

What is the order of $2d$ escape transition?

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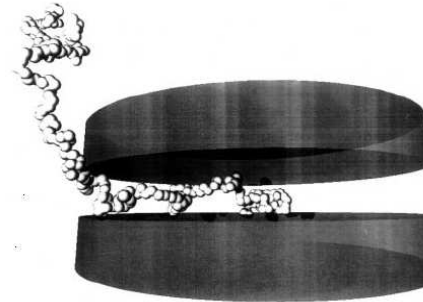
NTZ-Workshop CompPhys06 in Leipzig, Nov. 29 - Dec. 02, 2006

Motivations

- **3d** escape transition of a compressed end-grafted chain

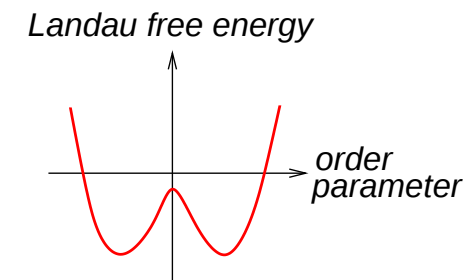
- Polymer chain: **first-order**
Milchev et. al., MC simulation

Europhys. Lett. 47, 675 (1999)



- Gaussian chain: **first-order**
Skvortsov et. al., Landau theory

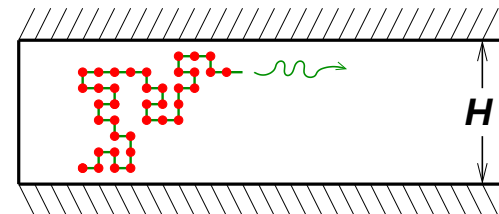
Europhys. Lett. 58, 292 (2002)



- **2d** long polymer chains confined in a strip

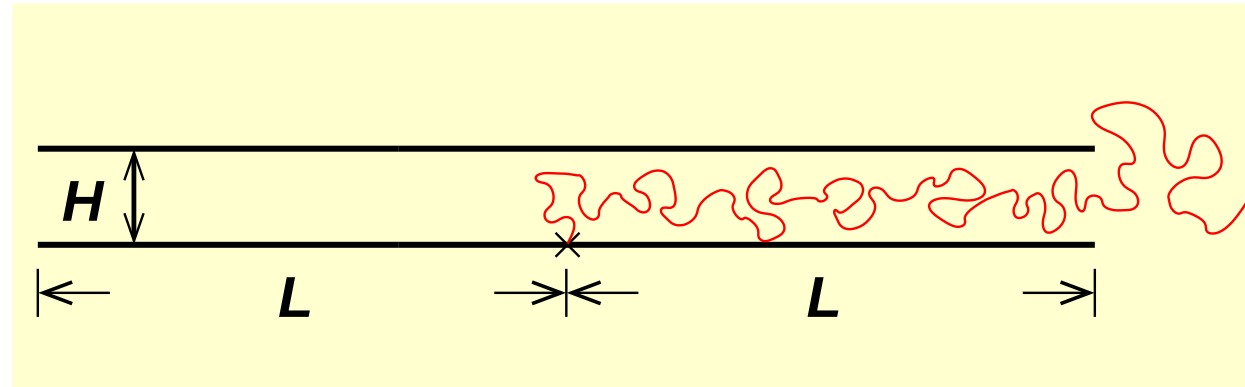
Hsu & Grassberger, PERM

Eur. Phys. J. B36, 209 (2003)



- What is the order of escape transition in **2d** ???

Outline



- Theoretical predictions

- Blob picture
- Landau theory

- Monte Carlo simulation:

Algorithm: Pruned-enriched Rosenbluth method **PERM**

- Conclusions

Blob picture

● Imprisoned chains

● End-to-end distance:

$$R = n_b(2r_b) = n_b H$$

within a blob, SAW

$$H = a n^\nu = 2r_b, \quad \nu = 3/4$$

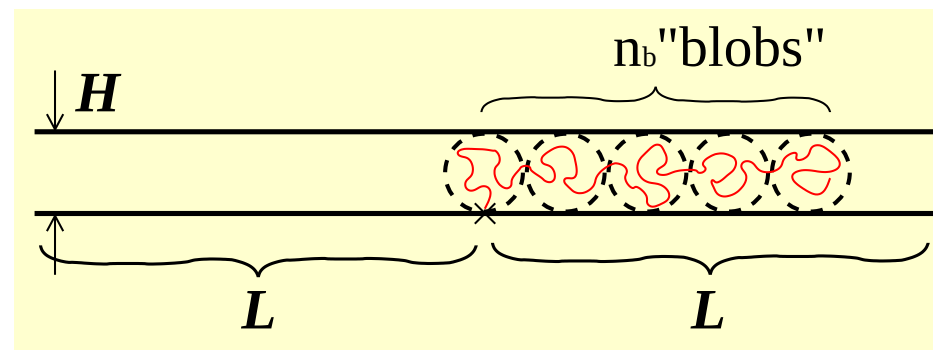
$$\Rightarrow n = (H/a)^{4/3} \text{ and } n = N/n_b$$

$$R/a = N(H/a)^{-1/3}$$

$$H \sim R_{\text{free}} = a N^\nu \Rightarrow R/a = N^\nu \text{ (mushroom)}$$

● Free energy: each blob contributes $k_B T$

$$F_{\text{impr}} = n_b [k_B T] = N(H/a)^{-4/3}$$



- Escape transition at $N = N^*$, $R = L$

- Escaped chains

- End-to-end distance:

$$R_{\text{esc}} = L + a(N - N^*)^{3/4}$$

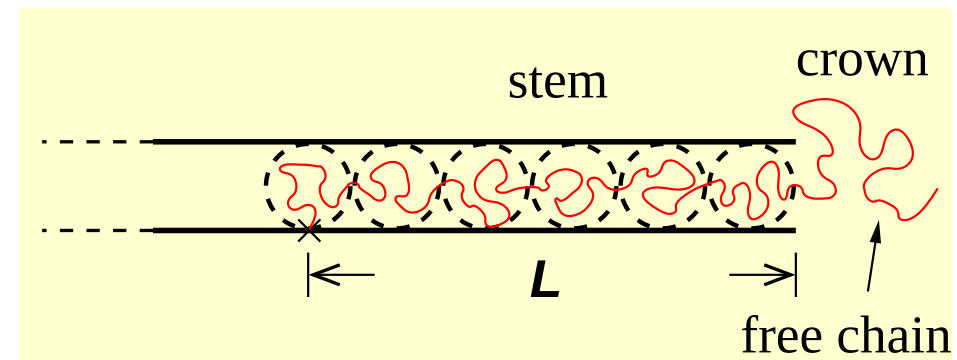
stem: N^* monomers

crown: $N - N^*$ monomers

- Free energy: $F_{\text{impr}} = N^*(H/a)^{-4/3}$

- Landau order parameter for the stretching, s

$$s = \begin{cases} R/Na & , \text{ imprisoned chain } (N < N^*) \\ L/N^*a & , \text{ escaped chain } (N > N^*) \end{cases}$$

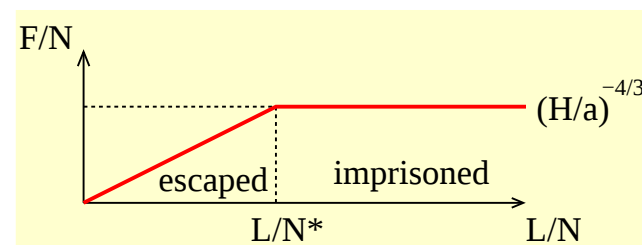


Skvortsov et. al., Europhys. Lett. 58, 292 (2002)

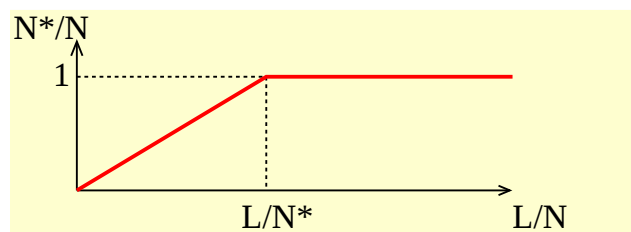
In the thermodynamic limit

$L \rightarrow \infty, N \rightarrow \infty,$
 L/N finite

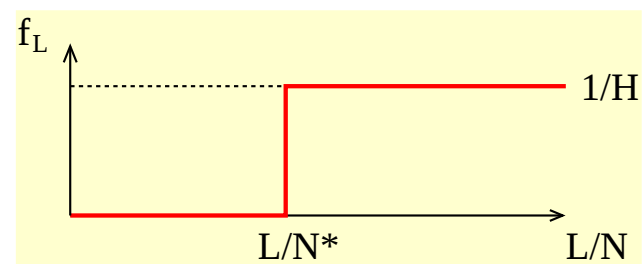
Free energy F



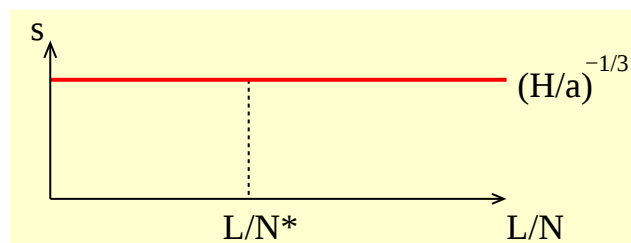
Imprisoned monomers N^*



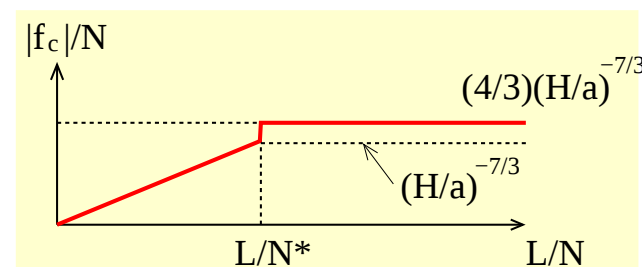
Lateral force $f_L = \left(\frac{\partial F}{\partial L}\right)_H$



Landau order parameter s



Compression force $f_c = \left(\frac{\partial F}{\partial H}\right)_L$



Landau theory

- Partition sum: $Z = \exp(-F) = \int ds \exp(-\Phi(s))$

$\Phi(s)$: free energy function, $\frac{d\Phi(s)}{ds} = 0 \Rightarrow s_{eq}$

- Imprisoned chains ($s = \frac{R}{Na} < \frac{L}{Na}$)

$$\Phi_{\text{impr}}(s) = N \left[A \left(\frac{a}{sH} \right)^2 + Bs^4 \right]$$

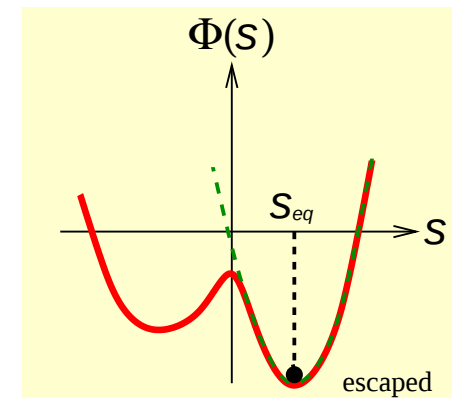
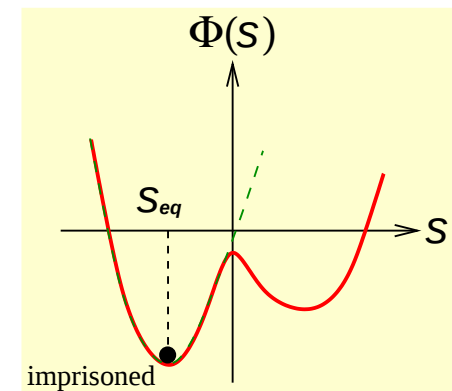
A and B : coefficients

- Escaped chains ($s = \frac{L}{N^*} > \frac{L}{Na}$)

$$\Phi_{\text{esc}}(s) = N^* \left[A \left(\frac{a}{sH} \right)^2 + Bs^4 \right]$$

$$= \frac{L}{a} \left[A \left(\frac{a}{H} \right)^2 s^{-3} + Bs^3 \right]$$

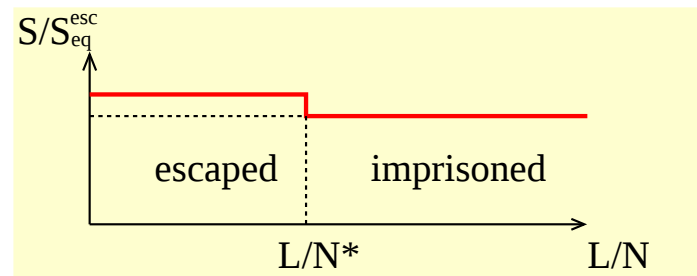
N^* : imprisoned monomers



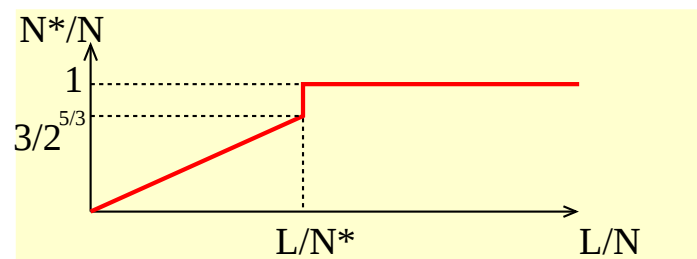
$$L \rightarrow \infty, N \rightarrow \infty, L/N \text{ finite}$$

- Landau order parameter $s = s_{eq}$:

$$s = \begin{cases} R/N = (A/2B)^{1/6} (a/H)^{1/3} & , \text{ imprisoned chain} \\ L/N^* = (A/B)^{1/6} (a/H)^{1/3} & , \text{ escaped chain} \end{cases}$$



- Imprisoned monomers N^* :



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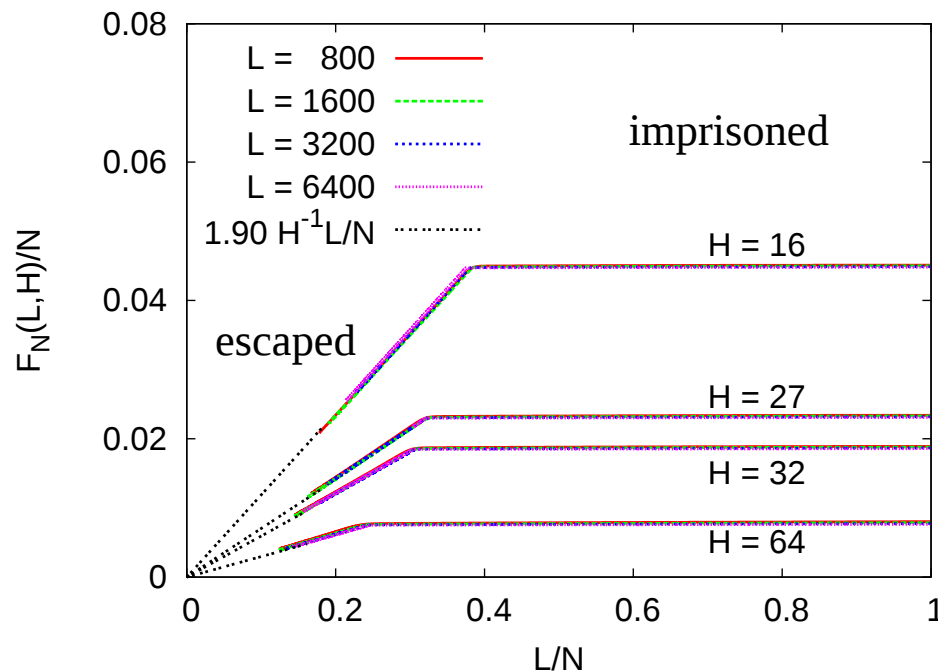
The diagram illustrates a quantum system in a one-dimensional potential well. Two horizontal black lines represent the boundaries of the well, with a vertical double-headed arrow on the right labeled H indicating the height. A green zigzag line with an arrow pointing right represents a wave function, with a horizontal double-headed arrow below it labeled L indicating the wavelength. A vertical line with a double-headed arrow indicates a position. A green path of dots starts from a point marked with a black 'x' on the bottom boundary, moves vertically up, then horizontally right, then vertically up again, then horizontally right, then vertically up, then horizontally right, and finally vertically up to the top boundary. The dots are red on the vertical segments and green on the horizontal segments.

- 2d escape transition – p.9

Free energy $F_N(L, H)$

● $F_N(L, H)/N$ vs. L/N

$$\frac{F_N(L, H)}{N} = \begin{cases} 1.93(2)H^{-4/3}, & \text{imprisoned chain} \\ 1.90(5)\frac{L}{N}H^{-1}, & \text{escaped chain} \end{cases}$$



Blob picture ✓

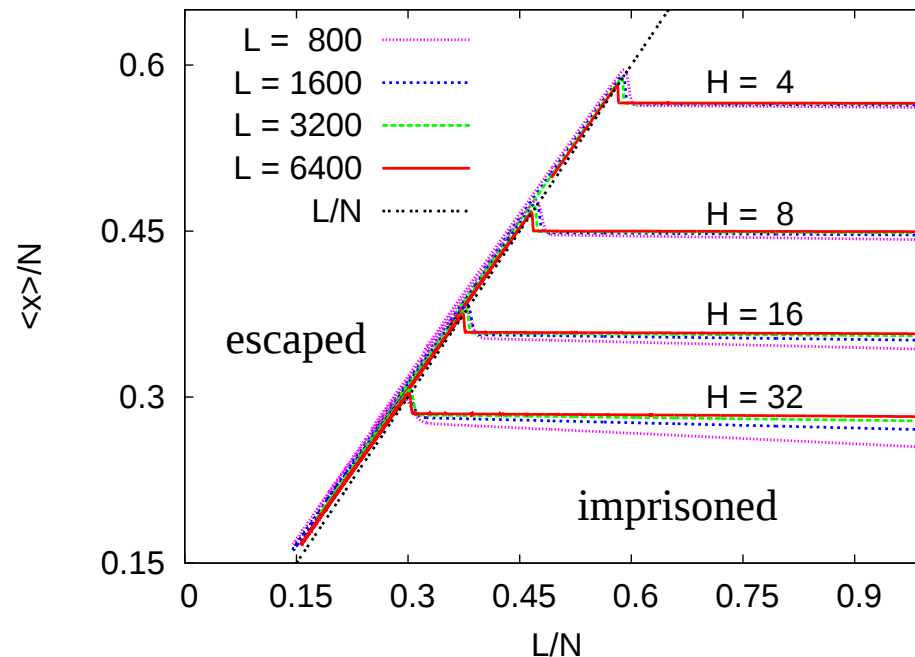
Transition point:

$$F_N^{\text{esc}}(L, H) = F_N^{\text{impr}}(L, H)$$

End-to-end distance $\langle x \rangle$

● $\langle x \rangle / N$ vs. L/N

$$\frac{\langle x \rangle}{N} = \begin{cases} 0.915 H^{-1/3} & , \text{ imprisoned chain} \\ \frac{L}{N} + \frac{a(N-N^*)}{N} H^{3/4} & , \text{ escaped chain} \end{cases}$$



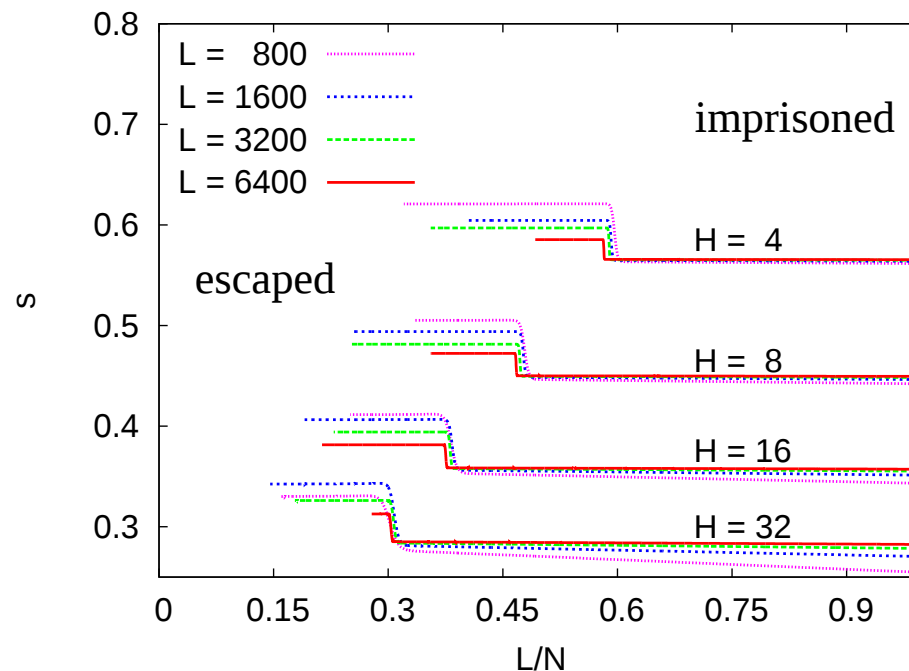
Blob picture ✕

N^* : imprisoned monomers

Landau order parameter s

● $\langle s \rangle$ vs. L/N

$$s = \begin{cases} R/N & , \text{imprisoned chain} \\ L/N^* & , \text{escaped chain} \end{cases}$$



Blob picture ✗

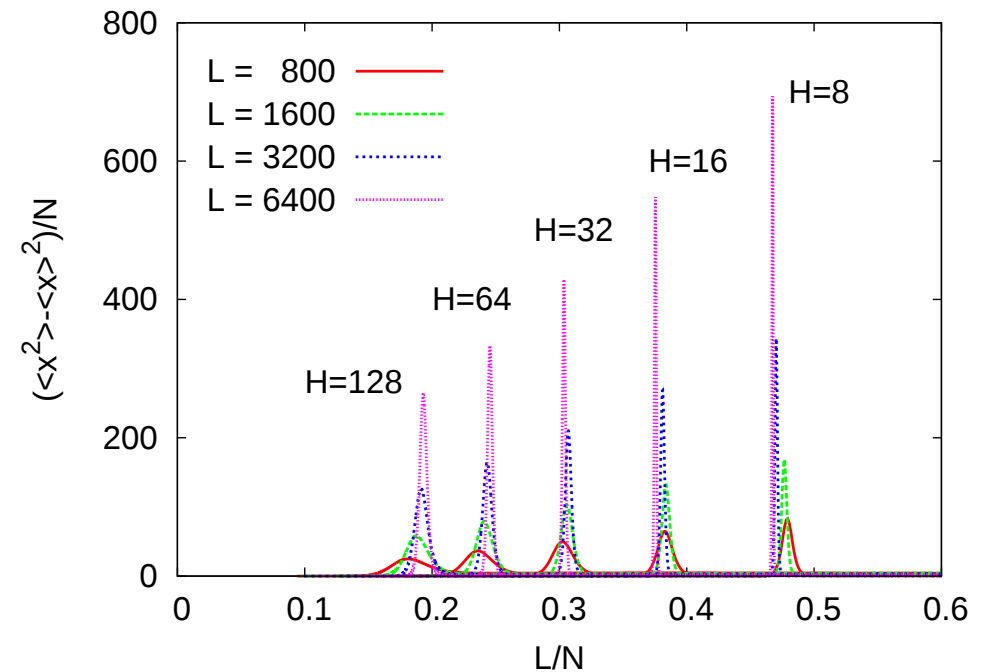
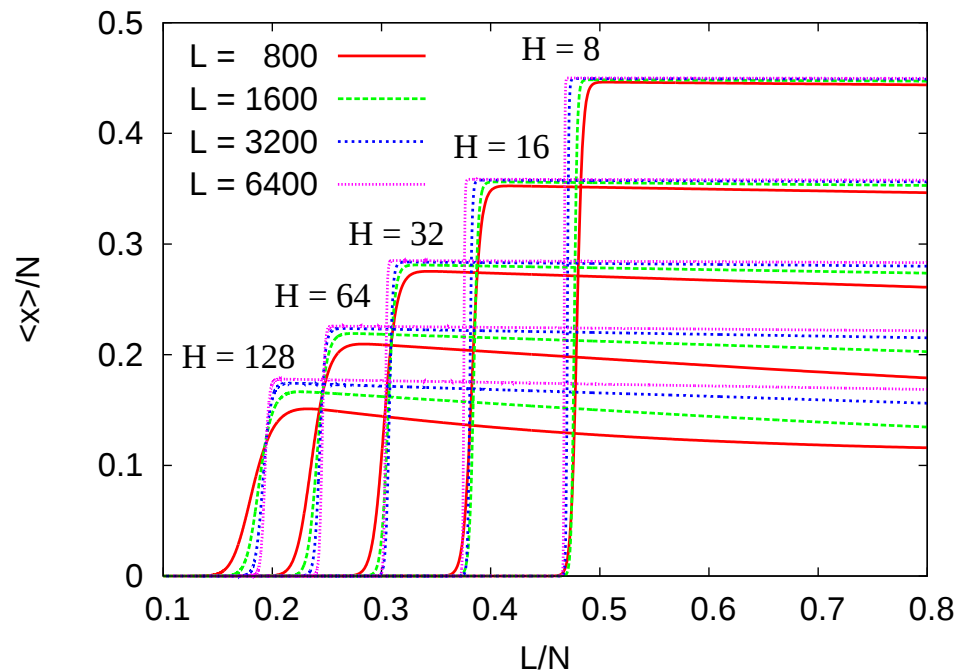
Landau theory ✓

Escape transition points

● End-to-end distance of imprisoned chains, x

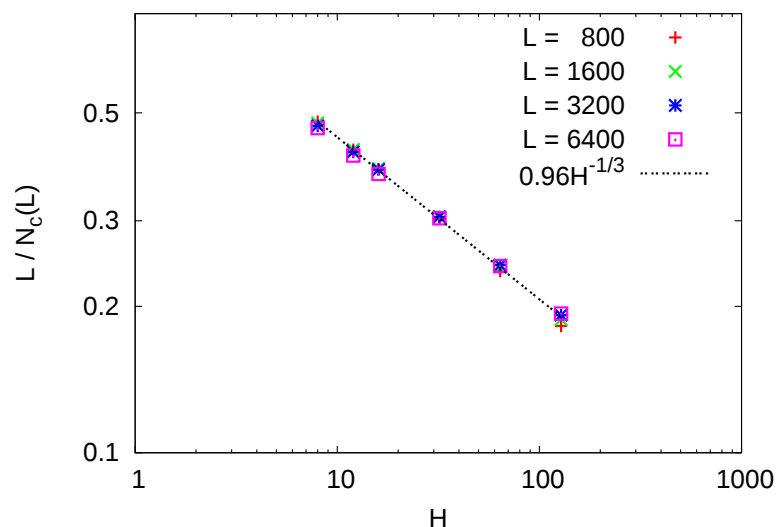
● $\langle x \rangle / N$ vs. L/N , $\langle x \rangle / N \sim H^{-1/3}$

● $(\langle x^2 \rangle - \langle x \rangle^2) / N$ vs. L/N



Escape transition points

- Transition points:
 - Peaks of variances of x



$$\Rightarrow L/N \sim 0.96(3)H^{-1/3}$$

- Free energy:

$$F_N^{\text{impr}}(L, H) = F_N^{\text{esc}}(L, H) \Rightarrow L/N \sim 1.02(4)H^{-1/3}$$

WHAT IS THE ORDER?

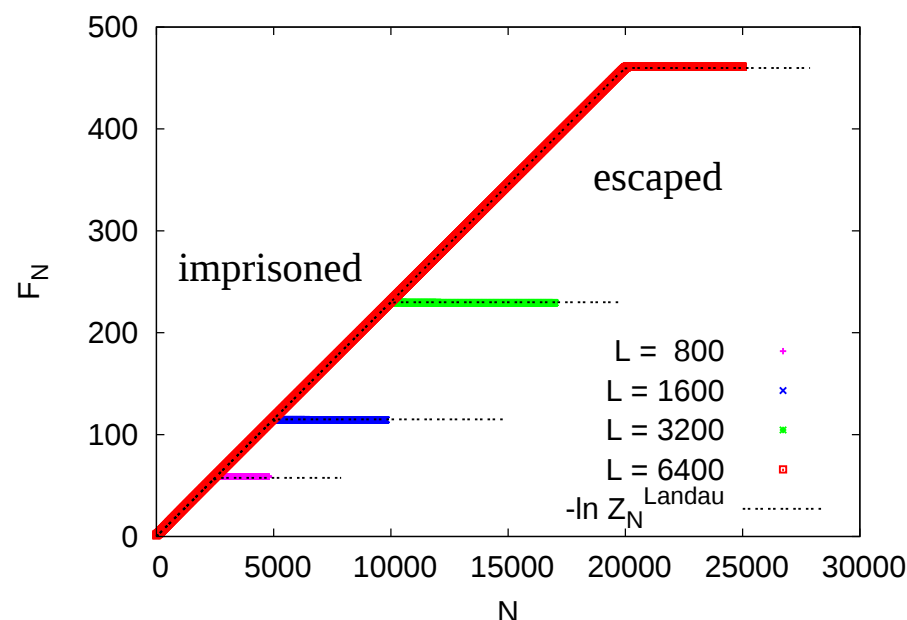
Landau free energy $\Phi(s)$

- Theoretical prediction of the partition sum: $Z = \int ds \exp(-\Phi(s))$

$$Z_N^{Landau}(L, H) = \exp(-\Phi_N^{\text{impr}}(s_{\text{eq}}^{\text{impr}})) + \exp(-\Phi_N^{\text{esc}}(s_{\text{eq}}^{\text{esc}}))$$

$$= \exp \left[-3 \left(\frac{A^2 B}{4} \right)^{1/3} N \left(\frac{a}{H} \right)^{4/3} \right] + \exp \left[-2(AB)^{1/2} \frac{L}{H} \right]$$

- Free energy $F_N(L, H) = -\ln\left(\frac{Z_N(L, H)}{Z_N}\right)$ vs. $N \Rightarrow A, B$



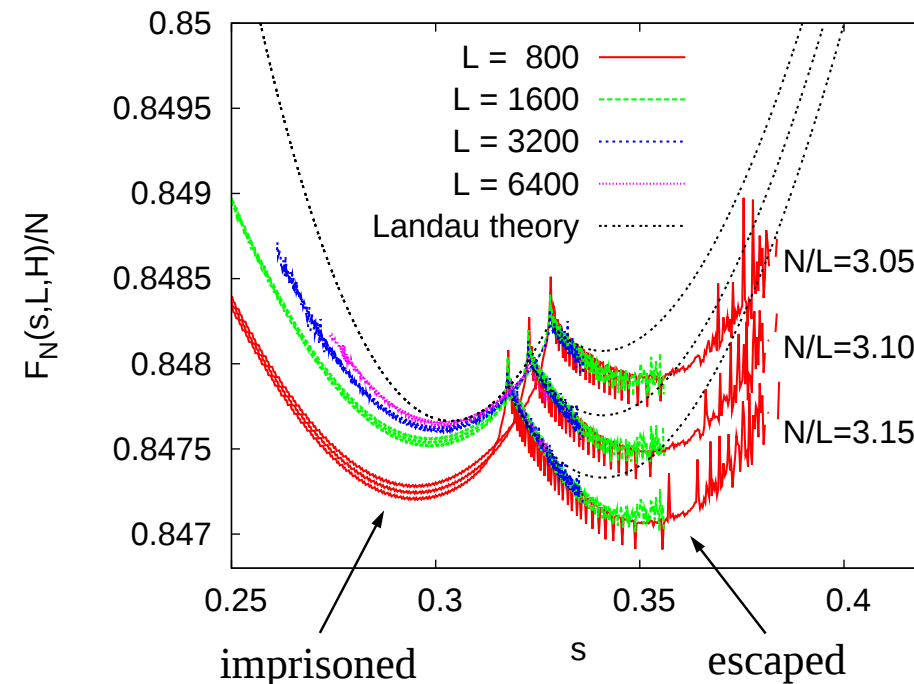
$$H = 27$$

● Theoretical prediction:

$$\Phi_N(s) = \begin{cases} N \left[A \left(\frac{a}{sH} \right)^2 + Bs^4 \right] & , \text{imprisoned chain} \\ \frac{L}{a} \left[A \left(\frac{a}{H} \right)^2 s^{-3} + Bs^3 \right] & , \text{escaped chain} \end{cases}$$

● MC simulations:

$F_N(s, L, H)/N$ vs. s



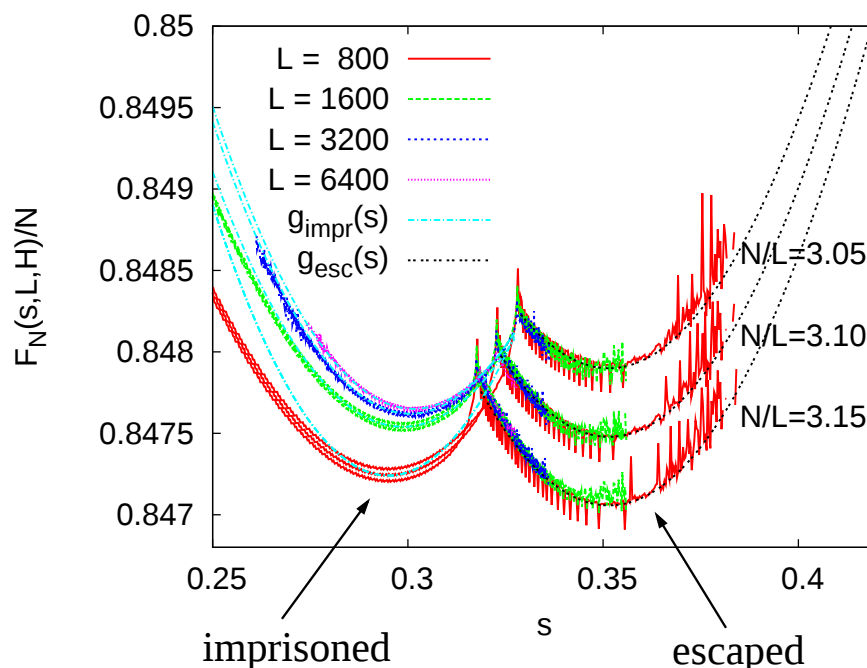
$F_N(s, L, H = 27)$ vs. s

- Min. of F_N is determined by fitting a parabola

- $g_{\text{impr}}(s) = a_{1,L}(s - s_{\text{eq}}^{\text{impr}}(L))^2 + c_{1,L}$

- $g_{\text{esc}}(s) = a_2(s - s_{\text{eq}}^{\text{esc}}(L))^2 + c_2 + d_2 \frac{N}{L}$

- Transition points: $g_{\text{impr}}(s = s_{\text{eq}}^{\text{impr}}(L)) = g_{\text{esc}}(s = s_{\text{eq}}^{\text{esc}}(L))$



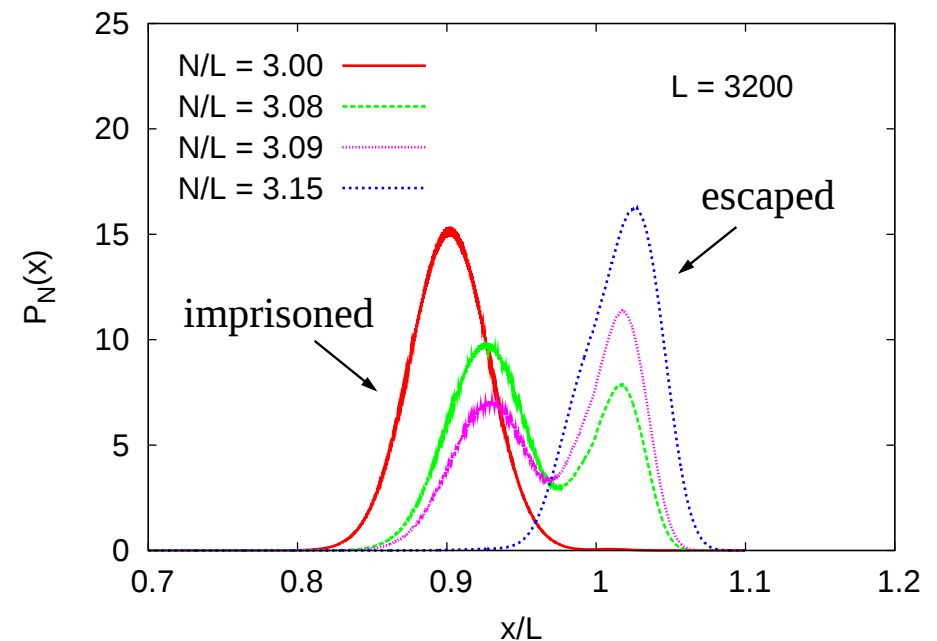
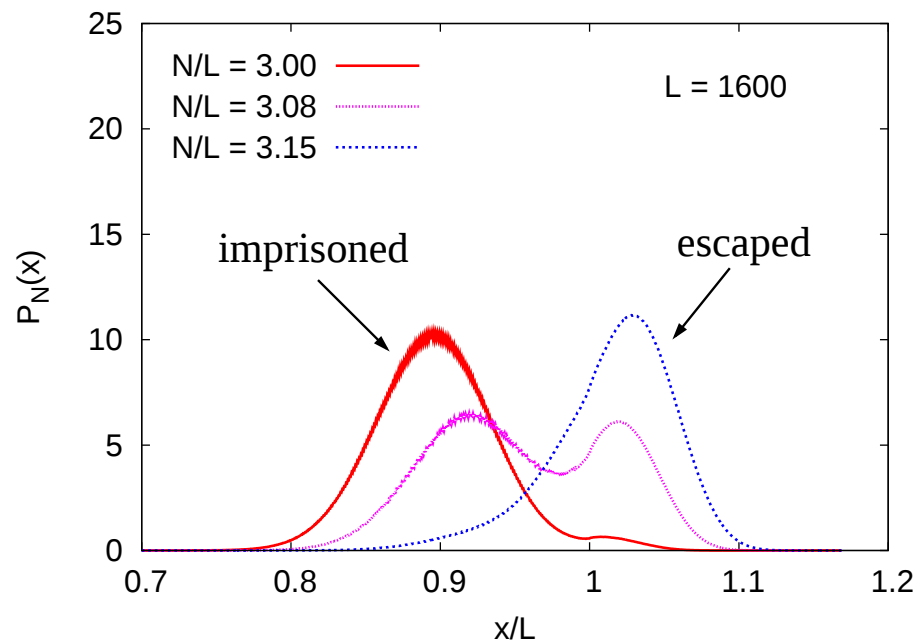
L	N^*/L
800	3.128
1600	3.093
3200	3.083
6400	3.080

$$N/L \sim 1.03(5)H^{1/3} \sim 3.09(15)$$

Histogram of x , $P_N(x)$

● $P_N(x) = \sum_{walks} \delta_{x',x}$

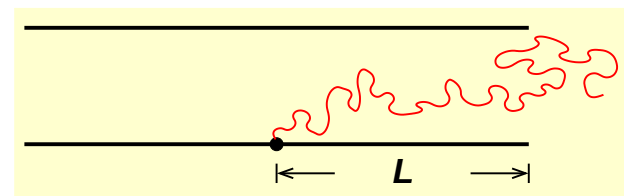
● $P_N(x)$ vs. x/L $(\sum_x P_N(x) = L)$



Conclusions

- What is the order of **2d** escape transition?
first-order phase transition

- Landau theory
- MC simulations (PERM)



- Transition points: $L/N \sim 1.02(4)H^{-1/3}$
- Does the Landau theory also explain **3d** escape transition? **No!**
 - Poor definition of the order parameter

$$s = \begin{cases} R/Na & , N < N^* \\ L/N^*a & , N > N^* \end{cases}$$

R : end-to-end distance

