

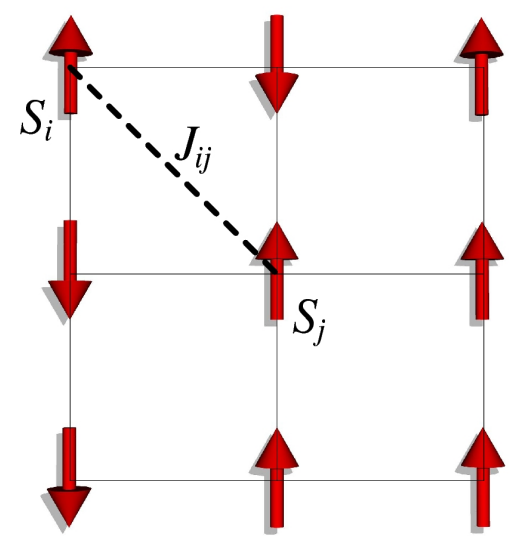
Critical Binder cumulant revisited

W. Selke (RWTH Aachen) and L. N. Shchur (Landau Institute, Chernogolovka)

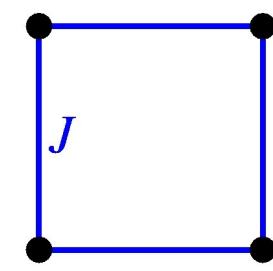
Ising model

$$\mathcal{H} = -\sum J_{ij} S_i S_j$$

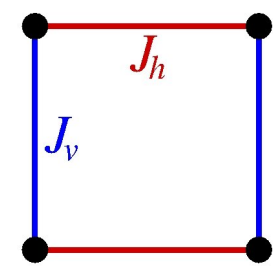
$$S_i, S_j = \pm 1$$



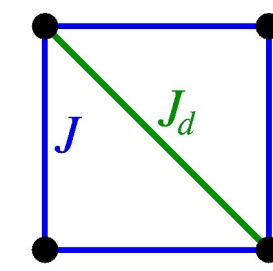
Square lattice



isotropic

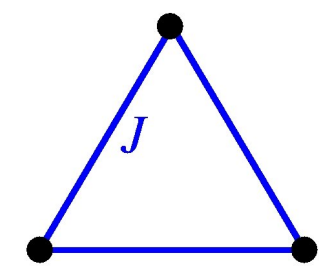


anisotropic
rectangular
symmetry



nnn interactions
along one diagonal

Triangular lattice

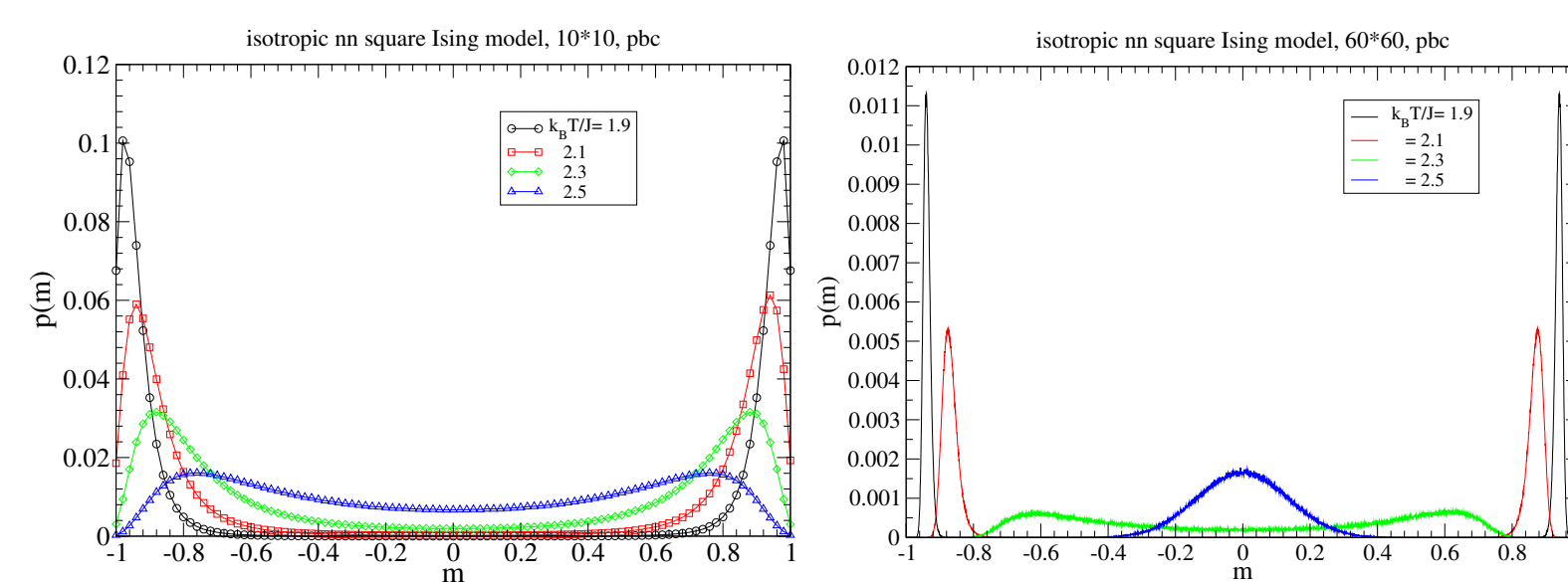


$\cong$ square lattice
with $J_d = J$

Magnetization histograms

$$\text{Distribution function of the magnetization } p(m) = \frac{1}{Z} \sum_K e^{-\mathcal{H}(K)/k_B T}$$

K : configurations with magnetization m



isotropic nn model, $L \times L$ sites, periodic boundary conditions, near critical temperature

$k_B T_c / J = 2.269...$

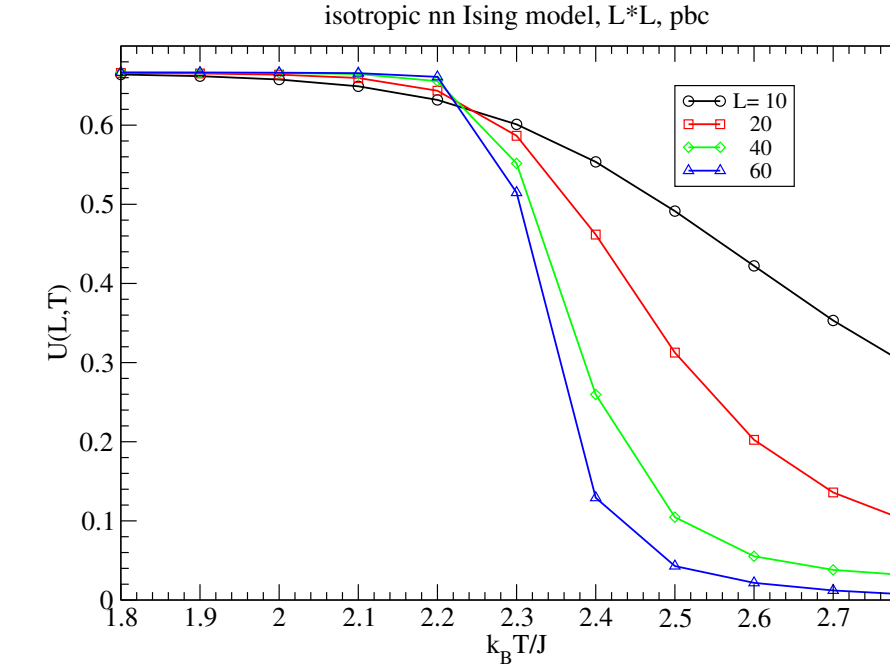
Observations: $T < T_c$: two-peak structure
 $T > T_c$: $p(m)$ approaching Gaussian function for larger L

Characterization of distributions by moments and/or cumulants:

Binder cumulant

Fourth-order cumulant of the distribution function of the magnetization

$$U = 1 - \frac{\langle M^4 \rangle}{3 \langle M^2 \rangle^2} \quad (\text{Binder, 1981})$$



Binder cumulant in the
thermodynamic limit:

$$U(T) = \begin{cases} \frac{2}{3} & T < T_c \\ U^* & T = T_c \\ 0 & T > T_c \end{cases}$$

From crossing points of $U(L_1) = U(L_2)$, one may estimate conveniently T_c :

U^* : **critical Binder cumulant** (universal?)

$U^* = 0.61069...$ (Kamieniarz + Blöte, 1993)

isotropic nn Ising model, square shape, periodic boundary conditions

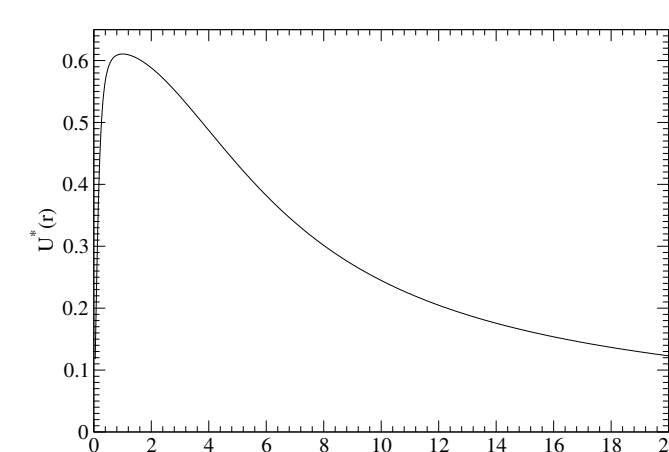
PREVIOUS results: Isotropic nn square lattice Ising model with square shape

- U^* has been found to be, employing periodic bc, **INDEPENDENT** of:
 - + **spin value**
 - $S=1/2, S=1$ (Nicolaidis and Bruce, 1988)
 - + **discrete or continuous nature of (Ising-type) spin variable**
 - Nicolaidis and Bruce, 1988; Kamieniarz and Blöte, 1993
- On the other hand, for the isotropic nn square Ising model, U^* has been observed to **DEPEND** on
 - + **boundary conditions**
 - periodic, free, fixed,...boundaries
 - K.Binder, D.W. Heermann, W.Janke, D.P.Landau, A. Milchev,...(scattered results)

Varying shape, lattice type, and anisotropy

With periodic bc, U^* has been found/argued to **DEPEND** on

- **Shape of the lattice**



$U^*(r)$: isotropic, square nn Ising model with rectangular shape, aspect ratio r ; exact calculations augmented by finite-size extrapolations

Kamieniarz and Blöte (K+B), 1993

- **Lattice type(?)**

Slightly different values of U^* for isotropic nn Ising models on **square lattice with square shape** and on **triangular lattice with rhombus shape** (K+B, 1993)

- **Anisotropy of nn couplings, J_v/J_h :**

$$U^* = U^*(J_v/J_h)$$

for square shapes, $r = 1$; with a mapping onto the isotropic nn Ising model with rectangular shape so that

$$U^*(r = 1, J_v/J_h) = U^*(r, J_v/J_h = 1) \text{ where}$$

$$\sinh(2J_h/k_B T_c(J_v/J_h)) = r,$$

(which follows from setting $r = \xi_v/\xi_h$; ξ correlation length) (K+B, 1993)

- **Anisotropy of nnn couplings, J_d :**

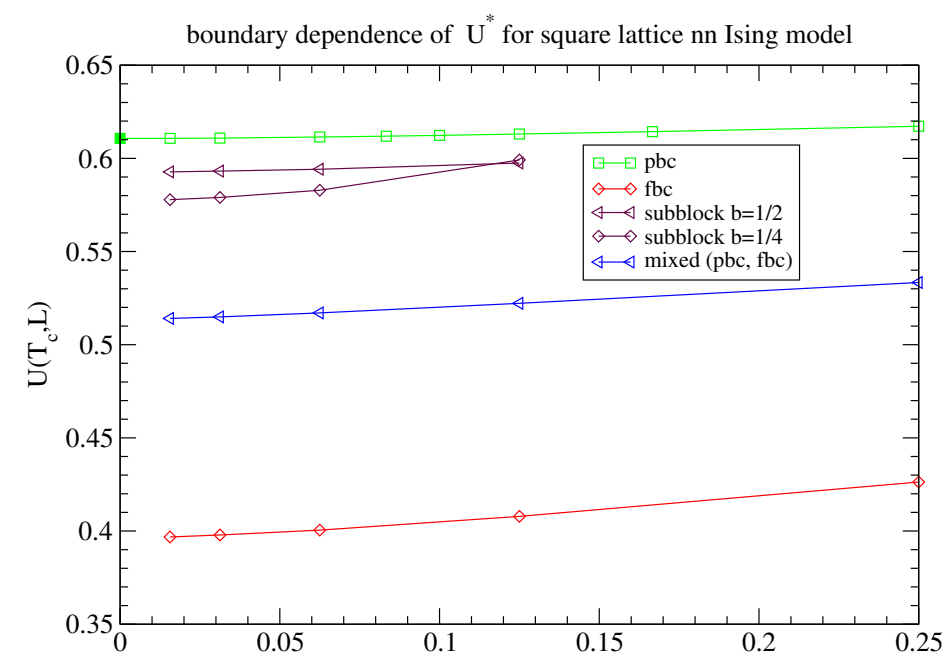
$$U^* = U^*(J_d/J)$$

but there is no mapping

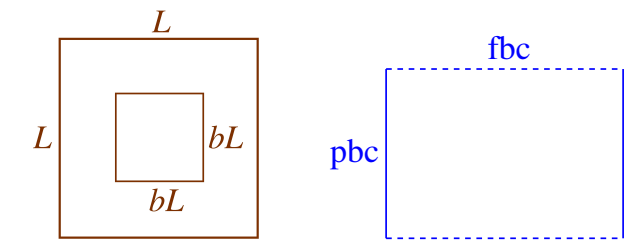
$$U^*(J_d/J, r = 1) = U^*(0, r), \text{ which would keep rectangular symmetry (Chen+Dohm, 2004)}$$

Our aim: Monte Carlo study on those (non)universal aspects of U^*

PRESENT results: isotropic nn Ising model on square lattice – various boundary conditions

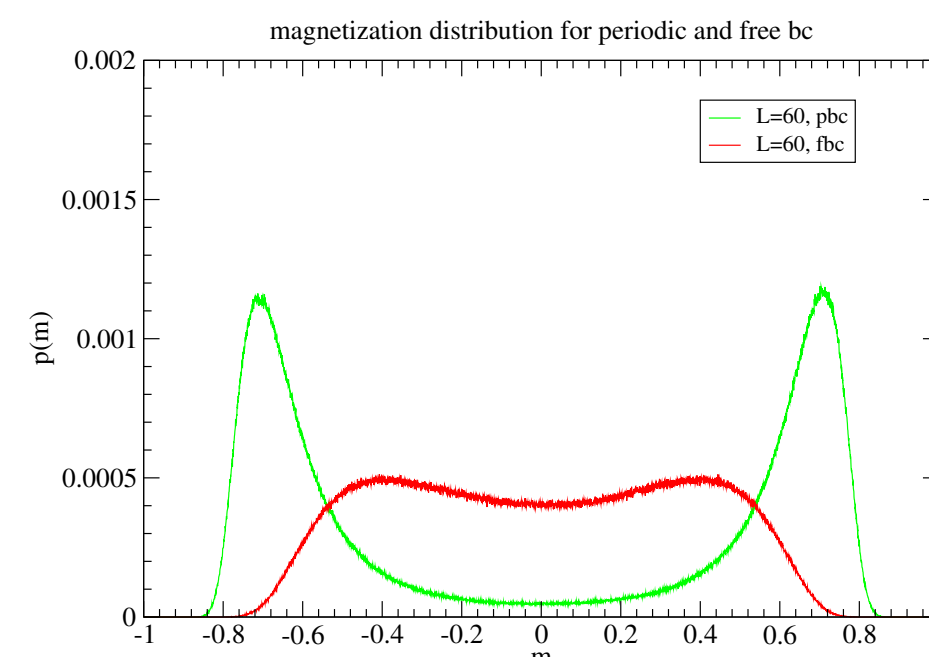


Cumulant at T_c for square lattice with **free, periodic, mixed** and other boundary conditions



Subblocks (Binder, 1981): squares of size $bL \times bL$, embedded in square of $L \times L$ sites; 'heat bath bc' when $b \rightarrow 0$, with $U^* = 0.560 \pm 0.002$

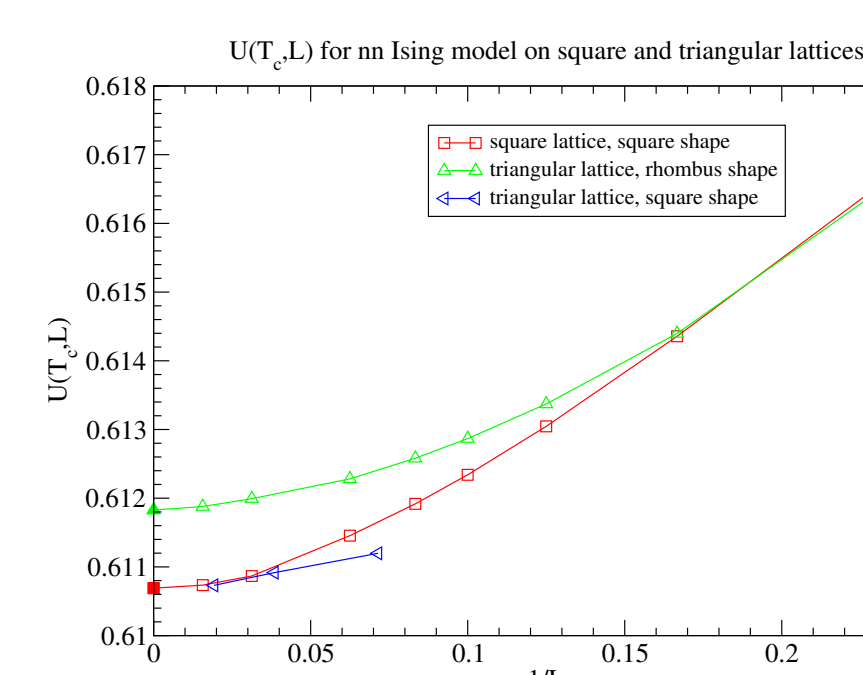
Mixed bc: pbc for two opposite sides, fbc for the other two opposite sites of squares of size L^2



Cumulant at T_c for square lattice, $L = 60$, with **periodic** and **free** boundary conditions

Note: less pronounced two-peak-distribution for **free boundary conditions**, and U^* is smaller than in the case of **periodic boundary conditions**

Isotropic nn Ising model on triangular and square lattices

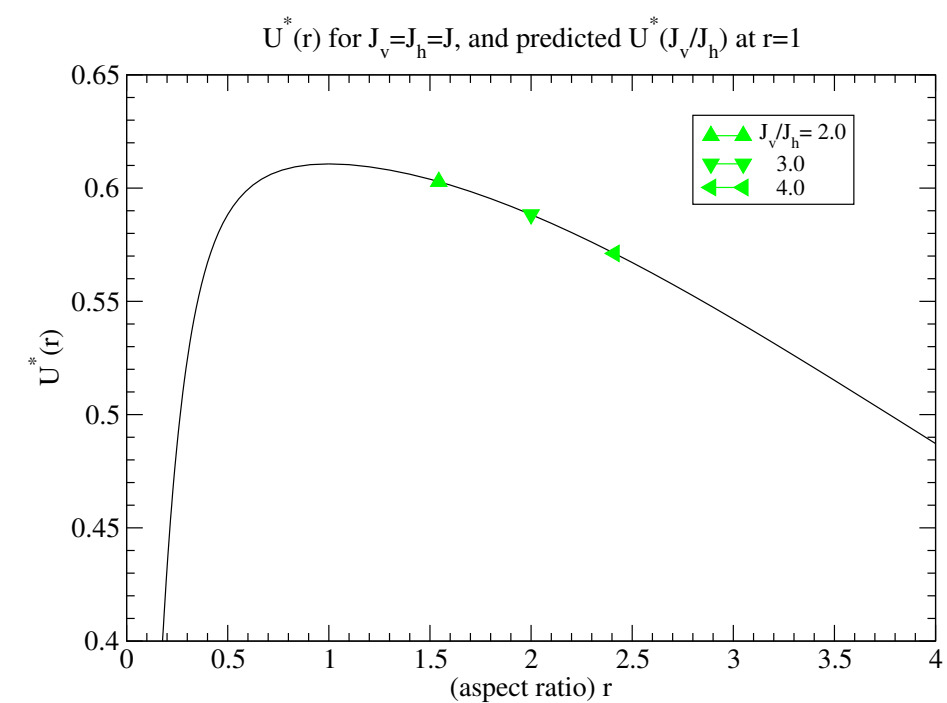
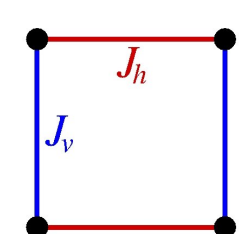


$U(T_c, L)$ for isotropic nn Ising model on **square lattice (square shape)**, **triangular lattice (rhombus shape)**, and **triangular lattice (square shape)**

W.S., E. P.J.B, 2006; J.Stat.Mech, 2007

Suggestion: For given shape, isotropic models lead to the same U^* (checked, in addition, for other rectangular shapes)

Square lattice with different nn (horizontal and vertical) couplings



PREDICTION:

$$U^*(J_v = J_h, r) =$$

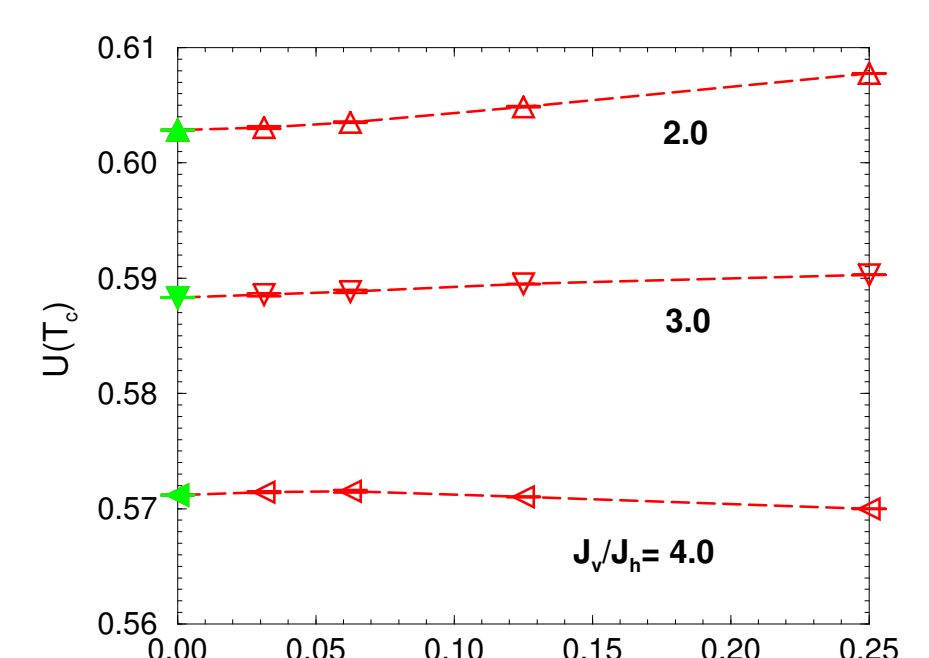
$$U^*(J_v/J_h, r = 1) \text{ such that:}$$

$$\sinh(2J_h/k_B T_c) = r$$

(K+B, 1993)

To be checked by MC simulations:

Checking the prediction

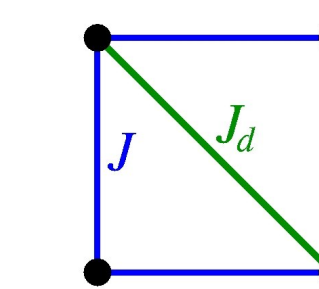


Cumulant $U(T_c, L)$ for **anisotropic nn Ising model, L^2 spins ($r = 1$)**, and various J_v/J_h , as compared to previous findings on $U^*(J_v/J_h = 1, r)$ in the isotropic case, presuming the mapping

Our findings confirm the predicted mapping

$$U^*(J_v/J_h, r = 1) = U^*(J_v/J_h = 1, r) \text{ with } \sinh(2J_h/k_B T_c) = r$$

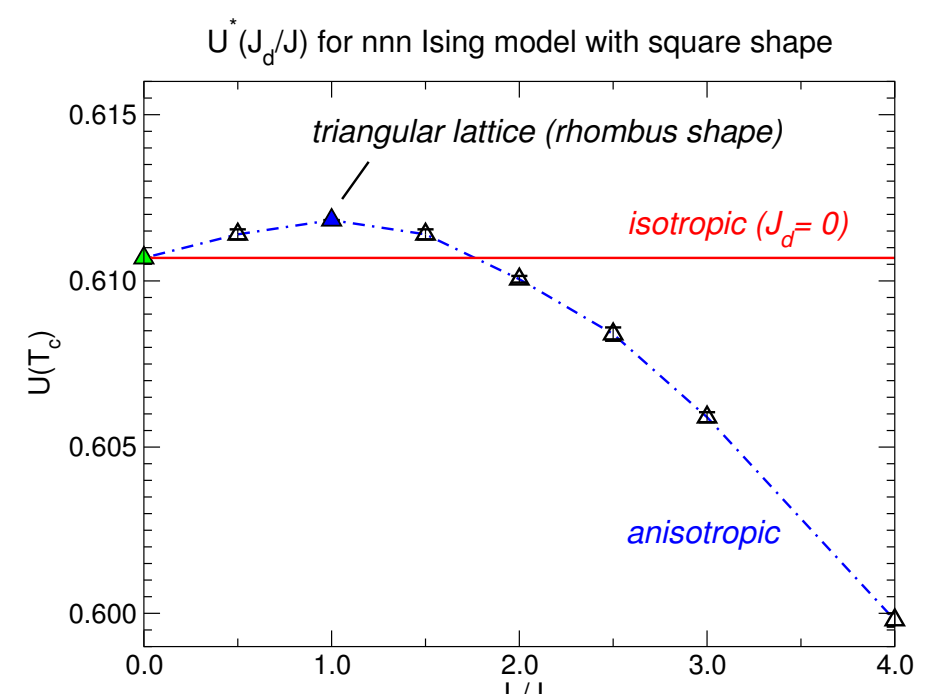
Square lattice with anisotropic nnn interactions



Statements of Chen and Dohm (2004):

- $U^*(J_d/J, r = 1)$ varies continuously with J_d/J
- Keeping rectangular symmetry, there is no mapping of U^* onto the isotropic case such that $U^*(r = 1, J_d/J) = U^*(r, J_d/J = 0)$
- In general: Violation of 'two-scale-factor universality' due to anisotropy

Checking by simulations



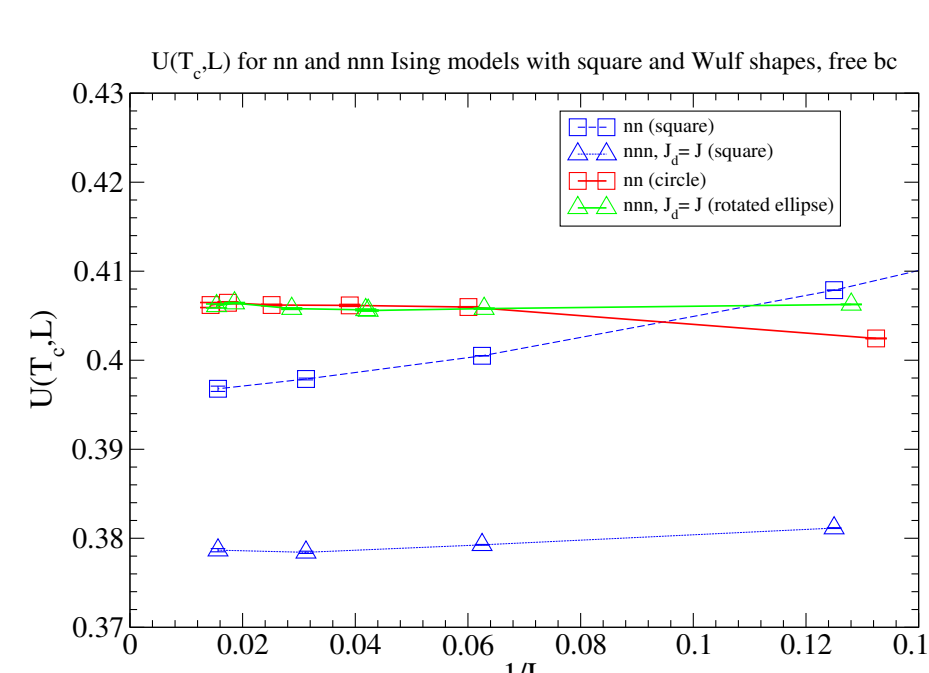
Simulated U^* for nnn Ising model on square lattice with $r=1$, including **nn isotropic case, $J_d = 0$**

Overshooting: No mapping with $U^*(J_d/J, r = 1) = U^*(J_d/J = 0, r)$

Note: U^* of square lattice nnn Ising model, $J_d = J$, with square shape is identical to that of the triangular lattice nn Ising model with rhombus shape (similarly for other ratios of J_d/J (Dohm, 2006)).

Question: Can dependence of U^* on anisotropy be transcribed, in general, into dependence on shape?

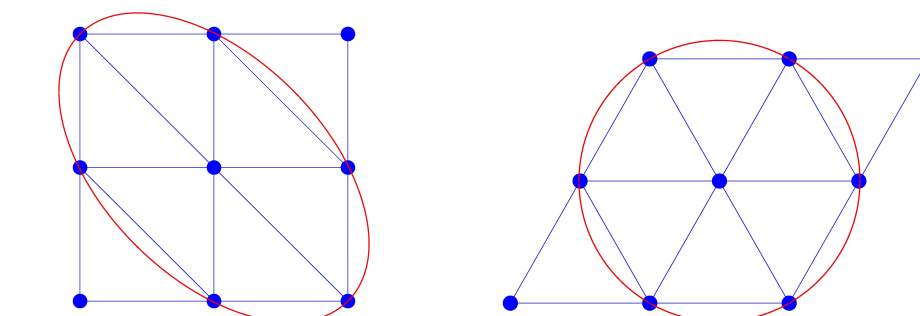
Critical Binder cumulant in (an)isotropic systems and Wulf shape



Wulf shapes at T_c with free bc: U^* for **nn isotropic square Ising model with circle shape** and **nnn anisotropic, $J_d = J$, case with rotated ellipse shape**, having L^2 sites.- For comparison: nn and nnn cases, free bc, with **square shapes, L^2 sites**.

Recall: The equilibrium Wulf shape, at T_c , of an Ising droplet results from the orientational dependence of the surface free energy at criticality, reflecting the interactions.

Note:



There are the same spins in the **rotated ellipse** for the anisotropic nnn, $J_d = J$, Ising model on the square lattice and in the **circle** for the nn Ising model on the **triangular lattice**; thence $U^*(nnn, sq, ellipse) = U^*(iso nn, tria, circle) = U^*(iso nn, sq, circle)$

Question(suggestion):

Does U^* take a generic/unique value when one considers systems (free bc) with their Wulf shape at criticality?

Conclusions:

- The critical Binder cumulant U^* in 2d Ising models **depends on boundary conditions, system shapes, anisotropy of interactions**.
- For **isotropic** models, U^* depends on shape and boundary conditions, but not on details of interactions and lattice type.
- For given boundary condition, the dependence of U^* on **ANISOTROPY** may be mapped onto a dependence on the **SHAPE**: verification for the nn anisotropic case, keeping rectangular symmetry; evidence for the nnn anisotropic Ising model, considering rhombus (parallelogram) shapes.
- **Question**: Can a generic/unique value of U^* be obtained for Ising models with a **shape following from the Wulf construction at criticality, using, e.g., free boundary conditions**?