

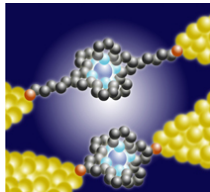
Switching the current through molecular wires with ultra-short laser pulses

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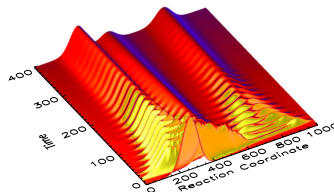
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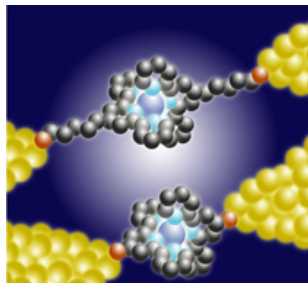
01. December 2006



Outline

- 1 Introduction
- 2 The quantum master equation
- 3 Short Gaussian laser pulses
- 4 Optimal control of the current
- 5 Summary

Molecular wires



- first reproducible experiments on molecular wires
 - break junctions, STM setups, quantum dot geometries, . . .
- influence of laser light on molecular wires gives an opto-electronic coupling
- with femtosecond laser pulses: high spatial as well as temporal resolution

Density matrices

- time-dependent Schrödinger equation

$$i\hbar \frac{d}{dt} |\Psi(x, t)\rangle = H(x, t) |\Psi(x, t)\rangle$$

- pure state: $\sigma = |\Psi\rangle\langle\Psi|$
- mixed state: $\sigma = \sum_n W_n |\Psi_n\rangle\langle\Psi_n|$
- time propagation

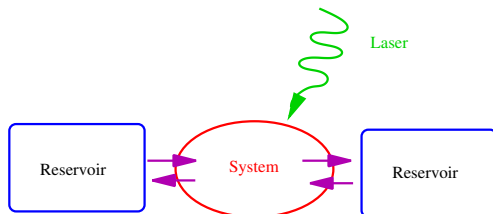
$$\begin{aligned} i\hbar \frac{d\sigma(t)}{dt} &= i\hbar \frac{d|\Psi\rangle\langle\Psi|}{dt} = i\hbar \left(\frac{d|\Psi\rangle}{dt} \langle\Psi| + |\Psi\rangle \frac{d\langle\Psi|}{dt} \right) \\ &= H(t)\sigma(t) - \sigma(t)H(t) = [H(t), \sigma(t)] \end{aligned}$$

- Observables:

$$\begin{aligned} \langle Q \rangle &= \langle \Psi | Q | \Psi \rangle = \sum_n \langle \Psi | n \rangle \langle n | Q | \Psi \rangle \\ &= \sum_n \langle n | Q | \Psi \rangle \langle \Psi | n \rangle = \text{tr}(Q\sigma) \end{aligned}$$

Reduced density matrix formalism

- Goal: Description of **ultra-fast (fs) processes** in the molecular wire
- **full quantum dynamics** including dephasing, energy dissipation but also coherences, not only steady state, temperature dependence
- splitting in relevant system and electronic reservoirs

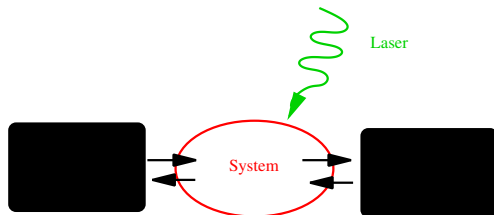


- σ - density matrix of the full system
(relevant system + bath)

$$i\hbar \frac{d\sigma(t)}{dt} = [H(t), \sigma(t)]$$

Reduced density matrix formalism

- Goal: Description of **ultra-fast (fs) processes** in the molecular wire
- **full quantum dynamics** including dephasing, energy dissipation but also coherences, not only steady state, temperature dependence
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- reduced density-matrix:
 $\rho = \text{tr}_R(\sigma)$ - density matrix of the relevant system

$$i\hbar \frac{d\rho(t)}{dt} = [H_S(t), \rho(t)] + \mathcal{D}(t)\rho(t)$$

System-bath coupling

Hamiltonian

$$H = H_S + H_R + H_{SR}$$

- every reservoir degree of freedom only slightly distorted
⇒ modeled by **harmonic oscillators**
- how strongly does the reservoir absorb energy?
⇒ **spectral density $J(\omega)$**
- **perturbation theory** in the system-bath coupling H_{SB}
(2nd order or higher)

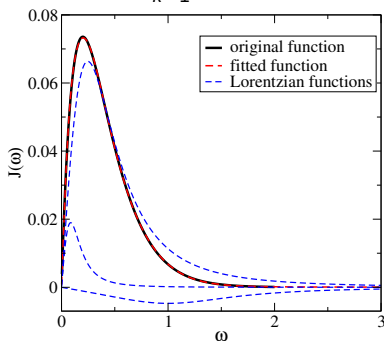
Decomposition of the spectral density

- information on the frequencies of the reservoir modes and their coupling to the system

$$J(\omega) = \frac{\pi}{2} \sum_i \frac{c_i^2}{m_i \omega_i} \delta(\omega - \omega_i)$$

- numerical decomposition into Lorentzians

$$J(\omega) = \sum_{k=1}^n \frac{p_k}{4\Omega_k} \frac{1}{(\omega - \Omega_k)^2 + \Gamma_k^2}$$



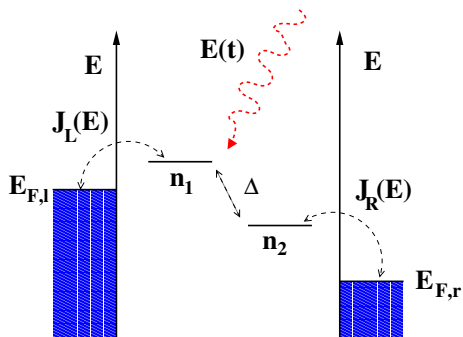
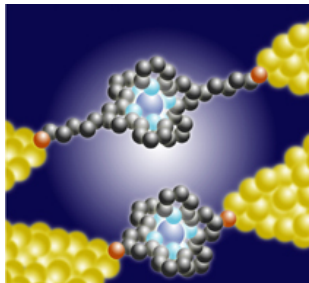
Overview of the theory

- reservoir correlation function $C(t) = \int_{-\infty}^{\infty} \frac{d\omega}{\pi} J(\omega) \frac{e^{i\omega t}}{e^{\beta\omega} - 1}$
 $\Rightarrow$ sum of exponentials in t
- this allows further analytical treatment
- definition of **auxiliary density matrices** (time-nonlocal approach) or **auxiliary operators** $\Lambda_{12}^k(t)$ (time-local approach)
- instead of one quantum master equation we get **one master equation of the system and several auxiliary master equations**
- Matsubara expansion $\Rightarrow$ many master equations for **low temperatures**
- **no further approximation in the light-matter coupling** (beyond semi-classical approximation)

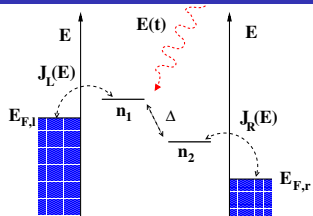
U. Kleinekathöfer, J. Chem. Phys. **121**, 2505 (2004).

S. Welack, M. Schreiber, U. Kleinekathöfer, J. Chem. Phys. **124**, 044712(2006).

The model

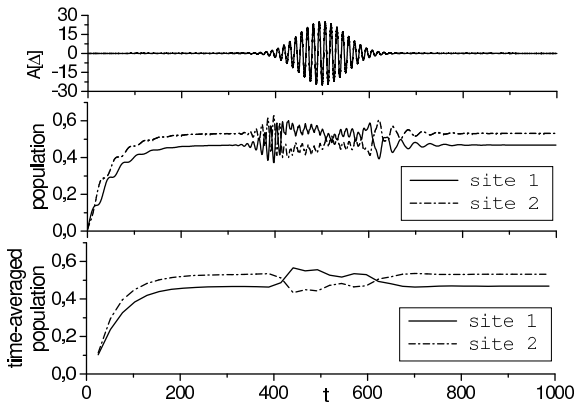


Population dynamics



$$U_n(t) = E(t)\delta_{1n} - E(t)\delta_{2n}$$

$$E_1 = E_2 = E_{F,r} + \Delta = E_{F,l} - \Delta$$



Determination of the current

- using electron number operator $N_L = \sum_q c_q^\dagger c_q$

$$I_L(t) = e \frac{d}{dt} \text{tr} \{ N_L \sigma(t) \} = -ie \text{tr} \{ [N_L, H(t)] \sigma(t) \}$$

- defining auxiliary operators

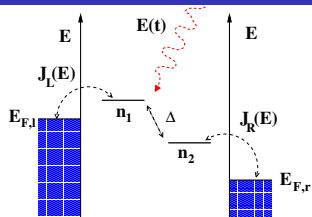
$$\Lambda_{12}(t) = \int_{t_0}^t dt' C_{12}(t-t') U_S(t, t') c_1 \rho_S(t') = \sum_k^{m+m'} \Lambda_{12}^k(t)$$

$$\Lambda_{21}(t) = \int_{t_0}^t dt' C_{21}^*(t-t') U_S^\dagger(t, t') c_1^\dagger \rho_S(t') = \sum_k^{m+m'} \Lambda_{21}^k(t)$$

- equations of motion for auxiliary operators $\Lambda_{12}^k(t)$ and $\Lambda_{21}^k(t)$
- final current equation

$$I_I(t) = 2e \text{Re} \left(\text{tr}_S \{ c_1^\dagger \Lambda_{12}(t) \} - \text{tr}_S \{ c_1 \Lambda_{21}(t) \} \right)$$

Coherent destruction of tunneling (CDT)

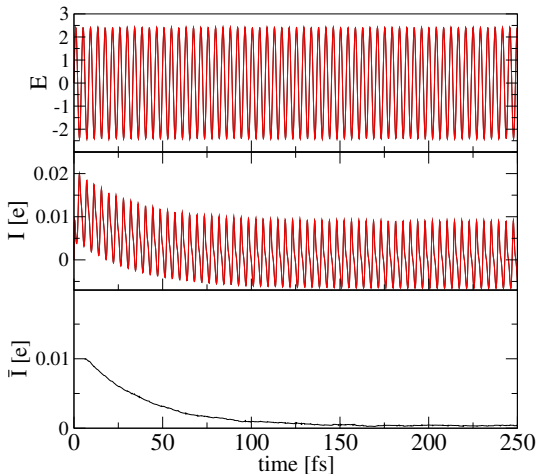


$$U_n(t) = E(t)\delta_{1n} - E(t)\delta_{2n}$$

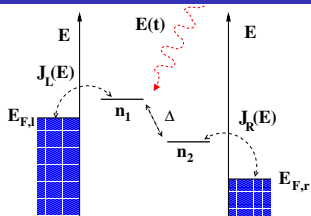
$$E(t) = E_0 \sin(\omega t)$$

$$\Delta = 0.1 \text{ eV}$$

$$1 [e] \approx 2.4 \times 10^{-4} \text{ A}$$

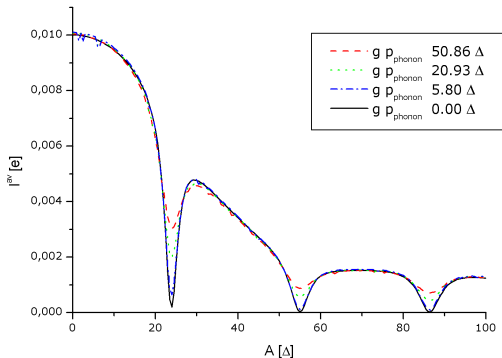


Coherent destruction of tunneling (CDT)

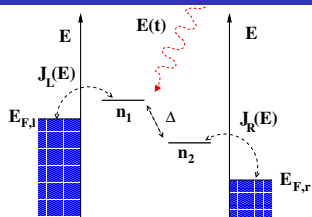


$$U_n(t) = E(t)\delta_{1n} - E(t)\delta_{2n}$$

$$E(t) = E_0 \sin(\omega t)$$

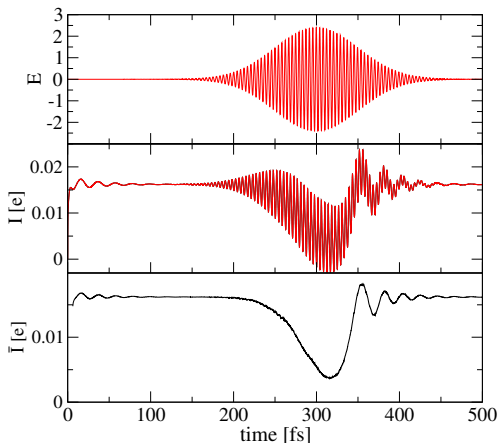


CDT: Short laser pulse

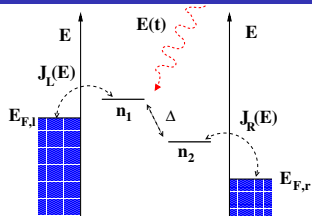


$$U_n(t) = E(t)\delta_{1n} - E(t)\delta_{2n}$$

$$E(t) = E_0 \exp\left(\frac{-(t-T)^2}{\sigma^2}\right)$$

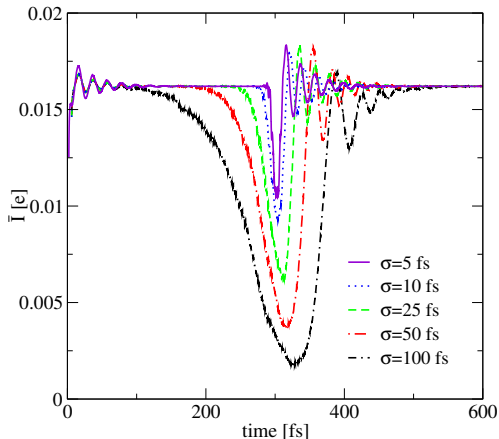


CDT: Pulse length σ

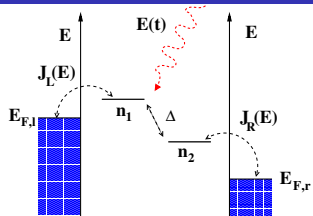


$$U_n(t) = E(t)\delta_{1n} - E(t)\delta_{2n}$$

$$E(t) = E_0 \exp\left(\frac{-(t - T)^2}{\sigma^2}\right)$$

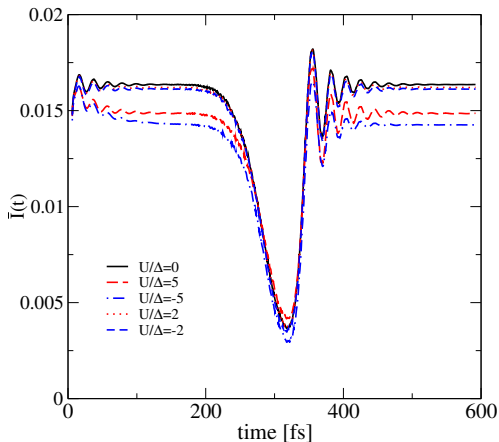


CDT: Electron correlation for spinless fermions

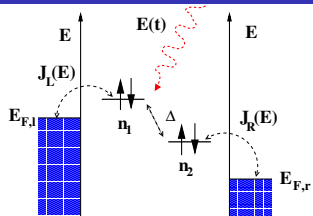


$$H_{corr} = U \sum_n n_n n_{n+1}$$

$$n_n = c_n^\dagger c_n$$

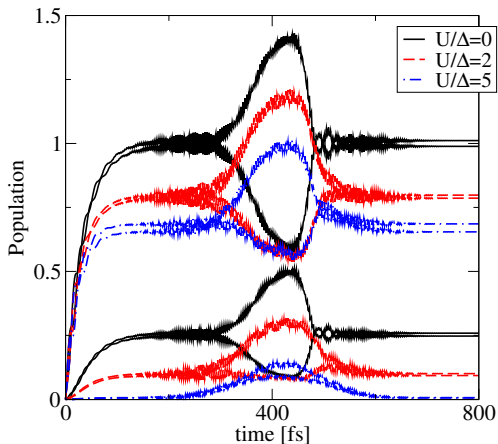


CDT: Electron correlation for fermions with spin

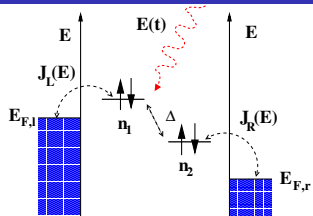


$$H_{corr} = U \sum_n n_{n\uparrow} n_{n\downarrow}$$

$$n_{n\sigma} = c_{n\sigma}^\dagger c_{n\sigma}$$

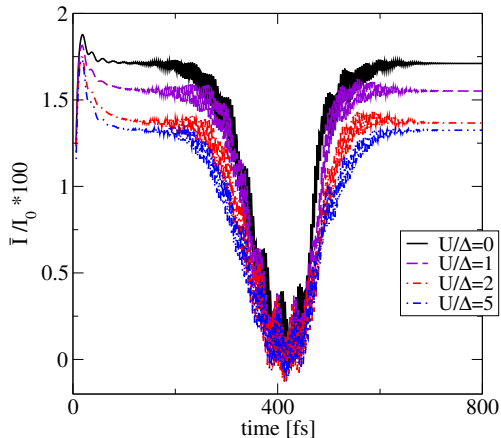


CDT: Electron correlation for fermions with spin



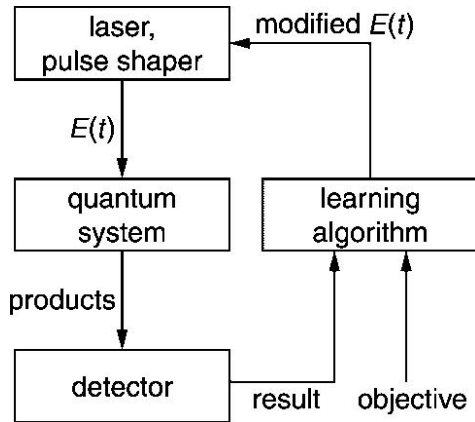
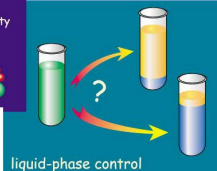
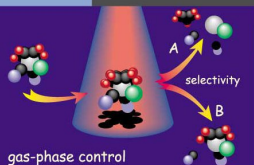
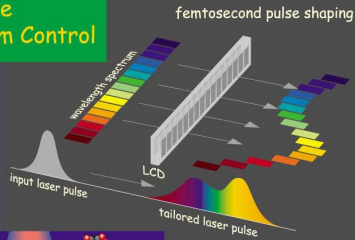
$$H_{corr} = U \sum_n n_{n\uparrow} n_{n\downarrow}$$

$$n_{n\sigma} = c_{n\sigma}^\dagger c_{n\sigma}$$



Coherent control with femtosecond laser pulses

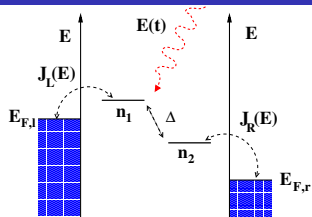
Adaptive Quantum Control



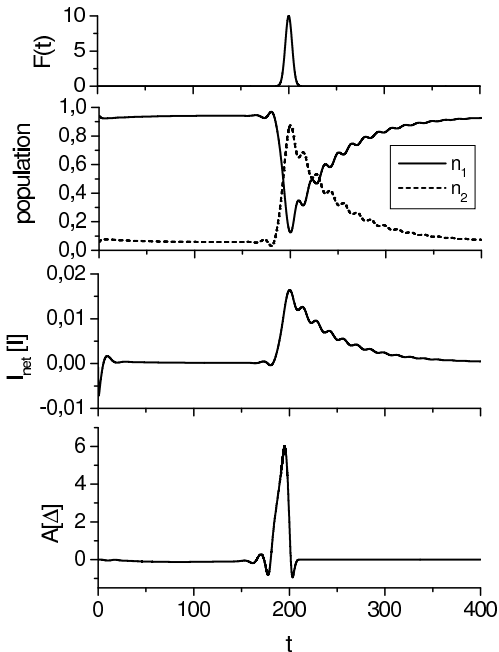
T. Brixner und G. Gerber,
ChemPhysChem **4**, 419 (2003).

R.S. Judson and H. Rabitz
Phys. Rev. Lett. **68**, 1500 (1992).

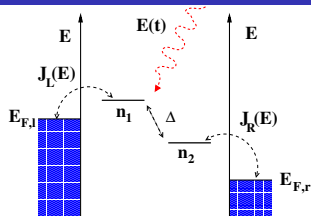
CDT: Optimal control of current with 2 sites



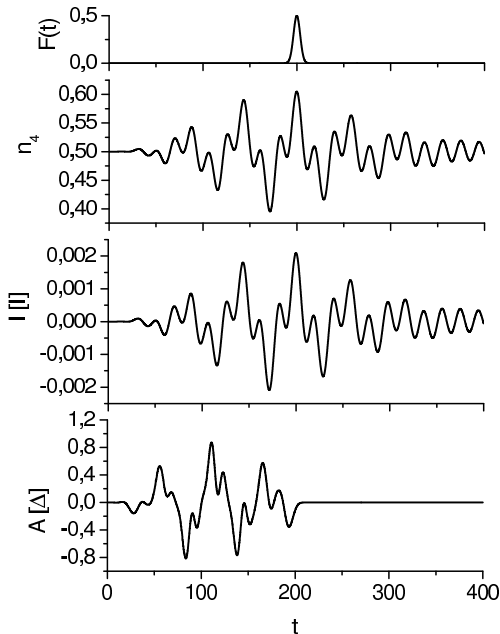
- maximal current at peak of goal function $F(t)$
- $E_1 - E_2 = -4\Delta$
- $E_{F,l} - E_{F,r} = 2\Delta$



CDT: Optimal control of current with 4 sites



- maximal current at peak of goal function $F(t)$
- $E_1 = E_2 = E_3 = E_4$
- $E_{F,l} = E_{F,r}$

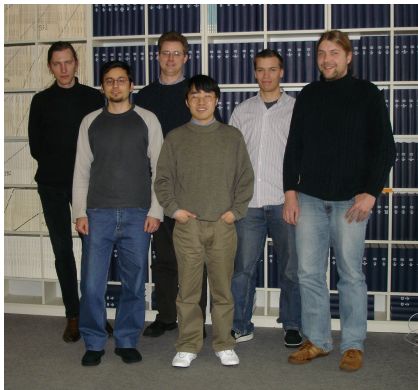


Summary

- density matrix approach to the electron transport through molecular wires
- numerical decomposition of the spectral density
- no approximation in the light-matter coupling beyond the semi-classical approximation
- calculation of population dynamics and current
- coherent destruction of tunneling also for rather short laser pulses
- first tests of optimal control of the current

Acknowledgments: The Group

- Quantum and Molecular Dynamics subgroup



- A. Novikov (now Rochester), S. Pezeshki, U. Kleinekathöfer (soon Bremen), G.-Q. Li, S. Welack (now Hong Kong), M. Schröder, J. Liebers (not shown)
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