# The critical behaviour of 3D Ising spin glass models: universality and scaling corrections

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### Overview

M. Hasenbusch, A. Pelissetto, E. Vicari, arXiv:cond-mat/0710.1980

- Definition of the Model
- Summary of results given in the literature
- Corrections to scaling and the Renormalization Group
- Our Monte Carlo Simulations
- Conclusions

We study the  $\pm J$  model on a simple cubic lattice with periodic boundary conditions. The Hamiltonian is given by

$$H = -\sum_{\langle xy \rangle} J_{\langle xy \rangle} s_x s_y$$

with  $s_x \in \{-1, 1\}$  and  $J_{\langle xy \rangle} \in \{-1, 1\}$ The distribution of the bonds: P(1) = 1 - P(-1) = p

We have simulated p = 0.5 and p = 0.7.

 $\pm J$  model with bond dilution:  $J_{\langle xy \rangle} \in \{-1, 0, 1\}$  with probabilities  $P(-1) = (1 - p)p_b$ ,  $P(0) = (1 - p_b)$ , and  $P(1) = pp_b$ .

We have simulated p = 0.5 and  $p_b = 0.45$ .

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### The observables

Overlapp variables:  $q_x = s_x^{(1)} s_x^{(2)}$ 

Second moment correlation length over the linear lattice size L:

 $R_{\xi} = \xi/L$ 

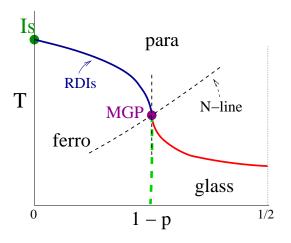
Cumulants:

$$U_4 = \frac{[\mu_4]}{[\mu_2]^2}, \quad U_{22} = \frac{[\mu_2^2] - [\mu_2]^2}{[\mu_2]^2}$$

with  $\mu_k = \langle (\sum_x q_x)^k \rangle$ 

[]: average over all bond configurations (samples)  $\{J_{\langle xy \rangle}\}$ 

Phase diagram of the  $\pm J$  model



 $p_{MGP} = 0.76820(4), 2p - 1 = \tanh(J/T)$ 

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Results for  $\pm J$  model at p = 0.5 taken from H. Katzgraber, M. Körner, and A. P. Young, Phys. Rev. B **73**, 224432 (2006).

authors	year	T <sub>c</sub>	ν	$\eta$
Ogielski, Morgenstern	1985	1.20(5)	1.2(1)	
Ogielski	1985	1.175(25)	1.3(1)	-0.22(5)
Singh, Chakravarty	1986	1.2(1)	1.3(2)	
Bhatt, Young	1985	1.2(2)	1.3(3)	-0.3(2)
Kawashima, Young	1996	1.11(4)	1.7(3)	-0.35(5)
Bernardi et al	1996	1.165(10)		-0.245(20)
Berg, Janke	1998	1.12(1)		-0.37(4)
Palassini, Caracciolo	1999	1.156(15)	1.8(2)	-0.26(4)
Mari, Campbell	1999	1.20(1)		-0.21(2)
Ballesteros et al.	2000	1.138(10)	2.15(15)	-0.337(15)
Mari, Campbell	2001	1.190(15)		-0.20(2)
Mari, Campbell	2002	1.195(15)	1.35(10)	-0.225(25)
Nakamura et al	2003	1.17(4)	1.5(3)	-0.4(1)
Pleimling, Campbell	2005	1.19(1)		-0.22(2)
Katzgraber et al.	2006	1.120(4)	2.39(5)	-0.395(17)
our result	2007	1.101(5)	2.53(8)	-0.384(9)

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Estimates for  $\nu$  from different quantities differ quite a lot:

E.g. FSS (finite size scaling), Katzgraber et al. (Using power law ansätze without corrections):

- ► 2.39(5) from <u>*ξ*/L</u>
- ▶ 2.79(11) from the Binder Cumulant
- 1.57(3) from the slope of  $\chi$

#### $\Rightarrow$ Have to understand corrections to scaling

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### The singular part of the free energy

 $f(\beta, h, L) = L^{-d} f_s(u_t L^{y_t}, u_h L^{y_h}, u_3 L^{y_3}) + f_{ns}(\beta, h)$ 

where  $u_t$ ,  $u_h$  and  $u_3$  are non-linear scaling fields:  $u_t = t + ah^2 + ...$  and  $u_h = h + bth + ...$  where  $t = \beta - \beta_c$ . For h = 0:

$$\chi = \frac{\partial^2 f}{\partial h^2} = L^{-d} (f_s^{(0,2)} L^{2y_h} (1 + 2bt + ...) + ...)) + f_{ns}^{(0,2)}$$

$$\frac{\partial \chi}{\partial \beta} = L^{-d} (f^{(1,2)} L^{2y_h + y_t} (1 + \dots) + 2b f^{(0,2)} L^{2y_h} (1 + \dots) + \dots) + f^{(1,2)}_{ns}$$

$$\propto L^{2y_h+y_t-d}(1+cL^{-y_t}+...)$$

The derivatives of  $\xi/L$  and the Binder cumulant with respect to  $\beta$  at  $\beta_c$  do not suffer from corrections  $\propto L^{-y_t}$ 

## The simulation

- Local Metropolis updates
- Multispin coding implementation (64 systems run in parallel, same random number for all systems)
- Random exchange method (Parallel tempering)
- Avoid bias in quantities like  $[\langle A \rangle \langle B \rangle]$  ...
- Careful check of equilibration: double lenght of the equilibration time until results are consistent within errors
- We compute the Taylor expansion of the observables up to 2<sup>nd</sup> order

In total about 30 years of a single core 2.4 GHz Opteron CPU.

The Models

L	samples/64	MC sweeps	Nt	CPU-time/days
8	100000	19200	5	4
9	110850	48000	8	27
10	100681	72000	8	50
11	109779	144000	10	183
12	106812	192000	10	308
13	38282	288000	10	210
14	31600	480000	10	361
16	24331	480000	20	831
20	1542	1920000	32	658
24	717	3000000	32	826
28	285	7200000	20	782

Effort grows roughly  $\propto L^9$ .

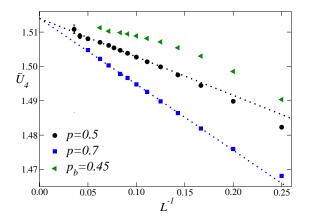
Compare with Katzgraber et al: Number of samples: 30000,15807,11360,9408,8416 for L = 8, 12, 16, 20, 24 respectively.

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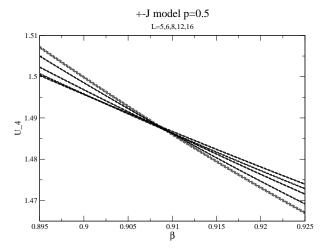
(A) (A) (B) (A)

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 $U_4$  at  $\xi/L = 0.63$  fixed ploted vs 1/L



Conclusion from various fits:  $\omega = 1.0(1)$ 

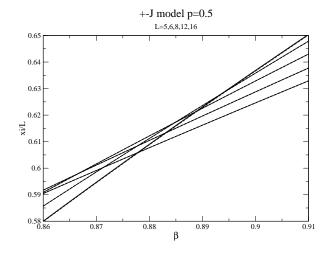


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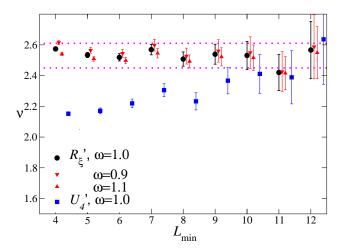
The Models	Results from the Literature	Analytic Corrections	Numerical Results	Conclusions
► V	ery small corrections	in the case of $U_4$		
		b = c = c = c = c = c = c = c = c = c =		

Strong corrections in the case of ξ/L; crossing points move to larger values of β and ξ/L as L increases
Including corrections with ω = 1 we get β<sub>c</sub> consistent with that obtained from U<sub>4</sub>

Our final results (From data with  $L \ge 8$ ):

 $U_4^* = 1.490(7), \ \xi/L^* = 0.654(7), \ \beta_c = 0.908(4)$  for p = 0.5.

Ansatz:  $R' = L^{1/\nu} (1 + cL^{-\omega})$ 



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## Conclusions

- In some quantities  $1/\nu$  corrections are leading!
- $U_4$  at  $\xi/L$  fixed is the same for all three models
- $\omega = 1.0(1)$
- ▶  $\nu = 2.53(8), \ \eta = -0.384(9)$

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