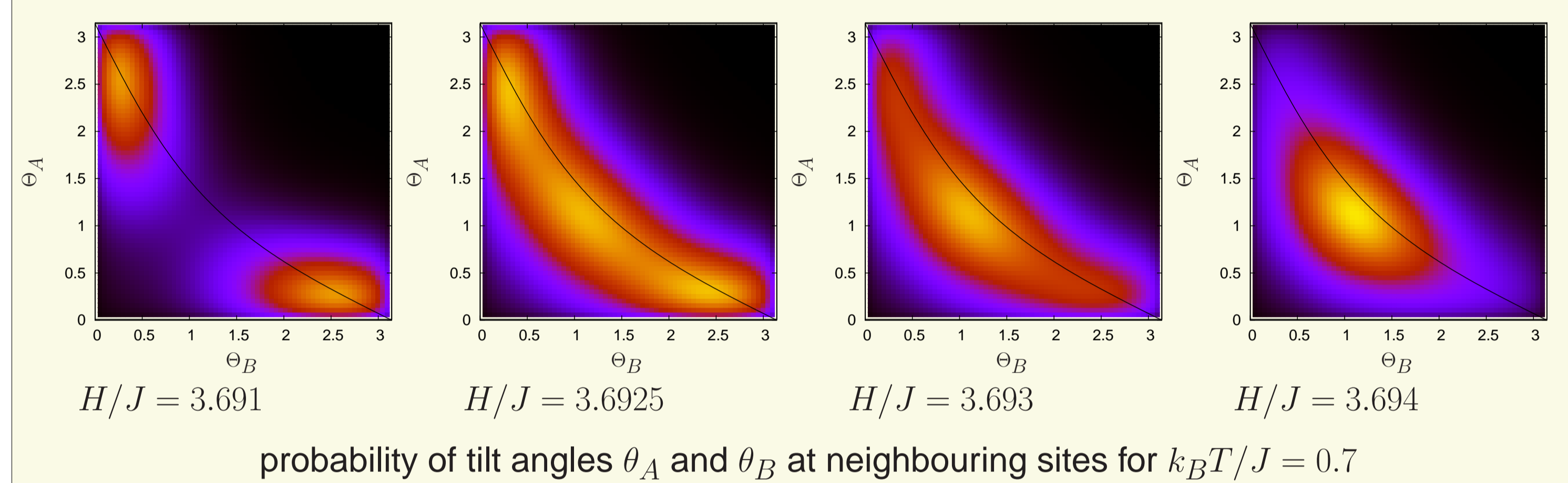
- spin-flop (SF) phase
  - XY-like phase
  - large magnetization
- antiferromagnetic (AF) phase
  - Ising-like phase
  - small magnetization
- **ground state degeneracy:**
  - AF, SF, and **biconical (BC)** structures

## Phase diagram ( $\Delta = 0.8$ )



## At low temperatures:

first order transition, biconical fluctuations



relation of tilt angles  $\theta_A, \theta_B$  in degenerate ground state (see solid black line)

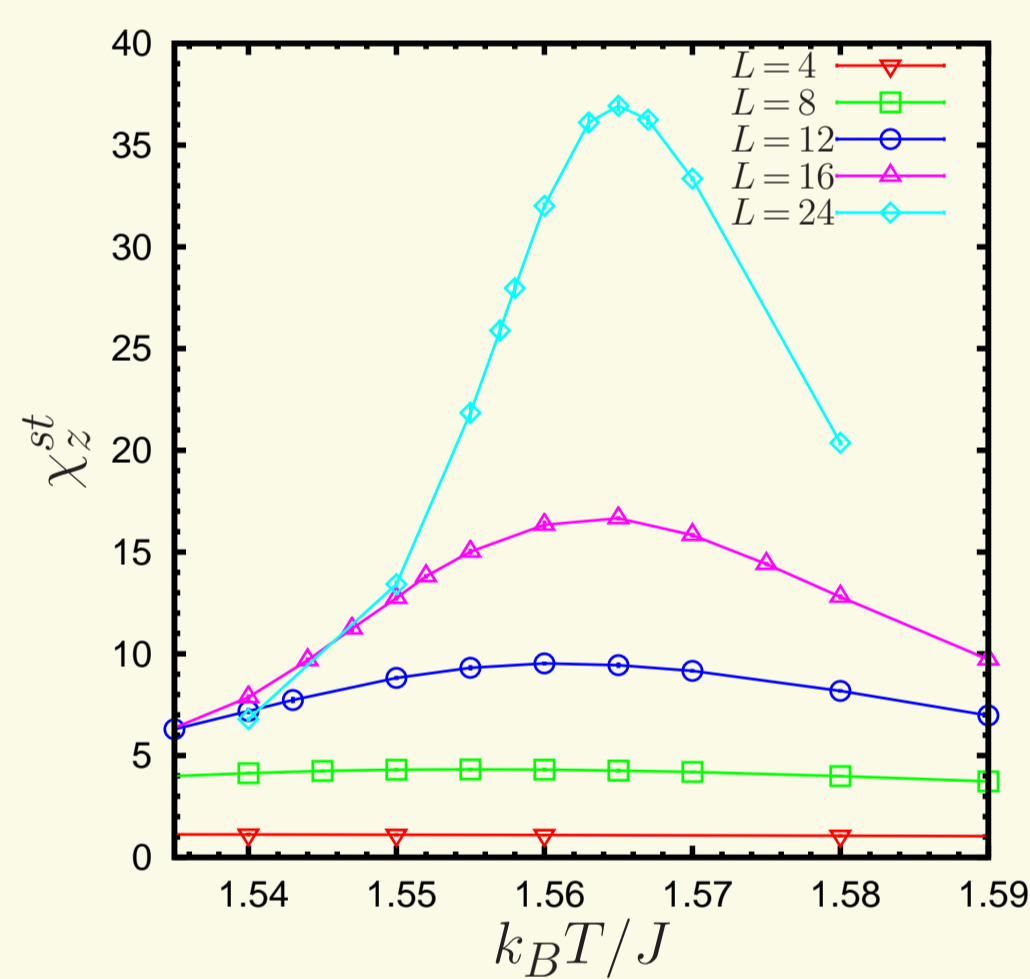
$$\theta_B = \arccos\left(\frac{\sqrt{1-\Delta^2} - \cos\theta_A}{1 - \sqrt{1-\Delta^2}\cos\theta_A}\right)$$

Evidence for coexistence of AF and SF phases at  $H/J = 3.6925$ , i.e. first order transitions (in agreement with analyses of critical exponents)

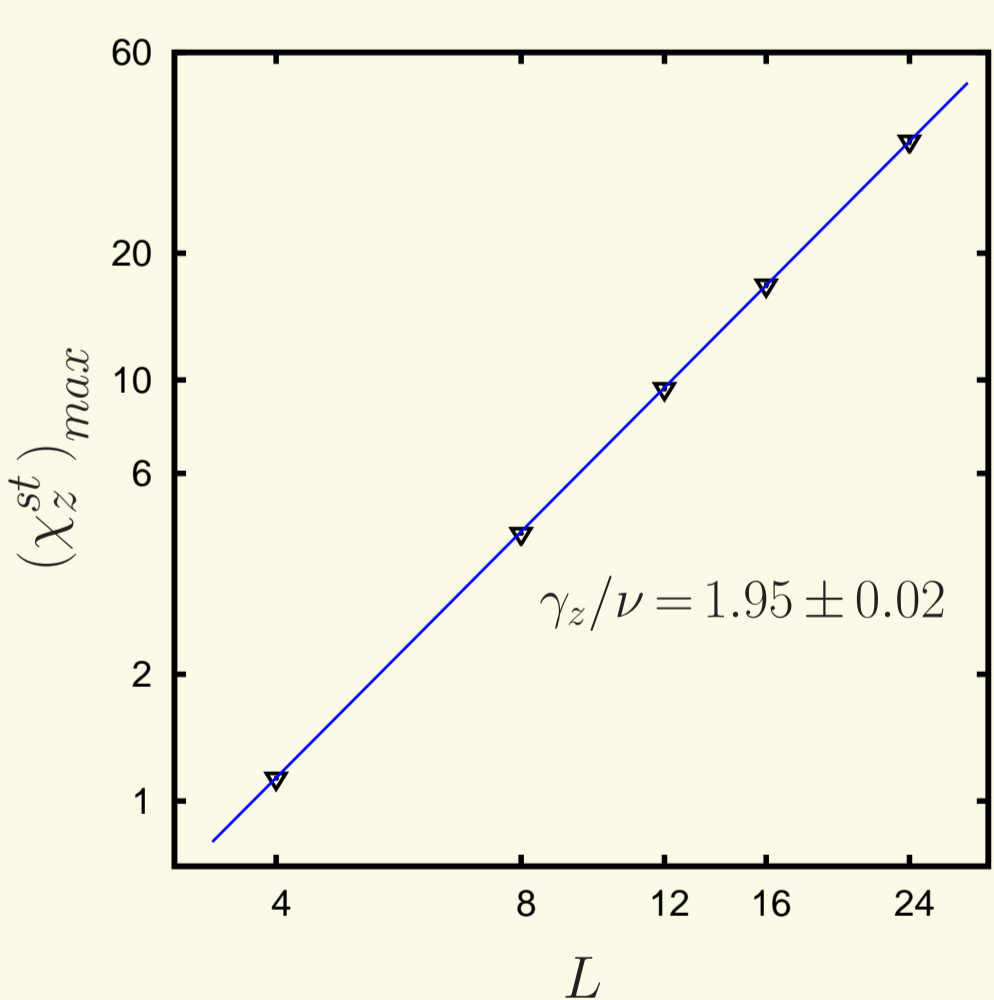
## Critical properties of AF and SF phase boundaries

Example: Finite size behaviour of the staggered susceptibilities at selected fields

AF phase boundary ( $H/J = 0$ )



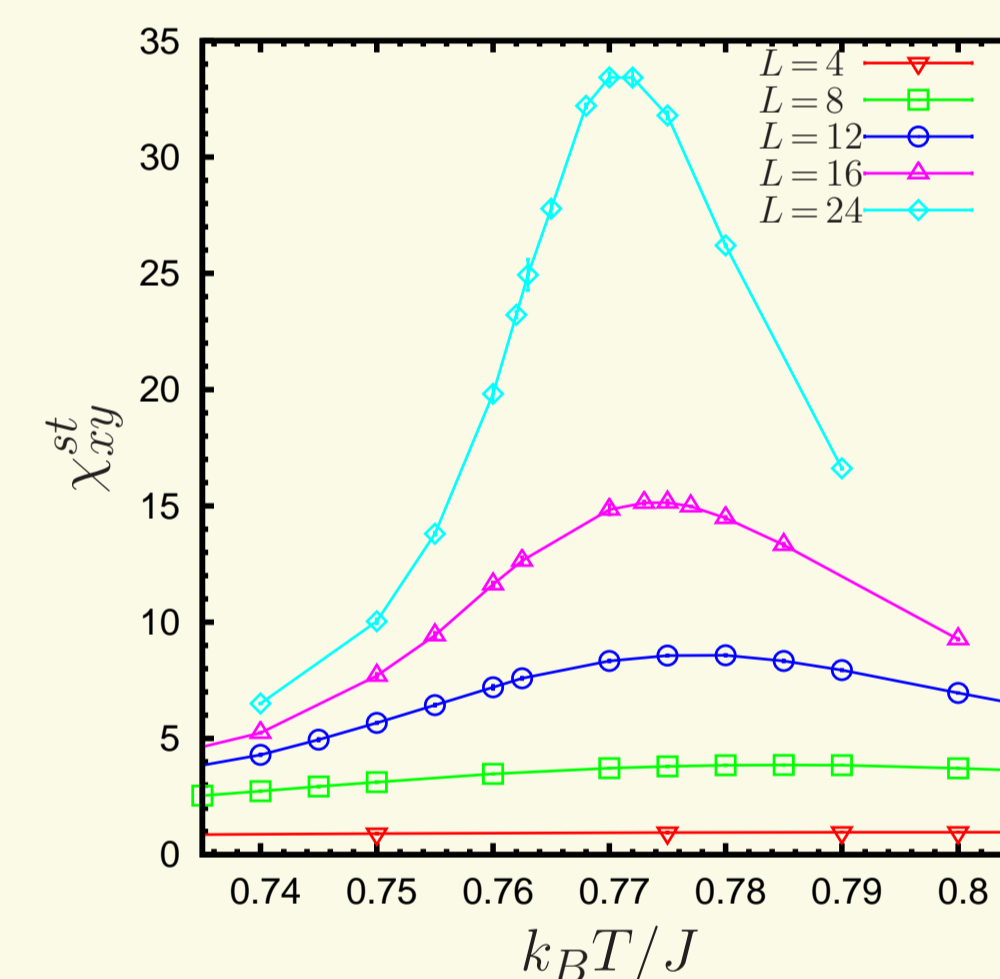
$$(\chi_z^{st})_{max} \sim L^{\frac{\gamma_z}{\nu}}$$



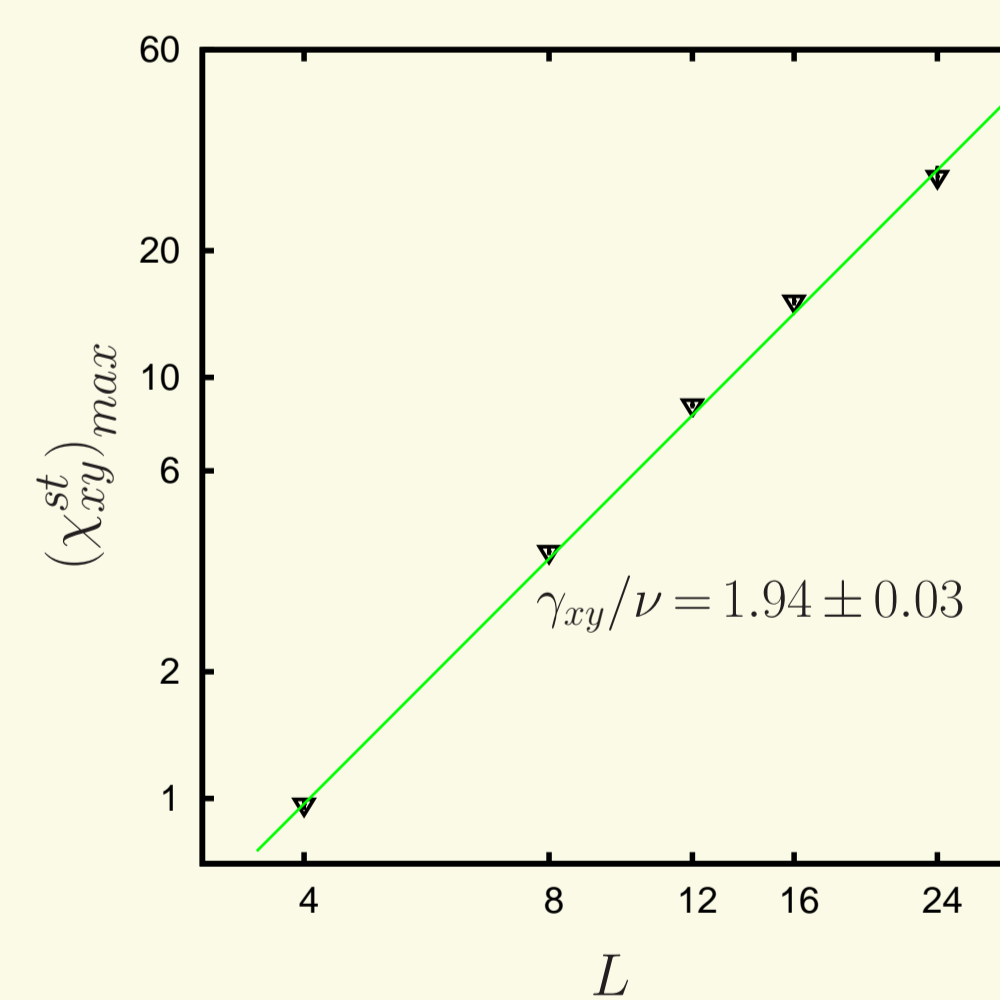
$\gamma_z/\nu = 1.95 \pm 0.01$  in agreement with **Ising universality** ( $\gamma/\nu = 1.9635\dots$ )

Confirmed by analyses of specific heat, Binder cumulant, ... for various cases

SF phase boundary ( $H/J = 7$ )



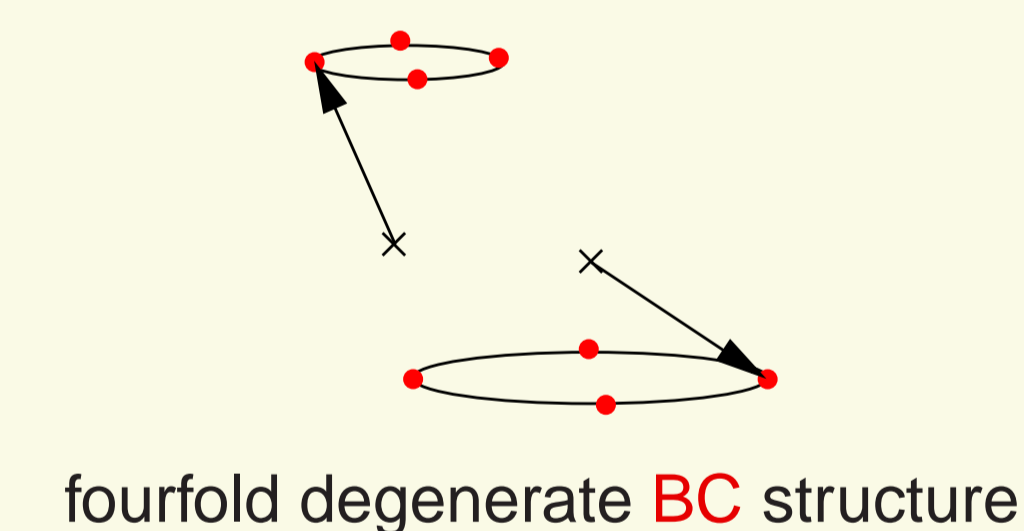
$$(\chi_{xy}^{st})_{max} \sim L^{\frac{\gamma_{xy}}{\nu}}$$



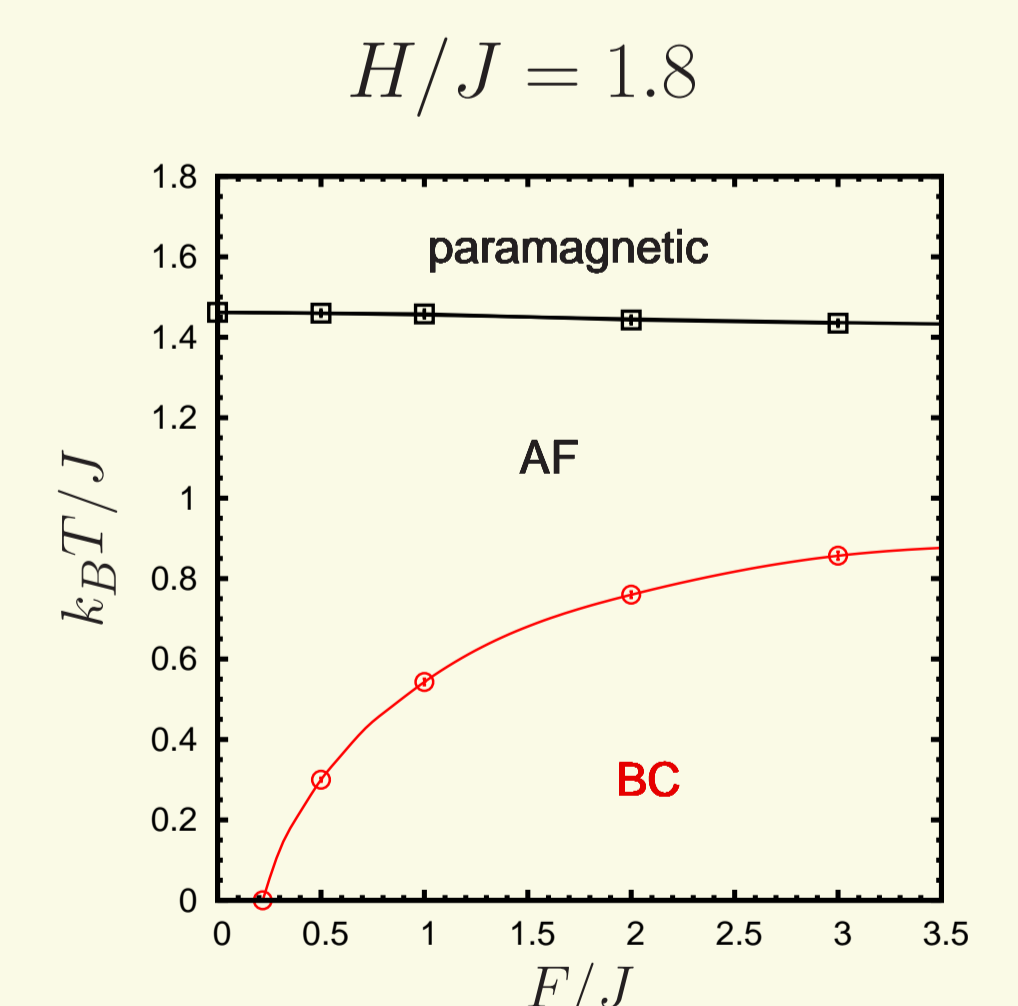
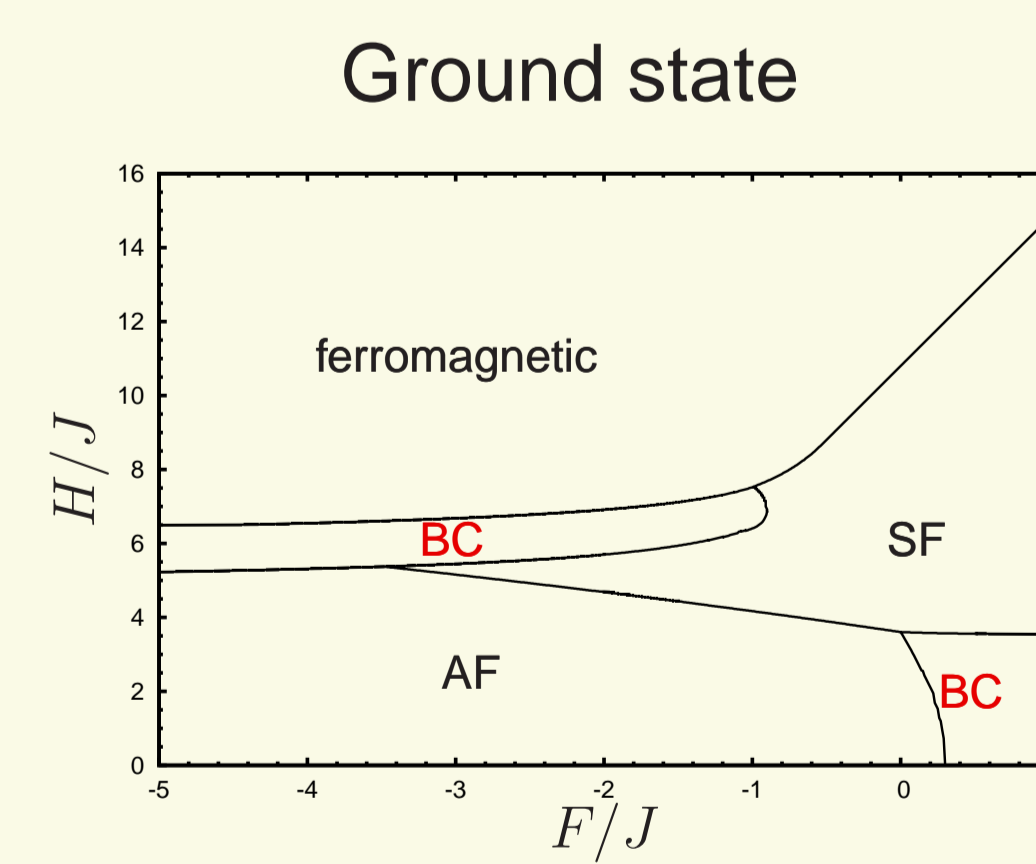
$\gamma_{xy}/\nu = 1.94 \pm 0.03$  in agreement with **XY universality** ( $\gamma/\nu = 1.9613\dots$ )

## Extension: XXZ antiferromagnet + cubic symmetry

$$\mathcal{H} = \mathcal{H}_{xxz} - F \sum_i ((S_i^x)^4 + (S_i^y)^4 + (S_i^z)^4)$$

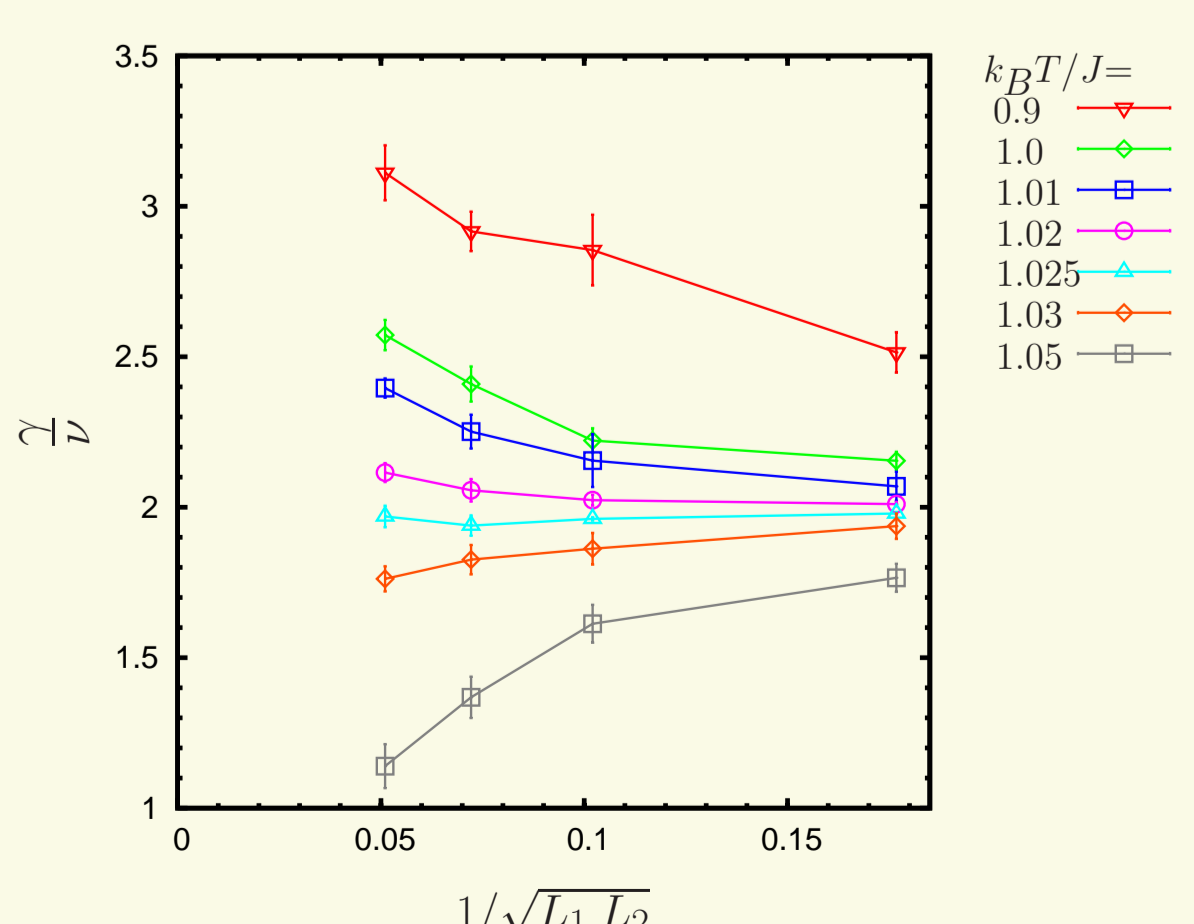


- $F > 0$ : Spins tend to align in direction of x-, y- and z-axis
- $F < 0$ : Spins tend to align along the diagonals of cubic lattice



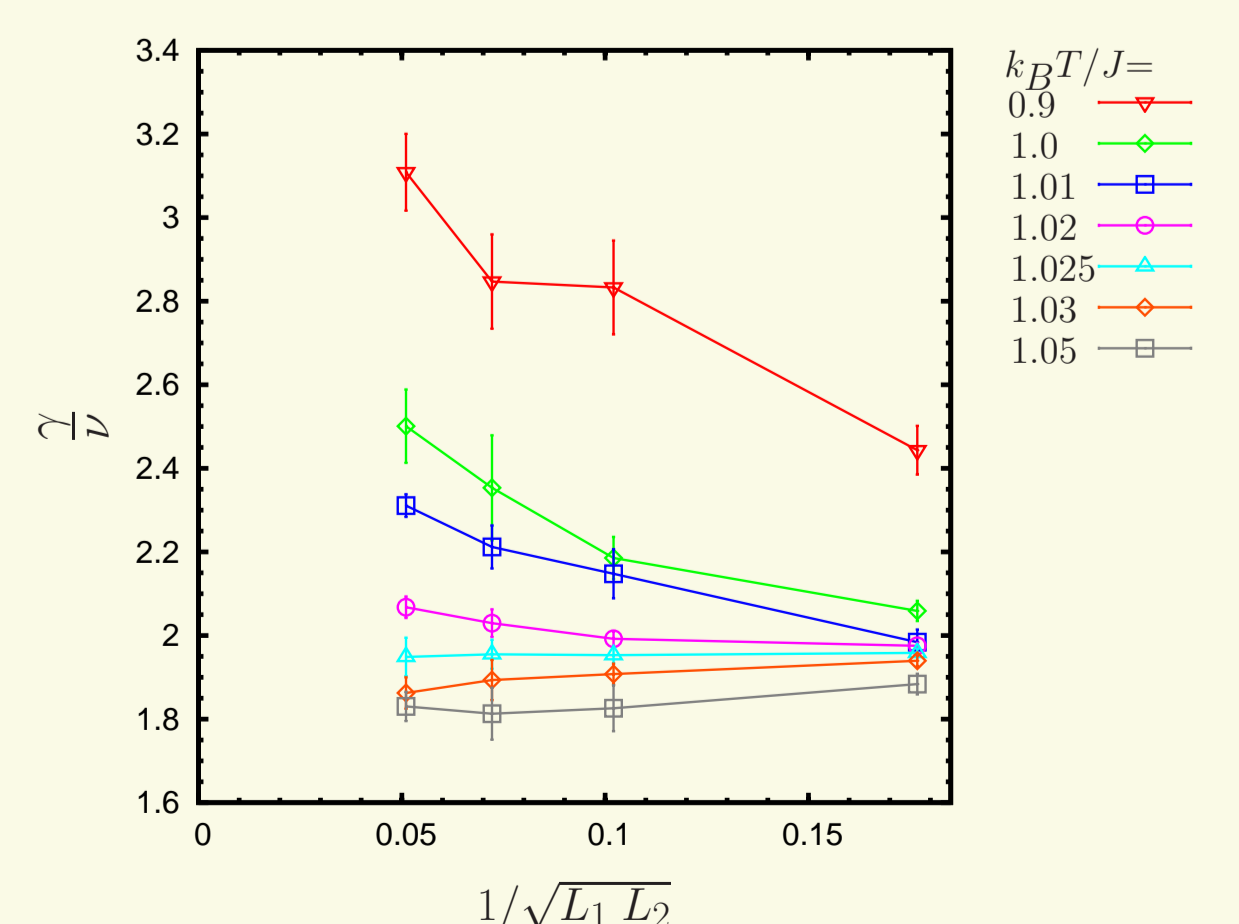
## Locating the triple point $k_B T_{tp}/J, H_{tp}/J$

effective exponents for  $\chi_{xy}^{st}$



Location:  $k_B T_{tp}/J = 1.025 \pm 0.015$   
 $H_{tp}/J = 3.90 \pm 0.03$

effective exponents for  $\chi_z^{st}$



For comparison: D. P. Landau, K. Binder, PRB (1978)

See also: R. Folk, Yu Holovatch, G. Moser, PRE (2008); P. Calabrese, A. Pelissetto, E. Vicari, PRB (2003)

## Conclusions

- At low temperatures, XXZ model in  $d = 3$  displays a first order phase transition between AF and SF phases, in contrast to  $2d$  version; BC fluctuations are observed.
- Triple point, at which AF, SF, and disordered phases meet, has been accurately located.
- Introducing a cubic symmetry term, discretized BC phases and structures are stabilized.

WS, M.Holtzschneider, R. Leidl, S. Wessel, and GB, Computer Simulations Studies in Condensed Matter Physics XXI (2008)

## Outlook, in progress

- Quantum-Heisenberg antiferromagnets with spin-liquid (spin-flop) and supersolid (biconical) phases, found using DMRG calculations.